

Assignment #1 Textbook Problems

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2. A free electron has wave function of the form:

$$\psi(z, t) = A e^{ikz} e^{i\omega t}$$

calculating ω :

$$\begin{aligned}\omega &= \frac{E}{\hbar} \\ &= \frac{3 \text{ eV}}{\hbar}\end{aligned}$$

calculating k :

$$\begin{aligned}k &= \frac{2\pi}{\lambda} \\ \lambda &= \frac{h}{p} = \frac{h}{mv} = \frac{h}{m_e 10^5} = 7.274 \times 10^{-9} \\ k &= \frac{2\pi}{7.274 \times 10^{-9}} = 8.64 \times 10^8 \text{ m}^{-1}\end{aligned}$$

putting it all together:

$$\psi(z, t) = A e^{i(8.64 \times 10^8)z} e^{i\left(\frac{3 \text{ eV}}{\hbar}\right)t}$$

To calculate the P.E.:

$$E = E_{PE} + E_{KE}$$

$$= V + E_{KE}$$

isolating and solving for V :

$$V = E - E_{KE}$$

$$= 3\text{eV} - \frac{1}{2} m_e v^2$$

$$= 3\text{eV} - \frac{1}{2} m_e (10^5)^2$$

$$= \boxed{2.97 \text{ eV}}$$

Verifying with Schrodinger time-dependent:

$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} + V \right) \psi$$

$$V\psi = i\hbar \dot{\psi} + \frac{\hbar^2}{2m} \psi''$$

assuming ψ is separable: $\psi = f(z) g(t)$

$$V f(z) g(t) = i\hbar f(z) \dot{g}(t) + \frac{\hbar^2}{2m} f''(z) g(t)$$

$$V = i\hbar \frac{\dot{g}(t)}{g(t)} + \frac{\hbar^2}{2m} \frac{f''(z)}{f(z)}$$

from earlier in the question, we know $f(z)$ and $g(t)$ so:

$$\begin{aligned}
 f(z) &= e^{ikz} & g(t) &= e^{-i\omega t} \\
 f'(z) &= ik e^{ikz} & g'(t) &= -i\omega e^{-i\omega t} \\
 f''(z) &= -k^2 e^{ikz}
 \end{aligned}$$

$$\begin{aligned}
 V &= i\hbar(-i\omega) + \frac{\hbar^2}{2m}(-k^2) \\
 &= 3\text{eV} + \frac{\hbar^2}{2m_e}(-(8.64 \times 10^8)^2) \\
 &= \boxed{2.97 \text{ eV}}
 \end{aligned}$$

3. We begin by simply analyzing the equation in question:

$$\begin{aligned}
 V &= \lambda f \\
 &= \frac{h}{p} \frac{\omega}{2\pi} = \frac{h}{p} \frac{E}{h} = \frac{E}{p} = \frac{\frac{1}{2}mv^2}{mv} = \frac{v}{2}
 \end{aligned}$$

this equation gives phase velocity which differs from the velocity of the electron. Therefore the equation does not hold for electrons. A more accurate approach would involve using

$$E = \frac{1}{2}mv^2$$

4. We can approximate this problem to be an electron in an infinite well.

Energy for an electron in an infinite well for a given state n is:

$$E_n = \frac{h^2 n^2}{2m_e L^2}$$

solving for n :

$$\begin{aligned} n &= \sqrt{\frac{E_n 2m_e L^2}{h^2}} \\ &= \sqrt{\frac{(5\text{eV}) 2m_e (10)^2}{h^2}} \\ &= \boxed{3.64 \times 10^{10}} \end{aligned}$$

approximate energy diff:

$$\begin{aligned} E_{211} - E_{111} &= \frac{h^2 n^2}{2m_e (10)^2} \left[(2^2 + 1^2 + 1^2) - (1^2 + 1^2 + 1^2) \right] \\ &= \frac{h^2 n^2}{2m_e (100)} (3) \\ &= \boxed{1.8 \times 10^{-39} \text{ J}} \end{aligned}$$

5. Exact same procedure as #4, only difference being $L = 10^{-9}$

$$\begin{aligned} n &= \sqrt{\frac{(5\text{eV}) 2m_e (10^{-9})^2}{h^2}} \\ &= \boxed{3.65} \end{aligned}$$

$$E_{2,1} - E_{1,1} = \frac{h^2 n^2}{2m_e 10^{-18}} \quad (3)$$

$$= \boxed{1.8 \times 10^{-19} \text{ J}}$$

6. Discretization occurs when the particle is confined to a space equal to or smaller than its wavelength. So we must calculate λ

$$E = \frac{1}{2} m v^2 \rightarrow v = \sqrt{\frac{2E}{m}}$$

$$\lambda = \frac{h}{mv} = \frac{h}{m} \sqrt{\frac{m}{2E}}$$

$$\lambda_e = \frac{h}{m_e} \sqrt{\frac{m_e}{2(2.5 \text{ eV})}} \quad \lambda_p = \frac{h}{m_p} \sqrt{\frac{m_p}{2(2.5 \text{ eV})}}$$

$$= \boxed{0.776 \text{ nm}} \quad = \boxed{0.018 \text{ nm}}$$

The electron must be confined to a space of λ_e or smaller, while the proton must be confined to λ_p or smaller to observe energy discretization

8. Probability distribution of a particle is given by:

$$P = \int_0^L \psi^* \psi dx$$

To find the particle in the range
 $L/2 \leq x \leq L$:

$$\begin{aligned}
 P &= \int_{L/2}^L \psi^* \psi dx \\
 &= \frac{2}{L} \int_{L/2}^L \sin\left(\frac{n\pi x}{L}\right) e^{+iE_n t/\hbar} \sin\left(\frac{n\pi x}{L}\right) e^{-iE_n t/\hbar} dx \\
 &= \frac{2}{L} \int_{L/2}^L \sin^2\left(\frac{n\pi x}{L}\right) dx \\
 &= \frac{2}{L} \int_{u_1}^{u_2} \sin^2(u) du \left(\frac{L}{n\pi}\right) \\
 &= \frac{2}{n\pi} \int_{u_1}^{u_2} \sin^2(u) du \\
 &= \frac{2}{n\pi} \left[\frac{u}{2} - \frac{\cos(u)\sin(u)}{2} \right]_{u_1}^{u_2} \\
 &= \frac{2}{n\pi} \left[\frac{n\pi x}{2L} - \frac{\sin\left(\frac{2n\pi x}{L}\right)}{2} \right]_{L/2}^L \\
 &= \frac{2}{n\pi} \left[\left(\frac{n\pi}{2}\right) - \left(\frac{n\pi}{4}\right) \right] \\
 &= \frac{2}{n\pi} \left(\frac{n\pi}{4} \right) \\
 &= \boxed{\frac{1}{2}}
 \end{aligned}$$

9. We know the energy for such an electron can be written as:

$$E_n = \frac{\hbar^2 n^2 \pi^2}{2 m_e L^2}$$

$$E_2 = \frac{\hbar^2 n^2 \pi^2}{2 m_e (2.3 \times 10^{-9})^2}$$

$$= 0.284 \text{ eV}$$

to calculate velocity:

$$v = \sqrt{\frac{2E}{m_e}}$$

$$= \sqrt{\frac{2(0.284 \text{ eV})}{m_e}}$$

$$= 316060 \text{ m/s}$$

11a. A normalized wave function Ψ obeys the following equation:

$$\int_0^L |\Psi|^2 dx = 1$$

plugging $\Psi(x)$ into the above equation we can solve for A :

$$\int_0^L A^2 (x(L-x))^2 dx = 1$$

$$A^2 \int_0^L x^2 (L^2 - 2xL + x^2) dx = 1$$

$$A^2 \int_0^L L^2 x^2 - 2x^3 L + x^4 dx = 1$$

$$A^2 \int_0^L x^2 (L^2 - 2xL + x^2) dx = 1$$

$$A^2 \int_0^L L^2 x^2 - 2x^3 L + x^4 dx = 1$$

$$A^2 \left[\frac{L^2 x^3}{3} - \frac{x^4 L}{2} + \frac{x^5}{5} \right]_0^L = 1$$

$$A^2 \left[\frac{L^5}{3} - \frac{L^5}{2} + \frac{L^5}{5} \right] = 1$$

$$A^2 \frac{L^5}{30} = 1$$

$$A = \sqrt{30/L^5}$$

$$11b. \quad \sqrt{\frac{30}{L^5}} (x(L-x)) = \sum_{n=1}^{\infty} c_n \psi_n(x)$$

We know that the eigenstates of an infinite well are orthogonal to each other such that:

$$\langle \psi_m | \psi_n \rangle = \delta_{mn}$$

We can use the above identity to determine c_n :

$$\begin{aligned} \langle \psi_m | \sum_{n=1}^{\infty} c_n \psi_n \rangle &= \sum_{n=1}^{\infty} c_n \langle \psi_m | \psi_n \rangle \\ &= c_m \end{aligned}$$

From the question, we know we can write $\psi(x)$ in terms of c_n, ψ_n such that:

$$\begin{aligned}
 c_n &= \int_0^L \psi_n \psi(x) dx \\
 &= \int_0^L \frac{2}{L} \sqrt{\frac{30}{L^5}} \int_0^L \sin\left(\frac{n\pi x}{L}\right) (x(L-x)) dx \\
 &= -\sqrt{\frac{2}{L}} \sqrt{\frac{30}{L^5}} L^3 \frac{2\cos(n\pi) - 2}{n^3 \pi^3} \\
 &= -\sqrt{60} \frac{L^3}{L^3} \frac{2(\cos(n\pi) - 1)}{n^3 \pi^3} \\
 &= \boxed{4\sqrt{15} \frac{1 - \cos(n\pi)}{n^3 \pi^3}}
 \end{aligned}$$

11c. The probability E_n is just the coefficient c_n 's magnitude:

$$\begin{aligned}
 P(E_n) &= |c_n|^2 \\
 &= \left| 4\sqrt{15} \left(\frac{1 - \cos(n\pi)}{n^3 \pi^3} \right) \right|^2 \\
 &= 240 \left(\frac{(1 - \cos(n\pi))^2}{n^6 \pi^6} \right)
 \end{aligned}$$

$$P(E_1) = \frac{240 (+2)^2}{1^6 \pi^6} = \boxed{0.999}$$

$$P(E_2) = 240 (0) = \boxed{0}$$

$$P(E_3) = \frac{240 (2)^2}{3^6 n^6} = \boxed{0.00137}$$

$$P(E_4) = 240 (0) = \boxed{0}$$

$$P(E_5) = \frac{240 (2)^2}{5^6 n^6} = \boxed{6.39 \times 10^{-5}}$$

$$P(E_6) = 240 (0) = \boxed{0}$$

It makes sense that only odd Energy eigenstates are possible since $\int |A(x(L-x))|^2 dx$ is an odd function.