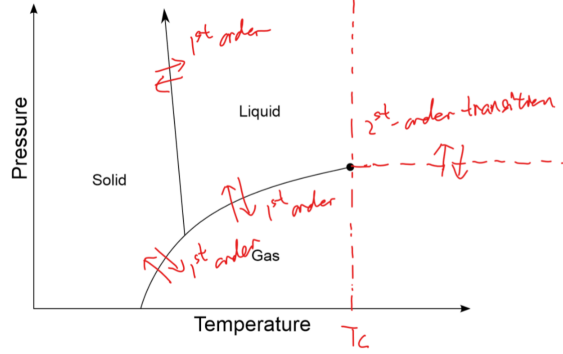


## 2d CFT

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### Week 1

**Exercies 0.0.1.** The first order transitions and second order transitions show in the diagram



**Exercies 0.0.2.** By the homogeneous relation

$$f(t, h) = b^{-d} f(b^{y_t} t, b^{y_h} h)$$

we have

$$f(t, h) = t^{\frac{d}{y_t}} g(\alpha)$$

where  $g(\alpha) = f(1, \alpha)$  and  $\alpha = t^{-\frac{y_h}{y_t}} h$ . It is easy to see that  $\alpha$  is invariance under scaling transformation  $x \rightarrow x/b$ . Hence we have

$$C(t, 0) = -T \frac{\partial^2 f}{\partial T^2} \Big|_{h=0} = -\frac{1}{T_c} t^{\frac{d}{y_t}-2} g''(0)$$

$$M(t, 0) = -\frac{\partial f}{\partial B} \Big|_{h=0} = t^{\frac{d-y_h}{y_t}} g'(0)$$

$$\chi(t, 0) = \frac{\partial^2 f}{\partial B^2} \Big|_{h=0} = t^{(d-2y_h)/y_t} g''(0)$$

As function with single variable  $h$ ,  $\lim_{t \rightarrow 0} M(t, h) \sim h^{\frac{1}{\delta}}$ , which implies that  $g'(\alpha) \sim \alpha^{\frac{1}{\delta}}$  since  $\alpha$  is linear function of  $h$ . Hence we have

$$\lim_{t \rightarrow 0} M = \lim_{t \rightarrow 0} t^{(d-y_h-\frac{y_h}{\delta})} h^{1/\delta}$$

since it is non-zero, we have  $d - y_h - y_h \frac{1}{\delta} = 0$ . Hence we have

$$\delta = \frac{y_h}{d - y_h}$$

**Exercies 0.0.3.** We have following relation

$$G_\sigma(\mathbf{r}; t, h) = t^{-2x_\sigma} G_\sigma\left(\frac{\mathbf{r}}{b}; b^{y_t} t, b^{y_h} h\right) \quad (1)$$

Let  $h = 0, K = b^{y_t} t$ ,

$$G_\sigma(\mathbf{r}; t, 0) = t^{2x_\sigma/y_t} G_\sigma\left(\frac{\mathbf{r}}{K t^{-1/y_t}}; K, 0\right)$$

Since  $G_\sigma(\mathbf{r}) \sim r^{-\tau} e^{-\frac{r}{\xi}}$ , we have  $\xi \sim t^{-1/y_t}$ . It implies  $\nu = 1/y_t$ . With relation 1, we have

$$\chi(t, h) = \frac{1}{T} \int d^d \mathbf{r} G_\sigma(\mathbf{r}; t, h) = t^{d-2x_\sigma} \chi(b^{y_t} t, b^{y_h} h)$$

So  $\gamma = (d - 2x_\sigma)/y_t$ . But we have  $\eta = 2x_\sigma + 2 - d$  for finite limit of  $G(r)$  when  $t \rightarrow 0$  and  $h = 0$ . Therefore, we get

$$\gamma = \nu(2 - \eta)$$

With scaling relations

$$\begin{aligned} \alpha + 2\beta + \gamma &= 2 \\ \alpha + \beta(1 + \delta) &= 2 \end{aligned}$$

and  $\alpha = 2 - d\nu$ , we have

$$\begin{aligned} \beta &= \frac{d\nu - 2\nu + \nu\eta}{2} \\ \delta &= \frac{d - \eta + 2}{d + \eta - 2} \end{aligned}$$

**Exercies 0.0.4.** By listed commutation relations, we have, for  $r, s > 0$ ,

$$\begin{aligned} [D, J_{rs}] &= [D, L_{rs}] = \frac{i}{2} [D, [K_r, P_s]] \\ &= -\frac{i}{2} ([P_s, [D, K_r]] + [K_r, [P_s, D]]) \\ &= \frac{1}{2} [P_s, K_r] - \frac{1}{2} [K_r, P_s] \\ &= 0 \end{aligned}$$

For  $r = -1, s = 0$ ,  $[D, J_{rs}] = [D, D] = 0$ . For  $r = -1, s \neq 0$ ,  $[D, J_{-1,s}] = [D, \frac{1}{2}(P_s - K_s)] = \frac{i}{2}(P_s + K_s)$ . For  $r = 0$ ,  $[D, J_{0s}] = \frac{i}{2}(P_s - K_s)$ . Hence (2,25) is satisfied when  $(m, n) = (-1, 0)$ .

If  $(m, n) = (-1, n)$ , then we have

$$[J_{mn}, J_{rs}] = \frac{1}{2} [P_n, J_{rs}] - \frac{1}{2} [K_n, J_{rs}]$$

With listed commutation relations, we can easily check it coincides with (2,25) respectively. Similarly check in the case of  $(m, n) = (0, n)$ .

## 2d CFT (Week 3)

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### 1 $SL_2(\mathbb{C})$

**Exercies 1.0.1.** We have  $\det X = t^2 - (x^2 + y^2 + z^2)$ . Since points in  $\mathbb{R}^{1,3}$  can be written with Pauli matrix as base. Elements in  $SO(1,3)$  can be viewed as action on  $M_2(\mathbb{C})$  with form  $P \mapsto PXP^*$ , which preserve  $\det$  of  $X$ . We have exact sequence of groups as follows:

$$0 \longrightarrow \mathbb{Z}_2 \longrightarrow SL_2(\mathbb{C}) \xrightarrow{sp} SO(1,3) \longrightarrow 1$$

where  $sp$  is map  $P \mapsto (X \mapsto PXP^*)$ . Since for  $P \in SL_2(\mathbb{C})$ ,  $\det(PXP^*) = \det(X) = t^2 - (x^2 + y^2 + z^2)$ ,  $sp$  is well-defined. Hence  $SO(1,3) \cong SL_2(\mathbb{C})/\mathbb{Z}_2$ .

**Exercies 1.0.2.** •

$$z \mapsto \frac{(z_2 - z_3)(z - z_1)}{(z_2 - z_1)(z - z_3)}$$

- We have

$$\begin{aligned} (w_1 - w_3) &= \frac{(az_1 + b)(cz_3 + d) - (az_2 + b)(cz_4 + d)}{(cz_1 + d)(cz_3 + d)} \\ &= \frac{z_1 - z_3}{(cz_1 + d)(cz_3 + d)} \end{aligned}$$

Hence we have  $[z_1, z_2, z_3, z_4] = [w_1, w_2, w_3, w_4]$ .

### 2 Three-point function

**Exercies 2.0.1.** Let  $\langle \phi_1(z_1)\phi_2(z_2)\phi_3(z_3) \rangle = f(z_{12}, z_{23}, z_{13})$ . Under scalar transformation  $z_i \mapsto \lambda z_i$ , we have

$$f(z_{12}, z_{23}, z_{13}) = \lambda^{h_1+h_2+h_3} f(\lambda z_{12}, \lambda z_{23}, \lambda z_{13})$$

Therefore,  $f$  is with form

$$f(z_{12}, z_{23}, z_{13}) = \frac{C_{123}}{z_{12}^a z_{23}^b z_{13}^c}$$

where  $a + b + c = h_1 + h_2 + h_3$ . Then under comformal transformation  $z_i \mapsto \frac{1}{z_i}$ , we have

$$z_1^{-2h_1} z_2^{-2h_2} z_3^{-2h_3} \frac{(z_1 z_2)^a (z_2 z_3)^b (z_1 z_3)^c}{z_{12}^a z_{23}^b z_{13}^c} = \frac{1}{z_{12}^a z_{23}^b z_{13}^c}$$

Hence  $a = h_1 + h_2 - h_3, b = h_2 + h_3 - h_1, c = h_1 + h_3 - h_2$ .

### 3 Energy-momentum tensor

**Exercies 3.0.1.** •

$$T^{\mu\nu} = -\eta^{\mu\nu} \partial_k \varphi \partial^k \varphi + 2\partial^\mu \varphi \partial^\nu \varphi$$

- We have

$$\delta \sqrt{g} = -\frac{1}{2} \sqrt{g} g_{\mu\nu} \delta g^{\mu\nu}$$

Therefore,

$$\begin{aligned} \tilde{T}^{\mu\nu} &= \frac{2}{\sqrt{g}} \frac{\delta S}{\delta g^{\mu\nu}} \\ &= \frac{1}{2} (-\delta_{\mu\nu} + 2) \partial_\mu \varphi \partial_\nu \varphi \end{aligned}$$

## 4 Derivations

### 4.1 Energy-momentum tensor in complex coordinate

Since

$$\begin{aligned}\partial_0\Phi &= \partial_z\Phi + \partial_{\bar{z}}\Phi \\ \partial_1\Phi &= i\partial_z\Phi - i\partial_{\bar{z}}\Phi\end{aligned}$$

we have

$$\begin{aligned}\partial_z\Phi &= \frac{1}{2}\partial_0\Phi - \frac{i}{2}\partial_1\Phi \\ \partial_{\bar{z}}\Phi &= \frac{1}{2}\partial_0\Phi + \frac{i}{2}\partial_1\Phi\end{aligned}$$

Also, since there are metric tensors in complex coordinates

$$g^{\alpha\beta} = \begin{pmatrix} 0 & 2 \\ 2 & 0 \end{pmatrix} \quad g_{\alpha\beta} = \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix}$$

, we have  $\partial^z\Phi = 2\partial_{\bar{z}}\Phi$  and  $\partial^{\bar{z}}\Phi = 2\partial_z\Phi$ . Therefore, from definition of energy-momentum tensor

$$T_{\alpha\beta} = -g_{\alpha\beta}\mathcal{L} + \frac{\partial\mathcal{L}}{\partial(\partial^\alpha\Phi)}\partial_\beta\Phi$$

we get expression of them in complex coordinates

$$\begin{aligned}T_{zz} &= \frac{1}{4}\left(\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_0\Phi - i\frac{\partial}{\partial_1\Phi}\partial_0\Phi - i\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_1\Phi - \frac{\partial\mathcal{L}}{\partial_1\Phi}\partial_1\Phi\right) \\ T_{\bar{z}\bar{z}} &= \frac{1}{4}\left(\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_0\Phi + i\frac{\partial}{\partial_1\Phi}\partial_0\Phi + i\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_1\Phi - \frac{\partial\mathcal{L}}{\partial_1\Phi}\partial_1\Phi\right) \\ T_{z\bar{z}} &= -\frac{1}{2}\mathcal{L} + \frac{1}{4}\left(\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_0\Phi - i\frac{\partial}{\partial_1\Phi}\partial_0\Phi + i\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_1\Phi + \frac{\partial\mathcal{L}}{\partial_1\Phi}\partial_1\Phi\right) \\ T_{\bar{z}z} &= -\frac{1}{2}\mathcal{L} + \frac{1}{4}\left(\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_0\Phi + i\frac{\partial}{\partial_1\Phi}\partial_0\Phi - i\frac{\partial\mathcal{L}}{\partial_0\Phi}\partial_1\Phi + \frac{\partial\mathcal{L}}{\partial_1\Phi}\partial_1\Phi\right)\end{aligned}$$

Hence

$$\begin{aligned}T_{zz} &= \frac{1}{4}(T_{00} - 2iT_{10} - T_{11}) \\ T_{\bar{z}\bar{z}} &= \frac{1}{4}(T_{00} + 2iT_{10} - T_{11}) \\ T_{z\bar{z}} &= T_{\bar{z}z} = \frac{1}{4}(T_{00} + T_{11})\end{aligned}$$

### 4.2 Schwarzian derivative

$$\tilde{T}(z + \epsilon(z)) = (1 + \partial_z\epsilon(z))^{-2} \left[ T(z) - \frac{c}{12} \left( \frac{\partial_z^3\epsilon(z)}{1 + \partial_z\epsilon(z)} - \frac{2}{3} \frac{\partial_z^2\epsilon(z)}{1 + \partial_z\epsilon(z)} \right) \right]$$

Since

$$\frac{1}{1 + \partial_z\epsilon(z)} = 1 - \partial_z\epsilon(z) + (\partial_z\epsilon)^2 + \dots$$

we have

$$\begin{aligned}\tilde{T}(z + \epsilon(z)) &\approx T(z)(1 - 2\partial_z \epsilon(z)) - \frac{c}{12}(\partial_z^3 \epsilon(z) - \frac{2}{3}\partial_z^2 \epsilon(z)) \\ &\approx T(z) - 2\partial_z \epsilon(z)T(z) - \frac{c}{12}\partial_z^3 \epsilon(z)\end{aligned}$$

Hence

$$\tilde{T}(z + \epsilon(z)) - [\epsilon(z)\partial_z T(z) + T(z)] \approx -\frac{c}{12}\partial_z^3 \epsilon(z) - 2\partial_z \epsilon(z)T(z) - \epsilon(z)\partial_z T(z)$$

It implies that  $\delta_\epsilon(T(z)) = -\frac{c}{12}\partial_z^3 \epsilon(z) - 2\partial_z \epsilon(z)T(z) - \epsilon(z)\partial_z T(z)$ .

### 4.3 Virasoro algebra

The Laurent expansion of  $z^{n+1}$  around  $\omega$  is

$$z^{n+1} = (z - \omega)^{n+1} + \binom{n+1}{1}\omega(z - \omega)^n + \cdots + \binom{n+1}{i}\omega^i(z - \omega)^{n+1-i} + \cdots + \omega^{n+1}$$

Hence we have following residues:

$$\begin{aligned}\text{Res}_\omega \frac{z^{n+1}}{(z - \omega)^4} &= 2\pi i \frac{(n+1)n(n-1)}{6} \omega^{n-2} \\ \text{Res}_\omega \frac{z^{n+1}}{(z - \omega)^2} &= 2\pi i (n+1) \omega^n \\ \text{Res}_\omega \frac{z^{n+1}}{z - \omega} &= 2\pi i \omega^{n+1}\end{aligned}$$

Hence we have

$$\begin{aligned}[L_n, L_m] &= \frac{1}{(2\pi i)^2} \oint_0 d\omega \omega^{m+1} \oint_\omega dz z^{n+1} \left( \frac{c}{2(z - \omega)^4} + \frac{2T(\omega)}{(z - \omega)^2} + \frac{\partial_\omega T(\omega)}{(z - \omega)} + \text{regular part} \right) \\ &= \frac{1}{2\pi i} \oint_0 d\omega \omega^{m+1} \left( \frac{c}{12}(n+1)n(n-1)\omega^{n-2} + 2(n+1)\omega^n T(\omega) + \omega^{n+1}\partial_\omega T(\omega) \right) \\ &= \frac{1}{2\pi i} \left\{ \oint_0 d\omega \left( \frac{c(n+1)n(n-1)}{12} \omega^{m+n-1} \right) - (m-n) \oint_0 d\omega \omega^{m+n+1} T(\omega) \right\} \\ &= \frac{c}{12} n(n^2 - 1) \delta_{n+m,0} - (m-n) L_{n+m}\end{aligned}$$

### 4.4 Commutation relations in free boson

We have

$$\varphi = \varphi_0 + \frac{4\pi}{l} \pi t + i \sum_{n \neq 0} \frac{1}{n} \left( a_n e^{-2\pi i n(t-x)/l} + \bar{a}_n e^{-2\pi i n(t+x)/l} \right)$$

on cylinder.

$$\Pi = \frac{\pi_0}{l} + \frac{1}{2l} \sum_{n \neq 0} \left( a_n e^{-2\pi i n(t-x)/l} + \bar{a}_n e^{-2\pi i n(t+x)/l} \right)$$

With  $[\Pi, \Pi] = 0$ , we can get  $[\pi_0, a_n] = 0$  and  $[\pi_0, \bar{a}_n] = 0$ . Furthermore, since  $[\varphi, \varphi] = 0$ , we have  $[\varphi_0, a_n] = [\varphi_0, \bar{a}_n] = 0$ .

$$\begin{aligned}i\delta(x-y) = [\varphi(x, t), \Pi(y, t)] &= \frac{1}{l} [\varphi_0, \pi_0] + \frac{i}{2l} \sum_{n \neq 0, m \neq 0} \frac{1}{n} \left( [a_n, \bar{a}_m] \exp(-2\pi i [(n+m)t - nx + my]/l) \right. \\ &\quad + [\bar{a}_n, a_m] \exp(-2\pi i [(n+m)t + nx - my]/l) \\ &\quad + [a_n, a_m] \exp(-2\pi i [(n+m)t - nx - my]/l) \\ &\quad \left. + [\bar{a}_n, \bar{a}_m] \exp(-2\pi i [(n+m)t + nx + my]/l) \right)\end{aligned}$$

Since  $t$  can be arbitrary,  $[a_n, \bar{a}_m] = [\bar{a}_n, a_m] = 0$  when  $n + m \neq 0$ . Then let  $t = 0$ , we get

$$\begin{aligned} i\delta(x-y) &= [\varphi(x, 0), \Pi(y, 0)] \\ &= \frac{1}{l}[\varphi_0, \pi_0] + \frac{i}{2l} \sum_{n \neq 0} \frac{1}{n} \left( [a_n, \bar{a}_{-n}] e^{-2\pi i(-nx-ny)/l} + [\bar{a}_n, a_{-n}] e^{-2\pi i(nx+ny)/l} + \text{other terms} \right) \\ &= \frac{1}{l}[\varphi_0, \pi_0] + \frac{i}{l} \sum_{n \neq 0} \left( \frac{2[a_n, \bar{a}_{-n}]}{n} \right) e^{2\pi i(n(x+y))/l} + \text{other terms} \end{aligned}$$

Since  $e^{2\pi i(n(x+y))/l}$  is independent of  $x - y$ , then its coefficient is zero. Hence  $[a_n, \bar{a}_{-n}] = 0$ . Take integral of both left and right side, we can get  $[\varphi_0, \pi_0] = i$  and  $[a_n, a_{-n}] = [\bar{a}_n, \bar{a}_{-n}] = 1$ .

#### 4.5 Action of free fermion

$$S = \frac{1}{4\pi} \int d^2x \psi^\dagger \gamma^0 (\gamma^0 \partial_0 \psi + \gamma^1 \partial_1 \psi)$$

But

$$\gamma^0 (\gamma^0 + \gamma^1) \psi = \begin{pmatrix} \partial_0 + i\partial_1 & 0 \\ 0 & \partial_0 - i\partial_1 \end{pmatrix} \psi$$

Write  $\psi$  as  $\begin{pmatrix} \varphi \\ \bar{\varphi} \end{pmatrix}$ , then we have

$$\gamma^0 (\gamma^0 + \gamma^1) \psi = \begin{pmatrix} 2\partial_{\bar{z}}\varphi \\ 2\partial_z\bar{\varphi} \end{pmatrix}$$

Hence

$$S = \frac{1}{2\pi} \int d^2x (\bar{\varphi} \partial_z \varphi + \varphi \partial_{\bar{z}} \bar{\varphi})$$

#### 4.6 $TT$ OPE in free fermion

By derivative, we get

$$\begin{aligned} \langle \psi(z), \partial_\omega \psi(\omega) \rangle &\sim \frac{1}{(z-\omega)^2} \\ \langle \partial_z \psi, \partial_\omega \psi \rangle &\sim \frac{-2}{(z-\omega)^3} \end{aligned}$$

Hence

$$\begin{aligned} T(z) \partial_\omega \psi(\omega) &= \frac{1}{2} : \psi(z) \partial_z \psi(z) : \partial_\omega \psi(\omega) \\ &\sim -\frac{\psi(\omega)}{(z-\omega)^3} - \frac{1}{2} \frac{\partial_\omega \psi(\omega)}{(z-\omega)^2} \end{aligned}$$

and

$$\begin{aligned} T(z)T(\omega) &= \frac{1}{4} : \psi(z) \partial_z \psi(z) :: \psi(\omega) \partial_\omega \psi(\omega) \\ &\sim \frac{1}{4} \left\{ -\frac{: \partial_z \psi(z) \partial_\omega \psi(\omega) :}{z-\omega} + \frac{2 : \psi(z) \psi(\omega) :}{(z-\omega)^3} - \frac{: \psi(z) \partial_\omega \psi(\omega) + \partial_z \psi(z) \psi(\omega) :}{(z-\omega)^2} + \frac{1}{(z-\omega)^4} \right\} \\ &\sim \frac{1}{4} \left\{ \frac{2\partial_\omega T(\omega)}{z-\omega} + \frac{4T(\omega)}{(z-\omega)^2} + \frac{1}{(z-\omega)^4} \right\} \end{aligned}$$

## 5 Vertex operator and OPE

If we write  $\varphi(z, \bar{z})$  into laurent series since  $\varphi$  is free boson, then we can find  $\exp(ik\varphi)$  is product of infinite exponential components which are commutative. Hence the normal ordering has taylor expansion form

$$:\exp(ik\varphi(z, \bar{z})):= \sum_{n=0}^{\infty} \frac{(ik)^n}{n!} : \varphi(z, \bar{z})^n :$$

To justify that  $:\exp(ik\varphi):$  is primary field, we calculate its OPE

$$\begin{aligned} T(z) : \exp ik\varphi(\omega, \bar{\omega}) : &= -\frac{1}{2} \sum_{n=0}^{\infty} : \partial\varphi(z) \partial\varphi(z) :: \varphi(\omega, \bar{\omega})^n \\ &\sim -\sum_{n=1}^{\infty} \frac{(ik)^n}{n!} n : \partial\varphi(z) \overbrace{\partial\varphi(z) \varphi(\omega, \bar{\omega})}^{n-1} \varphi(\omega, \bar{\omega})^{n-1} : \\ &\quad -\frac{1}{2} \sum_{n=2}^{\infty} \frac{(ik)^n}{n!} n(n-1) : \partial\varphi(z) \overbrace{\partial\varphi(z) \varphi(\omega, \bar{\omega}) \varphi(\omega, \bar{\omega})}^{n-2} \varphi(\omega, \bar{\omega})^{n-2} : \\ &\sim \frac{ik\partial_{\omega}\varphi(\omega)}{z-\omega} : \exp(ik\varphi) : + \frac{k^2}{2(z-\omega)^2} : \exp(ik\varphi) : \end{aligned}$$

This form implies that  $:\exp(ik\varphi):$  is primary field with conformal dimension  $\frac{k^2}{2}$ .

## 6 bc ghost system

$$\begin{aligned} T(z)b(\omega) &= (-2 : b(z)\partial c(z) : + : c(z)\partial b(z) :) b(\omega) \\ &\quad 2\frac{b(z)}{(z-\omega)^2} - \frac{\partial_z b(z)}{z-\omega} \end{aligned}$$

Take Taylor expansion of  $b(z)$  and  $\partial_z b(z)$  around  $\omega$ , we have

$$T(z)b(\omega) \sim 2\frac{b(\omega)}{(z-\omega)^2} + \frac{\partial_{\omega} b(\omega)}{z-\omega}$$

Hence the conformal dimension of  $b$  is 2.

Similarly, we have

$$\begin{aligned} T(z)c(\omega) &= (-2 : b(z)\partial c(z) : + : c(z)\partial b(z) :) c(\omega) \\ &\sim -\frac{c(z)}{(z-\omega)^2} + 2\frac{\partial_z c(z)}{z-\omega} \\ &\sim -\frac{c(\omega)}{(z-\omega)^2} + \frac{\partial_{\omega} c(\omega)}{z-\omega} \end{aligned}$$

Hence conformal dimension of  $c$  is -1.

$$\begin{aligned} T(z)T(\omega) &= 4 : b(z)\partial c(\omega) :: b(\omega)\partial c(\omega) : \\ &\quad -2 : c(z)\partial b(z) :: b(\omega)\partial c(\omega) : -2 : b(z)\partial c(z) :: c(\omega)\partial b(\omega) : \\ &\quad + c(z)\partial b(z) :: c(\omega)\partial b(\omega) : \end{aligned}$$

We will calculate it term by term

$$\begin{aligned}
& 4 : b(z) \partial c(z) :: b(\omega) \partial c(\omega) : \\
& \sim 4 \left( : \overbrace{b(z) \partial c(z) b(\omega) \partial c(\omega)} : + b(z) \overbrace{\partial c(z) b(\omega)} \partial c(\omega) + : \overbrace{b(z) \partial c(z) b(\omega)} \partial c(\omega) : \right) \\
& \sim \frac{4(- : \partial_z c(z) b(\omega) : + : b(z) \partial_\omega c(\omega) :)}{(z - \omega)^2} - \frac{4}{(z - \omega)^4} \\
& \sim -\frac{4}{(z - \omega)^4} + \frac{8b(\omega) \partial_\omega c(\omega)}{(z - \omega)^2} + \frac{-4 : \partial_\omega^2 c(\omega) b(\omega) : + 4 \partial_\omega b(z) \partial_\omega c(\omega)}{z - \omega}
\end{aligned}$$

and

$$\begin{aligned}
& 2 : c(z) \partial b(z) :: b(\omega) \partial c(\omega) : \\
& \sim 2 \left( - : \overbrace{c(z) \partial b(z) b(\omega)} \partial c(\omega) - : c(z) \overbrace{\partial b(z) b(\omega) \partial c(\omega)} : - : \overbrace{\partial b(z) c(z) b(\omega)} \partial c(\omega) \right) \\
& \sim \frac{4}{(z - \omega)^4} + \frac{4 : c(z) b(\omega) :}{(z - \omega)^3} - \frac{2 : \partial b(z) \partial c(\omega)}{(z - \omega)} \\
& \sim \frac{4}{(z - \omega)^4} + \frac{4 : c(\omega) b(\omega) :}{(z - \omega)^3} + \frac{4 : \partial_\omega c(\omega) b(\omega) :}{(z - \omega)^2} + \frac{2 : \partial_\omega^2 c(\omega) b(\omega) : - 2 : \partial_\omega b(\omega) \partial_\omega c(\omega) :}{z - \omega}
\end{aligned}$$

and symmetrically

$$\begin{aligned}
& 2 : b(z) \partial c(z) :: c(\omega) \partial b(\omega) : \\
& \sim \frac{4}{(z - \omega)^4} + \frac{4 : b(\omega) c(\omega) :}{(z - \omega)^3} + \frac{4 : \partial_\omega b(\omega) c(\omega) :}{(z - \omega)^2} + \frac{2 : \partial_\omega^2 b(\omega) c(\omega) : - 2 : \partial_\omega c(\omega) \partial_\omega b(\omega) :}{z - \omega}
\end{aligned}$$

and

$$\begin{aligned}
& c(z) \partial b(z) :: c(\omega) \partial b(\omega) : \\
& \sim \frac{2 : c(\omega) \partial_\omega b(\omega) :}{(z - \omega)^2} + \frac{-\partial_\omega^2 b(\omega) c(\omega) + \partial_\omega c(\omega) \partial_\omega b(\omega)}{z - \omega} - \frac{1}{(z - \omega)^4}
\end{aligned}$$

Hence we have

$$T(z)T(\omega) \sim -\frac{13}{(z - \omega)^4} + \frac{2T(\omega)}{(z - \omega)^2} + \frac{\partial T(\omega)}{z - \omega}$$

So the central charge is equal to  $-26$ .

## 7 Schwarzian derivatives

Let  $\omega = \frac{az+b}{cz+d}$ , then

$$\begin{aligned}
\frac{d\omega}{dz} &= \frac{ad - bc}{(cz + d)^2} \\
\frac{d^2\omega}{dz^2} &= \frac{-2c(ad - bc)}{(cz + d)^3} \\
\frac{d^3\omega}{dz^3} &= \frac{6c^2(ad - bc)}{(cz + d)^4}
\end{aligned}$$

Hence

$$\begin{aligned}
\{\omega, z\} &= \frac{6c^2}{(cz + d)^2} - \frac{3}{2} \left( \frac{-2c}{cz + d} \right)^2 \\
&= 0
\end{aligned}$$



Then

$$\begin{aligned}\left(\frac{a\omega + b}{c\omega + d}\right)'_z &= \frac{ad - bc}{(c\omega + d)^2} \omega'_z \\ \left(\frac{a\omega + b}{c\omega + d}\right)''_z &= \frac{-2c(ad - bc)}{(c\omega + d)^3} (\omega'_z)^2 + \frac{ad - bc}{(c\omega + d)^2} \omega''_z \\ \left(\frac{a\omega + b}{c\omega + d}\right)'''_z &= \frac{6c^2(ad - bc)}{(c\omega + d)^4} (\omega'_z)^3 + 3 \frac{-2c(ad - bc)}{(c\omega + d)^3} \omega''_z \omega'_z + \frac{ad - bc}{(c\omega + d)^2} \omega'''_z\end{aligned}$$

Hence we have

$$\left\{\frac{a\omega + b}{c\omega + d}, z\right\} = \frac{6c^2}{(c\omega + d)^2} (\omega'_z)^2 + 3 \frac{-2c}{c\omega + d} \omega''_z + \frac{\omega'''_z}{\omega'_z} - \frac{3}{2} \left( \frac{-2c}{c\omega + d} \omega'_z + \frac{\omega''_z}{\omega'_z} \right)^2$$

It is equal to  $\{\omega, z\}$ .

### 8 Modular invariant

$$Z_R(\tau, \bar{\tau}) = \frac{1}{|\eta(\tau)|^2} \sum_{m,n} q^{\frac{1}{2}(\frac{m}{R} + \frac{Rn}{2})^2} \bar{q}^{\frac{1}{2}(\frac{m}{R} - \frac{Rn}{2})^2}$$

### 9 Modular transformation

We have

$$\begin{aligned} \gamma \cdot \tau &= \frac{a\tau + b}{c\tau + d} \\ &= \frac{(a\tau + b)(c\bar{\tau} + d)}{|c\tau + d|^2} \\ &= \frac{ac|\tau|^2 + bc\bar{\tau} + ad\tau}{|c\tau + d|^2} \\ &= \frac{ac|\tau|^2 + (ad + bc) \operatorname{Re} \tau + i(ad - bc) \operatorname{Im} \tau}{|c\tau + d|^2} \end{aligned}$$

Since  $ad - bc = 1$ , we can conclude that

$$\operatorname{Im}(\gamma \cdot \tau) = \frac{\operatorname{Im} \tau}{|c\tau + d|^2}$$

In the upper-half plane, the gray region can be described as

$$\begin{aligned} -\frac{1}{2} &\leq \operatorname{Re} \tau \leq \frac{1}{2} \\ |\tau| &> 1 \end{aligned}$$

If  $S$  acts on the gray region, then we have  $S(z) = -\frac{1}{z} = -\frac{\bar{z}}{|z|^2}$ , hence it sends the region to the red region. And the blue region is the image of the gray origin under  $T$ . Finally, the transformation  $ST$ , maps to the green region.