# Quantum Chemistry by Levine

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# 1 Chapter 7 Theorems of Quantum Mechanics

#### 7.1

- (a) True
- (b) True
- (c) False

#### 7.2

c is real number

#### 7.3

(a) 
$$< m|n> = \int f_m^* f_n d\tau$$
  
 $< n|m> * = (\int f_n^* f_m d\tau)^* = \int f_m^* f_n d\tau$   
(b)  $< f|B|g> = \int f^* B g d\tau$   
 $< cf|B|g> = \int c^* f^* B g d\tau = c^* \int f^* B g d\tau = c^* < f|B|g>$ 

# 7.4

unity operator

## 7.5

if B is Hermitian

$$< f|B|g> = \int f^*Bg = < g|B|f>^* = (\int g^*Bf)^* = \int (Bf)^*g = < Bf|g>$$

# 7.6

Given  $\hat{A}$  is Hermitian:

(a) 
$$< f|A|g> = < g|A|f>^* < f|cA|g> = c < f|A|g> < f|cA|g> = c < f|A|g> < g|cA|f>^* = (\int g*cAf)^* = c* \int (Af)^*g = c^* < g|A|f>$$

(b) given 
$$< f|A|g>$$
 and  $< f|B|g>$   $< f|A+B|g> = \int f^*(A+B)g = \int f^*Ag + \int f^*Bg$ 

7.7
(a) verify that 
$$\frac{d^2}{dx^2}$$
 is Hermitian
$$\int f * \frac{d^2}{dx^2}g = f^* \frac{dg}{dx} \Big|_{-\inf}^{\inf} - \int \frac{df^*}{dx} \frac{dg}{dx}$$

$$= -\int \frac{df^*}{dx} \frac{dg}{dx} = \int \frac{d^2f}{dx^2}g$$
from 7.6
$$-\frac{\hbar^2}{2m} \text{ is real number}$$
(b)  $< T_x >= \int \Psi^* (-\frac{\hbar^2}{2m} \frac{d^2}{dx^2}) \Psi$ 

$$= -\frac{\hbar^2}{2m} (\Psi \frac{d\Psi}{dx} \Big|_{-\inf}^{\inf} - \int \frac{d\Psi^*}{dx} \frac{d\Psi}{dx})$$

$$= \frac{\hbar^2}{2m} \int \frac{d\Psi^*}{dx} \frac{d\Psi}{dx}$$
(c)  $< T >= < T_x > + < T_y > + < T_z >$ 
(d) Given each component nonnegative,  $< T >>=$ 

(d)Given each component nonnegative,  $\langle T \rangle = 0$ 

(a) 
$$\int f^* \frac{dg}{dx} = f^* g \Big|_{-\inf}^{\inf} - \int \frac{df^*}{dx} g$$
 there is an extra negative sign not Hermitian (b)  $\int f^* \frac{idg}{dx} = \int \frac{df^*}{dx} g$  so  $i \frac{d}{dx}$  is Hermitian (c)  $\frac{d^2}{dx^2}$  is Hermitian

- (d)  $i\frac{d^2}{dx^2}$  is not Hermitian

#### 7.9

representing a physical quantity means is Hermitian operator

- (a) false
- (b) false
- (c) true
- (d) true

Given 
$$\langle A^2 \rangle = \int \Psi^* A^2 \Psi = \int \Psi^* A(A\Psi)$$
  
then  $f = \Psi, g = A\Psi$  for  $\langle f|A|g \rangle$   
for Hermitian  $A$ :  
 $= \int (A\Psi) * (A\Psi) = \int |A\Psi|^2$ 

### 7.12

(a) given A and B are Hermitian

 $\int f^*ABg = \int (Af)*Bg = \int (BAf)*g$  if AB is also Hermitian:  $\int f^*ABg = \int (ABf)*g$  So we need AB = BA, which means A,B commute (b)  $\int f^*(AB+BA)g = \int f^*ABg + \int f^*BAg$  (c) both x and  $p_x$  are Hermitian but they do not commute due to uncertainty So  $xp_x$  not Hermitian (d)  $\frac{1}{2}(xp_x + p_xx)$  is Hermitian

7.13

(a) 
$$\int f^* \frac{dg}{dx} = f^* g|_{-\inf}^{\inf} - \int \frac{df^*}{dx} g = -\langle g| \frac{d}{dx} | f \rangle$$
  
(b)  $\int f^* (AB - BA) g = \int f^* AB g - \int f^* BA g = \int (Af)^* Bg - \int (Bf)^* Ag = \int (BAf)^* g - \int (ABf)^* g = \langle g|BA|f \rangle^* - \langle g|AB|f \rangle^* = -\langle g|(AB - BA)|f \rangle^*$   
So  $AB - BA$  is anti-Hermitian

#### 7.15

(a) the eigenfunctions are orthogonal given  $\hat{H}\Psi=E\Psi$  at stationary wave state:  $\hat{H}\Psi=(n+\frac{1}{2})hv\Psi$   $<\Psi|\hat{H}|\Psi>=(n+\frac{1}{2})hv\delta_{mn}$  (b) the eigenfunctions are orthogonal given  $\hat{H}\Psi=E\Psi$  at stationary wave state:  $\hat{H}\Psi=\frac{n^2h^2}{8ml^2}\Psi$   $<\Psi|\hat{H}|\Psi>=\frac{n^2h^2}{8ml^2}\delta_{mn}$ 

#### 7.16

$$<\Psi_2|H|f(x)>=< f(x)|H|\Psi_2>^*$$
 where  $H\Psi_2=(2+\frac{1}{2})hv\Psi$  the above equation becomes:  $=\frac{5}{2}hv< f(x)|\Psi_2>^*$ 

#### 7.19

given the simulation equation:  $f = \Sigma < g_i | f > g_i$  given  $g_i = \sqrt{\frac{2}{l}} sin(\frac{n\pi x}{l})$ 

given 
$$g_i = \sqrt{\frac{1}{l}} sin(\frac{m_i}{l})$$
  
so  $\langle g_i | f \rangle = \int_0^{l/2} g_i(-1) dx + \int_{l/2}^l g_i dx$ 

# 7.20

- (a) F.
- like  $Y_0^0$
- (b) F.

only a complete set of eigenfunctions.

(c) T.

 $L_z$  and H commute.

#### 7.21

if m is odd,  $\prod^{m} = \prod$  if m is even,  $\prod^{m} = \hat{1}$ 

#### 7.22

- (a) s has the wave function like a sphere as even function
- (b)  $2p_x$  is odd function.
- (c)  $2s + 2p_x$  are linear combination of  $\hat{H}$  under the same energy level but 2s is even parity and  $2p_x$  is odd parity, so they are not under the same eigenvalue of  $\prod$

#### 7.23

if  $m! = n \ \psi_m$  and  $\psi_n$  are orthogonal

$$\int \psi_m^* \prod \psi_n = 0$$

if m = n

since vibration wave function is either an odd or even function

 $\prod$  has eigenvalue as +1or-1

# 7.24

(a)  $<2s|x|2p_x>$  since  $\hat{x}2p_x$  is even, so when it integrate with 2s, the result is an even function

integral not zero

- (b)  $<2s|x^2|2p_x>$  since both are odd function, integral is zero
- (c)  $< 2p_y|x|2p_x >$  is an odd function, integral is zero

#### 7.25

same as calculation of  $\prod$  given  $\hat{R}f = rf$   $\hat{R}^n f = r^n f = f$ then  $r^n = 1$ 

# 7.26

prove that  $\prod$  is both linear and Hermitian (a)  $\prod (f(x)+g(x))=f(-x)+g(-x)=\prod f(x)+\prod g(x)$  (b)  $< f(x)|\prod |g(x)>=\int f(x)^*\prod g(x)=\int f(-x)^*\prod g(-x)d(-x)=\int (\prod f(x))^*g(x)dx$ 

#### 7.27

intuitively,  $\prod f = +1f$  so f is even function  $\prod g = -1g$  so g is odd function then  $\int f * g = 0$  is an odd function

#### 7.28

 $< v_1 | x | v_2 >$  where  $v_1$  and  $v_2$  are either odd or even according to quantum number so if  $v_1$  and  $v_2$  both odd, or both even, integral is zero otherwise, not zero

# 7.29

(a) do not need to calculate, can just imagine in 3-d cartisan coordinate.  $r=\sqrt{x^2+y^2+z^2}$  does not change  $\theta=\pi-\theta$   $\phi=\phi+\pi$  (b)  $\prod e^{im\phi}=e^{im\phi+\pi}=e^{im\phi}e^{im\pi}=(-1)^me^{im\phi}$ 

7.30 
$$\int \int \dots \left( \int f(q_1, q_2, \dots, q_m) dq_1 dq_2 \dots dq_k \right) \dots dq_m = 0$$

#### 7.32

the prob of getting  $L_z=\hbar$  is  $|\frac{1}{\sqrt{6}}|^2+|\frac{1}{\sqrt{3}}|^2=\frac{1}{2}$  the prob of getting  $L_z=0$  is  $|\frac{1}{\sqrt{3}}|^2=\frac{1}{3}$  so the average  $< L_z>=\frac{1}{2}\hbar$ 

# 7.33

The  $p_0$  and  $p_1$  state has  $L^2 = 2\hbar^2$ The  $d_0$  has  $L^2 = 2 * 3\hbar^2$  $< L^2 >= (\frac{1}{6} + \frac{1}{2}) * 2\hbar + \frac{1}{3} * 6\hbar = \frac{10}{3}\hbar$ 

## 7.34

for both p state:

$$E_1 = -\frac{e^2}{4\pi\epsilon_0 8a}$$
 for  $d$  state:  

$$E_2 = -\frac{e^2}{4\pi\epsilon_0 18a}$$
 So  $< E > = (\frac{1}{6} + \frac{1}{2})E_1 + \frac{1}{3}E_2$ 

# 7.35

Given  $L^2 = 2\hbar$ 

So l=1

then m = -1, 0, 1

So the measurement of  $L_x$  is  $-\hbar, 0, \hbar$ 

# 7.36

the first term is n = 1 stationary state the second term is n=2 stationary state So  $\frac{1}{4}E_{n=1}+\frac{3}{4}E_{n=2}=\frac{13h^2}{32ml^2}$ 

# 7.37

Given the wave function  $g_i = \sqrt{\frac{2}{l}} sin(\frac{n\pi x}{l})$ any non-stationary can be written as its combination So each index, i.e. possibility is:  $|c_i^2| = \int g_i^* \Psi$ 

### 7.38

Shown as above:

$$|c_i^2|=\int g_i^*\Psi$$

# 7.42

- (a)1
- (b)0
- (c)1
- (d)0

7.43
$$\int_{-\infty}^{\infty} \delta(x-a) dx = \int_{-\infty}^{\infty} \delta(x)^2 dx$$

$$= \delta(0) \int_{-\infty}^{\infty} \delta(x) dx = \inf$$

7.44  $\int_0^\infty f(x)\delta(x) \text{ is not the same as } \int_{-\infty}^\infty f(x)\delta(x)$ 

$$\begin{array}{l} \int_{0}^{\infty} f(x)\delta(x) = f(x)H(x)|_{0}^{\infty} - \int_{0}^{\infty} f(x)H(x)dx \\ = f(\infty) - \frac{1}{2}f(0) - f(x)|_{0}^{\infty} = \frac{1}{2}f(0) \end{array}$$

$$\begin{array}{c}
 7.49 \\
 (a) \begin{pmatrix} 6 & 2 \\
 -12 & -12 \end{pmatrix} \\
 (b) \begin{pmatrix} 2 & 4 \\
 8 & -8 \end{pmatrix} \\
 (c) \begin{pmatrix} 3 & 0 \\
 4 & 1 \end{pmatrix} \\
 (d) \begin{pmatrix} 6 & 3 \\
 0 & -9 \end{pmatrix} \\
 (e) \begin{pmatrix} -2 & 5 \\
 -16 & -19 \end{pmatrix}$$

# 7.50

$$CD = \begin{cases} 5i & 10 & 5 \\ 0 & 0 & 0 \\ -i & -2 & -1 \end{cases}$$

$$DC = (5i - 1)$$

#### 7.51

all the unity orthogonal vector

7.52 
$$< f_i|P|f_j> = < f_i|sC|f_j> = c < f_i|C|f_j>$$
 i.e.  $P_{ij} = sC_{ij}$ 

# 7.53

Given  $\langle f_i|A|f_j \rangle = a_i\delta_{ij}$ Given  $f_i$  is a complete orthogonal set Let  $Af_j = \Sigma_k c_k f_k$   $\langle f_i|A|f_j \rangle = \Sigma_k c_k \langle f_i|f_k \rangle = \Sigma_k c_k \delta_{ik} = c_i$ So:  $Af_j = \Sigma_k \langle f_k|A|f_j \rangle f_k = \Sigma_k a_i \delta_{ij} f_k = a_i f_j$ So  $a_i$  is the eigenvalue.

7.63

Given Hermitian:

$$< f_i|A|f_j> = < Af_i|f_j>$$
  
 $\int f_i^* Af_j = \int (Af_i) * f_j$ 

So A has eigenvalue a has to be real

#### 7.64

(a)F.

It can be a linear combination of time term prod stationary wave function

- (b) T
- (c) F.

needs to be a stationary function

(d) F.

These eigenfunctions need to have the same eigenvalue

(e) F.

The measurement has to be a eigenvalue.

(f) T.

 $|\Psi|^2$  is independent of time, which is what stationary means

(g) F.

They can have one common eigenfunction, but not a common complete set of eigenfunction.

(h) F.

They can have the same eigenfunction under different eigenvalue

(i) F.

It does not ensure they are non-degenerate

- (j) F.
- (k) T.
- (l) F.

It does not have a fixed unit.

(m) F.

 $<\Psi|A|\Psi>$  is real does not ensure  $\Psi$  is real

- (n) T.
- (o) T
- (p) F.

It holds for all well-behaved functions as long as B is Hermitian

- (q) T.
- (r) F.

only for stationary states