Bessel function

From Wikipedia, the free encyclopedia

Bessel functions, first defined by the mathematician Daniel Bernoulli and then generalized by Friedrich Bessel, are the canonical solutions $y(x)$ of **Bessel's differential equation**

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - \alpha^2)y = 0$$

for an arbitrary complex number α (the order of the Bessel function). Although α and −α produce the same differential equation for real α, it is conventional to define different Bessel functions for these two values in such a way that the Bessel functions are mostly smooth functions of α.

The most important cases are for α an integer or half-integer. Bessel functions for integer α are also known as **cylinder functions** or the **cylindrical harmonics** because they appear in the solution to Laplace's equation in cylindrical coordinates. **Spherical Bessel functions** with half-integer α are obtained when the Helmholtz equation is solved in spherical coordinates.

Contents

Applications of Bessel functions

Bessel's equation arises when finding separable solutions to Laplace's equation and the Helmholtz equation in cylindrical or spherical coordinates. Bessel functions are therefore especially important for many problems of wave propagation and static potentials. In solving problems in cylindrical coordinate systems, one obtains Bessel functions of integer order ($\alpha = n$); in spherical problems, one obtains half-integer orders ($\alpha = n+1/2$). For example:

- Electromagnetic waves in a cylindrical waveguide
- Pressure amplitudes of inviscid rotational flows
- Heat conduction in a cylindrical object
- Modes of vibration of a thin circular (or annular) acoustic membrane (such as a drum or other membranophone)
- Diffusion problems on a lattice
- Solutions to the radial Schrödinger equation (in spherical and cylindrical coordinates) for a free particle
- Solving for patterns of acoustical radiation
- Frequency-dependent friction in circular pipelines
- Dynamics of floating bodies
- Angular resolution

Bessel functions also appear in other problems, such as signal processing (e.g., see FM synthesis, Kaiser window, or Bessel filter).

Definitions

Because this is a second-order differential equation, there must be two linearly independent solutions. Depending upon the circumstances, however, various formulations of these solutions are convenient. Different variations are summarized in the table below, and described in the following sections.

Type	First kind	Second kind
Bessel functions	J_α	Y_α
modified Bessel functions	I_α	K_α
Hankel functions	$H_\alpha^{(1)} = J_\alpha + iY_\alpha$	$H_\alpha^{(2)} = J_\alpha - iY_\alpha$
Spherical Bessel functions	j_n	y_n
Spherical Hankel functions	$h_n^{(1)} = j_n + iy_n$	$h_n^{(2)} = j_n - iy_n$

Bessel functions of the first kind: J_α

Bessel functions of the first kind, denoted as $J_\alpha(x)$, are solutions of Bessel's differential equation that are finite at the origin ($x = 0$) for integer or positive α, and diverge as x approaches zero for negative non-integer α. It is possible to define the function by its series expansion around $x = 0$, which can be found by applying the Frobenius method to Bessel's equation:[1]

$$J_\alpha(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m!\,\Gamma(m+\alpha+1)} \left(\frac{x}{2}\right)^{2m+\alpha}$$

where $\Gamma(z)$ is the gamma function, a shifted generalization of the factorial function to non-integer values. The Bessel function of the first kind is an entire function if α is an integer, otherwise it is a multivalued function with singularity at zero. The graphs of Bessel functions look roughly like oscillating sine or cosine functions that decay proportionally to $1/\sqrt{x}$ (see also their asymptotic forms below), although their roots are not generally periodic, except asymptotically for large x. (The series indicates that $-J_1(x)$ is the derivative of $J_0(x)$, much like $-\sin(x)$ is the derivative of $\cos(x)$; more generally, the derivative of $J_n(x)$ can be expressed in terms of $J_{n\pm1}(x)$ by the identities below.)

For non-integer α, the functions $J_\alpha(x)$ and $J_{-\alpha}(x)$ are linearly independent, and are therefore the two solutions of the differential equation. On the other hand, for integer order α, the following relationship is valid (note that the Gamma function has simple poles at each of the non-positive integers):[2]

$$J_{-n}(x) = (-1)^n J_n(x).$$

This means that the two solutions are no longer linearly independent. In this case, the second linearly independent solution is then found to be the Bessel function of the second kind, as discussed below.

Bessel's integrals

Another definition of the Bessel function, for integer values of n, is possible using an integral representation:[3]

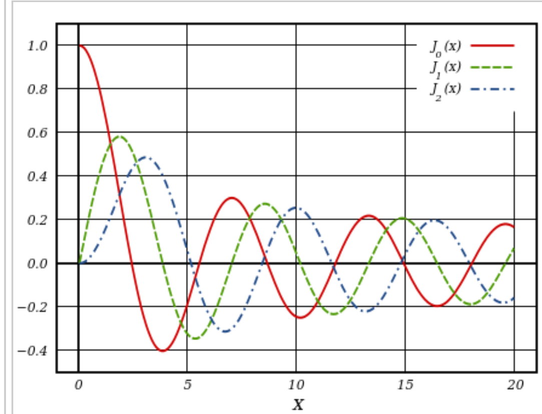
$$J_n(x) = \frac{1}{\pi} \int_0^\pi \cos(n\tau - x\sin(\tau))\,d\tau.$$

Another integral representation is:[3]

$$J_n(x) = \frac{1}{2\pi} \int_{-\pi}^\pi e^{i(n\tau - x\sin(\tau))}\,d\tau.$$

This was the approach that Bessel used, and from this definition he derived several properties of the function. The definition may be extended to non-integer orders by (for Re(x) > 0), one of Schläfli's integrals:[3]

$$J_\alpha(x) = \frac{1}{\pi} \int_0^\pi \cos(\alpha\tau - x\sin\tau)\,d\tau - \frac{\sin(\alpha\pi)}{\pi} \int_0^\infty e^{-x\sinh(t) - \alpha t}\,dt.$$ [4][5][6][7]



Plot of Bessel function of the first kind, $J_\alpha(x)$, for integer orders α = 0, 1, 2

Relation to hypergeometric series

The Bessel functions can be expressed in terms of the generalized hypergeometric series as[8]

$$J_\alpha(x) = \frac{(\frac{x}{2})^\alpha}{\Gamma(\alpha+1)}\ {}_0F_1(\alpha+1; -\tfrac{x^2}{4}).$$

This expression is related to the development of Bessel functions in terms of the Bessel–Clifford function.

Relation to Laguerre polynomials

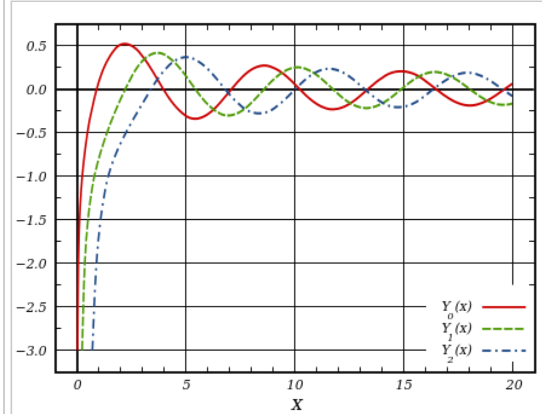
In terms of the Laguerre polynomials L_k and arbitrarily chosen parameter t, the Bessel function can be expressed as[9]

$$\frac{J_\alpha(x)}{\left(\frac{x}{2}\right)^\alpha} = \frac{e^{-t}}{\Gamma(\alpha+1)} \sum_{k=0}^\infty \frac{L_k^{(\alpha)}\left(\frac{x^2}{4t}\right)}{\binom{k+\alpha}{k}} \frac{t^k}{k!}.$$

Bessel functions of the second kind: Y_α

The Bessel functions of the second kind, denoted by $Y_\alpha(x)$, occasionally denoted instead by $N_\alpha(x)$, are solutions of the Bessel differential equation that have a singularity at the origin ($x = 0$) and are multivalued. These are sometimes called **Weber functions** as they were introduced by H. Weber (1873), and also **Neumann functions** after Carl Neumann.[10]

For non-integer α, $Y_\alpha(x)$ is related to $J_\alpha(x)$ by:



Plot of Bessel function of the second kind, $Y_\alpha(x)$, for integer orders α = 0, 1, 2.

$$Y_\alpha(x) = \frac{J_\alpha(x)\cos(\alpha\pi) - J_{-\alpha}(x)}{\sin(\alpha\pi)}.$$

In the case of integer order n, the function is defined by taking the limit as a non-integer α tends to n:

$$Y_n(x) = \lim_{\alpha\to n} Y_\alpha(x).$$

There is also a corresponding integral formula (for $\mathrm{Re}(x) > 0$),[11]

$$Y_n(x) = \frac{1}{\pi}\int_0^\pi \sin(x\sin\theta - n\theta)\,d\theta - \frac{1}{\pi}\int_0^\infty \left[e^{nt} + (-1)^n e^{-nt}\right] e^{-x\sinh t}\,dt.$$

$Y_\alpha(x)$ is necessary as the second linearly independent solution of the Bessel's equation when α is an integer. But $Y_\alpha(x)$ has more meaning than that. It can be considered as a 'natural' partner of $J_\alpha(x)$. See also the subsection on Hankel functions below.

When α is an integer, moreover, as was similarly the case for the functions of the first kind, the following relationship is valid:

$$Y_{-n}(x) = (-1)^n Y_n(x).$$

Both $J_\alpha(x)$ and $Y_\alpha(x)$ are holomorphic functions of x on the complex plane cut along the negative real axis. When α is an integer, the Bessel functions J are entire functions of x. If x is held fixed at a non-zero value, then the Bessel functions are entire functions of α.

The Bessel functions of the second kind when α is an integer is an example of the second kind of solution in Fuchs's theorem.

Hankel functions: $H_\alpha^{(1)}$, $H_\alpha^{(2)}$

Another important formulation of the two linearly independent solutions to Bessel's equation are the **Hankel functions of the first and second kind**, $H_\alpha^{(1)}(x)$ and $H_\alpha^{(2)}(x)$, defined by:[12]

$$H_\alpha^{(1)}(x) = J_\alpha(x) + iY_\alpha(x)$$
$$H_\alpha^{(2)}(x) = J_\alpha(x) - iY_\alpha(x)$$

where i is the imaginary unit. These linear combinations are also known as **Bessel functions of the third kind**; they are two linearly independent solutions of Bessel's differential equation. They are named after Hermann Hankel.

The importance of Hankel functions of the first and second kind lies more in theoretical development rather than in application. These forms of linear combination satisfy numerous simple-looking properties, like asymptotic formulae or integral representations. Here, 'simple' means an appearance of the factor of the form $e^{if(x)}$. The Bessel function of the second kind then can be thought to naturally appear as the imaginary part of the Hankel functions.

The Hankel functions are used to express outward- and inward-propagating cylindrical wave solutions of the cylindrical wave equation, respectively (or vice versa, depending on the sign convention for the frequency).

Using the previous relationships they can be expressed as:

$$H_\alpha^{(1)}(x) = \frac{J_{-\alpha}(x) - e^{-\alpha\pi i} J_\alpha(x)}{i \sin(\alpha\pi)}$$
$$H_\alpha^{(2)}(x) = \frac{J_{-\alpha}(x) - e^{\alpha\pi i} J_\alpha(x)}{-i \sin(\alpha\pi)}.$$

If α is an integer, the limit has to be calculated. The following relationships are valid, whether α is an integer or not:[13]

$$H_{-\alpha}^{(1)}(x) = e^{\alpha\pi i} H_\alpha^{(1)}(x)$$
$$H_{-\alpha}^{(2)}(x) = e^{-\alpha\pi i} H_\alpha^{(2)}(x).$$

In particular, if $\alpha = m + 1/2$ with m a nonnegative integer, the above relations imply directly that

$$J_{-(m+\frac{1}{2})}(x) = (-1)^{m+1} Y_{m+\frac{1}{2}}(x)$$
$$Y_{-(m+\frac{1}{2})}(x) = (-1)^m J_{m+\frac{1}{2}}(x).$$

These are useful in developing the spherical Bessel functions (below).

The Hankel functions admit the following integral representations for $\mathrm{Re}(x) > 0$:[14]

$$H_\alpha^{(1)}(x) = \frac{1}{\pi i} \int_{-\infty}^{+\infty+i\pi} e^{x \sinh t - \alpha t}\, dt,$$
$$H_\alpha^{(2)}(x) = -\frac{1}{\pi i} \int_{-\infty}^{+\infty-i\pi} e^{x \sinh t - \alpha t}\, dt,$$

where the integration limits indicate integration along a contour that can be chosen as follows: from −∞ to 0 along the negative real axis, from 0 to ±iπ along the imaginary axis, and from ±iπ to +∞±iπ along a contour parallel to the real axis.[15]

Modified Bessel functions: I_α, K_α

The Bessel functions are valid even for complex arguments x, and an important special case is that of a purely imaginary argument. In this case, the solutions to the Bessel equation are called the **modified Bessel functions** (or occasionally the **hyperbolic Bessel functions**) **of the first and second kind**, and are defined by:[16]

$$I_\alpha(x) = i^{-\alpha} J_\alpha(ix) = \sum_{m=0}^{\infty} \frac{1}{m!\,\Gamma(m+\alpha+1)} \left(\frac{x}{2}\right)^{2m+\alpha}$$
$$K_\alpha(x) = \frac{\pi}{2} \frac{I_{-\alpha}(x) - I_\alpha(x)}{\sin(\alpha\pi)},$$

when α is not an integer; when α is an integer, then the limit is used. These are chosen to be real-valued for real and positive arguments x. The series expansion for $I_\alpha(x)$ is thus similar to that for $J_\alpha(x)$, but without the alternating $(-1)^m$ factor.

If $-\pi < \arg(x) \leq \pi/2$, $K_\alpha(x)$ can be expressed as a Hankel function of the first kind:

$$K_\alpha(x) = \frac{\pi}{2} i^{\alpha+1} H_\alpha^{(1)}(ix),$$

and if $\pi/2 < \arg(x) \leq \pi$, it can be expressed as a Hankel function of the second kind:

$$K_\alpha(x) = \frac{\pi}{2} (-i)^{\alpha+1} H_\alpha^{(2)}(-ix).$$

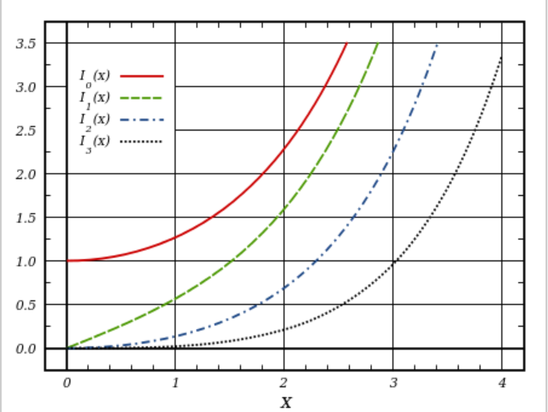
We can express the first and second Bessel functions in terms of the modified Bessel functions (these are valid if $-\pi < \arg(z) \leq \pi/2$):

$$J_\alpha(iz) = e^{\frac{\alpha i \pi}{2}} I_\alpha(z)$$
$$Y_\alpha(iz) = e^{\frac{(\alpha+1) i \pi}{2}} I_\alpha(z) - \frac{2}{\pi} e^{-\frac{\alpha i \pi}{2}} K_\alpha(z).$$

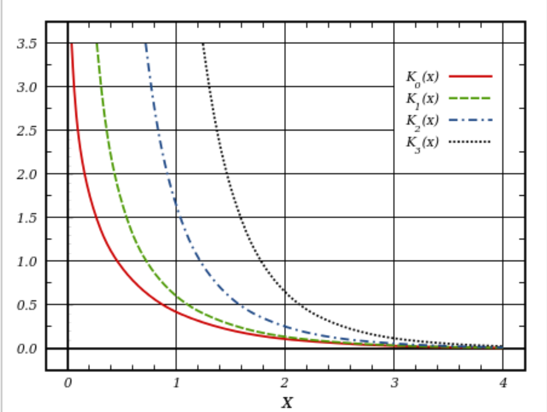
$I_\alpha(x)$ and $K_\alpha(x)$ are the two linearly independent solutions to the **modified Bessel's equation**:[17]

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} - (x^2 + \alpha^2) y = 0.$$

Unlike the ordinary Bessel functions, which are oscillating as functions of a real argument, I_α and K_α are exponentially growing and decaying functions, respectively. Like the ordinary Bessel function J_α, the function I_α goes to zero at $x = 0$ for $\alpha > 0$ and is finite at $x = 0$ for $\alpha = 0$. Analogously, K_α diverges at $x = 0$ with the singularity being of logarithmic type.[18]



Modified Bessel functions of 1st kind, $I_\alpha(x)$, for α = 0, 1, 2, 3



Modified Bessel functions of 2nd kind, $K_\alpha(x)$, for α = 0, 1, 2, 3

Two integral formulas for the modified Bessel functions are (for Re(x) > 0):[19]

$$I_\alpha(x) = \frac{1}{\pi}\int_0^\pi \exp(x\cos(\theta))\cos(\alpha\theta)\,d\theta - \frac{\sin(\alpha\pi)}{\pi}\int_0^\infty \exp(-x\cosh t - \alpha t)\,dt,$$
$$K_\alpha(x) = \int_0^\infty \exp(-x\cosh t)\cosh(\alpha t)\,dt.$$

Modified Bessel functions $K_{1/3}$ and $K_{2/3}$ can be represented in terms of rapidly converged integrals[20]

$$K_{\frac{1}{3}}(\xi) = \sqrt{3}\,\int_0^\infty \exp\left[-\xi\left(1+\frac{4x^2}{3}\right)\sqrt{1+\frac{x^2}{3}}\right]\,dx$$
$$K_{\frac{2}{3}}(\xi) = \frac{1}{\sqrt{3}}\,\int_0^\infty \frac{3+2x^2}{\sqrt{1+\frac{x^2}{3}}}\exp\left[-\xi\left(1+\frac{4x^2}{3}\right)\sqrt{1+\frac{x^2}{3}}\right]\,dx.$$

The **modified Bessel function of the second kind** has also been called by the now-rare names:

- **Basset function** after Alfred Barnard Basset
- **Modified Bessel function of the third kind**
- **Modified Hankel function**[21]
- **Macdonald function** after Hector Munro Macdonald

Spherical Bessel functions: j_n, y_n

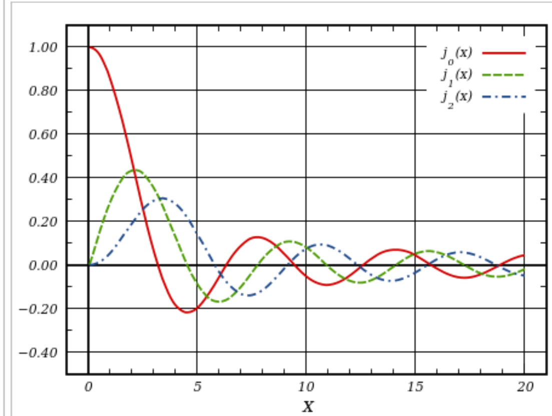
When solving the Helmholtz equation in spherical coordinates by separation of variables, the radial equation has the form:

$$x^2\frac{d^2y}{dx^2}+2x\frac{dy}{dx}+[x^2-n(n+1)]y=0.$$

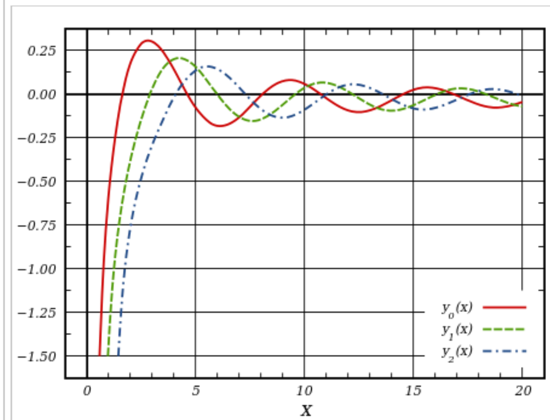
The two linearly independent solutions to this equation are called the **spherical Bessel functions** j_n and y_n, and are related to the ordinary Bessel functions J_n and Y_n by:[22]

$$\begin{aligned} j_n(x)&=\sqrt{\frac{\pi}{2x}}J_{n+\frac{1}{2}}(x),\\ y_n(x)&=\sqrt{\frac{\pi}{2x}}Y_{n+\frac{1}{2}}(x)=(-1)^{n+1}\sqrt{\frac{\pi}{2x}}J_{-n-\frac{1}{2}}(x). \end{aligned}$$

y_n is also denoted n_n or η_n; some authors call these functions the **spherical Neumann functions**.



Spherical Bessel functions of 1st kind, $j_n(x)$, for $n = 0, 1, 2$



Spherical Bessel functions of 2nd kind, $y_n(x)$, for $n = 0, 1, 2$

The spherical Bessel functions can also be written as (**Rayleigh's formulas**):[23]

$$\begin{aligned} j_n(x)&=(-x)^n\left(\frac{1}{x}\frac{d}{dx}\right)^n\frac{\sin(x)}{x},\\ y_n(x)&=-(-x)^n\left(\frac{1}{x}\frac{d}{dx}\right)^n\frac{\cos(x)}{x}. \end{aligned}$$

The first spherical Bessel function $j_0(x)$ is also known as the (unnormalized) sinc function. The first few spherical Bessel functions are:

$$\begin{aligned} j_0(x)&=\frac{\sin(x)}{x}\\ j_1(x)&=\frac{\sin(x)}{x^2}-\frac{\cos(x)}{x}\\ j_2(x)&=\left(\frac{3}{x^2}-1\right)\frac{\sin(x)}{x}-\frac{3\cos(x)}{x^2} \quad [24]\\ j_3(x)&=\left(\frac{15}{x^3}-\frac{6}{x}\right)\frac{\sin(x)}{x}-\left(\frac{15}{x^2}-1\right)\frac{\cos(x)}{x}, \end{aligned}$$

and

$$y_0(x) = -j_{-1}(x) = -\frac{\cos(x)}{x}$$
$$y_1(x) = j_{-2}(x) = -\frac{\cos(x)}{x^2} - \frac{\sin(x)}{x}$$
$$y_2(x) = -j_{-3}(x) = \left(-\frac{3}{x^2}+1\right)\frac{\cos(x)}{x} - \frac{3\sin(x)}{x^2}$$ [25]
$$y_3(x) = j_{-4}(x) = \left(-\frac{15}{x^3}+\frac{6}{x}\right)\frac{\cos(x)}{x} - \left(\frac{15}{x^2}-1\right)\frac{\sin(x)}{x}.$$

Generating function

The spherical Bessel functions have the generating functions [26]

$$\frac{1}{z}\cos\left(\sqrt{z^2-2zt}\right) = \sum_{n=0}^{\infty}\frac{t^n}{n!}j_{n-1}(z),$$
$$\frac{1}{z}\sin\left(\sqrt{z^2+2zt}\right) = \sum_{n=0}^{\infty}\frac{(-t)^n}{n!}y_{n-1}(z).$$

Differential relations

In the following f_n is any of $j_n, y_n, h_n^{(1)}, h_n^{(2)}$ for $n = 0, \pm1, \pm2, \ldots$ [27]

$$\left(\frac{1}{z}\frac{d}{dz}\right)^m \left(z^{n+1} f_n(z)\right) = z^{n-m+1} f_{n-m}(z),$$
$$\left(\frac{1}{z}\frac{d}{dz}\right)^m \left(z^{-n} f_n(z)\right) = (-1)^m z^{-n-m} f_{n+m}(z).$$

Spherical Hankel functions: $h_n^{(1)}$, $h_n^{(2)}$

There are also spherical analogues of the Hankel functions:

$$h_n^{(1)}(x) = j_n(x) + i y_n(x)$$
$$h_n^{(2)}(x) = j_n(x) - i y_n(x).$$

In fact, there are simple closed-form expressions for the Bessel functions of half-integer order in terms of the standard trigonometric functions, and therefore for the spherical Bessel functions. In particular, for non-negative integers n:

$$h_n^{(1)}(x) = (-i)^{n+1}\frac{e^{ix}}{x}\sum_{m=0}^{n}\frac{i^m}{m!(2x)^m}\frac{(n+m)!}{(n-m)!}$$

and $h_n^{(2)}$ is the complex-conjugate of this (for real x). It follows, for example, that $j_0(x) = \sin(x)/x$ and $y_0(x) = -\cos(x)/x$, and so on.

The spherical Hankel functions appear in problems involving spherical wave propagation, for example in the multipole expansion of the electromagnetic field.

Riccati–Bessel functions: S_n, C_n, ξ_n, ζ_n

Riccati–Bessel functions only slightly differ from spherical Bessel functions:

$$\begin{aligned}
S_n(x) &= x j_n(x) = \sqrt{\frac{\pi x}{2}}\, J_{n+\frac{1}{2}}(x) \\
C_n(x) &= -x y_n(x) = -\sqrt{\frac{\pi x}{2}}\, Y_{n+\frac{1}{2}}(x) \\
\xi_n(x) &= x h_n^{(1)}(x) = \sqrt{\frac{\pi x}{2}}\, H_{n+\frac{1}{2}}^{(1)}(x) = S_n(x) - iC_n(x) \\
\zeta_n(x) &= x h_n^{(2)}(x) = \sqrt{\frac{\pi x}{2}}\, H_{n+\frac{1}{2}}^{(2)}(x) = S_n(x) + iC_n(x).
\end{aligned}$$

They satisfy the differential equation:

$$x^2\frac{d^2y}{dx^2} + [x^2 - n(n+1)]y = 0.$$

For example, this kind of differential equation appears in Quantum Mechanics while solving the radial component of the Schrodinger's Equation with hypothetical cylindrical infinite potential barrier. For reference see p. 154, Introduction to Quantum Mechanics by Griffiths, 2nd Edition. This differential equation, and the Riccati–Bessel solutions, also arises in the problem of scattering of electromagnetic waves by a sphere, known as Mie scattering after the first published solution by Mie (1908). See e.g., Du (2004)[28] for recent developments and references.

Following Debye (1909), the notation ψ_n, χ_n is sometimes used instead of S_n, C_n.

Asymptotic forms

The Bessel functions have the following asymptotic forms. For small arguments[1] $0 < z \ll \sqrt{\alpha + 1}$, one obtains, when α is not a negative integer:

$$J_\alpha(z) \sim \frac{1}{\Gamma(\alpha+1)}\left(\frac{z}{2}\right)^\alpha$$

When α is a negative integer, we have:

$$J_\alpha(z) \sim \frac{(-1)^\alpha}{(-\alpha)!}\left(\frac{2}{z}\right)^\alpha$$

For the Bessel function of the second kind we have three cases:

$$Y_\alpha(z) \sim \begin{cases} \frac{2}{\pi}\left(\ln\left(\frac{z}{2}\right)+\gamma\right) & \text{if } \alpha = 0 \\ -\frac{\Gamma(\alpha)}{\pi}\left(\frac{2}{z}\right)^\alpha + \frac{1}{\Gamma(\alpha+1)}\left(\frac{z}{2}\right)^\alpha \cot(\alpha\pi) & \text{if } \alpha \text{ is not a non-positive integer (one term dominates unless } \alpha \text{ is imaginary)} \\ -\frac{(-1)^\alpha \Gamma(-\alpha)}{\pi}\left(\frac{z}{2}\right)^\alpha & \text{if } \alpha \text{ is a negative integer} \end{cases}$$

where γ is the Euler–Mascheroni constant (0.5772...).

For large real arguments $x \gg \left|\alpha^2 - \frac{1}{4}\right|$, one cannot write a true asymptotic form for Bessel functions of the first and second kind (unless α is half-integer) because they have zeros all the way out to infinity which would have to be matched exactly by any asymptotic expansion. However, for a given value of arg(z) one can write an equation containing a term of order $|z|^{-1}$:[29]

$$J_\alpha(z) = \sqrt{\frac{2}{\pi z}}\left(\cos\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right) + e^{|\mathrm{Im}(z)|}O(|z|^{-1})\right) \quad \text{for } |\arg z| < \pi$$

$$Y_\alpha(z) = \sqrt{\frac{2}{\pi z}}\left(\sin\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right) + e^{|\mathrm{Im}(z)|}O(|z|^{-1})\right) \quad \text{for } |\arg z| < \pi.$$

(For α = 1/2 the last terms in these formulas drop out completely; see the spherical Bessel functions above.) Even though these equations are true, better approximations may be available for complex z. For example, $J_0(z)$ when z is near the negative real line is approximated better by

$$J_0(z) \approx \sqrt{\frac{-2}{\pi z}}\cos\left(z + \frac{\pi}{4}\right)$$

than by

$$J_0(z) \approx \sqrt{\frac{2}{\pi z}}\cos\left(z - \frac{\pi}{4}\right).$$

The asymptotic forms for the Hankel functions are:

$$H_\alpha^{(1)}(z) \sim \sqrt{\frac{2}{\pi z}} \exp\left(i\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right)\right) \qquad \text{for } -\pi < \arg z < 2\pi$$
$$H_\alpha^{(2)}(z) \sim \sqrt{\frac{2}{\pi z}} \exp\left(-i\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right)\right) \qquad \text{for } -2\pi < \arg z < \pi$$

These can be extended to other values of arg(z) using equations relating $H_\alpha^{(1)}(ze^{im\pi})$ and $H_\alpha^{(2)}(ze^{im\pi})$ to $H_\alpha^{(1)}(z)$ and $H_\alpha^{(2)}(z)$.[30] It is interesting that although the Bessel function of the first kind is the average of the two Hankel functions, $J_\alpha(z)$ is not asymptotic to the average of these two asymptotic forms when z is negative (because one or the other will not be correct there, depending on the arg(z) used). But the asymptotic forms for the Hankel functions permit us to write asymptotic forms for the Bessel functions of first and second kinds for ***complex*** (non-real) z so long as $|z|$ goes to infinity at a constant phase angle arg z (using the square root having positive real part):

$$J_\alpha(z) \sim \frac{1}{\sqrt{2\pi z}} \exp\left(i\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right)\right) \qquad \text{for } -\pi < \arg z < 0$$
$$J_\alpha(z) \sim \frac{1}{\sqrt{2\pi z}} \exp\left(-i\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right)\right) \qquad \text{for } 0 < \arg z < \pi$$
$$Y_\alpha(z) \sim -i\frac{1}{\sqrt{2\pi z}} \exp\left(i\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right)\right) \qquad \text{for } -\pi < \arg z < 0$$
$$Y_\alpha(z) \sim -i\frac{1}{\sqrt{2\pi z}} \exp\left(-i\left(z - \frac{\alpha\pi}{2} - \frac{\pi}{4}\right)\right) \qquad \text{for } 0 < \arg z < \pi$$

For the modified Bessel functions, Hankel developed asymptotic expansions as well:

$$I_\alpha(z) \sim \frac{e^z}{\sqrt{2\pi z}}\left(1 - \frac{4\alpha^2-1}{8z} + \frac{(4\alpha^2-1)(4\alpha^2-9)}{2!(8z)^2} - \frac{(4\alpha^2-1)(4\alpha^2-9)(4\alpha^2-25)}{3!(8z)^3} + \cdots\right) \text{ for } |\arg z| < \tfrac{\pi}{2},$$ [31]
$$K_\alpha(z) \sim \sqrt{\frac{\pi}{2z}}e^{-z}\left(1 + \frac{4\alpha^2-1}{8z} + \frac{(4\alpha^2-1)(4\alpha^2-9)}{2!(8z)^2} + \frac{(4\alpha^2-1)(4\alpha^2-9)(4\alpha^2-25)}{3!(8z)^3} + \cdots\right) \text{ for } |\arg z| < \tfrac{3\pi}{2}.$$ [32]

When $\alpha = 1/2$ all the terms except the first vanish and we have

$$I_{\frac{1}{2}}(z) = \sqrt{\frac{2}{\pi z}}\sinh(z) \sim \frac{e^z}{\sqrt{2\pi z}} \qquad \text{for } |\arg z| < \tfrac{\pi}{2},$$
$$K_{\frac{1}{2}}(z) = \sqrt{\frac{\pi}{2z}}e^{-z}$$

For small arguments $0 < |z| \ll \sqrt{\alpha+1}$, we have:

$$I_\alpha(z) \sim \frac{1}{\Gamma(\alpha+1)}\left(\frac{z}{2}\right)^\alpha$$

$$K_\alpha(z) \sim \begin{cases} -\ln\left(\frac{z}{2}\right) - \gamma & \text{if } \alpha = 0 \\ \frac{\Gamma(\alpha)}{2}\left(\frac{2}{z}\right)^\alpha & \text{if } \alpha > 0. \end{cases}$$

Properties

For integer order α = n, J_n is often defined via a Laurent series for a generating function:

$$e^{\left(\frac{x}{2}\right)(t-1/t)} = \sum_{n=-\infty}^{\infty} J_n(x)t^n,$$

an approach used by P. A. Hansen in 1843. (This can be generalized to non-integer order by contour integration or other methods.) Another important relation for integer orders is the *Jacobi–Anger expansion*:

$$e^{iz\cos(\phi)} = \sum_{n=-\infty}^{\infty} i^n J_n(z) e^{in\phi},$$

and

$$e^{\pm iz\sin(\phi)} = J_0(z) + 2\sum_{n=1}^{\infty} J_{2n}(z)\cos(2n\phi) \pm 2i\sum_{n=0}^{\infty} J_{2n+1}(z)\sin([2n+1]\phi),$$

which is used to expand a plane wave as a sum of cylindrical waves, or to find the Fourier series of a tone-modulated FM signal.

More generally, a series

$$f(z) = a_0^\nu J_\nu(z) + 2\cdot\sum_{k=1} a_k^\nu J_{\nu+k}(z)$$

is called Neumann expansion of f. The coefficients for ν = 0 have the explicit form

$$a_k^0 = \frac{1}{2\pi i}\int_{|z|=c} f(z) O_k(z)\,dz,$$

where O_k is Neumann's polynomial.[33]

Selected functions admit the special representation

$$f(z) = \sum_{k=0} a_k^\nu J_{\nu+2k}(z)$$

with

$$a_k^\nu = 2(\nu + 2k) \int_0^\infty f(z) \frac{J_{\nu+2k}(z)}{z} \, dz$$

due to the orthogonality relation

$$\int_0^\infty J_\alpha(z) J_\beta(z) \frac{dz}{z} = \frac{2}{\pi} \frac{\sin\left(\frac{\pi}{2}(\alpha - \beta)\right)}{\alpha^2 - \beta^2}.$$

More generally, if f has a branch-point near the origin of such a nature that

$$f(z) = \sum_{k=0} a_k J_{\nu+k}(z),$$

then

$$\mathcal{L}\left\{\sum_{k=0} a_k J_{\nu+k}\right\}(s) = \frac{1}{\sqrt{1+s^2}} \sum_{k=0} \frac{a_k}{(s+\sqrt{1+s^2})^{\nu+k}}$$

or

$$\sum_{k=0} a_k \xi^{\nu+k} = \frac{1+\xi^2}{2\xi} \mathcal{L}\{f\} \left(\frac{1-\xi^2}{2\xi}\right),$$

where $\mathcal{L}\{f\}$ is f's Laplace transform.[34]

Another way to define the Bessel functions is the Poisson representation formula and the Mehler-Sonine formula:

$$\begin{aligned} J_\nu(z) &= \frac{\left(\frac{z}{2}\right)^\nu}{\Gamma\left(\nu + \frac{1}{2}\right)\sqrt{\pi}} \int_{-1}^1 e^{izs} (1 - s^2)^{\nu - \frac{1}{2}} \, ds, \\ &= \frac{2}{\left(\frac{z}{2}\right)^\nu \cdot \sqrt{\pi} \cdot \Gamma\left(\frac{1}{2} - \nu\right)} \int_1^\infty \frac{\sin(zu)}{(u^2 - 1)^{\nu + \frac{1}{2}}} \, du, \end{aligned}$$

where $\nu > -1/2$ and $z \in \mathbf{C}$.[35] This formula is useful especially when working with Fourier transforms.

The functions J_α, Y_α, $H_\alpha^{(1)}$, and $H_\alpha^{(2)}$ all satisfy the recurrence relations:[36]

$$\frac{2\alpha}{x}Z_\alpha(x) = Z_{\alpha-1}(x) + Z_{\alpha+1}(x)$$
$$2\frac{dZ_\alpha}{dx} = Z_{\alpha-1}(x) - Z_{\alpha+1}(x)$$

where Z denotes J, Y, $H^{(1)}$, or $H^{(2)}$. (These two identities are often combined, e.g. added or subtracted, to yield various other relations.) In this way, for example, one can compute Bessel functions of higher orders (or higher derivatives) given the values at lower orders (or lower derivatives). In particular, it follows that:[37]

$$\left(\frac{1}{x}\frac{d}{dx}\right)^m \left[x^\alpha Z_\alpha(x)\right] = x^{\alpha-m} Z_{\alpha-m}(x),$$
$$\left(\frac{1}{x}\frac{d}{dx}\right)^m \left[\frac{Z_\alpha(x)}{x^\alpha}\right] = (-1)^m \frac{Z_{\alpha+m}(x)}{x^{\alpha+m}}.$$

Modified Bessel functions follow similar relations :

$$e^{\left(\frac{x}{2}\right)(t+1/t)} = \sum_{n=-\infty}^{\infty} I_n(x) t^n,$$

and

$$e^{z\cos(\theta)} = I_0(z) + 2\sum_{n=1}^{\infty} I_n(z)\cos(n\theta).$$

The recurrence relation reads

$$C_{\alpha-1}(x) - C_{\alpha+1}(x) = \frac{2\alpha}{x}C_\alpha(x)$$
$$C_{\alpha-1}(x) + C_{\alpha+1}(x) = 2\frac{dC_\alpha}{dx}$$

where C_α denotes I_α or $e^{\alpha\pi i}K_\alpha$. These recurrence relations are useful for discrete diffusion problems.

Because Bessel's equation becomes Hermitian (self-adjoint) if it is divided by x, the solutions must satisfy an orthogonality relationship for appropriate boundary conditions. In particular, it follows that:

$$\int_0^1 x J_\alpha(x u_{\alpha,m}) J_\alpha(x u_{\alpha,n})\, dx = \frac{\delta_{m,n}}{2}\left[J_{\alpha+1}(u_{\alpha,m})\right]^2 = \frac{\delta_{m,n}}{2}\left[J'_\alpha(u_{\alpha,m})\right]^2,$$

where $\alpha > -1$, $\delta_{m,n}$ is the Kronecker delta, and $u_{\alpha,\mathrm{m}}$ is the m-th zero of $J_\alpha(x)$. This orthogonality relation can then be used to extract the coefficients in the Fourier–Bessel series, where a function is expanded in the basis of the functions $J_\alpha(x\, u_{\alpha,\mathrm{m}})$ for fixed α and varying m.

An analogous relationship for the spherical Bessel functions follows immediately:

$$\int_0^1 x^2 j_\alpha(xu_{\alpha,m}) j_\alpha(xu_{\alpha,n})\, dx = \frac{\delta_{m,n}}{2}[j_{\alpha+1}(u_{\alpha,m})]^2.$$

If one defines a boxcar function of x that depends on a small parameter ε as:

$$f_\epsilon(x) = \epsilon \operatorname{rect}\left(\frac{x-1}{\epsilon}\right)$$

(where rect() is the rectangle function) then the Hankel transform of it (of any given order α greater than −1/2), $g_\varepsilon(k)$, approaches $J_\alpha(k)$ as ε approaches zero, for any given k. Conversely, the Hankel transform (of the same order) of $g_\varepsilon(k)$ is $f_\varepsilon(x)$:

$$\int_0^\infty kJ_\alpha(kx)g_\epsilon(k)dk = f_\epsilon(x)$$

which is zero everywhere except near 1. As ε approaches zero, the right-hand side approaches $\delta(x-1)$, where δ is the Dirac delta function. So by abuse of language (or "formally"), one says that

$$\int_0^\infty kJ_\alpha(kx)J_\alpha(k)dk = \delta(x-1)$$

even though the integral on the left is not actually defined. A change of variables then yields the *closure equation*:[38]

$$\int_0^\infty xJ_\alpha(ux)J_\alpha(vx)\, dx = \frac{1}{u}\delta(u-v)$$

for $\alpha > -1/2$. The Hankel transform can express a fairly arbitrary function as an integral of Bessel functions of different scales. For the spherical Bessel functions the orthogonality relation is:

$$\int_0^\infty x^2 j_\alpha(ux) j_\alpha(vx)\, dx = \frac{\pi}{2u^2}\delta(u-v)$$

for $\alpha > -1$. Again, this is a useful formal equation whose left-hand side is not actually defined.

Another important property of Bessel's equations, which follows from Abel's identity, involves the Wronskian of the solutions:

$$A_\alpha(x)\frac{dB_\alpha}{dx} - \frac{dA_\alpha}{dx}B_\alpha(x) = \frac{C_\alpha}{x},$$

where A_α and B_α are any two solutions of Bessel's equation, and C_α is a constant independent of x (which depends on α and on the particular Bessel functions considered). In particular,

$$J_\alpha(x)\frac{dY_\alpha}{dx} - \frac{dJ_\alpha}{dx}Y_\alpha(x) = \frac{2}{\pi x},$$

and

$$I_\alpha(x)\frac{dK_\alpha}{dx} - \frac{dI_\alpha}{dx}K_\alpha(x) = -\frac{1}{x}.$$

for $\alpha > -1$. For $\alpha > -1$, the even entire function of genus 1, $x^{-\alpha}J_\alpha(x)$, has only real zeros. Let

$$0 < j_{\alpha,1} < j_{\alpha,2} < \ldots < j_{\alpha,n} < \ldots$$

be all its positive zeros, then

$$J_\alpha(z) = \frac{\left(\frac{z}{2}\right)^\alpha}{\Gamma(\alpha+1)} \prod_{n=1}^{\infty}\left(1 - \frac{z^2}{j_{\alpha,n}^2}\right).$$

(There are a large number of other known integrals and identities that are not reproduced here, but which can be found in the references.)

Multiplication theorem

The Bessel functions obey a multiplication theorem

$$\lambda^{-\nu}J_\nu(\lambda z) = \sum_{n=0}^{\infty}\frac{1}{n!}\left(\frac{(1-\lambda^2)z}{2}\right)^n J_{\nu+n}(z)$$

where λ and ν may be taken as arbitrary complex numbers, see.[39][40] The above expression also holds if J is replaced by Y. The analogous identities for modified Bessel functions are

$$\lambda^{-\nu}I_\nu(\lambda z) = \sum_{n=0}^{\infty}\frac{1}{n!}\left(\frac{(\lambda^2-1)z}{2}\right)^n I_{\nu+n}(z)$$

and

$$\lambda^{-\nu}K_\nu(\lambda z) = \sum_{n=0}^{\infty}\frac{(-1)^n}{n!}\left(\frac{(\lambda^2-1)z}{2}\right)^n K_{\nu+n}(z).$$

Bourget's hypothesis

Bessel himself originally proved that for non-negative integers n, the equation $J_n(x) = 0$ has an infinite number of solutions in x.[41] When the functions $J_n(x)$ are plotted on the same graph, though, none of the zeros seem to coincide for different values of n except for the zero at $x = 0$. This phenomenon is known as **Bourget's hypothesis** after the nineteenth century French mathematician who studied Bessel functions. Specifically it states that for any integers $n \geq 0$ and $m \geq 1$, the functions $J_n(x)$ and $J_{n+m}(x)$ have no common zeros other than the one at $x = 0$. The hypothesis was proved by Carl Ludwig Siegel in 1929.[42]

Selected identities[43]

$$K_{\frac{1}{2}}(z) = \sqrt{\frac{\pi}{2}} e^{-z} z^{-\frac{1}{2}}, \qquad z > 0$$

$$I_{-\frac{1}{2}}(z) = \sqrt{\frac{2}{\pi z}} \cosh(z)$$

$$I_{\frac{1}{2}}(z) = \sqrt{\frac{2}{\pi z}} \sinh(z)$$

$$I_\nu(z) = \sum_{k=0} \frac{z^k}{k!} J_{\nu+k}(z)$$

$$J_\nu(z) = \sum_{k=0} (-1)^k \frac{z^k}{k!} I_{\nu+k}(z)$$

$$I_\nu(\lambda z) = \lambda^\nu \sum_{k=0} \frac{\left((\lambda^2 - 1)\frac{z}{2}\right)^k}{k!} I_{\nu+k}(z)$$

$$I_\nu(z_1 + z_2) = \sum_{k=-\infty}^{\infty} I_{\nu-k}(z_1) I_k(z_2)$$

$$J_\nu(z_1 \pm z_2) = \sum_{k=-\infty}^{\infty} J_{\nu \mp k}(z_1) J_k(z_2)$$

$$I_\nu(z) = \tfrac{z}{2\nu} \left(I_{\nu-1}(z) - I_{\nu+1}(z)\right)$$

$$J_\nu(z) = \tfrac{z}{2\nu} \left(J_{\nu-1}(z) + J_{\nu+1}(z)\right)$$

$$J'_\nu(z) = \begin{cases} \frac{1}{2}\left(J_{\nu-1}(z) - J_{\nu+1}(z)\right) & \nu \neq 0 \\ J\,(z) & \nu = 0 \end{cases}$$

$$-J_1(z) \qquad \nu = 0$$

$$I'_\nu(z) = \begin{cases} \frac{1}{2}\left(I_{\nu-1}(z) + I_{\nu+1}(z)\right) & \nu \neq 0 \\ I_1(z) & \nu = 0 \end{cases}$$

$$\left(\frac{z}{2}\right)^\nu = \Gamma(\nu) \sum_{k=0} I_{\nu+2k}(z)(\nu+2k)\binom{-\nu}{k} = \Gamma(\nu) \sum_{k=0} (-1)^k J_{\nu+2k}(z)(\nu+2k)\binom{-\nu}{k} = \Gamma(\nu+1) \sum_{k=0} \frac{\left(\frac{z}{2}\right)^k}{k!} J_{\nu+k}(z)$$

$$1 = \sum_{n=0}^{\infty} (2n+1) j_n(z)^2$$

$$\frac{\sin(2z)}{2z} = \sum_{n=0}^{\infty} (-1)^n (2n+1) j_n(z)^2$$

See also

- Anger function
- Bessel–Clifford function
- Bessel–Maitland function
- Bessel polynomials
- Fourier–Bessel series
- Hahn–Exton q-Bessel function
- Hankel transform
- Jackson q-Bessel function
- Kelvin functions
- Lerche–Newberger sum rule
- Lommel function
- Lommel polynomial
- Neumann polynomial
- Sonine formula
- Struve function
- Vibrations of a circular drum
- Weber function

Notes

1. Abramowitz and Stegun, p. 360, 9.1.10 (http://www.math.sfu.ca/~cbm/aands/page_360.htm).
2. Abramowitz and Stegun, p. 358, 9.1.5 (http://www.math.sfu.ca/~cbm/aands/page_358.htm).
3. Temme, Nico M. (1996). *Special functions : an introduction to the classical functions of mathematical physics* (2. print. ed.). New York [u.a.]: Wiley. pp. 228–231. ISBN 0471113131.
4. Watson, p. 176 (http://books.google.com/books?id=Mlk3FrNoEVoC&pg=PA176)

5. http://www.math.ohio-state.edu/~gerlach/math/BVtypset/node122.html
6. http://www.nbi.dk/~polesen/borel/node15.html
7. Arfken & Weber, exercise 11.1.17.
8. Abramowitz and Stegun, p. 362, 9.1.69 (http://www.math.sfu.ca/~cbm/aands/page_362.htm).
9. Szegö, G. Orthogonal Polynomials, 4th ed. Providence, RI: Amer. Math. Soc., 1975.
10. http://www.mhtlab.uwaterloo.ca/courses/me755/web_chap4.pdf
11. Watson, p. 178 (http://books.google.com/books?id=Mlk3FrNoEVoC&pg=PA178).
12. Abramowitz and Stegun, p. 358, 9.1.3, 9.1.4 (http://www.math.sfu.ca/~cbm/aands/page_358.htm).
13. Abramowitz and Stegun, p. 358, 9.1.6 (http://www.math.sfu.ca/~cbm/aands/page_358.htm).
14. Abramowitz and Stegun, p. 360, 9.1.25 (http://www.math.sfu.ca/~cbm/aands/page_360.htm).
15. Watson, p. 178 (http://books.google.com/books?id=Mlk3FrNoEVoC&pg=PA178)
16. Abramowitz and Stegun, p. 375, 9.6.2, 9.6.10, 9.6.11 (http://www.math.sfu.ca/~cbm/aands/page_375.htm).
17. Abramowitz and Stegun, p. 374, 9.6.1 (http://www.math.sfu.ca/~cbm/aands/page_374.htm).
18. Quantum electrodynamics. Greiner, Walter and Reinhardt, Joachim. 2009 Springer. pg. 72
19. Watson, p. 181 (http://books.google.com/books?id=Mlk3FrNoEVoC&pg=PA181).
20. M.Kh.Khokonov. *Cascade Processes of Energy Loss by Emission of Hard Photons*, JETP, V.99, No.4, pp. 690-707 (2004). Derived from formulas sourced to I. S. Gradshteĭn and I. M. Ryzhik, *Table of Integrals, Series, and Products* (Fizmatgiz, Moscow, 1963; Academic, New York, 1980).
21. Referred to as such in: Teichroew, D. *The Mixture of Normal Distributions with Different Variances*, The Annals of Mathematical Statistics. Vol. 28, No. 2 (Jun., 1957), pp. 510–512
22. Abramowitz and Stegun, p. 437, 10.1.1 (http://www.math.sfu.ca/~cbm/aands/page_437.htm).
23. Abramowitz and Stegun, p. 439, 10.1.25, 10.1.26 (http://www.math.sfu.ca/~cbm/aands/page_439.htm);
24. Abramowitz and Stegun, p. 438, 10.1.11 (http://www.math.sfu.ca/~cbm/aands/page_438.htm).
25. Abramowitz and Stegun, p. 438, 10.1.12 (http://www.math.sfu.ca/~cbm/aands/page_438.htm);
26. Abramowitz and Stegun, p. 439, 10.1.39 (http://www.math.sfu.ca/~cbm/aands/page_439.htm).
27. Abramowitz and Stegun, p. 439, 10.1.23, 10.1.24 (http://people.math.sfu.ca/~cbm/aands/page_439.htm).
28. Hong Du, "Mie-scattering calculation," *Applied Optics* **43** (9), 1951–1956 (2004)
29. Abramowitz and Stegun, p. 364, 9.2.1 (http://www.math.sfu.ca/~cbm/aands/page_364.htm);
30. NIST Digital Library of Mathematical Functions, Section 10.11 (http://dlmf.nist.gov/10.11#E1).
31. Abramowitz and Stegun, p. 377, 9.7.1 (http://www.math.sfu.ca/~cbm/aands/page_377.htm);
32. Abramowitz and Stegun, p. 378, 9.7.2 (http://www.math.sfu.ca/~cbm/aands/page_378.htm);
33. Abramowitz and Stegun, p. 363, 9.1.82 (http://www.math.sfu.ca/~cbm/aands/page_363.htm) ff.
34. E. T. Whittaker, G. N. Watson, A course in modern Analysis p. 536 (http://books.google.com/books?id=Mlk3FrNoEVoC&lpg=PA522&dq=bessel%20neumann%20series&pg=PA536#v=onepage&q=bessel%20neumann%20series&f=false)
35. I.S. Gradshteyn (И.С. Градштейн), I.M. Ryzhik (И.М. Рыжик); Alan Jeffrey, Daniel Zwillinger, editors. *Table of Integrals, Series, and Products*, seventh edition. Academic Press, 2007. ISBN 978-0-12-373637-6. Equation 8.411.10
36. Abramowitz and Stegun, p. 361, 9.1.27 (http://people.math.sfu.ca/~cbm/aands/page_361.htm).
37. Abramowitz and Stegun, p. 361, 9.1.30 (http://people.math.sfu.ca/~cbm/aands/page_361.htm).
38. Arfken & Weber, section 11.2
39. Abramowitz and Stegun, p. 363, 9.1.74 (http://www.math.sfu.ca/~cbm/aands/page_363.htm).
40. C. Truesdell, "On the Addition and Multiplication Theorems for the Special Functions (http://www.pnas.org/cgi/reprint/36/12/752.pdf)", *Proceedings of the National Academy of Sciences, Mathematics*, (1950) pp.752–757.
41. F. Bessel, *Untersuchung des Theils der planetarischen Störungen*, Berlin Abhandlungen (1824), article 14.
42. Watson, pp. 484–5
43. See, for example, Lide DR. CRC handbook of chemistry and physics: a ready-reference book of chemical CRC Press, 2004, ISBN 0-8493-0485-7, p. A-95

References

- Abramowitz, Milton; Stegun, Irene A., eds. (December 1972) [1964]. "Chapter 9". *Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables*. Applied Mathematics Series **55** (10 ed.). New York, USA: United States Department of Commerce, National Bureau of Standards; Dover Publications. pp. 355, 435. ISBN 978-0-486-61272-0. LCCN 64-60036. MR 0167642. See also chapter 10 (http://www.math.sfu.ca/~cbm/aands/page_435.htm).
- Arfken, George B. and Hans J. Weber, *Mathematical Methods for Physicists*, 6th edition (Harcourt: San Diego, 2005). ISBN 0-12-059876-0.
- Bayin, S.S. *Mathematical Methods in Science and Engineering*, Wiley, 2006, Chapter 6.
- Bayin, S.S., *Essentials of Mathematical Methods in Science and Engineering*, Wiley, 2008, Chapter 11.
- Bowman, Frank *Introduction to Bessel Functions* (Dover: New York, 1958). ISBN 0-486-60462-4.
- G. Mie, "Beiträge zur Optik trüber Medien, speziell kolloidaler Metallösungen", *Ann. Phys. Leipzig* **25** (1908), p. 377.
- Olver, F. W. J.; Maximon, L. C. (2010), "Bessel function", in Olver, Frank W. J.; Lozier, Daniel M.; Boisvert, Ronald F.; Clark, Charles W., *NIST Handbook of Mathematical Functions*, Cambridge University Press, ISBN 978-0521192255, MR 2723248
- Press, WH; Teukolsky, SA; Vetterling, WT; Flannery, BP (2007), "Section 6.5. Bessel Functions of Integer Order", *Numerical Recipes: The Art of Scientific Computing* (3rd ed.), New York: Cambridge University Press, ISBN 978-0-521-88068-8
- B Spain, M.G. Smith, *Functions of mathematical physics*, Van Nostrand Reinhold Company, London, 1970. Chapter 9 deals with Bessel functions.
- N. M. Temme, *Special Functions. An Introduction to the Classical Functions of Mathematical Physics*, John Wiley and Sons, Inc., New York, 1996. ISBN 0-471-11313-1. Chapter 9 deals with Bessel functions.
- Watson, G.N., *A Treatise on the Theory of Bessel Functions, Second Edition*, (1995) Cambridge University Press. ISBN 0-521-48391-3.
- Weber, H. (1873), "Ueber eine Darstellung willkürlicher Functionen durch Bessel'sche Functionen", *Mathematische Annalen* **6** (2): 146–161, doi:10.1007/BF01443190

External links

- Lizorkin, P.I. (2001), "Bessel functions", in Hazewinkel, Michiel, *Encyclopedia of Mathematics*, Springer, ISBN 978-1-55608-010-4
- Karmazina, L.N.; Prudnikov, A.P. (2001), "Cylinder function", in Hazewinkel, Michiel, *Encyclopedia of Mathematics*, Springer, ISBN 978-1-55608-010-4
- Rozov, =N.Kh. (2001), "Bessel equation", in Hazewinkel, Michiel, *Encyclopedia of Mathematics*, Springer, ISBN 978-1-55608-010-4
- Wolfram function pages on Bessel J (http://functions.wolfram.com/Bessel-TypeFunctions/BesselJ/) and Y (http://functions.wolfram.com/Bessel-TypeFunctions/BesselY/) functions, and modified Bessel I (http://functions.wolfram.com/Bessel-TypeFunctions/BesselI/) and K (http://functions.wolfram.com/Bessel-TypeFunctions/BesselK/) functions. Pages include formulas, function evaluators, and plotting calculators.
- Wolfram Mathworld – Bessel functions of the first kind (http://mathworld.wolfram.com/BesselFunctionoftheFirstKind.html)
- Bessel functions J_ν (http://www.librow.com/articles/article-11/appendix-a-34), Y_ν (http://www.librow.com/articles/article-11/appendix-a-35), I_ν (http://www.librow.com/articles/article-11/appendix-a-36) and K_ν (http://www.librow.com/articles/article-11/appendix-a-37) in Librow Function handbook (http://www.librow.com/articles/article-11).
- F. W. J. Olver, L. C. Maximon, Bessel Functions (http://dlmf.nist.gov/10) (chapter 10 of the Digital Library of Mathematical Functions).

Retrieved from "https://en.wikipedia.org/w/index.php?title=Bessel_function&oldid=696605545"

Categories: Special hypergeometric functions | Fourier analysis

---