

Statistical Mechanics

Problem Sheet 2

Gases and Black Body Radiation

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Gases and black body radiation

3.1

For particles to be distinguishable, we need to use some parameter (position, spin, colour, etc.) to label each particle and tell one from another. For example, particles in a solid with lattice structure can be distinguished by their positions. We can contrast this with an ideal gas, where particles do not have fixed positions and are indistinguishable.

In an ideal gas, we treat each particle using particle in a box model, so that they have the wave function:

$$\psi(x, y, z) = A \sin(k_x x) \sin(k_y y) \sin(k_z z) \quad (1)$$

where the wave numbers satisfy $k_i = n_i \pi / L_i$.

The energy levels are quantised according to the relation:

$$E_{n_x, n_y, n_z} = \frac{\hbar^2}{2m} \left(\frac{n_x^2}{L_x^2} + \frac{n_y^2}{L_y^2} + \frac{n_z^2}{L_z^2} \right) \equiv \frac{\hbar^2}{2m} \mathbf{k}^2 \quad (2)$$

where we have defined the wave vector $\mathbf{k} = (k_x, k_y, k_z)$.

We may write the partition function for a single particle as:

$$\begin{aligned} Z_1 &= \sum e^{-\beta E_{n_x, n_y, n_z}} \\ &= \frac{V}{\pi^3} \sum \exp\left(-\beta \frac{\hbar^2}{2m} k^2\right) \frac{\pi}{L_x} \frac{\pi}{L_y} \frac{\pi}{L_z} \\ &\approx \frac{V}{\pi^3} \int \exp\left(-\beta \frac{\hbar^2}{2m} k^2\right) dk_x dk_y dk_z \\ &= \frac{V}{\pi^3} \int \exp\left(-\beta \frac{\hbar^2}{2m} k^2\right) k^2 \sin \theta dk d\theta d\phi \\ &\equiv \int_0^\infty g(k) \exp\left(-\beta \frac{\hbar^2}{2m} k^2\right) dk \\ &= \frac{V}{(2\pi)^3} \left(\frac{2\pi m}{\beta \hbar^2} \right)^{3/2} \\ &\equiv \frac{V}{\lambda_{\text{th}}^3} \end{aligned} \quad (3)$$

where we have defined the density of states $g(k) \equiv V k^2 / 2\pi^2$ and the thermal wavelength $\lambda_{\text{th}} \equiv \sqrt{2\pi \hbar^2 / m k_B T}$.

For a system of N non-interacting, indistinguishable particles that sparsely populate the energy

levels (i.e., each energy level is occupied by at most one particle), the partition function of the system is:

$$Z = \frac{Z_1^N}{N!} = \frac{V^N}{\lambda_{\text{th}}^{3N} N!} \quad (4)$$

Using Stirling's approximation, we have:

$$F = -k_B T \ln Z = -Nk_B T \left[\ln \left(\frac{V}{N\lambda_{\text{th}}^3} \right) + 1 \right] \quad (5)$$

which leads to the entropy:

$$S = \left(\frac{\partial F}{\partial T} \right)_{V,N} = Nk_B \left[\frac{5}{2} - \ln \left(\frac{V}{N\lambda_{\text{th}}^3} \right) \right] \quad (6)$$

Also note the equation of state:

$$P = - \left(\frac{\partial F}{\partial V} \right)_{T,N} = \frac{Nk_B T}{V} \quad (7)$$

Consider two gases of equal temperature, volume and pressure. By the equation of state, the number of particles in each gas must be equal. Suppose that they are made of particles of masses m_1 and m_2 , and that the thermal wavelength of the first gas is $\lambda_{\text{th},1}$ and that of the second gas is $\lambda_{\text{th},2}$. Then we have the expression for their entropies:

$$S_i = Nk_B \left[\frac{5}{2} - \ln \left(\frac{V}{N\lambda_{\text{th},i}^3} \right) \right] \quad (8)$$

The action of mixing the two gases is to increase the volume of the system to $2V$, and the number of particles to $2N$. The entropy of the mixed system becomes:

$$S_{\text{mix}} = 2Nk_B \left[\frac{5}{2} - \ln \left(\frac{V}{N\lambda_{\text{th},f}^3} \right) \right] \quad (9)$$

where we define the thermal wavelength of the final system via an average particle mass:

$$\lambda_{\text{th},f} = \sqrt{\frac{2\pi\hbar^2}{(m_1 + m_2)k_B T/2}} \quad (10)$$

Apparently, if $m_1 = m_2$, i.e., the two gases are identical (at least in term of their thermal properties), then $S_{\text{mix}} = S_1 + S_2$ so that there is no entropy change upon mixing. If, however, $m_1 \neq m_2$, we

need to compare:

$$\begin{aligned}
 \Delta S &= S_{\text{mix}} - S_1 - S_2 \\
 &= Nk_B \left[-2 \ln \left(\frac{V}{N\lambda_{\text{th},f}^3} \right) + \ln \left(\frac{V}{N\lambda_{\text{th},1}^3} \right) + \ln \left(\frac{V}{N\lambda_{\text{th},2}^3} \right) \right] \\
 &= Nk_B \ln \left(\frac{\lambda_{\text{th},f}^6}{\lambda_{\text{th},1}^3 \lambda_{\text{th},2}^3} \right) \\
 &= 3Nk_B \ln \left[\frac{\sqrt{m_1 m_2}}{(m_1 + m_2)/2} \right]
 \end{aligned} \tag{11}$$

The function inside the logarithm has maximum of unity along the line $m_1 = m_2$, so that $\Delta S \geq 0$ for any $m_1 \neq m_2$. There is an increase in entropy upon mixing two gases of different masses since information about the identity of the particles is lost in the process.

3.2

(i) Treating the nitrogen gas as an ideal diatomic gas, we have the single-particle partition function:

$$Z_1 = Z_{\text{trans}} Z_{\text{vib}} Z_{\text{rot}} \tag{12}$$

By equipartition theorem, the heat capacity of the gas approaches $C_V = 7k_B/2$ at high temperature. In the temperature range $T_r < T < T_v$, the heat capacity increases from $5k_B/2$ to $7k_B/2$ as the vibrational modes become excited. From the data given, we have approximately $T_r = 170 \text{ K}$. In this range, the increase in C_V is mostly explained by the vibrational modes, which gives the heat capacity:

$$C_{V,\text{vib}} = k_B \left(\frac{\hbar\omega}{k_B T} \right)^2 \frac{\exp(\beta\hbar\omega)}{[\exp(\beta\hbar\omega) - 1]^2} \tag{13}$$

3.3

In the range $\theta_{\text{rot}} \ll T \ll \theta_{\text{vib}}$, the single-particle partition function of a diatomic gas is:

$$Z_1 = \left(\frac{V}{\lambda_{\text{th}}^3} \right) \left(\frac{T}{2\theta_{\text{rot}}} \right) \frac{\exp(-\theta_{\text{vib}}/2T)}{1 - \exp(-\theta_{\text{vib}}/2T)} \tag{14}$$

where the last term does not vary much with temperature as long as $T \ll \theta_{\text{vib}}$.

Since $\lambda_{\text{th}} \propto T^{-1/2}$, we see that $Z_1 \propto VT^{5/2}$. In this regime, the free energy is:

$$F_1 = -k_B T \ln Z_1 \propto -\frac{5}{2} T \ln (V^{2/5} T) \quad (15)$$

and the entropy is:

$$S_1 = - \left(\frac{\partial F_1}{\partial T} \right)_V \propto \frac{5}{2} \ln (V^{2/5} T) \propto \ln (VT^{5/2}) \quad (16)$$

Therefore, along an adiabat, we need $VT^{5/2}$ to be constant. But the equation of state is:

$$P = - \left(\frac{\partial F_1}{\partial V} \right)_T \propto \frac{T}{V} \quad (17)$$

We can then replace T with PV to demand that $V^{7/2} P^{5/2}$ is constant, which is equivalent to:

$$PV^{7/5} = \text{const} \quad (18)$$

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3.4

The Gibbs function can be written as:

$$G = U - TS + PV = F + PV = F - V \left(\frac{\partial F}{\partial V} \right)_T \quad (19)$$

This means that G can be expressed in terms of the partition function as:

$$G = k_B T \left[V \left(\frac{\partial \ln Z}{\partial V} \right)_T - \ln Z \right] \quad (20)$$

For an ideal gas, this becomes:

$$G = F + PV = F + Nk_B T \quad (21)$$

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3.5

We treat photons as a group of spin-one bosons who are not allowed to have zero spin. Since photons are not conserved, we set the chemical potential to zero. The density of states in k -space is:

$$g(k) dk = (2s + 1) \frac{V}{(2\pi)^3} 4\pi k^2 dk = \frac{V}{\pi^2} k^2 dk \quad (22)$$

where we replace the spin degeneracy factor $(2s + 1)$ with two since zero spin is not allowed.

Transforming to ω -space, we have:

$$g(\omega) d\omega = \frac{V}{\pi^2} \frac{\omega^2}{c^3} d\omega \quad (23)$$

We know that the energy of a photon is given by $E_\omega = \hbar\omega$ and the occupation number of a mode is:

$$n_\omega = \frac{1}{\exp(\beta\hbar\omega) - 1} \quad (24)$$

This means we can write the internal energy as:

$$\frac{U}{V} = \frac{1}{V} \int_0^\infty g(\omega) E_\omega n_\omega d\omega \equiv \int_0^\infty \rho_\omega d\omega \quad (25)$$

where we define the spectral energy density:

$$\rho_\omega \equiv \frac{\hbar}{\pi^2 c^3} \frac{\omega^3}{\exp(\beta\hbar\omega) - 1} \quad (26)$$

The maximum of ρ_ω can be numerically found to be $2.82k_B T/\hbar$. We can also define the number density of photons:

$$n_\omega \equiv \frac{1}{V} g(\omega) \frac{1}{\exp(\beta\hbar\omega) - 1} = \frac{1}{\pi^2 c^3} \frac{\omega^2}{\exp(\beta\hbar\omega) - 1} \quad (27)$$

whose maximum is at $1.59k_B T/\hbar$.

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3.6

Instead of treating the photons as a gas of bosons, we can treat them as a collection of harmonic oscillators. The single-particle partition function is:

$$Z_1 = \frac{1}{1 - \exp(-\beta\hbar\omega)} = \frac{1}{1 - \exp(-\hbar\nu/k_B T)} \quad (28)$$

where we have ignored the zero-point energy.

(i) For a collection of N non-interacting harmonic oscillators each labelled by a frequency ν_j , the partition function is:

$$\ln Z = \sum_j \ln Z_j = - \sum_{j=1}^{\infty} \ln [1 - \exp(-h\nu_j/k_B T)] \quad (29)$$

If ν_j are very closely spaced, we can replace the sum with an integral:

$$\ln Z = - \int_0^{\infty} g(\nu) \ln [1 - \exp(-h\nu/k_B T)] d\nu \quad (30)$$

where we define the density of states $g(\nu) d\nu$ as the number of modes in the frequency range ν to $\nu + d\nu$.

(ii) For photons, the energy levels depend on the frequency ν via $E = h\nu$. On the other hand, the wave numbers of photons are quantised according to boundary conditions:

$$k_i = \frac{2\pi\nu_i}{c} = \frac{\pi}{L_i} n_i \quad (31)$$

which means that the frequency levels are quantised as:

$$\nu_i = \frac{c}{2L_i} n_i \quad (32)$$

with the smallest frequency change $\Delta\nu_i = c/2L_i$.

We can then write the sum as an integral:

$$\begin{aligned} \ln Z &= -2 \sum_{j=1}^{\infty} \ln [1 - \exp(-h\nu_j/k_B T)] \\ &= -2 \left(\frac{8L_x L_y L_z}{c^3} \right) \sum_{j=1}^{\infty} \ln [1 - \exp(-h\nu_j/k_B T)] \Delta\nu_x \Delta\nu_y \Delta\nu_z \\ &\approx -16 \frac{V}{c^3} \int \ln [1 - \exp(-h\nu/k_B T)] d^3\nu \\ &= -16 \frac{V}{c^3} \int_0^{\pi/2} \int_0^{\pi/2} \int_0^{\infty} \ln [1 - \exp(-h\nu/k_B T)] \nu^2 \sin \theta d\nu d\theta d\phi \\ &= -\frac{8\pi V}{c^3} \int_0^{\infty} \ln [1 - \exp(-h\nu/k_B T)] \nu^2 d\nu \end{aligned} \quad (33)$$

where the extra factor of 2 comes from the two polarisations of the photon.

(iii) The internal energy is:

$$U = - \left(\frac{\partial \ln Z}{\partial \beta} \right)_V = \frac{8\pi V}{c^3} \int_0^\infty \frac{h\nu^3 \exp(-h\nu/k_B T)}{1 - \exp(-h\nu/k_B T)} d\nu \quad (34)$$

Consider the change of variable $x \equiv h\nu/k_B T$, we have:

$$\begin{aligned} U &= \frac{8\pi V}{c^3} h \left(\frac{k_B T}{h} \right)^4 \int_0^\infty \frac{x^3}{e^x - 1} dx \\ &= \frac{8\pi V}{c^3} h \left(\frac{k_B T}{h} \right)^4 \frac{\pi^4}{15} \\ &= \frac{8\pi^5 (k_B T)^4}{15c^3 h^3} V \end{aligned} \quad (35)$$

(iv) Returning to the expression for $\ln Z$, we can use integration by parts to write:

$$\begin{aligned} \ln Z &= -\frac{8\pi V}{c^3} \int_0^\infty \ln[1 - \exp(-h\nu/k_B T)] \nu^2 d\nu \\ &= -\frac{8\pi V}{c^3} \left[\frac{1}{3} \nu^3 \ln[1 - \exp(-h\nu/k_B T)] \right]_0^\infty - \frac{1}{3} \int_0^\infty \frac{(h/k_B T) \nu^3 \exp(-h\nu/k_B T)}{1 - \exp(-h\nu/k_B T)} d\nu \\ &= \frac{8\pi V}{c^3} \frac{1}{3k_B T} \int_0^\infty \frac{h\nu^3 \exp(-h\nu/k_B T)}{1 - \exp(-h\nu/k_B T)} d\nu \\ &= \frac{1}{3k_B T} U \\ &= \frac{8\pi^5 (k_B T)^3}{45c^3 h^3} V \end{aligned} \quad (36)$$

The equation of state is then:

$$P = - \left(\frac{\partial F}{\partial V} \right)_T = k_B T \left(\frac{\partial \ln Z}{\partial V} \right)_T = \frac{1}{3} \frac{U}{V} = \frac{8\pi^5 (k_B T)^4}{45c^3 h^3} \quad (37)$$

where we identify the Stefan's constant $\sigma = 8\pi^5 k_B^4 / (15c^3 h^3)$.

3.7

At $T = 2 \times 10^7$ K, the radiation pressure is:

$$P = \frac{4\sigma}{3c} T^4 = 4.03 \times 10^{13} \text{ N/m}^2 = 4.03 \times 10^8 \text{ atm} \quad (38)$$

which is only 0.1% of the gas pressure at the centre of the sun.

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