行星运动的求解

先证一个引理:

圆锥曲线在顶点处的曲率半径等于该顶点与(最近的)焦点的距离. 证明:

在极坐标系下,设圆锥曲线的方程为 $r = \frac{r_0}{1 + \varepsilon \cos \theta}$. 顶点坐标为 $\left(\frac{r_0}{1+\varepsilon},0\right)$.

$$\kappa = \lim_{\Delta s \to 0} \frac{\Delta \phi}{\Delta s} = \frac{\mathrm{d}\theta}{\sqrt{r^2 + \left(\frac{\mathrm{d}r}{\mathrm{d}\theta}\right)^2} \mathrm{d}\theta} = \frac{1}{\sqrt{r^2 + \left(\frac{\mathrm{d}r}{\mathrm{d}\theta}\right)^2}}$$

设恒星的质量为 M,万有引力常数为 G. 轨道的极坐标方程可表示为 $r=\frac{r_0}{1+\varepsilon\cos(\theta-\theta_0)}$,其中 ε 为轨道离心率.

令 $\cos(\theta - \theta_0) = 1$ 得,近日点与恒星的距离 $r_m = \frac{r_0}{1+\epsilon}$.

近日点速率
$$v_m$$
 满足 $\frac{v_m^2}{r_m} = \frac{GM}{r_m^2}$.
$$\therefore v_m = \sqrt{\frac{GM}{r_m}} = \sqrt{\frac{(1+\varepsilon)GM}{r_0}}$$

设 $h = v_m r_m = \sqrt{\frac{GMr_0}{1+\varepsilon}}$,由角动量守恒,行星在轨道上运动时 h = vr 不变.

单位时间行星与恒星连线扫过的面积

$$dS = \frac{1}{2}hdt = \frac{1}{2}\sqrt{\frac{GMr_0}{1+\varepsilon}}dt = \frac{1}{2}r^2d\theta$$

若以行星位于近日点时为计时起点,则行星运动到 (r,θ) 处的时间

$$t = \sqrt{\frac{1+\varepsilon}{GMr_0}} \int_{\theta_0}^{\theta} r^2 d\theta = r_0^2 \sqrt{\frac{1+\varepsilon}{GMr_0}} \int_0^{\theta-\theta_0} \frac{dx}{(1+\varepsilon\cos x)^2}$$

经计算:

$$\int_{0}^{\theta} \frac{\mathrm{d}x}{(1+\varepsilon\cos x)^{2}} = \begin{cases} -\frac{\varepsilon\sin\theta}{(1-\varepsilon^{2})(\varepsilon\cos\theta+1)} + \frac{2\arctan\left(\sqrt{\frac{1-\varepsilon}{1+\varepsilon}\tan\frac{\theta}{2}}\right)}{(1-\varepsilon^{2})^{\frac{3}{2}}}, \ 0<\varepsilon<1\\ \frac{1}{2}\tan\frac{\theta}{2} + \frac{1}{6}\tan^{3}\frac{\theta}{2}, \ \varepsilon=1\\ \frac{\varepsilon\sin\theta}{(\varepsilon^{2}-1)(\varepsilon\cos\theta+1)} - \frac{2\arctan\left(\sqrt{\frac{\varepsilon-1}{1+\varepsilon}\tan\frac{\theta}{2}}\right)}{(\varepsilon^{2}-1)^{\frac{3}{2}}}, \ \varepsilon>1 \end{cases}$$

于是已知 θ 、r、t 中的任意一个, 可求另外两者.