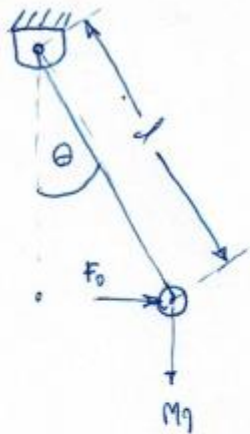


- 3-8. A simple pendulum of mass m and length l is initially motionless at the bottom equilibrium position, $\theta(0) = 0$. Then a constant horizontal force $F_0 = \frac{1}{2}mg$ is applied to the particle. Find θ_{\max} for the motion which follows.



Potential energy term.

$$V = -mgl \cos \theta - \frac{1}{2} mgl \sin \theta$$

$$T_1 = \frac{1}{2} m l^2 \dot{\theta}^2$$

$$T_1 + V_1 = T_2 + V_2$$

from conservation of energy

$$\frac{1}{2} m l^2 \dot{\theta}^2 - mgl (\cos \theta + \frac{1}{2} \sin \theta) = -mgl$$

set $\dot{\theta} = 0$ and $\theta = \theta_{\max}$, so $\cos \theta + \frac{1}{2} \sin \theta = 1$

But $\cos \theta = \cos^2 \frac{\theta}{2} - \sin^2 \frac{\theta}{2}$ and $\sin \theta = 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2}$

Hence $\sin \frac{\theta}{2} \cos \frac{\theta}{2} = 2 \sin^2 \frac{\theta}{2}$, so $\tan \frac{\theta_{\max}}{2} = \frac{1}{2}$

$$\theta_{\max} = 2 \tan^{-1} \frac{1}{2} = 53.13^\circ \quad \underline{\text{Ans}}$$

3-9. A pendulum consists of a particle of mass m and a massless string of length $2R$. As the deflection angle θ increases, the string wraps around one of two fixed cylinders of radius R adjacent to the support point O .

- (a) Obtain the differential equation for θ , where θ is assumed to be positive.
 (b) Assuming the initial conditions $\theta(0) = 0$, $\dot{\theta}(0) = (g/2R)^{1/2}$, find θ_{\max} .

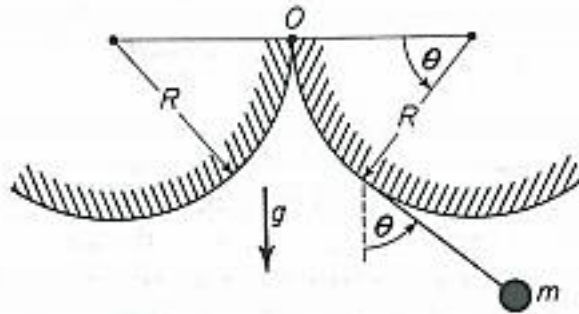
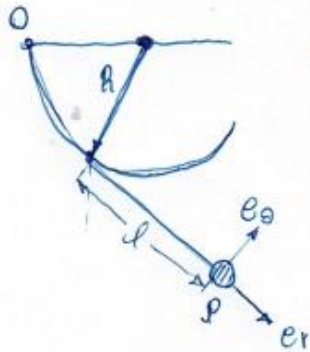


Figure P3-9



a: from $l = 2R - R\theta = (2 - \theta)R$

Using tangential and normal components.

$$\vec{a}_p = \ddot{s} \vec{e}_\theta - \frac{\dot{s}^2}{\rho} \vec{e}_r$$

Where $\dot{s} = l\dot{\theta} = (2 - \theta)R\dot{\theta}$ and $\rho = l$

$$\ddot{s} = (2 - \theta)R\ddot{\theta} - R\dot{\theta}^2$$

$$\sum F_\theta = m a_\theta$$

$$-mg \sin \theta = m \left(R(2 - \theta)\ddot{\theta} - R\dot{\theta}^2 \right) \text{ and then } R(2 - \theta)\ddot{\theta} - R\dot{\theta}^2 + g \sin \theta = 0$$

b. Use conservation of energy. set $\theta(0) = 0$, $\dot{\theta}(0) = \sqrt{\frac{g}{2R}}$

$$T + V = \frac{1}{2} m (2 - \theta)^2 R^2 \dot{\theta}^2 - mgR (\sin \theta + (2 - \theta) \cos \theta) = mgR - 2mgR$$

$$= -mgR$$

$$\theta_{\max} \text{ at } \dot{\theta} = 0 \quad \text{so} \quad \sin \theta + (2 - \theta) \cos \theta = 1 \quad \text{or} \quad \theta_{\max} = \frac{\pi}{2} \quad \underline{\text{Ans}}$$

- 3-10. A massless disk of radius R has an embedded particle of mass m at a distance $R/2$ from the center. The disk is released from rest in the position shown and rolls without slipping down the fixed inclined plane. Find: (a) $\dot{\theta}$ as a function of θ ; (b) a position θ at which the speed of the particle is momentarily constant.

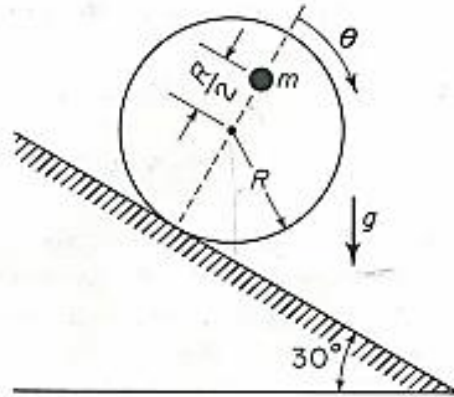


Figure P3-10

Conservative system

a. from $v^2 = \left[R^2 + \left(\frac{R}{2} \right)^2 + 2R \left(\frac{R}{2} \right) \cos \theta \right] \dot{\theta}^2$

$$T = \frac{1}{2} m v^2 = \frac{1}{2} m R^2 \dot{\theta}^2 \left(\frac{5}{4} + \cos \theta \right)$$

$$V = -\frac{1}{2} m g R \theta + \frac{1}{2} m g R \cos(\theta + 30^\circ)$$

Using conservation of energy.

$$T + V = \frac{1}{2} m R^2 \dot{\theta}^2 \left(\frac{5}{4} + \cos \theta \right) + \frac{1}{2} m g R \times$$

$$[\cos(\theta + 30^\circ) - \cos \theta] = \frac{\sqrt{3}}{4} m g$$

solving for $\dot{\theta}$, $\dot{\theta} = \sqrt{\frac{g}{R} \left[\frac{2\sqrt{3} + 4\theta - 4\cos(\theta + 30^\circ)}{5 + 4\cos \theta} \right]}$ Ans

- b. The speed of the particle is constant when potential energy constant.

$$\frac{dV}{d\theta} = -\frac{1}{2} m g R - \frac{1}{2} m g R \sin(\theta + 30^\circ) = 0, \text{ so } \theta = 240^\circ$$

3-13. Given that the force acting on a particle has the following components: $F_x = -x + y$, $F_y = x - y + y^2$, $F_z = 0$. Solve for the potential energy V .

Given $\vec{F} = F_x \hat{i} + F_y \hat{j} + F_z \hat{k}$ where $F_x = -x + y$, $F_y = x - y + y^2$
 $F_z = 0$

from $\frac{\partial F_x}{\partial y} = \frac{\partial F_y}{\partial x} = 1$, so $(-x + y) dx + (x - y + y^2) dy = -dV$

from the exact. we have

$-\frac{\partial V}{\partial x} = -x + y$ integrates to $-V = -\frac{1}{2}x^2 + xy + f_1(y)$

$-\frac{\partial V}{\partial y} = x - y + y^2$ integrates to $-V = xy - \frac{1}{2}y^2 + \frac{1}{3}y^3 + f_2(x)$

result $V = \frac{1}{2}x^2 - xy + \frac{1}{2}y^2 - \frac{1}{3}y^3 + C$ Ans

3-15. Initially the spring has its unstressed length l_0 and the particle has a velocity v_0 in maximum length $4l_0/3$. Assuming no gravity, solve for the spring stiffness k as a function of m , l_0 , and v_0 .

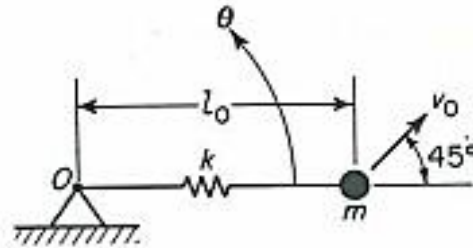


Figure P3-15

Use polar coordinates.

$$T = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2)$$

$$V = \frac{1}{2} k (r - l_0)^2$$

Conservative system.

$$\frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) + \frac{1}{2} k (r - l_0)^2 = \frac{1}{2} m v_0^2$$

At max stretch $r = \frac{4l_0}{3}$, $\dot{r} = 0$ we obtain

$$\frac{1}{2} m \left(\frac{4l_0}{3} \right)^2 \dot{\theta}^2 + \frac{1}{2} k \left(\frac{l_0}{3} \right)^2 = \frac{1}{2} m v_0^2$$

Using conservation of angular momentum.

$$m r^2 \dot{\theta} = \frac{m l_0}{\sqrt{2}} v_0 \quad \text{of } \dot{\theta} = \frac{l_0 v_0}{\sqrt{2} r^2} \text{ at } r = \text{max}$$

$$\text{Then } \frac{8}{9} m l_0^2 \left(\frac{9 v_0}{16 \sqrt{2} l_0} \right)^2 + \frac{k l_0^2}{18} = \frac{1}{2} m v_0^2$$

$$\text{or } k = 18 \frac{m v_0^2}{l_0^2} \left(\frac{1}{2} - \frac{9}{64} \right) = \frac{207 m v_0^2}{32 l_0^2} = 6.4668 \frac{m v_0^2}{l_0^2} \underline{\underline{\text{ans}}}$$