

European Transport Simulator:

*transport equations, variables
and Fortran implementation*

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Table of Contents

General comments	4
<i>Grids:</i>	4
<i>Form of transport equations and adiabatic compression terms:</i>	4
<i>Generalized form of transport equations:</i>	6
Current diffusion equation.....	7
<i>Variables and units:</i>	7
<i>Equation:</i>	7
<i>Boundary Conditions:</i>	8
<i>Definitions for quantities derived from current equation:</i>	9
<i>Translation to generalized form for interface with numerical solver:</i>	10
<i>Variables used inside fortran routines:</i>	11
Ion density.....	14
<i>Variables and units:</i>	14
<i>Equation:</i>	14
<i>Boundary Conditions:</i>	15
<i>Definitions for quantities derived from ion density equation:</i>	16
<i>Translation to generalized form for interface with numerical solver:</i>	16
<i>Variables used inside fortran routines:</i>	18
Electron density.....	21
<i>Variables and units:</i>	21
<i>Equation:</i>	21
<i>Boundary Conditions:</i>	22
<i>Definitions for quantities derived from electron density equation:</i>	23
<i>Translation to generalized form for interface with numerical solver:</i>	23
<i>Variables used inside fortran routines:</i>	25
Quasi-neutrality condition	27
<i>Variables and units</i>	27
<i>Equations:</i>	27
<i>Definitions for quantities derived from quasi-neutrality conditions:</i>	27
<i>Variables used inside fortran routines:</i>	29
Ion temperatures.....	30
<i>Variables and units:</i>	30
<i>Equation:</i>	30
<i>Boundary Conditions:</i>	31
<i>Definitions for quantities derived from ion temperature equation:</i>	32
<i>Translation to generalized form for interface with numerical solver:</i>	33
<i>Variables used inside fortran routines:</i>	35
Electron energy transport equation	38
<i>Variables and units:</i>	38
<i>Equation:</i>	38
<i>Boundary Conditions:</i>	39
<i>Definitions for quantities derived from electron temperature equation:</i>	40
<i>Translation to generalized form for interface with numerical solver:</i>	41

<i>Variables used inside fortran routines:</i>	43
Rotation transport equation.....	46
<i>Variables and units:</i>	46
<i>Equation:</i>	46
<i>Boundary Conditions:</i>	47
<i>Definitions for quantities derived from rotation equation:</i>	48
<i>Translation to generalized form for interface with numerical solver:</i>	49
<i>Variables used inside fortran routines:</i>	52

General comments

Grids:

ETS uses two types of grids for the toroidal flux coordinate, not normalized ρ and normalized x , which are defined as:

$$\rho = \sqrt{\frac{\Phi}{\pi B_0}} \quad [\text{m}];$$

Φ [Wb] is the toroidal flux and B_0 [T] is the magnetic field measured at the characteristic major radius of the device R_0 [m];

$x = \frac{\rho}{\rho_b}$ [-], where ρ_b is the coordinate of plasma magnetic boundary.

$$\frac{\partial}{\partial \rho} = \frac{\partial}{\partial x} \frac{\partial x}{\partial \rho} = \frac{1}{\rho_b} \frac{\partial}{\partial x}$$

All interfaces between different modules and subroutines use ρ as the primary coordinate, internally, transport equations are solved using x coordinate.

Use of the normalized coordinate x introduces an artificial non-physical convection term in all transport equations. Any change in the boundary value of $\rho_b(t)$ results in rearrangement of all plasma characteristics throughout the whole cross-section. In many circumstances, as the initial phase of discharge, L-H transition, strong additional heating, etc., the correction is not negligible.

Form of transport equations and adiabatic compression terms:

Consider the equation:

$$\frac{\partial}{\partial t} \left(\frac{\partial V}{\partial \rho} \cdot \langle f \rangle \right) + \frac{\partial}{\partial \rho} \left(\frac{\partial V}{\partial \rho} \cdot \langle |\nabla \rho|^2 \rangle \cdot \Gamma_\rho \right) = \frac{\partial V}{\partial \rho} \cdot S, \quad (0.1)$$

where:

- $\Gamma_\rho = \langle f (\vec{u} - \vec{u}_\rho) \nabla \rho \rangle / \langle |\nabla \rho|^2 \rangle$ is the flux through the surface $\rho = \text{const}$. One shall note, that it is different to the diffusive flux $\Gamma_D = \langle f (\vec{u} - \vec{u}_\Phi) \nabla \rho \rangle / \langle |\nabla \rho|^2 \rangle$. The difference between Γ_ρ and Γ_D can be highlighted by noting that, in absence of collisions, $\Gamma_D \equiv 0$ whereas Γ_ρ can be anything. On the other hand under the condition of $\dot{B}_0 = \frac{dB_0}{dt} = 0$, $\Gamma_D \equiv \Gamma_\rho$.
- here \vec{u} is the plasma mass velocity and \vec{u}_ρ is the velocity of the surface $\rho = \text{const}$
- S is the source density.

For negligible variations of f on a flux-surface, the difference between Γ_ρ and Γ_D is proportional to

$$\langle (\vec{u}_\rho - \vec{u}_\Phi) \cdot \nabla \rho \rangle = \frac{\partial \rho}{\partial t} \Big|_\Phi = \frac{\dot{B}_0}{2B_0} \rho.$$

Thus the transport equation can be re-written in the form:

$$\frac{\partial}{\partial t} \left|_{\rho} \left(\frac{\partial V}{\partial \rho} \langle f \rangle \right) + \frac{\partial}{\partial \rho} \left(\frac{\partial V}{\partial \rho} \left[\langle |\nabla \rho|^2 \rangle \Gamma_D - \frac{\dot{B}_0}{2B_0} \rho \langle f \rangle \right] \right) = \frac{\partial V}{\partial \rho} \cdot S,$$

or as

$$\left(\frac{\partial}{\partial t} \left|_{\rho} - \frac{\dot{B}_0}{2B_0} \frac{\partial}{\partial \rho} \rho \right) \cdot \left(\frac{\partial V}{\partial \rho} \langle f \rangle \right) + \frac{\partial}{\partial \rho} \left(\frac{\partial V}{\partial \rho} \langle |\nabla \rho|^2 \rangle \left(-D \frac{\partial}{\partial \rho} \langle f \rangle + \nu \langle f \rangle \right) \right) = \frac{\partial V}{\partial \rho} \cdot S, \quad (0.2)$$

on the assumption of

$$\Gamma_D = -D \frac{\partial}{\partial \rho} \langle f \rangle + \nu \langle f \rangle.$$

This is a standard representation of the transport equation used in ETS for the fixed ρ grid

If we introduce following notations:

$$k_{\Phi}(t) = \frac{\dot{\Phi}_b}{2\Phi_b} = k_B(t) + k_{\rho}(t), \quad k_B(t) = \frac{\dot{B}_0}{2B_0} = \frac{1}{2B_0} \cdot \frac{dB}{dt}, \quad k_{\rho}(t) = \frac{\dot{\rho}_b}{\rho_b} = \frac{1}{\rho_b} \cdot \frac{d\rho_b}{dt} \quad (0.3)$$

The equation can be rewritten as:

$$\frac{\partial}{\partial t} \left|_{\rho} \cdot \left(\frac{\partial V}{\partial \rho} \langle f \rangle \right) + \frac{\partial}{\partial \rho} \left(\frac{\partial V}{\partial \rho} \left[\langle |\nabla \rho|^2 \rangle \left(-D \frac{\partial}{\partial \rho} \langle f \rangle + \nu \langle f \rangle \right) - k_B \rho \langle f \rangle \right] \right) = \frac{\partial V}{\partial \rho} \cdot S. \quad (0.4)$$

Using (0.3) and

$$\frac{\partial a}{\partial t} \left|_{\rho} = \frac{\partial a}{\partial t} \left|_x + \frac{\partial a}{\partial x} \cdot \frac{\partial x}{\partial t} \right|_{\rho},$$

equation (0.4) can be rewritten for the normalised coordinate, x :

$$\frac{\partial}{\partial t} \left|_x \cdot \left(\frac{\partial V}{\partial \rho} \langle f \rangle \right) + \frac{\partial}{\partial \rho} \left(\frac{\partial V}{\partial \rho} \left[\langle |\nabla \rho|^2 \rangle \left(-D \frac{\partial}{\partial \rho} \langle f \rangle + \nu \langle f \rangle \right) - k_{\Phi} \rho \langle f \rangle \right] \right) = \frac{\partial V}{\partial \rho} \cdot (S - k_{\rho} \langle f \rangle). \quad (0.5)$$

It is also useful to define some transformations as:

$$\begin{aligned} \frac{3}{2} \left(\frac{\partial V}{\partial \rho} \right)^{-5/3} \cdot \left(\frac{\partial}{\partial t} \left|_{\rho} - \frac{\dot{B}_0}{2B_0} \frac{\partial}{\partial \rho} \rho \right) \cdot \left[\left(\frac{\partial V}{\partial \rho} \right)^{5/3} \langle f \rangle \right] = \\ \frac{3}{2} \left(\frac{\partial V}{\partial \rho} \right)^{-5/3} \frac{\partial}{\partial t} \left|_x \left[\left(\frac{\partial V}{\partial \rho} \right)^{5/3} \langle f \rangle \right] - \left(\frac{\partial V}{\partial \rho} \right)^{-1} \frac{\partial}{\partial \rho} \cdot \left[\left(\frac{\partial V}{\partial \rho} \right) \cdot \frac{3}{2} \frac{\dot{\Phi}_b}{2\Phi_b} \rho \cdot \langle f \rangle \right] + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \rho \frac{\partial}{\partial \rho} \ln \frac{\partial V}{\partial \rho} \right] \langle f \rangle \end{aligned} \quad (0.6)$$

and:

$$\rho \frac{\partial}{\partial \rho} \ln \frac{\partial V}{\partial \rho} = 1 + \rho \frac{\partial}{\partial \rho} \ln \left(\frac{1}{\rho} \frac{\partial V}{\partial \rho} \right) = 1 + \rho \frac{\partial}{\partial \rho} \ln \left(\frac{4\pi^2 B_0}{F \left\langle \frac{1}{R^2} \right\rangle} \right) = 1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \left\langle \frac{1}{R^2} \right\rangle \right) \quad (0.7)$$

Generalized form of transport equations:

All numerical solvers used within the ETS are adapted to solve the transport equation in the form:

$$\frac{a(x) \cdot Y(x, t) - b(x) \cdot Y(x, t-1)}{h} + \frac{1}{c(x)} \frac{\partial}{\partial x} \left(-d(x) \cdot \frac{\partial Y(x, t)}{\partial x} + e(x) \cdot Y(x, t) \right) = f(x) - g(x) \cdot Y(x, t) \quad (0.8)$$

where Y is the function, t is the time and x is the radial coordinate.

The boundary conditions shall be provided in the form:

$$v(x_{bnd}) \cdot \frac{\partial Y(x, t)}{\partial x} \Big|_{bnd} + u(x_{bnd}) \cdot Y(x_{bnd}, t) = w(x_{bnd})$$

Therefore the following set of quantities is expected by the numerical solver:

$x, a(x), b(x), c(x), d(x), e(x), f(x), g(x), h, Y(x, t-1)$

$u(1:2), v(1:2), w(1:2),$

where index 1 refers to the plasma axis and 2 to the plasma boundary.

The solver will return:

$$Y(x, t), \frac{\partial Y(x, t)}{\partial x}$$

Current diffusion equation

Variables and units:

- ψ - flux function [Wb]
- j_{\parallel} - parallel current density [A/m²]
- j_{tor} - toroidal current density [A/m²]
- Q_{OH} - ohmic heating power [W]
- q - safety factor [-]
- E_{\parallel} - parallel electric field [V/m]
- σ_{\parallel} - parallel conductivity [Ohm⁻¹ m⁻¹]
- μ_0 - permeability of free space ;
- R_0 - characteristic major radius of the device [m]
- B_0 - magnetic field measured at R_0 [T]
- F - diamagnetic function [m·T]

Equation:

For a constant ρ grid, the equation

$$\text{is } \sigma_{\parallel} \left(\frac{\partial \Psi}{\partial \alpha} \Big|_{\rho} - \frac{\rho \dot{B}_0}{2B_0} \frac{\partial \Psi}{\partial \rho} \right) = \frac{F^2}{\mu_0 B_0 \rho} \frac{\partial}{\partial \rho} \left[\frac{V'}{4\pi^2} \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle \frac{1}{F} \frac{\partial \Psi}{\partial \rho} \right] - \frac{V'}{2\pi \rho} (j_{ni,exp} + j_{ni,imp} \cdot \Psi)$$

Where the diamagnetic function, $F = RB_{\phi}$, the geometry coefficients V' and $\left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle$ are provided by the equilibrium solver.

For a constant normalized grid, this becomes

$$\sigma_{\parallel} \frac{\partial \Psi}{\partial \alpha} \Big|_x = \frac{F^2}{\mu_0 B_0 \rho} \frac{\partial}{\partial \rho} \left[\frac{V'}{4\pi^2} \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle \frac{1}{F} \frac{\partial \Psi}{\partial \rho} \right] - \frac{V'}{2\pi \rho} (j_{ni,exp} + j_{ni,imp} \cdot \Psi) + \rho \sigma_{\parallel} \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \frac{\partial \Psi}{\partial \rho} \quad (1.1)$$

Equation (1.1) can be reformulated as:

$$\begin{aligned} \sigma_{\parallel} \frac{\partial \Psi}{\partial \alpha} \Big|_x = & \frac{F^2}{\mu_0 B_0 \rho} \frac{\partial}{\partial \rho} \left[\frac{V'}{4\pi^2} \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle \frac{1}{F} \frac{\partial \Psi}{\partial \rho} + \frac{\mu_0 B_0 \rho^2 \sigma_{\parallel}}{F^2} \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \Psi \right] - \frac{V'}{2\pi \rho} (j_{ni,exp} + j_{ni,imp} \cdot \Psi) \\ & + \sigma_{\parallel} \left(2\rho \cdot \frac{1}{F} \cdot \frac{\partial F}{\partial \rho} - \rho \cdot \frac{1}{\sigma_{\parallel}} \cdot \frac{\partial \sigma_{\parallel}}{\partial \rho} - 2 \right) \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \Psi \end{aligned} \quad (1.2)$$

non inductive current includes contributions in generic form from external sources (*computed by other modules, like ICRH, ECRH, NBI, neutrals and etc.*), in the form of explicit and implicit terms:

$$j_{ni,exp} = \sum_{isource=1}^{nsource} m_{source} \cdot j_{isource,exp} \quad (1.3)$$

$$j_{ni,imp} = \sum_{isource=1}^{nsource} m_{source} \cdot j_{isource,imp} \cdot \quad (1.4)$$

This summation is done in the workflow by [Source Combiner](#), where individual weights, m_{source} , can be adjusted.

Boundary Conditions:

On axis, at $\rho = 0$, ETS assumes:

$$\left. \frac{\partial \psi}{\partial \rho} \right|_{\rho=0} = 0 \quad (1.5)$$

At the edge, at $\rho = \rho_{bnd}$, following options (distinguished by boundary condition type) are available:

type=1 (value)

$$\psi|_{\rho=\rho_{bnd}} = \psi_{bnd}, \quad (1.6)$$

where ψ_{bnd} is the value of poloidal flux at the outer boundary.

type=2 (total current inside $\rho = \rho_{bnd}$)

$$\left. \frac{\partial \psi}{\partial \rho} \right|_{\rho=\rho_{bnd}} = \frac{4\pi^2 \mu_0}{V' \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle} I_{bnd} \quad (1.7)$$

where I_{bnd} is the value of plasma current inside the boundary.

type=3 (loop voltage at $\rho = \rho_{bnd}$)

$$\left. \frac{\partial \psi}{\partial t} \right|_{\rho=\rho_{bnd}} = U_{loop,bnd} - \left. \frac{\partial \Psi}{\partial \rho} \right|_{\rho=\rho_{bnd}} \cdot \frac{\partial \rho_b}{\partial t} \quad (1.8)$$

where $U_{loop,bnd}$ is the plasma loop voltage at the boundary.

type=4 (generic)

$$v_{gen} \left. \frac{\partial \psi}{\partial \rho} \right|_{\rho=\rho_{bnd}} + u_{gen} \psi|_{\rho=\rho_{bnd}} = w_{gen} \quad (1.9)$$

Definitions for quantities derived from current equation:

Poloidal components of the magnetic field and flux function:

$$B_{pol} = \frac{|\nabla \rho|}{2\pi R_0} \frac{\partial \Psi}{\partial \rho} \quad (1.10)$$

Safety factor:

$$q = \frac{2\pi B_0 \rho}{(\partial \Psi / \partial \rho)} \quad (1.12)$$

Total current:

$$\text{(toroidal)} \quad j_{tor} = \frac{2\pi R_0}{\mu_0 V'} \cdot \frac{\partial}{\partial \rho} \left(H \frac{\partial \Psi}{\partial \rho} \right); \quad H = \frac{V'}{4\pi^2} \cdot \left\langle \left(\frac{\nabla \rho}{R} \right)^2 \right\rangle \quad (1.13)$$

$$\text{(parallel)} \quad j_{\parallel} = \frac{2\pi}{\mu_0 R_0 V'} \cdot \left(\frac{F}{B_0} \right)^2 \frac{\partial}{\partial \rho} \left(\frac{R_0 B_0}{F} H \frac{\partial \Psi}{\partial \rho} \right) \quad (1.14)$$

Ohmic heating & parallel electric field:

$$Q_{OH} = \sigma_{\parallel} E_{\parallel}^2 \quad E_{\parallel} = \frac{1}{\sigma_{\parallel}} (j_{\parallel} - j_{ni,exp} - j_{ni,imp} \cdot \Psi) \quad (1.15)$$

If equation (1.2) is not solved, ets will extrapolate the quantities using safety factor profile as primary quantity:

$$\Psi = \int_0^{\rho} \frac{2\pi B_0}{q} \rho \partial \rho; \quad (1.16)$$

$$j_{tor} = \frac{2\pi R_0}{\mu_0 V'} \cdot \frac{\partial}{\partial \rho} \left(H \frac{2\pi B_0 \rho}{q} \right); \quad (1.17)$$

$$j_{\parallel} = \frac{2\pi}{\mu_0 R_0 V'} \cdot \left(\frac{F}{B_0} \right)^2 \frac{\partial}{\partial \rho} \left(\frac{R_0 B_0}{F} H \frac{2\pi B_0 \rho}{q} \right) \quad (1.18)$$

$$Q_{OH} = \frac{1}{\sigma_{\parallel}} (j_{\parallel} - j_{ni,exp} - j_{ni,imp} \cdot \Psi)^2 \quad (1.19)$$

Translation to generalized form for interface with numerical solver:

Rewrite the equation (1.2) using the time discretization: $\frac{\Psi - \Psi^-}{\tau}$, where Ψ^- is the value of the poloidal flux at previous time and τ is the time step.

$$\sigma_{//} \frac{\Psi - \Psi^-}{\tau} \Big|_x + \frac{F^2}{\mu_0 B_0 \rho \rho_b} \frac{\partial}{\partial x} \left[-\frac{V'}{4\pi^2} \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle \frac{1}{F \rho_b} \frac{\partial \Psi}{\partial x} - \frac{\mu_0 B_0 \rho^2 \sigma_{//}}{F^2} \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \Psi \right] = -\frac{V'}{2\pi \rho} j_{ni,exp} - \left[\frac{V'}{2\pi \rho} j_{ni,imp} + \sigma_{//} \left(2 - 2\rho \cdot \frac{1}{F} \cdot \frac{\partial F}{\partial \rho} + \rho \cdot \frac{1}{\sigma_{//}} \cdot \frac{\partial \sigma_{//}}{\partial \rho} \right) \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \right] \cdot \Psi \quad (1.20)$$

Then, using the equation (0.8), coefficients for the interface with numerical solver are defined as:

$$\begin{aligned} a(x) &= \sigma_{//} \\ b(x) &= \sigma_{//} \\ c(x) &= \frac{\mu_0 B_0 \rho \rho_b}{F^2} \\ d(x) &= \frac{V'}{4\pi^2} \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle \frac{1}{F \rho_b} \\ e(x) &= -\frac{\mu_0 B_0 \rho^2 \sigma_{//}}{F^2} \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \\ f(x) &= -\frac{V'}{2\pi \rho} j_{ni,exp} \\ g(x) &= \frac{V'}{2\pi \rho} j_{ni,imp} + \sigma_{//} \left(2 - 2\rho \cdot \frac{1}{F} \cdot \frac{\partial F}{\partial \rho} + \rho \cdot \frac{1}{\sigma_{//}} \cdot \frac{\partial \sigma_{//}}{\partial \rho} \right) \cdot \frac{\dot{\Phi}_b}{2\Phi_b} \\ h &= \tau \\ Y^{t-1}(\rho) &= \Psi^- \end{aligned} \quad (1.21)$$

boundary conditions

at the axis:

$$v(1) = 1; \quad u(1) = 0; \quad w(1) = 0$$

At the edge, $x = 1$:

(1.22)

type=1 (value)

$$\psi|_{x=1} - \psi_{bnd} = 0$$

$$v(2) = 0; \quad u(2) = 1; \quad w(2) = \psi_{bnd}$$

type=2 (total current inside $x = 1$)

$$\frac{\partial \psi}{\partial x} \Big|_{x=1} = \frac{4\pi^2 \mu_0}{V' \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle} I_{bnd} \cdot \rho_b$$

$$v(2) = 1; \quad u(2) = 0; \quad w(2) = \frac{4\pi^2 \mu_0}{V' \left\langle \left| \frac{\nabla \rho}{R} \right|^2 \right\rangle} I_{bnd} \cdot \rho_b$$

type=3 (loop voltage at $x = 1$)

$$\psi|_{x=1} - (\tau U_{loop,bnd} + \psi_{bnd}^-) = 0$$

$$v(2) = 0; \quad u(2) = 1; \quad w(2) = \tau U_{loop,bnd} + \psi_{bnd}^-$$

type=4 (generic)

$$\frac{v_{gen}}{\rho_b} \frac{\partial \Psi}{\partial x} \Big|_{bnd} + u_{gen} \Psi_{bnd} = w_{gen}$$

$$v(2) = \frac{v_{gen}}{\rho_b}; \quad u(2) = u_{gen}; \quad w(2) = w_{gen}$$

Variables used inside fortran routines:

(Routine: MAIN_PLASMA, Subroutine: CURRENT)

Variable	TYPE%NAME used in ETS data flow	Internal name used in CURRENT routine	Units
τ	EVOLUTION%TAU	TAU	s
\dot{B}_0	GEOMETRY%BTPRIME	BTPRIME	T/s
B_0	GEOMETRY%BT	BT	T
R_0	GEOMETRY%RO	RO	m
ρ	GEOMETRY%RHO	RHO	m

x	GEOMETRY%RHO_NORM	RHO_NORM	-
V'	GEOMETRY%VPR	VPR	m^2
$\left\langle \left \frac{\nabla \rho}{R} \right ^2 \right\rangle$	GEOMETRY%G3	G3	m^{-2}
F	GEOMETRY%FDIA	FDIA	T·m
H	--	H	
π	--	PI	
μ_0	--	MU0	H/m
Ψ	PROFILES%PSI	PSI	V·s
Ψ^-	EVOLUTION%PSIM	PSIM	V·s
q	PROFILES%QSF	QSF	--
j_{tor}	PROFILES%CURR_TOR	CURR_TOR	A/m ²
j_{\parallel}	PROFILES %CURR_PAR	CURR_PAR	A/m ²
$j_{ni,exp}$	--	CURR_NI_EXP	A/m ²
$j_{ni,imp}$	--	CURR_NI_IMP	A/(V·s·m ²)
$j_{source1}$	SOURCES%CURR_NI_EXP	--	A/m ²
$j_{source2}$	SOURCES %CURR_NI_IMP	--	A/(V·s·m ²)
Q_{OH}	SOURCES%QOH	QOH	W/m ³
σ_{\parallel}	TRANSPORT%SIGMA or SOURCES%SIGMA or COLLISIONS%SIGMA	SIGMA	(Ohm·m) ⁻¹
E_{\parallel}	PROFILES %E_PAR	E_PAR	V/m
a	SOLVER%A	A	--
b	SOLVER%B	B	--
c	SOLVER%C	C	--
d	SOLVER%D	D	--
e	SOLVER%E	E	--
f	SOLVER%F	F	--
g	SOLVER%G	G	--
h	SOLVER%H	H	--
$v(1:2)$	SOLVER%V	V	--
$u(1:2)$	SOLVER%U	U	--

$w(1:2)$	SOLVER%W	W	--
Y^{i-1}	SOLVER%YM	W	--
<i>solution</i>	SOLVER%Y	Y	--
<i>derivative of solution</i>	SOLVER%DY	DY	--
Functions		Internal name used in CURRENT routine	
$\frac{\sigma_{//}\mu_0\rho^2}{F^2}$		FUN1	
$\frac{\partial}{\partial\rho} \frac{\sigma_{//}\mu_0\rho^2}{F^2}$		DFUN1	
$\frac{V'}{4\pi^2} \left\langle \left \frac{\nabla\rho}{R} \right ^2 \right\rangle$		FUN2	
$\frac{2\pi B_0}{q}$		FUN3	
$H \frac{\partial\psi}{\partial\rho}$		FUN4	
$\frac{\partial}{\partial\rho} \left(H \frac{\partial\psi}{\partial\rho} \right)$		DFUN4	
$\frac{R_0 B_0}{F} H \frac{\partial\psi}{\partial\rho}$		FUN5	
$\frac{\partial}{\partial\rho} \left(\frac{R_0 B_0}{F} H \frac{\partial\psi}{\partial\rho} \right)$		DFUN5	

Ion density

Variables and units:

n_i - ion density [m^{-3}]

Γ_i - ion flux [s^{-1}]

γ_i - ion flux contributing to heat transport [A/m^2]

Γ_{Si} - integral source (approximation for the flux if equation is not solved) [s^{-1}]

$n_{i,\text{int}}$ - interpretative ion density [m^{-3}], required if equation is not solved

D_i - ion diffusion coefficient [m^2/s]

V_i^{pinch} - ion pinch velocity [m/s]

$S_{i,\text{exp}}$ - explicit part of ion source density [$\text{s}^{-1} \text{m}^{-3}$]

$S_{i,\text{imp}}$ - implicit part of ion source density [s^{-1}];

Equation:

$$\left(\frac{\partial}{\partial \rho} \right)_{\rho} - \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial}{\partial \rho} \rho \left(V' n_i \right) + \frac{\partial}{\partial \rho} \Gamma_i = V' (S_{i,\text{exp}} - S_{i,\text{imp}} \cdot n_i) \quad (2.1)$$

where the total flux defined as:

$$\Gamma_i = V' \left(|\nabla \rho|^2 \right) \left(-D_i \frac{\partial n_i}{\partial \rho} + n_i V_i^{\text{pinch}} \right); \quad (2.2)$$

where total transport coefficients, D_i and V_i^{pinch} , are defined as a sum of individual contributions from different transport models weighted with coefficients m_{model} :

$$D_i = \sum_{i \text{ model}=1}^{n \text{ model}} m_{\text{model}} \cdot D_{i,i \text{ model}}; \quad (2.3)$$

$$V_i^{\text{pinch}} = \sum_{i \text{ model}=1}^{n \text{ model}} m_{\text{model}} \cdot V_{i,i \text{ model}}^{\text{pinch}}. \quad (2.4)$$

(Combination of individual contributions, $D_{i,i \text{ model}}$ and $V_{i,i \text{ model}}^{\text{pinch}}$, into D_i and V_i^{pinch} is done outside of the ETS equation solver in **Transport combiner**)

Sources in equation (2.3) include contributions in generic form from NBI, recycling and puffed neutrals, ripple, ergodization, and other possible sources, in the form of explicit and implicit terms:

$$S_{i,exp} = \sum_{source=1}^{nsource} m_{source} \cdot S_{i,source1} \quad (2.5)$$

$$S_{i,imp} = \sum_{source=1}^{nsource} m_{source} \cdot S_{i,source2} \quad (2.6)$$

This summation is done in the workflow by [Source Combiner](#), where individual weights, m_{source} , can be adjusted.

Equation (2.1) can be rewritten for the normalized grid coordinate as:

$$\frac{\partial}{\partial x} (V' n_i) + \frac{\partial}{\partial \rho} \left[V' \left(\langle |\nabla \rho|^2 \rangle \left(-D_i \frac{\partial n_i}{\partial \rho} + n_i V_i^{pinch} \right) - \frac{\Phi_b}{2\Phi_b} \cdot \rho n_i \right) \right] = V' \left[S_{i,exp} - \left(S_{i,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \cdot n_i \right]. \quad (2.7)$$

Boundary Conditions:

NOTE! Specified positive values for $\nabla n_{i,bnd}$ and L_{ni} correspond to “normal” profile with density decreasing towards the edge

On axis, at $\rho = 0$, ETS assumes:

$$\left. \frac{\partial n_i}{\partial \rho} \right|_{x=0} = 0 \quad \text{and} \quad V_i^{pinch} = 0 \quad (2.8)$$

At the edge, at $\rho = \rho_{bnd}$, following options (distinguished by boundary condition type) are available:

type=1 (value)

$$n_i \Big|_{x=1} = n_{i,bnd} \quad (2.9)$$

type=2 (gradient)

$$\left. \frac{\partial n_i}{\partial \rho} \right|_{x=1} = -\nabla n_{i,bnd} \quad (2.10)$$

type=3 (scale length)

$$\left. \frac{1}{(\partial \ln n_i / \partial \rho)} \right|_{x=1} = -L_{ni} \quad (2.11)$$

type=4 (flux)

$$V' \left\langle |\nabla \rho|^2 \right\rangle \left(-D_i \frac{\partial n_i}{\partial \rho} + n_i V_i^{pinch} \right) \Big|_{x=1} = \Gamma_{i,bnd} \quad (2.12)$$

type=5 (generic)

$$v_{gen} \frac{\partial n_i}{\partial \rho} \Big|_{x=1} + u_{gen} n_i \Big|_{x=1} = w_{gen} \quad (2.13)$$

Definitions for quantities derived from ion density equation:

Ion diffusion equation provides the flux contributing to ion heat transport, this is produced based on the input from attached transport models:

$$\gamma_i = \sum_{i\ model=1}^{n\ model} c_{1,i\ model} V' \langle |\nabla \rho|^2 \rangle \left(-D_{i,i\ model} \frac{\partial n_i}{\partial \rho} + n_i V_{i,i\ model}^{pinch} \right), \quad (2.14)$$

where $c_{1,i\ model}$ are the coefficients 1, 3/2 or 5/2 also provided by transport models.

If equation (2.7) is not solved ($NI_BND_TYPE=0$), interpretative density, $n_{i,int}$, should be specified and flux is calculated from the integral of sources:

$$n_i = n_{i,int}; \quad (2.15)$$

$$\Gamma_i = \Gamma_{Si}; \quad (2.16)$$

$$\gamma_i = \frac{3}{2} \Gamma_{Si}; \quad (2.17)$$

where τ is the time step, V'^- and $n_{i,int}^-$ are taken from the previous time step, and

$$\Gamma_{Si} = \frac{\dot{\Phi}_b}{2\Phi_b} \cdot (\rho V' n_{i,int}) + \rho_b \int_0^x \left(V' S_{i,exp} + \frac{V'^- n_{i,int}^-}{\tau} - n_{i,int} V' \left(\frac{1}{\tau} + S_{i,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \right) dx; \quad (2.18)$$

Translation to generalized form for interface with numerical solver:

Rewrite the equation (2.7) using the time discretization: $\frac{\partial V' n_i}{\partial t} = \frac{V' n_i - V'^- n_i^-}{\tau}$, where V'^- and n_i^- are taken at the previous time and τ is the time step:

$$\begin{aligned} \frac{V' n_i - V'^- n_i^-}{\tau} + \frac{1}{\rho_b} \frac{\partial}{\partial x} \left[V' \left(\langle |\nabla \rho|^2 \rangle \left(-\frac{D_i}{\rho_b} \frac{\partial n_i}{\partial x} + n_i V_i^{pinch} \right) - \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \rho n_i \right) \right] = \\ = V' S_{i,exp} - \left(S_{i,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \cdot V' n_i \end{aligned} \quad (2.19)$$

Then, using the equation (0.8), coefficients for the interface with numerical solver are defined as:

$$\begin{aligned}
 a(x) &= V' \\
 b(x) &= V'^{-} \\
 c(x) &= \rho_b \\
 d(x) &= V' \left\langle |\nabla \rho|^2 \right\rangle \frac{D_i}{\rho_b} \\
 e(x) &= V' \left\langle |\nabla \rho|^2 \right\rangle V_i^{pinch} - \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \rho V' \\
 f(x) &= V' S_{i,exp} \\
 g(x) &= V' \left(S_{i,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \\
 h &= \tau \\
 Y^{i-1}(\rho) &= n_i^-
 \end{aligned} \tag{2.20}$$

boundary conditions:

at the axis, $x = 0$:

the choice of the boundary condition is connected with the choice on numerical solver, two options are available:

$$\left. \frac{\partial n_i}{\partial \rho} \right|_{x=0} = 0 \tag{2.21}$$

$$v(1) = 1; \quad u(1) = 0; \quad w(1) = 0$$

Or:

$$\Gamma_i \big|_{x=0} = 0$$

$$v(1) = -\frac{D_i \big|_{x=0}}{\rho_b}; \quad u(1) = V_i^{pinch} \big|_{x=0}; \quad w(1) = 0$$

at the edge, $x = 1$:

type=1 (value)

$$n_i \big|_{x=1} = n_{i,bnd}$$

$$u(2) = 0; \quad v(2) = 1; \quad w(2) = n_{i,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial n_i}{\partial x} \right|_{x=1} = -\nabla n_{i,bnd} \cdot \rho_b$$

$$v(2)=1; \quad u(2)=0; \quad w(2)=-\nabla n_{i,bnd} \cdot \rho_b$$

type=3 (scale lenght)

$$\left. \frac{\rho_b}{(\partial \ln n_i / \partial x)} \right|_{x=1} = -L_{ni}$$

$$v(2)=1; \quad u(2)=\rho_b / L_{ni}; \quad w(2)=0$$

type=4 (flux)

$$V' \left\langle |\nabla \rho|^2 \right\rangle \left(-\frac{D_i}{\rho_b} \frac{\partial n_i}{\partial x} + n_i V_i^{pinch} \right) \Big|_{x=1} = \Gamma_{i,bnd}$$

$$v(2) = -V' \left\langle |\nabla \rho|^2 \right\rangle \frac{D_i}{\rho_b}; \quad u(2) = V' \left\langle |\nabla \rho|^2 \right\rangle V_i^{pinch}; \quad w(2) = \Gamma_{i,bnd}$$

type=5 (generic)

$$\frac{v_{gen}}{\rho_b} \left. \frac{\partial n_i}{\partial x} \right|_{bnd} + u_{gen} n_{i,bnd} = w_{gen}$$

$$v(2) = \frac{v_{gen}}{\rho_b}; \quad u(2) = u_{gen}; \quad w(2) = w_{gen}$$

Variables used inside fortran routines:

(Routine: MAIN_PLASMA, Subroutine: ION_DENSITY)

Variable	TYPE%NAME used in ETS data flow	Internal name used in ION_DENSITY routine	Units
τ	EVOLUTION%TAU	TAU	s
\dot{B}_0	GEOMETRY%BTPRIME	BTPRIME	T/s
B_0	GEOMETRY%BT	BT	T
ρ	GEOMETRY%RHO	RHO	m
V'	GEOMETRY%VPR	VPR	m ²

V^r	EVOLUTION%VPRM	VPRM	m^2
$\langle \nabla \rho ^2 \rangle$	GEOMETRY%G1	G1	--
n_i	PROFILES%NI	NI	m^{-3}
n_i^-	EVOLUTION%NI	NIM	m^{-3}
Γ_i	PROFILES%FLUX_NI	FLUX	1/s
γ_i	PROFILES%FLUX_NI_CONV	FLUX_NI_CONV	1/s
Γ_{Si}	PROFILES%INT_SOURCE_NI	INT_SOURCE	1/s
$S_{i,exp}$	-	SI_EXP	$m^{-3}s^{-1}$
$S_{i,imp}$	-	SI_IMP	1/s
$S_{i,source,1}$	SOURCES%SI_EXP	-	$m^{-3}s^{-1}$
$S_{i,source,2}$	SOURCES%SI_IMP	-	1/s
D_i	--	DIFF	m^2/s
V_i^{pinch}	--	VCONV	m/s
$D_{i,model}$	TRANSPORT%DIFF_NI	DIFF_MOD	m^2/s
$V_{i,model}^{pinch}$	TRANSPORT%VCONV_NI	VCONV_MOD	m/s
$c_{1,model}$	TRANSPORT%C1	C1	--
a	SOLVER%A	A	--
b	SOLVER%B	B	--
c	SOLVER%C	C	--
d	SOLVER%D	D	--
e	SOLVER%E	E	--
f	SOLVER%F	F	--
g	SOLVER%G	G	--
h	SOLVER%H	H	--
$v(1:2)$	SOLVER%V	V	--
$u(1:2)$	SOLVER%U	U	--
$w(1:2)$	SOLVER%W	W	--
Y^{r-1}	SOLVER%YM	W	--
$solution$	SOLVER%Y	Y	--
$derivative\ of\ solution$	SOLVER%DY	DY	--

Functions		Internal name used in ION_DENSITY routine	
$\frac{1}{\rho} \left(V' S_{i,\text{exp}} + \frac{V'^{-} n_{i,\text{interpretative}}^{-}}{\tau} - n_{i,\text{interpretative}} V' \left(\frac{1}{\tau} + S_{i,\text{imp}} \right) \right)$		FUN1	
$\int_0^{\rho} \left(V' S_{i,\text{exp}} + \frac{V'^{-} n_{i,\text{interpretative}}^{-}}{\tau} - n_{i,\text{interpretative}} V' \left(\frac{1}{\tau} + S_{i,\text{imp}} \right) \right) \partial \rho$		INTFUN1	

Electron density

Variables and units:

n_e - electron density [m^{-3}]

Γ_e - electron flux [s^{-1}]

γ_e - electron flux contributing to heat transport [A/m^2]

Γ_{Se} - integral source (approximation for the flux if equation is not solved) [s^{-1}]

$n_{e,\text{int}}$ - interpretative electron density [m^{-3}], required if equation is not solved

D_e - electron diffuse electron coefficient [m^2/s]

V_e^{pinch} - electron pinch velocity [m/s]

$S_{e,\text{exp}}$ - explicit part of electron source density [$\text{s}^{-1} \text{m}^{-3}$]

$S_{e,\text{imp}}$ - implicit part of electron source density [s^{-1}];

Equation:

$$\left(\frac{\partial}{\partial t} \bigg|_{\rho} - \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial}{\partial \rho} \rho \right) (V' n_e) + \frac{\partial}{\partial \rho} \Gamma_e = V' (S_{e,\text{exp}} - S_{e,\text{imp}} \cdot n_e) \quad (3.1)$$

where the total flux defined as:

$$\Gamma_e = V' \left\langle |\nabla \rho|^2 \right\rangle \left(-D_e \frac{\partial n_e}{\partial \rho} + n_e V_e^{\text{pinch}} \right); \quad (3.2)$$

where total transport coefficients, D_e and V_e^{pinch} , are defined as a sum of individual contributions from different transport models weighted with coefficients m_{model} :

$$D_e = \sum_{i \text{ model}=1}^{n \text{ model}} m_{\text{model}} \cdot D_{e,i \text{ model}}; \quad (3.3)$$

$$V_e^{\text{pinch}} = \sum_{i \text{ model}=1}^{n \text{ model}} m_{\text{model}} \cdot V_{e,i \text{ model}}^{\text{pinch}}. \quad (3.4)$$

(Combination of individual contributions, $D_{e,i \text{ model}}$ and $V_{e,i \text{ model}}^{\text{pinch}}$, into D_e and V_e^{pinch} is done outside of the ETS equation solver in **Transport combiner**)

Sources in equation (3.3) include contributions in generic form from NBI, recycling and puffed neutrals and other possible sources, in the form of explicit and implicit terms:

$$S_{e,exp} = \sum_{isource=1}^{nsource} m_{source} \cdot S_{e,isource,exp} \quad (3.5)$$

$$S_{e,imp} = \sum_{isource=1}^{nsource} m_{source} \cdot S_{e,isource,imp} \quad (3.6)$$

This summation is done in the workflow by [Source Combiner](#), where individual weights, m_{source} , can be adjusted.

Equation (3.1) can be rewritten for the normalized grid coordinate as:

$$\frac{\partial}{\partial x} (V' n_e) + \frac{\partial}{\partial \rho} \left[V' \left(\langle |\nabla \rho|^2 \rangle \left(-D_e \frac{\partial n_e}{\partial \rho} + n_e V_e^{pinch} \right) - \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \rho n_e \right) \right] = V' \left[S_{e,exp} - \left(S_{e,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \cdot n_e \right]. \quad (3.7)$$

Boundary Conditions:

NOTE! Specified positive values for $\nabla n_{e,bnd}$ and L_{ne} correspond to “normal” profile with density decreasing towards the edge

On axis, at $\rho = 0$, ETS assumes:

$$\left. \frac{\partial n_e}{\partial \rho} \right|_{x=0} = 0 \quad \text{and} \quad V_e^{pinch} = 0 \quad (3.8)$$

At the edge, at $\rho = \rho_{bnd}$, following options (distinguished by boundary condition type) are available:

type=1 (value)

$$n_e|_{x=1} = n_{e,bnd} \quad (3.9)$$

type=2 (gradient)

$$\left. \frac{\partial n_e}{\partial \rho} \right|_{x=1} = -\nabla n_{e,bnd} \quad (3.10)$$

type=3 (scale lenght)

$$\left. \frac{1}{(\partial \ln n_e / \partial \rho)} \right|_{x=1} = -L_{ne} \quad (3.11)$$

type=4 (flux)

$$V' \left\langle |\nabla \rho|^2 \right\rangle \left(-D_e \frac{\partial n_e}{\partial \rho} + n_e V_e^{pinch} \right) \Big|_{x=1} = \Gamma_{e,bnd} \quad (3.12)$$

type=5 (generic)

$$v_{gen} \frac{\partial n_e}{\partial \rho} \Big|_{x=1} + u_{gen} n_e = w_{gen} \quad (3.13)$$

Definitions for quantities derived from electron density equation:

If equation (3.7) is not solved, interpretative density, $n_{e,int}$, should be specified and flux is calculated from the integral of sources:

$$n_e = n_{e,int}; \quad (3.14)$$

$$\Gamma_e = \Gamma_{Se}; \quad (3.15)$$

$$\gamma_e = \frac{3}{2} \Gamma_{Se}; \quad (3.16)$$

where τ is the time step, V'^- and $n_{e,int}^-$ are taken from the previous time step, and

$$\Gamma_{Se} = \frac{\dot{\Phi}_b}{2\Phi_b} \cdot (\rho V' n_{e,int}) + \rho_b \int_0^x \left(V' S_{e,exp} + \frac{V'^- n_{e,int}^-}{\tau} - n_{e,int} V' \left(\frac{1}{\tau} + S_{e,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \right) \partial x; \quad (3.17)$$

Translation to generalized form for interface with numerical solver:

Rewrite the equation (3.7) using the time discretization: $\frac{\partial V' n_e}{\partial t} = \frac{V' n_e - V'^- n_e^-}{\tau}$, where V'^- and n_e^- are taken at the previous time and τ is the time step:

$$\begin{aligned} \frac{V' n_e - V'^- n_e^-}{\tau} + \frac{1}{\rho_b} \frac{\partial}{\partial x} \left[V' \left(\langle |\nabla \rho|^2 \rangle \left(-\frac{D_e}{\rho_b} \frac{\partial n_e}{\partial x} + n_e V_e^{pinch} \right) - \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \rho n_e \right) \right] = \\ = V' S_{e,exp} - \left(S_{e,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \cdot V' n_e \end{aligned} \quad (3.18)$$

Then, using the equation (0.8), coefficients for the interface with numerical solver are defined as:

$$a(x) = V'$$

$$b(x) = V'^-$$

$$c(x) = \rho_b$$

$$\begin{aligned}
d(x) &= V' \langle |\nabla \rho|^2 \rangle \frac{D_e}{\rho_b} \\
e(x) &= V' \langle |\nabla \rho|^2 \rangle V_e^{pinch} - \frac{\dot{\Phi}_b}{2\Phi_b} \cdot \rho V' \\
f(x) &= V' S_{e,exp} \\
g(x) &= V' \left(S_{e,imp} + \frac{\dot{\rho}_b}{\rho_b} \right) \\
h &= \tau \\
Y^{i-1}(\rho) &= n_e^-
\end{aligned} \tag{3.19}$$

boundary conditions:

at the axis, $x = 0$:

the choice of the boundary condition is connected with the choice on numerical solver, two options are available:

$$\left. \frac{\partial n_e}{\partial x} \right|_{x=0} = 0 \tag{3.20}$$

$$v(1) = 1; \quad u(1) = 0; \quad w(1) = 0$$

Or:

$$\Gamma_e|_{x=0} = 0$$

$$v(1) = -\frac{D_e}{\rho_b} \Big|_{x=0}; \quad u(1) = V_e^{pinch} \Big|_{x=0}; \quad w(1) = 0$$

at the edge, $x = 1$:

type=1 (value)

$$n_e|_{x=1} = n_{e,bnd}$$

$$u(2) = 0; \quad v(2) = 1; \quad w(2) = n_{e,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial n_e}{\partial x} \right|_{x=1} = -\nabla n_{e,bnd} \cdot \rho_b$$

$$v(2) = 1; \quad u(2) = 0; \quad w(2) = -\nabla n_{e,bnd} \cdot \rho_b$$

type=3 (scale lenght)

$$\left. \frac{\rho_b}{(\partial \ln n_e / \partial x)} \right|_{x=1} = -L_{ne}$$

$$v(2)=1; \quad u(2)=\rho_b / L_{ne}; \quad w(2)=0$$

type=4 (flux)

$$V' \left\langle |\nabla \rho|^2 \right\rangle \left(-\frac{D_e}{\rho_b} \frac{\partial n_e}{\partial x} + n_e V_e^{pinch} \right) \bigg|_{x=1} = \Gamma_{e,bnd}$$

$$v(2) = -V' \left\langle |\nabla \rho|^2 \right\rangle \frac{D_e}{\rho_b}; \quad u(2) = V' \left\langle |\nabla \rho|^2 \right\rangle V_e^{pinch}; \quad w(2) = \Gamma_{e,bnd}$$

type=5 (generic)

$$\frac{v_{gen}}{\rho_b} \frac{\partial n_e}{\partial x} \bigg|_{bnd} + u_{gen} n_{e,bnd} = w_{gen}$$

$$v(2) = \frac{v_{gen}}{\rho_b}; \quad u(2) = u_{gen}; \quad w(2) = w_{gen}$$

Variables used inside fortran routines:

(Routine: MAIN_PLASMA, Subroutine: ELECTRON_DENSITY)

Variable	TYPE%NAME used in ETS data flow	Internal name used in ELECTRON_DENSITY routine	Units
τ	EVOLUTELECTRON%TAU	TAU	s
\dot{B}_0	GEOMETRY%BTPRIME	BTPRIME	T/s
B_0	GEOMETRY%BT	BT	T
ρ	GEOMETRY%RHO	RHO	m
V'	GEOMETRY%VPR	VPR	m ²
V'^{-}	EVOLUTELECTRON%VPRM	VPRM	m ²
$\left\langle \nabla \rho ^2 \right\rangle$	GEOMETRY%G1	G1	--
n_i	PROFILES%NI	NI	m ⁻³
n_i^{-}	EVOLUTELECTRON%NI	NIM	m ⁻³
Γ_i	PROFILES%FLUX_NI	FLUX	1/s

γ_i	PROFILES%FLUX_NI_CONV	FLUX_NI_CONV	1/s
Γ_{Si}	PROFILES%INT_SOURCE_NI	INT_SOURCE	1/s
$S_{i,exp}$	-	SI_EXP	$m^{-3}s^{-1}$
$S_{i,imp}$	-	SI_IMP	1/s
$S_{i,source,1}$	SOURCES%SI_EXP	-	$m^{-3}s^{-1}$
$S_{i,source,2}$	SOURCES%SI_IMP	-	1/s
D_i	--	DIFF	m^2/s
V_i^{pinch}	--	VCONV	m/s
$D_{i,model}$	TRANSPORT%DIFF_NI	DIFF_MOD	m^2/s
$V_{i,model}^{pinch}$	TRANSPORT%VCONV_NI	VCONV_MOD	m/s
$C_{1,model}$	TRANSPORT%C1	C1	--
a	SOLVER%A	A	--
b	SOLVER%B	B	--
c	SOLVER%C	C	--
d	SOLVER%D	D	--
e	SOLVER%E	E	--
f	SOLVER%F	F	--
g	SOLVER%G	G	--
h	SOLVER%H	H	--
$v(1:2)$	SOLVER%V	V	--
$u(1:2)$	SOLVER%U	U	--
$w(1:2)$	SOLVER%W	W	--
Y^{t-1}	SOLVER%YM	W	--
$solution$	SOLVER%Y	Y	--
$derivative\ of\ solution$	SOLVER%DY	DY	--
Functions		Internal name used in ELECTRON_DENSITY routine	
$\frac{1}{\rho} \left(V' S_{i,exp} + \frac{V'^{-} n_{i,interpretative}^{-}}{\tau} - n_{i,interpretative} V' \cdot \left(\frac{1}{\tau} + S_{i,imp} \right) \right)$		FUN1	

$\int_0^{\rho} \left(V' S_{i,exp} + \frac{V' n_{i,interpretative}}{\tau} - n_{i,interpretative} V' \left(\frac{1}{\tau} + S_{i,imp} \right) \right) \partial \rho$	INTFUN1
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Quasi-neutrality condition

Variables and units:

- n_i - ion density [m^{-3}]
- Γ_i - ion flux [s^{-1}]
- γ_i - ion flux contributing to heat transport [A/m^2]
- Z_{ion} - ion charge [-]
- n_e - electron density [m^{-3}]
- Γ_e - electron flux [s^{-1}]
- γ_e - electron flux contributing to heat transport [A/m^2]
- n_{imp} - impurity density [m^{-3}]
- Γ_{imp} - impurity flux [s^{-1}]
- Z_{imp} - impurity charge [-]
- Z_{eff} - plasma effective charge [-]

Equations:

Plasma quasi-neutrality can be described by following conditions:

$$n_e = \sum_{ion} Z_{ion} \cdot n_{ion} + \sum_{imp} Z_{imp} \cdot n_{imp} , \quad (4.1)$$

$$\Gamma_e = \sum_{ion} Z_{ion} \cdot \Gamma_{ion} + \sum_{imp} Z_{imp} \cdot \Gamma_{imp} , \quad (4.2)$$

where indexes e , ion and imp refer to electrons, ions and impurity ions retrospectively. Impurity contribution is considered by standalone *Impurity module* and provided to the main transport solver as input.

Definitions for quantities derived from quasi-neutrality conditions:

Quasi-neutrality conditions in ETS can be used to calculate any particle plasma component (electrons or ions) or several of them (several ions).

Electrons computed from quasi-neutrality:

Density and the flux are obtained from equations (4.1) and (4.2), plasma effective charge is computed as:

$$Z_{eff} = \frac{\sum_{ion} Z_{ion}^2 \cdot n_{ion} + \sum_{imp} Z_{imp}^2 \cdot n_{imp}}{n_e} \quad (4.4)$$

Electron flux contributing to heat transport is obtained as:

$$\gamma_e = \sum_{ion} Z_{ion} \cdot \gamma_{ion} \quad (4.5)$$

Ins computed from quasi-neutrality:

Density and the flux are obtained from:

$$(n \cdot Z)_{sum} = n_e - \sum_{ion} Z_{ion} \cdot n_{ion} - \sum_{imp} Z_{imp} \cdot n_{imp} , \quad (4.6)$$

$$(\Gamma \cdot Z)_{sum} = \Gamma_e - \sum_{ion} Z_{ion} \cdot \Gamma_{ion} - \sum_{imp} Z_{imp} \cdot \Gamma_{imp} , \quad (4.7)$$

Here index *ion* refers to ion density and fluxes obtained from solving the diffusion equation (2.7), index *sum* refers to the sum of ions for which the predictive equation is not solved. Individual contributions for these ions are obtained even from:

$$n_{ion} = \frac{(n \cdot Z)_{sum}}{N_{ion}} \cdot \frac{1}{Z_{ion}} , \quad (4.8)$$

$$\Gamma_{ion} = \frac{(\Gamma \cdot Z)_{sum}}{N_{ion}} \cdot \frac{1}{Z_{ion}} , \quad (4.9)$$

where N_{ion} is the number of ions to be predicted from the quasi neutrality condition.

Or from defining the weights, w_{ion} , between contributions from different ions, then individual contributions are obtained as:

$$n_{ion} = \frac{(n \cdot Z)_{sum}}{\sum_{ion} w_{ion}} \cdot w_{ion} , \quad (4.10)$$

$$\Gamma_{ion} = \frac{(\Gamma \cdot Z)_{sum}}{\sum_{ion} w_{ion}} \cdot w_{ion} . \quad (4.11)$$

Variables used inside fortran routines:

(Routine: MAIN_PLASMA, Subroutine: QUASI_NEUTRALITY)

Variable	TYPE%NAME used in ETS data flow	Internal name used in QUASI_NEUTRALITY routine	Units
ρ	GEOMETRY%RHO	RHO	m
n_e	PROFILES%NE	NE	m ⁻³
Γ_e	PROFILES%FLUX_NE	FLUX_NE	1/s
γ_e	PROFILES%FLUX_NE_CONV	FLUX_NE_CONV	1/s
Z_{eff}	PROFILES%ZEFF	ZEFF	--
Z_{ion}	PROFILES%ZION	ZION	--
Z_{ion}^2	PROFILES%ZION2	ZION2	--
n_{ion}	PROFILES%NI	NI	m ⁻³
Γ_{ion}	PROFILES%FLUX_NI	FLUX_NI	1/s
γ_{ion}	PROFILES%FLUX_NI_CONV	FLUX_NI_CONV	1/s
Z_{imp}	IMPURITY%ZIMP	ZIMP	--
Z_{imp}^2	IMPURITY%ZIMP2	ZIMP2	--
n_{imp}	IMPURITY%NZ	NZ	m ⁻³
Γ_{imp}	IMPURITY %FLUX_NZ	FLUX_NZ	1/s

Ion temperatures

Variables and units:

T_i - ion temperature [eV]

q_i - ion conductive heat flux [W]

$T_i \cdot \gamma_i$ - ion convective heat flux [W]

H_i - total ion heat flux [W]

$H_{i,int}$ - total ion heat flux calculated as the integral of sources [W], if equation is not solved

χ_i - ion heat conductivity [m²/s]

V_{Ti}^{pinch} - ion heat pinch velocity [m/s]

$Q_{i,exp}$ - explicit part of ion heat source density [eV·m⁻³s⁻¹]

$Q_{i,imp}$ - implicit part of ion heat source density [m⁻³s⁻¹];

Q_{ei} - electron-ion exchange frequency [m⁻³s⁻¹];

Q_{zi} - impurity-ion exchange frequency [m⁻³s⁻¹];

$Q_{\gamma i}$ - flow exchange terms [m⁻³s⁻¹];

Equation:

$$\frac{3}{2} \left(\frac{\partial}{\partial t} - \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial}{\partial \rho} \right) \rho \left(n_i T_i V'^{\frac{5}{3}} \right) + V'^{\frac{2}{3}} \frac{\partial}{\partial \rho} (q_i + T_i \gamma_i) = V'^{\frac{5}{3}} [Q_{i,exp} - Q_{i,imp} \cdot T_i + Q_{ei} + Q_{zi} + Q_{\gamma i}] \quad (5.1)$$

where conductive and convective heat flux are defined as:

$$q_i = V' \left\langle |\nabla \rho|^2 \right\rangle \left[n_i \left(-\chi_i \frac{\partial T_i}{\partial \rho} + T_i V_{Ti}^{pinch} \right) \right] \quad \text{and} \quad T_i \gamma_i, \quad (5.2)$$

γ_i is defined by equation (2.14)

Total heat flux:

$$H_i = q_i + T_i \gamma_i \quad (5.3)$$

Total transport coefficients are defined as a sum of individual contributions from different transport models weighted with coefficients m_{model} :

$$\chi_i = \sum_{i model=1}^{n model} m_{model} \cdot \chi_{i,i model} \quad (5.4)$$

$$V_{Ti}^{pinch} = \sum_{i \text{ model}=1}^{n \text{ model}} m_{\text{model}} \cdot V_{Ti, i \text{ model}}^{pinch} \quad (5.5)$$

(Combination of individual contributions, $\chi_{i, i \text{ model}}$ and $V_{Ti, i \text{ model}}^{pinch}$, into χ_i and V_{Ti}^{pinch} is done outside of the ETS equation solver in **Transport combiner**)

Right hand side of equation (5.1) includes contributions in generic form from external sources (computed by other modules, like ICRH, ECRH, NBI, neutrals and etc.), in the form of explicit and implicit terms:

$$Q_{i, \text{exp}} = \sum_{i \text{ source}=1}^{n \text{ source}} m_{\text{source}} \cdot Q_{i, i \text{ source exp}} \quad (5.6)$$

$$Q_{i, \text{imp}} = \sum_{i \text{ source}=1}^{n \text{ source}} m_{\text{source}} \cdot Q_{i, i \text{ source imp}} \quad (5.7)$$

This summation is done in the workflow by **Source Combiner**, where individual weights, m_{source} , can be adjusted.

Number of source terms are computed by the ETS internally, this includes exchange terms with electrons, $Q_{ei} = \nu_{ei} (T_e - T_i)$, (5.8)

and impurity ions,

$$Q_{zi} = \sum_z \nu_{zi} (T_z - T_i) = \sum_z \nu_{zi} T_z - T_i \cdot \sum_z \nu_{zi}, \quad (5.9)$$

where ν_{ei} and $\sum_z \nu_{zi}$ are the exchange frequency.

The flow exchange terms are imported from the attached transport models:

$$Q_{ji} = \sum_{i \text{ model}=1}^{n \text{ model}} Q_{ji, i \text{ model}} \cdot \quad (5.10)$$

Making use of equations (0.6-0.7), equation (5.1) can be rewritten for the normalized coordinate x :

$$\begin{aligned} \frac{3}{2} \frac{\partial}{\partial x} \left[n_i T_i V^{5/3} + V^{2/3} \frac{\partial}{\partial \rho} \left[V \langle |\nabla \rho|^2 \rangle \left(n_i \left(-\chi_i \frac{\partial T_i}{\partial \rho} + T_i V_{Ti}^{pinch} \right) + T_i \gamma_i - V \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_i T_i \right) \right] \right] = \\ = V^{5/3} \left[Q_{i, \text{exp}} - Q_{i, \text{imp}} \cdot T_i + Q_{ei} + Q_{zi} + Q_{ji} - \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_i T_i \right] \end{aligned} \quad (5.11)$$

Boundary Conditions:

NOTE! Specified positive values for $\nabla T_{i, \text{bnd}}$ and L_{Ti} correspond to “normal” profile with density decreasing towards the edge

on axis, $x = 0$:

$$\left. \frac{\partial T_i}{\partial \rho} \right|_{x=0} = 0, \text{ assuming } V_{Ti}^{pinch} = 0 \quad (5.12)$$

at the edge, $x = 1$

type=1 (value)

$$T_i \Big|_{\rho=\rho_{bnd}} = T_{i,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial T_i}{\partial \rho} \right|_{\rho=\rho_{bnd}} = -\nabla T_{i,bnd}$$

type=3 (scale lenght)

$$\left. \frac{1}{(\partial \ln T_i / \partial \rho)} \right|_{\rho=\rho_{bnd}} = -L_{Ti}$$

type=4 (flux)

$$(q_i + T_i \gamma_i) \Big|_{\rho=\rho_{bnd}} = f_{Ti,bnd}$$

type=5 (generic)

$$v_{gen} \left. \frac{\partial T_i}{\partial \rho} \right|_{bnd} + u_{gen} T_{i,bnd} = w_{gen}$$

Definitions for quantities derived from ion temperature equation:

If equation (5.11) is not solved, interpretative temperature should be specified and various flux components are calculated from the integral of sources:

$$T_i = T_{i,interpretative} ; \quad (5.13)$$

$$T_i \gamma_i = \gamma_i \cdot T_{i,interpretative} \quad (5.14)$$

$$H_i = H_{i,int} \quad (5.15)$$

$$q_i = H_{i,int} - \gamma_i \cdot T_{i,interpretative} \quad (5.16)$$

where:

$$\begin{aligned}
H_{i,int} &= \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_i V' T_{i,interpretative} \\
&+ \rho_b \cdot \int_0^\rho V' \left[\frac{3}{2} \frac{n_i^- T_{i,interpretative}^-}{\tau} \left(\frac{V'^-}{V'} \right)^{\frac{5}{3}} + Q_{i,exp} + Q_{\gamma i} + v_{ei} T_e + \sum_z v_{zi} T_z \right] \partial x \\
&- \rho_b \cdot \int_0^\rho V' \left[\frac{3}{2} \frac{n_i}{\tau} + Q_{i,imp} + v_{ei} + \sum_z v_{zi} + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_i \right] \cdot T_{i,interpretative} \cdot \partial x
\end{aligned} \tag{5.17}$$

Translation to generalized form for interface with numerical solver:

Equation (5.11) can be rewritten using time discretization, $\frac{\partial}{\partial t} n_i T_i V'^{\frac{5}{3}} = \frac{n_i T_i V'^{\frac{5}{3}} - n_i^- T_i^- V'^{-\frac{5}{3}}}{\tau}$, as

$$\begin{aligned}
&\frac{3}{2} \frac{n_i T_i V'^{\frac{5}{3}} - n_i^- T_i^- V'^{-\frac{5}{3}}}{\tau} + V'^{\frac{2}{3}} \frac{\partial}{\partial \rho} \left[V' \langle |\nabla \rho|^2 \rangle \left(n_i \left(-\chi_i \frac{\partial T_i}{\partial \rho} + T_i V_{Ti}^{pinch} \right) \right) + T_i \gamma_i - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_i T_i \right] = \\
&= V'^{\frac{5}{3}} \left[Q_{i,exp} - Q_{i,imp} \cdot T_i + Q_{ei} + Q_{zi} + Q_{\gamma i} - \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_i T_i \right]
\end{aligned} \tag{5.18}$$

where τ is the time step, T_i^- , n_i^- and V'^- are taken from the previous time step.

To make it compatible with the form defined by the equation (0.8), it can be:

$$\begin{aligned}
&\frac{\left(\frac{3}{2} n_i V' \right) \cdot T_i - \left(\frac{3}{2} n_i^- \frac{(V'^-)^{\frac{5}{3}}}{V'^{\frac{2}{3}}} \right) \cdot T_i^-}{\tau} \\
&+ \frac{1}{\rho_b} \frac{\partial}{\partial x} \left[-V' \langle |\nabla \rho|^2 \rangle \frac{\chi_i n_i}{\rho_b} \cdot \frac{\partial T_i}{\partial x} + \left(V' \langle |\nabla \rho|^2 \rangle n_i V_{Ti}^{pinch} + \gamma_i - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_i \right) \cdot T_i \right] \\
&= V' \left[Q_{i,exp_i} + Q_{ei} + Q_{zi} + Q_{\gamma i} - \left[Q_{i,imp} + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_i \right] \cdot T_i \right]
\end{aligned} \tag{5.19}$$

Then, taking in account (5.8) and (5.9), coefficients for the interface with numerical solver can be used as:

$$a(x) = \frac{3}{2} V' n_i \tag{5.20}$$

$$b(x) = \frac{3}{2} n_i - \left(\frac{V'^{-\frac{5}{3}}}{V'^{\frac{2}{3}}} \right)$$

$$c(x) = \rho_b$$

$$d(x) = V' \langle |\nabla \rho|^2 \rangle \frac{n_i \chi_i}{\rho_b}$$

$$e(x) = V' \langle |\nabla \rho|^2 \rangle n_i V_{Ti}^{pinch} + \gamma_i - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_i$$

$$f(x) = V' \left[Q_{i,exp} + Q_{\gamma i} + v_{ei} T_e + \sum_z v_{zi} T_z \right]$$

$$g(x) = V' \left[Q_{i,imp} + v_{ei} + \sum_z v_{zi} + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_i \right]$$

$$h = \tau$$

$$Y^{t-1}(\rho) = T_i^-$$

boundary conditions:

at the axis, $x = 0$:

the choice of the boundary condition is connected with the choice on numerical solver, two options are available:

$$\left. \frac{\partial T_i}{\partial x} \right|_{x=0} = 0 \quad (5.21)$$

$$v(1) = 1; \quad u(1) = 0; \quad w(1) = 0$$

at the edge, $x = 1$:

type=1 (value)

$$T_i \big|_{x=1} = T_{i,bnd}$$

$$v(2) = 0; \quad u(2) = 1; \quad w(2) = T_{i,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial T_i}{\partial x} \right|_{x=1} = -\nabla T_{i,bnd} \cdot \rho_b$$

$$v(2) = 1; \quad u(2) = 0; \quad w(2) = -\nabla T_{i,bnd} \cdot \rho_b$$

type=3 (scale length)

$$\left. \frac{\rho_b}{(\partial \ln T_i / \partial x)} \right|_{x=1} = -L_{Ti}$$

$$v(2)=1; \quad u(2)=\rho_b / L_{Ti}; \quad w(2)=0$$

type=4 (flux)

$$(q_i + T_i \gamma_i) \Big|_{x=1} = f_{Ti,bnd}$$

$$v(2) = -\frac{n_i \chi_i V'}{\rho_b} \langle |\nabla \rho|^2 \rangle; \quad u(2) = n_i V_{Ti}^{pinch} V' \langle |\nabla \rho|^2 \rangle + \gamma_i; \quad w(2) = f_{Ti,bnd}$$

type=5 (generic)

$$\frac{v_{gen}}{\rho_b} \frac{\partial T_i}{\partial x} \Big|_{bnd} + u_{gen} T_{i,bnd} = w_{gen}$$

$$v(2) = \frac{v_{gen}}{\rho_b}; \quad u(2) = u_{gen}; \quad w(2) = w_{gen}$$

Variables used inside fortran routines:

Variable	TYPE%NAME used in ETS data flow	Internal name used in ION_TEMPERATURE routine	Units
τ	EVOLUTION%TAU	TAU	s
\dot{B}_0	GEOMETRY%BTPRIME	BTPRIME	T/s
B_0	GEOMETRY%BT	BT	T
ρ	GEOMETRY%RHO	RHO	m
V'	GEOMETRY%VPR	VPR	m ²
V''	EVOLUTION%VPRM	VPRM	m ²
$\frac{\partial V'}{\partial \rho}$	--	DVPR	m
$\langle \nabla \rho ^2 \rangle$	GEOMETRY%G1	G1	--

T_i	PROFILES%TI	TI	eV
T_i^-	EVOLUTION%TI	TIM	eV
n_i	PROFILES%NI	NI	m ⁻³
n_i^-	EVOLUTION%NI	NIM	m ⁻³
γ_i	PROFILES%FLUX_NI_CONV	FLUX_NI	1/s
H_i	PROFILES%FLUX_TI	FLUX_TI	W
$T_i\gamma_i$	PROFILES%FLUX_TI_CONV	FLUX_TI_CONV	W
q_i	PROFILES%FLUX_TI_COND	FLUX_TI_COND	W
$H_{i,int}$	PROFILES%INT_SOURCE_TI	INT_SOURCE	W
$Q_{i,exp}$	--	QI_EXP	eV·m ⁻³ s ⁻¹
$Q_{i,imp}$	--	QI_IMP	m ⁻³ s ⁻¹
$Q_{i,source,1}$	SOURCES%QI_EXP	--	eV·m ⁻³ s ⁻¹
$Q_{i,source,2}$	SOURCES%QI_IMP	--	m ⁻³ s ⁻¹
ν_{ei}	COLLISIONS%VEI	VEI	m ⁻³ s ⁻¹
$\nu_{ei}T_e$	COLLISIONS %QEI	QEI	eV·m ⁻³ s ⁻¹
$\sum_z \nu_{zi}$	COLLISIONS %VZI	VZI	m ⁻³ s ⁻¹
$\sum_z \nu_{zi}T_z$	COLLISIONS %QZI	QZI	eV·m ⁻³ s ⁻¹
$Q_{\gamma i}$	--	QGI	eV·m ⁻³ s ⁻¹
$Q_{\gamma i,model}$	TRANSPORT%QGI	--	eV·m ⁻³ s ⁻¹
χ_i	--	DIFF	m ² /s
V_{Ti}^{pinch}	--	VCONV	m/s
$\chi_{i,model}$	TRANSPORT%DIFF_TI	--	m ² /s
$V_{Ti,model}^{pinch}$	TRANSPORT%VCONV_TI	--	m/s
a	SOLVER%A	A	--
b	SOLVER%B	B	--
c	SOLVER%C	C	--
d	SOLVER%D	D	--
e	SOLVER%E	E	--
f	SOLVER%F	F	--

g	SOLVER%G	G	--
h	SOLVER%H	H	--
$v(1:2)$	SOLVER%V	V	--
$u(1:2)$	SOLVER%U	U	--
$w(1:2)$	SOLVER%W	W	--
Y^{t-1}	SOLVER%YM	W	--
<i>solution</i>	SOLVER%Y	Y	--
<i>derivative of solution</i>	SOLVER%DY	DY	--
Functions		Internal name used in ION_TEMPERATURE routine	
$\frac{V'}{\rho} \left(\frac{3}{2} \frac{n_i^- T_{i,interpretative}}{\tau} \left(\frac{V'^-}{V'} \right)^{\frac{5}{3}} + Q_{i,exp} + Q_{\gamma i} + v_{ei} T_e + \sum_z v_{zi} T_z - \right.$ $\left. - \left[\frac{3}{2} \frac{n_i}{\tau} + Q_{i,imp} + v_{ei} + \sum_z v_{zi} - \rho n_i \cdot \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial V'}{\partial \rho} \right] \cdot T_{i,interpretative} \right)$		FUN1	
$\int_0^\rho V' \left[\frac{3}{2} \frac{n_i^- T_{i,interpretative}}{\tau} \left(\frac{V'^-}{V'} \right)^{\frac{5}{3}} + Q_{i,exp} + Q_{\gamma i} + v_{ei} T_e + \sum_z v_{zi} T_z \right] \partial \rho -$ $- \int_0^\rho V' \left[\frac{3}{2} \frac{n_i}{\tau} + Q_{i,imp} + v_{ei} + \sum_z v_{zi} - \rho n_i \cdot \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial V'}{\partial \rho} \right] \cdot T_{i,interpretative} \cdot \partial \rho$		INTFUN1	

Electron energy transport equation

Variables and units:

- T_e - electron temperature [eV]
 q_e - electron conductive heat flux [W]
 $T_e \cdot \gamma_e$ - electron convective heat flux [W]
 H_e - total electron heat flux [W]
 $H_{e,int}$ - total electron heat flux calculated as the integral of sources [W], if equation is not solved
 χ_e - electron heat conductivity [m²/s]
 V_{Te}^{pinch} - electron heat pinch velocity [m/s]
 $Q_{e,exp}$ - explicit part of electron heat source density [eV·m⁻³s⁻¹]
 $Q_{e,imp}$ - implicit part of electron heat source density [m³s⁻¹];
 Q_{ie} - ion- electron exchange frequency [m⁻³s⁻¹];
 $Q_{\gamma i}$ - flow exchange terms [m⁻³s⁻¹];

Equation:

$$\frac{3}{2} \left(\frac{\partial}{\partial t} - \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial}{\partial \rho} \right) \rho \left(n_e T_e V^{\frac{5}{3}} \right) + V^{\frac{2}{3}} \frac{\partial}{\partial \rho} (q_e + T_e \gamma_e) = V^{\frac{5}{3}} [Q_{e,exp} - Q_{e,imp} \cdot T_e + Q_{ie} - Q_{\gamma i}] \quad (6.1)$$

where conductive and convective heat flux are defined as:

$$q_e = V \left\langle |\nabla \rho|^2 \right\rangle \left[n_e \left(-\chi_e \frac{\partial T_e}{\partial \rho} + T_e V_{Te}^{pinch} \right) \right] \quad \text{and} \quad T_e \gamma_e, \quad (6.2)$$

γ_e is defined by equation (3.14)

Total heat flux:

$$H_e = q_e + T_e \gamma_e \quad (6.3)$$

Total transport coefficients are defined as a sum of individual contributions from different transport models weighted with coefficients m_{model} :

$$\chi_e = \sum_{i=1}^{n_{model}} m_{model} \cdot \chi_{e,i model} \quad (6.4)$$

$$V_{Te}^{pinch} = \sum_{i=1}^{n_{model}} m_{model} \cdot V_{Te,i model}^{pinch} \quad (6.5)$$

(Combination of individual contributions, $\chi_{e,i\text{ model}}$ and $V_{Te,i\text{ model}}^{pinch}$, into χ_e and V_{Te}^{pinch} is done outside of the ETS equation solver in **Transport combiner**)

Right hand side of equation (6.1) includes contributions in generic form from external sources (computed by other modules, like ICRH, ECRH, NBI, neutrals and etc.), in the form of explicit and implicit terms:

$$Q_{e,\text{exp}} = \sum_{i\text{ source}=1}^{n\text{ source}} m_{\text{source}} \cdot Q_{e,i\text{ source},\text{exp}} \quad (6.6)$$

$$Q_{e,\text{imp}} = \sum_{i\text{ source}=1}^{n\text{ source}} m_{\text{source}} \cdot Q_{e,i\text{ source},\text{imp}} \quad (6.7)$$

This summation is done in the workflow by **Source Combiner**, where individual weights, m_{source} , can be adjusted.

Number of source terms are computed by the ETS internally, this includes exchange terms with electrons, $Q_{ie} = \sum_i \nu_{ei} (T_i - T_e)$, (6.8)

where ν_{ei} is the exchange frequency.

The flow exchange terms are imported from the attached transport models:

$$Q_{\gamma i} = \sum_{i\text{ model}=1}^{n\text{ model}} Q_{\gamma i,i\text{ model}} \cdot \quad (6.9)$$

Making use of equations (0.6-0.7), equation (5.1) can be rewritten for the normalized coordinate x :

$$\begin{aligned} \frac{3}{2} \frac{\partial}{\partial x} \left[n_e T_e V^{5/3} + V^{2/3} \frac{\partial}{\partial \rho} \left[V' \langle |\nabla \rho|^2 \rangle \left(n_e \left(-\chi_e \frac{\partial T_i}{\partial \rho} + T_e V_{Te}^{pinch} \right) \right) + T_e \gamma_e - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_e T_e \right] \right] = \\ = V^{1/3} \left[Q_{e,\text{exp}} - Q_{e,\text{imp}} \cdot T_e + Q_{ie} - Q_{\gamma i} - \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_e T_e \right] \end{aligned} \quad (6.10)$$

Boundary Conditions:

NOTE! Specified positive values for $\nabla T_{e,bnd}$ and L_{Te} correspond to “normal” profile with density decreasing towards the edge

on axis, $x = 0$:

$$\left. \frac{\partial T_e}{\partial \rho} \right|_{x=0} = 0, \text{ assuming } V_{Te}^{pinch} = 0 \quad (6.11)$$

at the edge, $x = 1$

type=1 (value)

$$T_e|_{\rho=\rho_{bnd}} = T_{e,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial T_e}{\partial \rho} \right|_{\rho=\rho_{bnd}} = -\nabla T_{e,bnd}$$

type=3 (scale lenght)

$$\left. \frac{1}{(\partial \ln T_e / \partial \rho)} \right|_{\rho=\rho_{bnd}} = -L_{Te}$$

type=4 (flux)

$$(q_e + T_e \gamma_e)|_{\rho=\rho_{bnd}} = H_{Te,bnd}$$

type=5 (generic)

$$v_{gen} \left. \frac{\partial T_e}{\partial \rho} \right|_{bnd} + u_{gen} T_{e,bnd} = w_{gen}$$

Definitions for quantities derived from electron temperature equation:

If equation (6.10) is not solved, interpretative temperature should be specified and various flux components are calculated from the integral of sources:

$$T_e = T_{e,interpretative} ; \quad (6.12)$$

$$T_e \gamma_e = \gamma_e \cdot T_{e,interpretative} \quad (6.13)$$

$$H_e = H_{e,int} \quad (6.14)$$

$$q_e = H_{e,int} - \gamma_e \cdot T_{e,interpretative} \quad (6.15)$$

where:

$$\begin{aligned} H_{e,int} = & \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_e V' \cdot T_{e,interpretative} \\ & + \rho_b \cdot \int_0^\rho V' \left[\frac{3}{2} \frac{n_e^- T_{e,interpretative}^-}{\tau} \left(\frac{V'^-}{V'} \right)^{\frac{5}{3}} + Q_{e,exp} - Q_{\tilde{\eta}} + \sum_i v_{ei} T_i \right] \partial x \\ & - \rho_b \cdot \int_0^\rho V' \left[\frac{3}{2} \frac{n_e}{\tau} + Q_{e,imp} + \sum_i v_{ei} + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_e \right] \cdot T_{e,interpretative} \cdot \partial x \end{aligned} \quad (6.16)$$

Translation to generalized form for interface with numerical solver:

Equation (5.11) can be rewritten using time discretization, $\frac{\partial}{\partial t} n_i T_i V'^{\frac{5}{3}} = \frac{n_i T_i V'^{\frac{5}{3}} - n_i^- T_i^- V'^{-\frac{5}{3}}}{\tau}$, as

$$\begin{aligned} & \frac{3}{2} \frac{n_i T_i V'^{\frac{5}{3}} - n_i^- T_i^- V'^{-\frac{5}{3}}}{\tau} + V'^{\frac{2}{3}} \frac{\partial}{\partial \phi} \left[V' \langle |\nabla \rho|^2 \rangle \left(n_i \left(-\chi_i \frac{\partial T_i}{\partial \phi} + T_i V_{T_i}^{pinch} \right) \right) + T_i \gamma_i - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_i T_i \right] = \\ & = V'^{\frac{5}{3}} \left[Q_{i,exp} - Q_{i,imp} \cdot T_i + Q_{ei} + Q_{zi} + Q_{\gamma i} - \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \phi} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_i T_i \right] \end{aligned} \quad (6.17)$$

where τ is the time step, T_e^- , n_e^- and V'^- are taken from the previous time step.

To make it compatible with the form defined by the equation (0.8), it can be:

$$\begin{aligned} & \frac{\left(\frac{3}{2} n_e V' \right) \cdot T_e - \left(\frac{3}{2} n_e^- \frac{V'^{-\frac{5}{3}}}{V'^{\frac{2}{3}}} \right) \cdot T_e^-}{\tau} \\ & + \frac{1}{\rho_b} \frac{\partial}{\partial x} \left[-V' \langle |\nabla \rho|^2 \rangle \frac{\chi_e n_e}{\rho_b} \cdot \frac{\partial T_e}{\partial x} + \left(V' \langle |\nabla \rho|^2 \rangle n_e V_{Te}^{pinch} + \gamma_e - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_e \right) \cdot T_e \right] \\ & = V' \left[Q_{e,exp_i} + Q_{ie} - Q_{\gamma i} - \left[Q_{e,imp} + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \phi} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_e \right] \cdot T_e \right] \end{aligned} \quad (6.18)$$

Then, taking in account (6.8) and (6.9), coefficients for the interface with numerical solver can be used as:

$$a(x) = \frac{3}{2} V' n_e \quad (6.19)$$

$$b(x) = \frac{3}{2} n_e^- \left(\frac{V'^{-\frac{5}{3}}}{V'^{\frac{2}{3}}} \right)$$

$$c(x) = \rho_b$$

$$d(x) = V' \langle |\nabla \rho|^2 \rangle \frac{n_e \chi_e}{\rho_b}$$

$$e(x) = V' \langle |\nabla \rho|^2 \rangle n_e V_{Te}^{pinch} + \gamma_e - V' \frac{3\dot{\Phi}_b}{4\Phi_b} \cdot \rho n_e$$

$$f(x) = V' \left[Q_{e,\text{exp}} - Q_{\gamma i} + \sum_i v_{ei} T_i \right]$$

$$g(x) = V' \left[Q_{e,\text{imp}} + \sum_i v_{ei} + \left[\frac{3}{2} \frac{\dot{\rho}_b}{\rho_b} - \frac{\dot{\Phi}_b}{2\Phi_b} \left(1 - \rho \frac{\partial}{\partial \rho} \ln \left(F \cdot \left\langle \frac{1}{R^2} \right\rangle \right) \right) \right] \cdot n_e \right]$$

$$h = \tau$$

$$Y^{i-1}(\rho) = T_e^-$$

boundary conditions:

at the axis, $x = 0$:

the choice of the boundary condition is connected with the choice on numerical solver, two options are available:

$$\left. \frac{\partial T_e}{\partial x} \right|_{x=0} = 0 \quad (6.20)$$

$$v(1) = 1; \quad u(1) = 0; \quad w(1) = 0$$

at the edge, $x = 1$:

type=1 (value)

$$T_e|_{x=1} = T_{e,bnd}$$

$$v(2) = 0; \quad u(2) = 1; \quad w(2) = T_{e,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial T_e}{\partial x} \right|_{x=1} = -\nabla T_{e,bnd} \cdot \rho_b$$

$$v(2) = 1; \quad u(2) = 0; \quad w(2) = -\nabla T_{e,bnd} \cdot \rho_b$$

type=3 (scale length)

$$\left. \frac{\rho_b}{(\partial \ln T_e / \partial x)} \right|_{x=1} = -L_{Te}$$

$$v(2) = 1; \quad u(2) = \rho_b / L_{Te}; \quad w(2) = 0$$

type=4 (flux)

$$(q_e + T_e \gamma_i)_{|x=1} = H_{Te,bnd}$$

$$v(2) = -\frac{n_e \chi_e V'}{\rho_b} \langle |\nabla \rho|^2 \rangle; \quad u(2) = n_e V_{Te}^{pinch} V' \langle |\nabla \rho|^2 \rangle + \gamma_e; \quad w(2) = H_{Te,bnd}$$

type=5 (generic)

$$\frac{v_{gen}}{\rho_b} \frac{\partial T_e}{\partial x} \Big|_{bnd} + u_{gen} T_{e,bnd} = w_{gen}$$

$$v(2) = \frac{v_{gen}}{\rho_b}; \quad u(2) = u_{gen}; \quad w(2) = w_{gen}$$

Variables used inside fortran routines:

Variable	TYPE%NAME used in ETS data flow	Internal name used in ELECTRON_TEMPERATURE routine	Units
τ	EVOLUTION%TAU	TAU	s
\dot{B}_0	GEOMETRY%BTPRIME	BTPRIME	T/s
B_0	GEOMETRY%BT	BT	T
ρ	GEOMETRY%RHO	RHO	m
V'	GEOMETRY%VPR	VPR	m ²
V'^{-}	EVOLUTION%VPRM	VPRM	m ²
$\frac{\partial V'}{\partial \rho}$	--	DVPR	m
$\langle \nabla \rho ^2 \rangle$	GEOMETRY%G1	G1	--
T_e	PROFILES%TE	TE	eV
T_e^{-}	EVOLUTION%TE	TEM	eV
n_e	PROFILES%NE	NE	m ⁻³
n_e^{-}	EVOLUTION%NE	NEM	m ⁻³
γ_e	PROFILES%FLUX_NE_CONV	FLUX_NE	1/s
H_e	PROFILES%FLUX_TE	FLUX_TE	W

$T_e \gamma_e$	PROFILES%FLUX_TE_CONV	FLUX_TE_CONV	W
q_e	PROFILES%FLUX_TE_COND	FLUX_TE_COND	W
$H_{e,int}$	PROFILES%INT_SOURCE_TE	INT_SOURCE	W
$Q_{e,exp}$	--	QE_EXP	$\text{eV} \cdot \text{m}^{-3} \text{s}^{-1}$
$Q_{e,imp}$	--	QE_IMP	$\text{m}^{-3} \text{s}^{-1}$
$Q_{e,source1}$	SOURCES%QE_EXP	--	$\text{eV} \cdot \text{m}^{-3} \text{s}^{-1}$
$Q_{e,source2}$	SOURCES%QE_IMP	--	$\text{m}^{-3} \text{s}^{-1}$
$\sum_i \nu_{ei}$	COLLISIONS%VIE	VIE	$\text{m}^{-3} \text{s}^{-1}$
$\sum_i \nu_{ei} T_i$	COLLISIONS %QIE	QIE	$\text{eV} \cdot \text{m}^{-3} \text{s}^{-1}$
$Q_{\gamma i}$	--	QGI	$\text{eV} \cdot \text{m}^{-3} \text{s}^{-1}$
$Q_{\gamma i,model}$	TRANSPORT%QGI	--	$\text{eV} \cdot \text{m}^{-3} \text{s}^{-1}$
χ_e	--	DIFF	m^2/s
V_{Te}^{pinch}	--	VCONV	m/s
$\chi_{e,model}$	TRANSPORT%DIFF_TE	--	m^2/s
$V_{Te,model}^{pinch}$	TRANSPORT%VCONV_TE	--	m/s
a	SOLVER%A	A	--
b	SOLVER%B	B	--
c	SOLVER%C	C	--
d	SOLVER%D	D	--
e	SOLVER%E	E	--
f	SOLVER%F	F	--
g	SOLVER%G	G	--
h	SOLVER%H	H	--
$v(1:2)$	SOLVER%V	V	--
$u(1:2)$	SOLVER%U	U	--
$w(1:2)$	SOLVER%W	W	--
Y^{t-1}	SOLVER%YM	W	--
$solution$	SOLVER%Y	Y	--
$derivative\ of\ solution$	SOLVER%DY	DY	--
Internal name used in			

Functions	ELECTRON_TEMPERATURE routine
$\frac{V}{\rho} \left[\frac{3}{2} \frac{n_e T_{e,interpretative}}{\tau} \left(\frac{V'}{V} \right)^{\frac{5}{3}} + Q_{e,exp} + \sum_i \nu_{ei} T_i - Q_{\gamma i} \right.$ $\left. - \left[-\frac{3}{2} \frac{n_e}{\tau} Q_{e,imp} + \sum_i \nu_{ei} - \rho n_e \cdot \frac{\dot{B}_0}{2B_0} \frac{1}{V'} \cdot \frac{\partial V'}{\partial \rho} \right] \cdot T_{e,interpretative} \right)$	FUN1
$\int_0^\rho V' \left[\frac{3}{2} \frac{n_e T_{e,interpretative}}{\tau} \left(\frac{V'}{V} \right)^{\frac{5}{3}} + Q_{e,exp} + \sum_i \nu_{ei} T_i - Q_{\gamma i} \right] \partial \rho -$ $- \int_0^\rho V' \left[-\frac{3}{2} \frac{n_e}{\tau} Q_{e,imp} + \sum_i \nu_{ei} - \rho n_e \cdot \frac{\dot{B}_0}{2B_0} \frac{1}{V'} \cdot \frac{\partial V'}{\partial \rho} \right] \cdot T_{e,interpretative} \partial \rho$	INTFUN1

Rotation transport equation

Variables and units:

$u_{i,\varphi}$	- ion toroidal rotation velocity [m/s]
$\omega_{i,\varphi}$	- ion angular toroidal velocity [1/s]
M_i	- ion toroidal momentum [kg*m/s]
Φ_i	- ion toroidal momentum flux [kg*m ² /s ²]
$\Phi_{i,conv}$	- convective component of ion toroidal momentum flux [kg*m ² /s ²]
$\Phi_{i,cond}$	- conductive component of ion toroidal momentum flux [kg*m ² /s ²]
M_{tot}	- total ion toroidal momentum [kg*m/s]
Φ_{tot}	- total ion toroidal momentum flux [kg*m ² /s ²]

Equation:

$$\left(\frac{\partial}{\partial \rho} \left| - \frac{\dot{B}_0}{2B_0} \cdot \frac{\partial}{\partial \rho} \rho \right| \right) (V' \langle R \rangle m_i n_i u_{i,\varphi}) + \frac{\partial}{\partial \varphi} \Phi_i = V' (U_{i,\varphi,exp} - U_{i,\varphi,imp} \cdot u_{i,\varphi} + U_{zi,\varphi}) \quad (7.1)$$

where total flux is defined as:

$$\Phi_i = V' \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \cdot \left(-\chi_{u\varphi,i} \frac{\partial u_{i,\varphi}}{\partial \varphi} + u_{i,\varphi} V_{u\varphi,i}^{pinch} \right) + m_i \langle R \rangle u_{i,\varphi} \Gamma_i \quad (7.2)$$

with convective and conductive components:

$$\Phi_{i,conv} = m_i \langle R \rangle u_{i,\varphi} \Gamma_i \quad (7.3)$$

$$\Phi_{i,cond} = V' \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \cdot \left(-\chi_{u\varphi,i} \frac{\partial u_{i,\varphi}}{\partial \varphi} + u_{i,\varphi} V_{u\varphi,i}^{pinch} \right) \quad (7.4)$$

Total transport coefficients are defined as a sum of individual contributions from different transport models weighted with coefficients m_{model} :

$$\chi_{u\varphi,i} = \sum_{i\ model=1}^{n\ model} m_{model} \chi_{u\varphi,i\ model} \quad (7.5)$$

$$V_{u\varphi}^{pinch} = \sum_{i\ model=1}^{n\ model} m_{model} V_{u\varphi,i\ model}^{pinch} \quad (7.6)$$

(Combination of individual contributions, $\chi_{u\varphi,i\text{ model}}$ and $V_{u\varphi,i\text{ model}}^{pinch}$, into $\chi_{u\varphi}$ and $V_{u\varphi}^{pinch}$ is done outside of the ETS equation solver in **Transport combiner**)

Right hand side of equation (7.1) includes contributions in generic form from external sources in the form of explicit and implicit terms:

$$U_{i,\varphi,\text{exp}} = \sum_{i\text{source}=1}^{n\text{source}} m_{\text{source}} U_{i,\varphi,i\text{source},\text{exp}} \quad (7.7)$$

$$U_{i,\varphi,\text{imp}} = \sum_{i\text{source}=1}^{n\text{source}} m_{\text{source}} U_{i,\varphi,i\text{source},\text{imp}} \quad (7.8)$$

This summation is done in the workflow by **Source Combiner**, where individual weights, m_{source} , can be adjusted.

momentum exchange with other ion components:

$$U_{zi,\varphi} = \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} (u_{z,\varphi} - u_{i,\varphi}) \quad (7.9)$$

various collisions quantities ($\langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}}$ and $\langle R \rangle \sum_z \frac{m_{zi} n_z u_{z,\varphi}}{\tau_{zi}}$) should be provided by stand alone

COLLISIONS module

Equation (7.1) can be rewritten in the form:

$$\begin{aligned} & \frac{\partial}{\partial t} \Big|_x (V \langle R \rangle m_i n_i u_{i,\varphi}) \\ & + \frac{\partial}{\partial \rho} \left(V \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \cdot \left(-\chi_{u\varphi,i} \frac{\partial u_{i,\varphi}}{\partial \rho} + u_{i,\varphi} V_{u\varphi,i}^{pinch} \right) + m_i \langle R \rangle u_{i,\varphi} \Gamma_i - \frac{\dot{\Phi}_b}{2\Phi_b} V m_i n_i \langle R \rangle u_{i,\varphi} \rho \right) \\ & = V \left(U_{i,\varphi,\text{exp}} + U_{zi,\varphi} - \left(U_{i,\varphi,\text{imp}} + \frac{\dot{\rho}_b}{\rho_b} m_i n_i \langle R \rangle \right) \cdot u_{i,\varphi} \right) \end{aligned} \quad (7.10)$$

Boundary Conditions:

NOTE! Specified positive values for $\nabla u_{\varphi,bnd}$ and $L_{u\varphi}$ correspond to “normal” profile with density decreasing towards the edge

on axis, $\rho = 0$:

$$\left. \frac{\partial u_{i,\phi}}{\partial \rho} \right|_{\rho=0} = 0 \quad (7.11)$$

at the edge, $\rho = \rho_{bnd}$:

type=1 (value)

$$u_{i,\phi} \Big|_{\rho=\rho_{bnd}} = u_{\phi,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial u_{i,\phi}}{\partial \rho} \right|_{\rho=\rho_{bnd}} = -\nabla u_{\phi,bnd}$$

type=3 (scale lenght)

$$\left. \frac{1}{(\partial \ln u_{i,\phi} / \partial \rho)} \right|_{\rho=\rho_{bnd}} = -L_{u\phi}$$

type=4 (flux)

$$\Phi \Big|_{\rho=\rho_{bnd}} = f_{u\phi,bnd}$$

type=5 (generic)

$$v_{gen} \frac{\partial u_{i,\phi,bnd}}{\partial \rho} \Big|_{bnd} + u_{gen} u_{i,\phi,bnd} = w_{gen}$$

Definitions for quantities derived from rotation equation:

total toroidal momentum:

$$M_i = m_i n_i \langle R \rangle u_{i,\phi} \quad (7.12)$$

angular velocity for plasma component i :

$$\omega_{i,\phi} = \frac{u_{i,\phi}}{\langle R \rangle} \quad (7.13)$$

total momentum and flux are defined as:

$$M_{tot} = \sum_i M_i ; \quad \Phi_{tot} = \sum_i \Phi_i \quad (7.14)$$

If equation (7.10) is not solved, interpretative toroidal velocity should be specified, momentum and flux are assumed:

$$u_{i,\varphi} = u_{i,\varphi,interpretative} ; \quad (7.15)$$

$$\omega_{i,\varphi} = \frac{u_{i,\varphi,interpretative}}{\langle R \rangle} ; \quad (7.16)$$

$$M_i = m_i n_i \langle R \rangle u_{i,\varphi,interpretative} ; \quad (7.17)$$

$$\Phi_i = \Phi_{i,int} \quad (7.18)$$

$$\Phi_{i,conv} = m_i \langle R \rangle \Gamma_i u_{i,\varphi,interpretative} \quad (7.19)$$

$$\Phi_{i,cond} = \Phi_{i,int} - m_i \langle R \rangle \Gamma_i u_{i,\varphi,interpretative} \quad (7.20)$$

$$M_{tot} = \sum_i M_i ; \quad \Phi_{tot} = \sum_i \Phi_i \quad (7.21)$$

where:

$$\begin{aligned} \Phi_{i,int} = & \frac{\dot{\Phi}_b}{2\Phi_b} V' m_i n_i \langle R \rangle \rho_b x u_{i,\varphi,interpretative} + \\ & + \rho_b \int_0^x V' \left(U_{i,\varphi,exp} + \langle R \rangle \sum_z \frac{m_{zi} n_z u_{z,\varphi}}{\tau_{zi}} + \frac{\langle R \rangle^- m_i n_i^- u_{i,\varphi,interpretative}^-}{\tau} \left(\frac{V'^-}{V'} \right) \right) \partial x \\ & - \rho_b \int_0^x V' \left(U_{i,\varphi,imp} - \langle R \rangle \frac{m_i n_i}{\tau} - \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} - \frac{\dot{\rho}_b}{\rho_b} m_i n_i \langle R \rangle \right) u_{i,\varphi,interpretative} \partial x \end{aligned} \quad (7.22)$$

Translation to generalized form for interface with numerical solver:

Equation (6.10) can be rewritten using time discretization,

$$\frac{\partial}{\partial x} \Big|_x \left(V' \langle R \rangle m_i n_i u_{i,\varphi} \right) = \frac{V' \langle R \rangle m_i n_i u_{i,\varphi} - V'^- \langle R \rangle^- m_i n_i^- u_{i,\varphi}^-}{\tau},$$

where V'^- , n_i^- , $\langle R \rangle^-$ and $u_{i,\varphi}^-$ are taken at the previous time step as:

$$\begin{aligned}
& m_i \frac{V' n_i \langle R \rangle u_{i,\varphi} - V'^- n_i^- \langle R \rangle^- u_{i,\varphi}^-}{\tau} \\
& + \frac{\partial}{\partial \rho} \left(V' \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \cdot \left(-\chi_{u\varphi,i} \frac{\partial u_{i,\varphi}}{\partial \rho} + u_{i,\varphi} V_{u\varphi,i}^{pinch} \right) + m_i \langle R \rangle u_{i,\varphi} \Gamma_i - \frac{\dot{\Phi}_b}{2\Phi_b} V' m_i n_i \langle R \rangle u_{i,\varphi} \rho \right) , \quad (7.23) \\
& = V' \left(U_{i,\varphi,\exp} + \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} (u_{z,\varphi} - u_{i,\varphi}) - \left(U_{i,\varphi,imp} + \frac{\dot{\rho}_b}{\rho_b} m_i n_i \langle R \rangle \right) \cdot u_{i,\varphi} \right)
\end{aligned}$$

or as:

$$\begin{aligned}
& m_i \frac{V' n_i \langle R \rangle u_{i,\varphi} - V'^- n_i^- \langle R \rangle^- u_{i,\varphi}^-}{\tau} \\
& + \frac{1}{\rho_b} \frac{\partial}{\partial x} \left(- \frac{V' \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \chi_{u\varphi,i}}{\rho_b} \frac{\partial u_{i,\varphi}}{\partial x} + \left(V' \langle |\nabla \rho|^2 \rangle n_i V_{u\varphi,i}^{pinch} + \Gamma_i - \frac{\dot{\Phi}_b}{2\Phi_b} V' n_i \rho \right) \cdot \langle R \rangle m_i \cdot u_{i,\varphi} \right) \quad (7.24) \\
& = V' \left(U_{i,\varphi,\exp} + \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} u_{z,\varphi} - \left(U_{i,\varphi,imp} + \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} + \frac{\dot{\rho}_b}{\rho_b} m_i n_i \langle R \rangle \right) \cdot u_{i,\varphi} \right)
\end{aligned}$$

Then, taking in account (6.8) and (6.9), coefficients for the interface with numerical solver can be used as:

$$a(x) = V' \langle R \rangle m_i n_i \quad (7.25)$$

$$b(x) = V'^- \langle R \rangle^- m_i n_i^-$$

$$c(x) = \rho_b$$

$$d(x) = \frac{V' \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \chi_{u\varphi,i}}{\rho_b}$$

$$e(\rho) = \left(V' \langle |\nabla \rho|^2 \rangle n_i V_{u\varphi,i}^{pinch} + \Gamma_i - \frac{\dot{\Phi}_b}{2\Phi_b} V' n_i \rho \right) \cdot \langle R \rangle m_i$$

$$f(\rho) = V' \left(U_{i,\varphi,\exp} + \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} u_{z,\varphi} \right)$$

$$g(\rho) = V' \left(U_{i,\varphi,imp} + \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} + \frac{\dot{\rho}_b}{\rho_b} m_i n_i \langle R \rangle \right)$$

$$h = \tau$$

$$Y^{t-1}(\rho) = u_{i,\varphi}^-$$

boundary conditions:

at the axis, $x = 0$:

$$\left. \frac{\partial u_{i,\varphi}}{\partial x} \right|_{x=0} = 0 \quad (7.26)$$

$$v(1) = 1; \quad u(1) = 0; \quad w(1) = 0$$

at the edge, $x = 1$:

type=1 (value)

$$u_{i,\varphi} \Big|_{x=1} = u_{\varphi,bnd}$$

$$v(2) = 0; \quad u(2) = 1; \quad w(2) = u_{\varphi,bnd}$$

type=2 (gradient)

$$\left. \frac{\partial u_{i,\varphi}}{\partial x} \right|_{x=1} = -\nabla u_{\varphi,bnd} \cdot \rho_b$$

$$v(2) = 1; \quad u(2) = 0; \quad w(2) = -\nabla u_{\varphi,bnd} \cdot \rho_b$$

type=3 (scale lenght)

$$\left. \frac{\rho_b}{(\partial \ln u_{i,\varphi} / \partial \rho)} \right|_{x=1} = -L_{u\varphi}$$

$$v(2) = 1; \quad u(2) = \frac{\rho_b}{L_{u\varphi}}; \quad w(2) = 0$$

type=4 (flux)

$$\Phi \Big|_{x=1} = f_{u\varphi,bnd}$$

$$\Phi_i = V' \langle |\nabla \rho|^2 \rangle m_i n_i \langle R \rangle \cdot \left(-\frac{\chi_{u\varphi,i}}{\rho_b} \frac{\partial u_{i,\varphi}}{\partial x} + u_{i,\varphi} V_{u\varphi,i}^{pinch} \right) + m_i \langle R \rangle u_{i,\varphi} \Gamma_i$$

$$v(2) = -V' \langle |\nabla \rho|^2 \rangle \langle R \rangle m_i n_i \frac{\chi_{u\varphi,i}}{\rho_b}; \quad u(2) = V' \langle |\nabla \rho|^2 \rangle \langle R \rangle m_i n_i V_{u\varphi,i}^{pinch} + m_i \langle R \rangle \Gamma_i; \quad w(2) = f_{u\varphi,bnd}$$

type=5 (generic)

$$\left. \frac{v_{gen}}{\rho_b} \frac{\partial u_{i,\varphi,bnd}}{\partial x} \right|_{bnd} + u_{gen} u_{i,\varphi,bnd} = w_{gen}$$

$$v(2) = \frac{v_{gen}}{\rho_b}; \quad u(2) = u_{gen}; \quad w(2) = w_{gen}$$

Variables used inside fortran routines:

Variable	TYPE%NAME used in ETS data flow	Internal name used in ROTATION routine	Units
τ	EVOLUTION%TAU	TAU	s
\dot{B}_0	GEOMETRY%BTPRIME	BTPRIME	T/s
B_0	GEOMETRY%BT	BT	T
ρ	GEOMETRY%RHO	RHO	m
V'	GEOMETRY%VPR	VPR	m ²
V'^{-}	EVOLUTION%VPRM	VPRM	m ²
$\langle \nabla \rho ^2 \rangle$	GEOMETRY%G1	G1	--
$\langle R \rangle$	GEOMETRY%G2	G2	m
$\langle R \rangle^{-}$	EVOLUTION%G2M	G2M	m
$u_{i,\varphi}$	PROFILES%VTOR	VTOR	m/s
$u_{i,\varphi}^{-}$	EVOLUTION%VTORM	VTORM	m/s
$\omega_{i,\varphi}$	PROFILES%WTOR	WTOR	1/s
M_i	PROFILES%MTOR	MTOR	kg*m/s
M_{tot}	PROFILES% MTOR_TOT	MTOR_TOT	kg*m/s
n_i	PROFILES%NI	NI	m ⁻³
n_i^{-}	EVOLUTION%NI	NIM	m ⁻³
Φ_i	PROFILES%FLUX_MTOR	FLUX_MTOR	kg*m ² /s ²
$\Phi_{i,conv}$	PROFILES%FLUX_MTOR_CONV	FLUX_MTOR_CONV	kg*m ² /s ²
$\Phi_{i,cond}$	PROFILES% FLUX_MTOR_COND	FLUX_MTOR_COND	kg*m ² /s ²
$\Phi_{i,int}$	PROFILES% INT_SOURCE_MTOR	INT_SOURCE	kg*m ² /s ²
Φ_{tot}	PROFILES% FLUX_MTOR_TOT	FLUX_MTOR_TOT	kg*m ² /s ²

$U_{i,\varphi,\text{exp}}$	--	UI_EXP	kg/m/s ²
$U_{i,\varphi,\text{imp}}$	--	UI_IMP	kg/m ² /s
$U_{i,\varphi,\text{source},1}$	SOURCES%UI_EXP	--	kg/m/s ²
$U_{i,\varphi,\text{source},2}$	SOURCES%UI_IMP	--	kg/m ² /s
$\langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}}$	COLLISIONS %WZI	WZI	kg/m ² /s
$\langle R \rangle \sum_z \frac{m_{zi} n_z u_{z,\varphi}}{\tau_{zi}}$	COLLISIONS %UZI	UZI	kg/m/s ²
$\chi_{u\varphi,i}$	--	DIFF	m ² /s
$V_{u\varphi,i}^{\text{pinch}}$	--	VCONV	m/s
$\chi_{u\varphi,i,\text{model}}$	TRANSPORT%DIFF_VTOR	--	m ² /s
$V_{u\varphi,i,\text{model}}^{\text{pinch}}$	TRANSPORT%VCONV_VTOR	--	m/s
a	SOLVER%A	A	--
b	SOLVER%B	B	--
c	SOLVER%C	C	--
d	SOLVER%D	D	--
e	SOLVER%E	E	--
f	SOLVER%F	F	--
g	SOLVER%G	G	--
h	SOLVER%H	H	--
$v(1:2)$	SOLVER%V	V	--
$u(1:2)$	SOLVER%U	U	--
$w(1:2)$	SOLVER%W	W	--
Y^{t-1}	SOLVER%YM	W	--
solution	SOLVER%Y	Y	--
$\text{derivative of solution}$	SOLVER%DY	DY	--
Functions		Internal name used in ROTATION routine	

$\frac{V'}{\rho} \left(U_{i,\varphi,\text{exp}} + \langle R \rangle \sum_z \frac{m_{zi} n_z u_{z,\varphi}}{\tau_{zi}} + \frac{\langle R \rangle^- m_i n_i^- u_{i,\varphi,\text{interpretative}}^-}{\tau} \left(\frac{V'^-}{V'} \right) \right. \\ \left. - \left(U_{i,\varphi,\text{imp}} - \langle R \rangle \frac{m_i n_i}{\tau} - \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} \right) u_{i,\varphi,\text{interpretative}} \right)$	FUN1
$\int_0^\rho V' \left(U_{i,\varphi,\text{exp}} + \langle R \rangle \sum_z \frac{m_{zi} n_z u_{z,\varphi}}{\tau_{zi}} + \frac{\langle R \rangle^- m_i n_i^- u_{i,\varphi,\text{interpretative}}^-}{\tau} \left(\frac{V'^-}{V'} \right) \right) \partial \rho \\ - \int_0^\rho V' \left(U_{i,\varphi,\text{imp}} - \langle R \rangle \frac{m_i n_i}{\tau} - \langle R \rangle \sum_z \frac{m_{zi} n_z}{\tau_{zi}} \right) u_{i,\varphi,\text{interpretative}} \partial \rho$	INTFUN1