Nodal Discontinuous Galerkin Method for the Euler Equations in GR

Samuel J Dunham[1], Eirik Endeve[2], et al.

Contents

1. Discontinuous Galerkin Scheme

We assume a spacetime metric

$$ds^2 = -\alpha^2\,dt^2 + \gamma_{ij}\,dx^i\,dx^j\,, \tag{1}$$

and consider the system of conservation laws with sources

$$\partial_t\big(\sqrt{\gamma}\,\boldsymbol{U}\big) + \sum_{i=1}^{d}\partial_i\big(\alpha\,\sqrt{\gamma}\,\boldsymbol{F}^i(\boldsymbol{U})\big) = \alpha\,\sqrt{\gamma}\,\boldsymbol{G}(\boldsymbol{U}), \tag{2}$$

where

$$\boldsymbol{U} = \big(D,\,S_j,\,\tau\big)^{\mathrm{T}} = \big(\rho\,W,\,\rho\,h\,W^2\,v_j,\,\rho\,W\,(h\,W-1)-p\big)^{\mathrm{T}}, \tag{3}$$

$$\boldsymbol{F}^i(\boldsymbol{U}) = \big(D\,v^i,\,\big)^{\mathrm{T}} \tag{4}$$

[2]Department of Astronomy, Vanderbilt University, 6301 Stevenson Science Center, Nashville TN, 37212, USA; samuel.j.dunham@vanderbilt.edu

[2]Computational and Applied Mathematics Group, Oak Ridge National Laboratory, Oak Ridge, TN 37831-6354, USA; endevee@ornl.gov

2. Bound-Preserving Methods Using First-Order DG Scheme

2.1. Cartesian Coordinates

This section closely follows Qin et al. (2016).

2.1.1. *Set of Admissible States*

We consider a one-dimensional system of conservation laws:

$$\partial_t \boldsymbol{U} + \partial_x \boldsymbol{F}(\boldsymbol{U}) = \boldsymbol{0}, \tag{5}$$

where $\boldsymbol{U}$ is a vector of conserved variables, defined as:

$$\boldsymbol{U} \longrightarrow \begin{pmatrix} D \\ S \\ \tau \end{pmatrix} = \begin{pmatrix} \rho\, W \\ \rho\, h\, W^2\, v \\ \rho\, W\,(h\, W - 1) - p \end{pmatrix}, \tag{6}$$

and $\boldsymbol{F}(\boldsymbol{U})$ are the fluxes of those conserved quantities:

$$\boldsymbol{F}(\boldsymbol{U}) \longrightarrow \begin{pmatrix} \rho\, W\, v \\ \rho\, h\, W^2\, v^2 + p \\ \rho\, h\, W^2\, v - D\, v \end{pmatrix}. \tag{7}$$

The physics leads us to define a set of admissible states, $\mathcal{G}_p$ (the subscript p stands for primitive), as:

$$\mathcal{G}_p \equiv \left\{ \boldsymbol{U} \middle| \rho > 0,\, p > 0,\, v^2 < 1 \right\}. \tag{8}$$

It is shown in Mignone & Bodo (2005) that $\mathcal{G}$ is a convex set[3] and can equivalently be written in terms of the conserved variables as:

$$\mathcal{G} \equiv \left\{ \boldsymbol{U} \middle| D > 0,\, \tau + D > \sqrt{D^2 + S^2} \right\}. \tag{9}$$

2.1.2. *Time-Step Derivation/CFL Condition*

For the first-order DG method using forward-Euler time-stepping, we evolve the vector of conserved variables as:

$$\overline{\boldsymbol{U}}_i^{n+1} = \overline{\boldsymbol{U}}_i^{n} - \eta_i \left[\hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}_i^{n}, \overline{\boldsymbol{U}}_{i+1}^{n} \right) - \hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}_{i-1}^{n}, \overline{\boldsymbol{U}}_i^{n} \right) \right], \tag{10}$$

where

$$\overline{\boldsymbol{U}}_i \equiv \frac{1}{\Delta x_i} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} \boldsymbol{U}_i\, dx, \tag{11}$$

$\eta_i \equiv \Delta t_i / \Delta x_i$, and $\hat{\boldsymbol{F}}$ is the numerical flux. In this document we use the local Lax-Friedrichs flux, defined as:

$$\hat{\boldsymbol{F}}(a, b) = \frac{1}{2} \left[\boldsymbol{F}(a) + \boldsymbol{F}(b) - \alpha_{ab}\,(b - a) \right], \tag{12}$$

[3] Convex in the sense that if $\boldsymbol{U}_1 \in \mathcal{G}$ and $\boldsymbol{U}_2 \in \mathcal{G}$, then $\alpha_1\, \boldsymbol{U}_1 + \alpha_2\, \boldsymbol{U}_2 \in \mathcal{G}$, where $\alpha_1,\, \alpha_2 \in [0, 1]$ and $\alpha_1 + \alpha_2 = 1$.

where a and b represent the state of the fluid in two different elements, α_{ab} is an estimate for the wave-speed:

$$\alpha_{ab} = \max\left[\alpha\left(a\right), \alpha\left(b\right)\right], \tag{13}$$

and α is the largest (in absolute value) eigenvalue of the flux-Jacobian:

$$\alpha = \left|\left|\frac{\partial \boldsymbol{F}}{\partial \boldsymbol{U}}\right|\right|. \tag{14}$$

Using this we define the following variables:

$$\alpha_{i+\frac{1}{2}} = \max\left[\alpha\left(\overline{\boldsymbol{U}}_i\right), \alpha\left(\overline{\boldsymbol{U}}_{i+1}\right)\right], \qquad \alpha_{i-\frac{1}{2}} = \max\left[\alpha\left(\overline{\boldsymbol{U}}_{i-1}\right), \alpha\left(\overline{\boldsymbol{U}}_i\right)\right]. \tag{15}$$

Substituting (12) with (15) into (10):

$$\begin{aligned}
\overline{\boldsymbol{U}}_i^{n+1} &= \overline{\boldsymbol{U}}_i^n - \frac{\eta_i}{2}\Big[\boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right) + \boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i+1}^n\right) - \alpha_{i+\frac{1}{2}}\left(\overline{\boldsymbol{U}}_{i+1}^n - \overline{\boldsymbol{U}}_i^n\right) \\
&\qquad\qquad - \boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right) - \boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i-1}^n\right) + \alpha_{i-\frac{1}{2}}\left(\overline{\boldsymbol{U}}_i^n - \overline{\boldsymbol{U}}_{i-1}^n\right)\Big] \\
&= \left[1 - \frac{\eta_i}{2}\left(\alpha_{i+\frac{1}{2}} + \alpha_{i-\frac{1}{2}}\right)\right]\overline{\boldsymbol{U}}_i^n + \frac{\eta_i}{2}\alpha_{i+\frac{1}{2}}\left[\overline{\boldsymbol{U}}_{i+1}^n - \frac{1}{\alpha_{i+\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i+1}^n\right)\right] \\
&\qquad\qquad + \frac{\eta_i}{2}\alpha_{i-\frac{1}{2}}\left[\overline{\boldsymbol{U}}_{i-1}^n + \frac{1}{\alpha_{i-\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i-1}^n\right)\right] \\
&= \left[1 - \frac{\eta_i}{2}\left(\alpha_{i+\frac{1}{2}} + \alpha_{i-\frac{1}{2}}\right)\right]\overline{\boldsymbol{U}}_i^n + \frac{\eta_i}{2}\alpha_{i+\frac{1}{2}}\boldsymbol{H}^-\left(\overline{\boldsymbol{U}}_{i+1}^n, \alpha_{i+\frac{1}{2}}\right) + \frac{\eta_i}{2}\alpha_{i+\frac{1}{2}}\boldsymbol{H}^+\left(\overline{\boldsymbol{U}}_{i-1}^n, \alpha_{i-\frac{1}{2}}\right),
\end{aligned} \tag{16}$$

where

$$\boldsymbol{H}^{\pm}\left(\overline{\boldsymbol{U}}, \alpha\right) \equiv \overline{\boldsymbol{U}} \pm \frac{1}{\alpha}\boldsymbol{F}\left(\overline{\boldsymbol{U}}\right). \tag{17}$$

The proof that $\boldsymbol{H}^{\pm} \in \mathcal{G}$ is given in Qin et al. (2016). Therefore, we see that with a restriction on $\alpha_{i\pm\frac{1}{2}}$ that (16) is a convex combination. The restriction is (recalling that $\eta_i = \Delta t_i / \Delta x_i$):

$$1 - \frac{\eta_i}{2}\left(\alpha_{i+\frac{1}{2}} + \alpha_{i-\frac{1}{2}}\right) > 0 \implies \frac{\eta_i}{2}\left(\alpha_{i+\frac{1}{2}} + \alpha_{i-\frac{1}{2}}\right) < 1 \implies \Delta t_i < \frac{2\,\Delta x_i}{\alpha_{i+\frac{1}{2}} + \alpha_{i-\frac{1}{2}}} \leq \frac{\Delta x_i}{\max\left(\alpha_{i\pm\frac{1}{2}}\right)}. \tag{18}$$

We want a time-step that is the same for all elements at a given time, so we tighten the restriction to:

$$\Delta t < \min_i\left(\frac{\Delta x_i}{\max\left(\alpha_{i\pm\frac{1}{2}}\right)}\right) = \frac{\Delta x}{\max_i\left(\alpha_{i\pm\frac{1}{2}}\right)}, \tag{19}$$

where the equality follows for a uniform mesh, i.e. $\Delta x_i = \Delta x\ \forall i$.

2.2. Curvilinear Coordinates

NOTE: We assume a conformally-flat, time-independent spatial three-metric.

2.2.1. *Set of Admissible States*

We again consider a one-dimensional system of conservation laws, but this time with a curvilinear metric:

$$\partial_t(\sqrt{\gamma}\,\boldsymbol{U}) + \partial_1(\sqrt{\gamma}\,\boldsymbol{F}) = \sqrt{\gamma}\,\boldsymbol{Q}, \tag{20}$$

where $\boldsymbol{U}$ is given by:

$$\boldsymbol{U} \longrightarrow \begin{pmatrix} D \\ S_1 \\ \tau \end{pmatrix} = \begin{pmatrix} \rho\,W \\ \rho\,h\,W^2\,v_1 \\ \rho\,W\,(h\,W-1)-p \end{pmatrix} = \begin{pmatrix} \rho\,W \\ \rho\,h\,W^2\,\gamma_{1j}\,v^j \\ \rho\,W\,(h\,W-1)-p \end{pmatrix}, \tag{21}$$

$\boldsymbol{F}\,(\boldsymbol{U})$ are the fluxes of those conserved quantities:

$$\boldsymbol{F}\,(\boldsymbol{U}) \longrightarrow \begin{pmatrix} D\,v^1 \\ S^1\,v_1 + p\,\delta^1{}_1 \\ S^1 - D\,v^1 \end{pmatrix} = \begin{pmatrix} \rho\,W\,v^1 \\ \rho\,h\,W^2\,v^1\,v_1 + p \\ \rho\,h\,W^2\,v^1 - D\,v^1 \end{pmatrix} = \begin{pmatrix} \rho\,W\,v^1 \\ \rho\,h\,W^2\,\gamma_{1j}\,v^1\,v^j + p \\ \rho\,h\,W^2\,v^1 - D\,v^1 \end{pmatrix}, \tag{22}$$

and $\boldsymbol{Q}$ is a source term:

$$\boldsymbol{Q} \longrightarrow \begin{pmatrix} 0 \\ \frac{1}{2}\,P^{kj}\,\partial_1\gamma_{kj} \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{1}{2}\left[P^{11}\,\partial_1\,\gamma_{11} + P^{22}\,\partial_1\,\gamma_{22} + P^{33}\,\partial_1\,\gamma_{33}\right] \\ 0 \end{pmatrix} \tag{23}$$

$$= \begin{pmatrix} 0 \\ P^{11}\,h_1\,\partial_1\,h_1 + P^{22}\,h_2\,\partial_1\,h_2 + P^{33}\,h_3\,\partial_1\,h_3 \\ 0 \end{pmatrix}, \tag{24}$$

where we have used the fact that $\gamma_{jj} = (h_j)^2$. The P^{kj} are components of the pressure tensor:

$$P^{kj} = S^k\,v^j + p\,\gamma^{kj} = \gamma^{kj}\left[S_a\,v^a + p\right]. \tag{25}$$

These definitions lead us to define the same set of admissible states as before, namely:

$$\mathcal{G}_p \equiv \left\{\boldsymbol{U}\middle|\rho > 0,\, p > 0,\, v^2 < 1\right\}, \tag{26}$$

the only difference being that v^2 now involves the metric:

$$v^2 = v^j\,v_j = \gamma_{kj}\,v^k\,v^j. \tag{27}$$

Before continuing, we show that the introduction of the metric doesn't affect the translation between $\mathcal{G}_p$ and $\mathcal{G}$...(SD: Need to do this)

2.2.2. *Time-Step Derivation/CFL Condition*

We start by integrating both sides over dx^1 and dividing by the volume, ΔV_i:

$$\frac{1}{\Delta V_i}\int_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} \partial_t(\sqrt{\gamma}\,\boldsymbol{U}_i)dx^1 + \frac{1}{\Delta V_i}\int_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} \partial_1(\sqrt{\gamma}\,\boldsymbol{F}\left(\boldsymbol{U}_i\right))dx^1 = \frac{1}{\Delta V_i}\int_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} \sqrt{\gamma}\,\boldsymbol{Q}_i\,dx^1, \tag{28}$$

where:

$$\Delta V_i = \int_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} dV = \int_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} \sqrt{\gamma}\,dx^1. \tag{29}$$

By defining the cell-average as:

$$\overline{\boldsymbol{W}}_i \equiv \frac{1}{\Delta V_i}\int_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} \boldsymbol{W}_i\,dV, \tag{30}$$

we have:

$$\frac{d\overline{\boldsymbol{U}}_i}{dt} + \frac{1}{\Delta V_i}\left(\sqrt{\gamma}\,\hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}\right)\right)\Big|_{x^1_{i-\frac{1}{2}}}^{x^1_{i+\frac{1}{2}}} = \overline{\boldsymbol{Q}}_i, \tag{31}$$

or, using the common notation of the time step being represented as a superscript and the spatial element represented by a subscript:

$$\overline{\boldsymbol{U}}_i^{n+1} = \overline{\boldsymbol{U}}_i^{n} - \frac{\Delta t_i}{\Delta V_i}\left[\sqrt{\gamma}_{i+\frac{1}{2}}\,\hat{\boldsymbol{F}}^n_{i+\frac{1}{2}} - \sqrt{\gamma}_{i-\frac{1}{2}}\,\hat{\boldsymbol{F}}^n_{i-\frac{1}{2}}\right] + \Delta t_i\,\overline{\boldsymbol{Q}}_i^n. \tag{32}$$

Now we define a parameter a la Zhang & Shu (2011): $\varepsilon \in (0,1)$, such that (NOTE: Zhang & Shu (2011) set $\varepsilon = 1/2$):

$$\overline{\boldsymbol{U}}_i^{n} = \varepsilon\,\overline{\boldsymbol{U}}_i^{n} + (1-\varepsilon)\,\overline{\boldsymbol{U}}_i^{n}. \tag{33}$$

We can use the first term to balance out the term in the square brackets and the second term to balance out the source term.

So, we get:

$$\begin{aligned}
\overline{\boldsymbol{U}}_i^{n+1} &= \varepsilon\left\{\overline{\boldsymbol{U}}_i^{n} - \frac{\Delta t_i}{\varepsilon\,\Delta V_i}\left[\sqrt{\gamma}_{i+\frac{1}{2}}\,\hat{\boldsymbol{F}}^n_{i+\frac{1}{2}} - \sqrt{\gamma}_{i-\frac{1}{2}}\,\hat{\boldsymbol{F}}^n_{i-\frac{1}{2}}\right]\right\} + (1-\varepsilon)\,\overline{\boldsymbol{U}}_i^{n} + \Delta t_i\,\overline{\boldsymbol{Q}}_i^n && (34)\\
&= \varepsilon\left\{\overline{\boldsymbol{U}}_i^{n} - \eta_i\left[\sqrt{\gamma}_{i+\frac{1}{2}}\,\hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}_{i+1}^{n},\overline{\boldsymbol{U}}_i^{n}\right) - \sqrt{\gamma}_{i-\frac{1}{2}}\,\hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}_i^{n},\overline{\boldsymbol{U}}_{i-1}^{n}\right)\right]\right\} + (1-\varepsilon)\,\overline{\boldsymbol{U}}_i^{n} + \Delta t_i\,\overline{\boldsymbol{Q}}_i^n && (35)\\
&= \varepsilon\,\boldsymbol{H}_1 + (1-\varepsilon)\,\boldsymbol{H}_2, && (36)
\end{aligned}$$

where

$$\boldsymbol{H}_1 \equiv \overline{\boldsymbol{U}}_i^{n} - \eta_i\left[\sqrt{\gamma}_{i+\frac{1}{2}}\,\hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}_{i+1}^{n},\overline{\boldsymbol{U}}_i^{n}\right) - \sqrt{\gamma}_{i-\frac{1}{2}}\,\hat{\boldsymbol{F}}\left(\overline{\boldsymbol{U}}_i^{n},\overline{\boldsymbol{U}}_{i-1}^{n}\right)\right], \tag{37}$$

$$\boldsymbol{H}_2 \equiv \overline{\boldsymbol{U}}_i^{n} + \frac{\Delta t_i}{1-\varepsilon}\,\overline{\boldsymbol{Q}}_i^n, \tag{38}$$

and

$$\eta_i \equiv \frac{\Delta t_i}{\varepsilon\,\Delta V_i}. \tag{39}$$

We proceed by focusing on each term individually, starting with the numerical flux term, $\boldsymbol{H}_1$.

2.2.3. *Numerical flux term*

We have to show that $\boldsymbol{H}_1 \in \mathcal{G}$. We again we use the Local-Lax-Friedrichs flux, (12), yielding for $\boldsymbol{H}_1$:

$$\overline{\boldsymbol{U}}_i^n - \frac{\eta_i}{2}\Big\{\sqrt{\gamma}_{i+\frac{1}{2}}\left[\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i+1}^n\right) + \boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right) - \alpha_{i+\frac{1}{2}}\left(\overline{\boldsymbol{U}}_{i+1}^n - \overline{\boldsymbol{U}}_i^n\right)\right] \tag{40}$$

$$- \sqrt{\gamma}_{i-\frac{1}{2}}\left[\boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right) + \boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i-1}^n\right) - \alpha_{i-\frac{1}{2}}\left(\overline{\boldsymbol{U}}_i^n - \overline{\boldsymbol{U}}_{i-1}^n\right)\right]\Big\} \tag{41}$$

$$= \left(1 - \frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}} - \frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\right)\overline{\boldsymbol{U}}_i^n \tag{42}$$

$$-\frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right) + \frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right) \tag{43}$$

$$+\frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\left[\overline{\boldsymbol{U}}_{i-1}^n + \frac{1}{\alpha_{i-\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i-1}^n\right)\right] + \frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}}\left[\overline{\boldsymbol{U}}_{i+1}^n - \frac{1}{\alpha_{i+\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i+1}^n\right)\right]. \tag{44}$$

Now we add and subtract $\frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}}\,\overline{\boldsymbol{U}}_i^n$ and $\frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\,\overline{\boldsymbol{U}}_i^n$, yielding:

$$\left(1 - \eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}} - \eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\right)\overline{\boldsymbol{U}}_i^n \tag{45}$$

$$+\frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}}\left[\overline{\boldsymbol{U}}_i^n - \frac{1}{\alpha_{i+\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right)\right] + \frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\left[\overline{\boldsymbol{U}}_i^n + \frac{1}{\alpha_{i-\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_i^n\right)\right] \tag{46}$$

$$+\frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\left[\overline{\boldsymbol{U}}_{i-1}^n + \frac{1}{\alpha_{i-\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i-1}^n\right)\right] + \frac{1}{2}\,\eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}}\left[\overline{\boldsymbol{U}}_{i+1}^n - \frac{1}{\alpha_{i+\frac{1}{2}}}\boldsymbol{F}\left(\overline{\boldsymbol{U}}_{i+1}^n\right)\right]. \tag{47}$$

All of the terms in square brackets are similar to the $\boldsymbol{H}$ quantities in Qin et al. (2016), and are therefore in $\mathcal{G}$. It can easily be seen that the sum of the coefficients is unity. The final condition is that the coefficient of $\overline{\boldsymbol{U}}_i^n > 0$, or (recalling that $\eta_i = \Delta t_i / \left(\varepsilon\,\Delta V_i\right)$):

$$1 - \eta_i\,\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}} - \eta_i\,\sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}} > 0 \implies \eta_i\left(\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}} + \sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}\right) < 1 \tag{48}$$

$$\implies \Delta t_i < \frac{\varepsilon\,\Delta V_i}{\sqrt{\gamma}_{i+\frac{1}{2}}\,\alpha_{i+\frac{1}{2}} + \sqrt{\gamma}_{i-\frac{1}{2}}\,\alpha_{i-\frac{1}{2}}} \leq \frac{\varepsilon\,\Delta V_i}{2\max\left(\sqrt{\gamma}_{i\pm\frac{1}{2}}\,\alpha_{i\pm\frac{1}{2}}\right)}. \tag{49}$$

Again we want a time-step that is the same for all elements at a given time, so:

$$\Delta t < \min_i\left(\frac{\varepsilon\,\Delta V_i}{2\max\left(\sqrt{\gamma}_{i\pm\frac{1}{2}}\,\alpha_{i\pm\frac{1}{2}}\right)}\right). \tag{50}$$

We close the numerical flux section by writing the explicit form of the time-step for spherical-polar coordinates.

Time-step for Spherical-Polar Coordinates

For spherical-polar coordinates in 1-D we have that $\Delta V_i = 1/3\left(r^3_{i+\frac{1}{2}} - r^3_{i-\frac{1}{2}}\right)$, and (assuming $\alpha_{i\pm\frac{1}{2}} = 1 \;\forall\; i$) $\max\left(\sqrt{\gamma}_{i\pm\frac{1}{2}}\,\alpha_{i\pm\frac{1}{2}}\right) = r^2_{i+\frac{1}{2}}$, so:

$$\begin{aligned}
\Delta t &< \min_i \left\{ \frac{\varepsilon\, 1/3 \left[r^3_{i+\frac{1}{2}} - r^3_{i-\frac{1}{2}}\right]}{2\, r^2_{i+\frac{1}{2}}} \right\} && (51)\\
&= \min_i \left\{ \frac{\varepsilon}{6}\, r_{i+\frac{1}{2}} \left[1 - \frac{r^3_{i-\frac{1}{2}}}{r^3_{i+\frac{1}{2}}}\right] \right\} && (52)\\
&= \min_i \left\{ \frac{\varepsilon}{6}\, r_{i+\frac{1}{2}} \left[1 - \left(1 - \frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)^3\right] \right\} && (53)\\
&= \min_i \left\{ \frac{\varepsilon}{6}\, r_{i+\frac{1}{2}} \left[1 - \left(1 + \left(\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)^2 - 2\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)\left(1 - \frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)\right] \right\} && (54)\\
&= \min_i \left\{ \frac{\varepsilon}{6}\, r_{i+\frac{1}{2}} \left[\left(\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)^3 - 3\left(\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)^2 + 3\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right] \right\} && (55)\\
&= \min_i \left\{ \frac{\varepsilon}{6}\, \Delta r_i \left[\left(\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right)^2 - 3\left(\frac{\Delta r_i}{r_{i+\frac{1}{2}}}\right) + 3\right] \right\}. && (56)
\end{aligned}$$

We know that $\Delta r_i / r_{i+\frac{1}{2}} \in [0, 1]$; the minimum value of the quadratic function in this domain is unity. So, we have that for spherical-polar coordinates:

$$\Delta t < \frac{\varepsilon}{6} \min\left(\Delta r_i\right). \tag{57}$$

Next we handle the source term.

2.2.4. Source term

For this section we drop the subscript i and the superscript n. We have to show that $\boldsymbol{H}_2 \in \mathcal{G}$, where

$$\boldsymbol{H}_2 = \begin{pmatrix} \overline{D} \\ \overline{S}_1 + \frac{\Delta t}{2(1-\varepsilon)}\overline{P^{kk}\,\partial_1\gamma_{kk}} \\ \overline{\tau} \end{pmatrix}, \quad (H_2)_1 > 0, \quad (H_2)_3 + (H_2)_1 > \sqrt{(H_2)_1\,(H_2)_1 + (H_2)_2\,(H_2)^2}. \tag{58}$$

It is clear that the first requirement for $\boldsymbol{H}_2$ is met, i.e. $\overline{D} > 0$. The second requirement is:

$$\overline{D} + \overline{\tau} > \sqrt{\overline{D}^2 + \left[\overline{S}_1 + \frac{\Delta t}{2\,(1-\varepsilon)}\overline{P^{kk}\,\partial_1\gamma_{kk}}\right]\left[\overline{S}^1 + \frac{\Delta t}{2\,(1-\varepsilon)}\overline{P^{kk}\,\partial^1\,\gamma_{kk}}\right]} \tag{59}$$

$$= \sqrt{\overline{D}^2 + \left[\overline{S}_1 + \frac{\Delta t}{2\,(1-\varepsilon)}\overline{P^{kk}\,\partial_1\gamma_{kk}}\right]\left[\overline{S}^1 + \gamma^{11}\frac{\Delta t}{2\,(1-\varepsilon)}\overline{P^{kk}\,\partial_1\,\gamma_{kk}}\right]} \tag{60}$$

$$= \sqrt{\overline{D}^2 + \overline{S}_1\,\overline{S}^1 + \gamma^{11}\left[\frac{\overline{P^{kk}\,\partial_1\,\gamma_{kk}}}{2\,(1-\varepsilon)}\right]^2(\Delta t)^2 + \frac{\overline{P^{kk}\,\partial_1\,\gamma_{kk}}}{2\,(1-\varepsilon)}\left[\overline{S}_1\,\gamma^{11} + \overline{S}^1\right]\Delta t} \tag{61}$$

$$= \sqrt{a\,(\Delta t)^2 + b\,\Delta t + c} \geq 0, \tag{62}$$

where

$$a = \left[\frac{\overline{P^{kk}\,\partial_1\,\gamma_{kk}}}{2\,\sqrt{\gamma_{11}}\,(1-\varepsilon)}\right]^2 \geq 0 \tag{63}$$

$$b = \frac{\overline{P^{kk}\,\partial_1\,\gamma_{kk}}}{\gamma_{11}\,(1-\varepsilon)}\overline{S}_1 = \frac{2\,\overline{S}_1}{\sqrt{\gamma_{11}}}\sqrt{a} \tag{64}$$

$$c = \overline{D}^2 + \overline{S}_1\,\overline{S}^1 \geq 0. \tag{65}$$

The condition for a quadratic function to be positive-definite is that $a > 0$ (concave-up) and $b^2 - 4\,a\,c < 0$ (no real roots). The first condition is clearly met. For the second condition:

$$\frac{4\left(\overline{S}_1\right)^2}{\gamma_{11}}a - 4\,a\left(\overline{D}^2 + \overline{S}_1\,\overline{S}^1\right) = 4\,\overline{S}_1\,\overline{S}^1\,a - 4\,a\overline{D}^2 - 4\,a\,\overline{S}_1\,\overline{S}^1 = -4\,a\,\overline{D}^2 < 0. \tag{66}$$

So, the square root is always positive.

With that we have:

$$\overline{D}^2 + \overline{\tau}^2 + 2\,\overline{D}\,\overline{\tau} > a\,(\Delta t)^2 + b\,\Delta t + c \tag{67}$$

$$\Longrightarrow\ a\,(\Delta t)^2 + b\,\Delta t + c' < 0, \tag{68}$$

where:

$$c' = c - \overline{D}^2 - \overline{\tau}^2 - 2\,\overline{D}\,\overline{\tau} = \overline{S}_1\,\overline{S}^1 - \overline{\tau}^2 - 2\,\overline{D}\,\overline{\tau}. \tag{69}$$

The solutions of this quadratic equation are:

$$\Delta t = \frac{-b}{2\,a} \pm \frac{1}{2\,a}\sqrt{b^2 - 4\,a\,c'} = \frac{-\overline{S}_1}{\sqrt{\gamma_{11}}\,\sqrt{a}} \pm \frac{1}{2\,a}\sqrt{\frac{4\overline{S}_1^2}{\gamma_{11}}a - 4\,a\,c'} \tag{70}$$

$$= \frac{-\overline{S}_1}{\sqrt{\gamma_{11}}\,\sqrt{a}} \pm \frac{1}{\sqrt{\gamma_{11}}\,\sqrt{a}}\sqrt{\overline{S}_1^2 - c'\,\gamma_{11}} = \frac{1}{\sqrt{\gamma_{11}}\,\sqrt{a}}\left[-\overline{S}_1 \pm \sqrt{\overline{S}_1^2 - c'\,\gamma_{11}}\right] \tag{71}$$

$$= \frac{1}{\sqrt{\gamma_{11}}\,\sqrt{a}}\left[-\overline{S}_1 \pm \sqrt{\gamma_{11}\left(\overline{\tau}^2 + 2\,\overline{D}\,\overline{\tau}\right)}\right] \tag{72}$$

$$= \frac{2\,(1-\varepsilon)}{\overline{P^{kk}\,\partial_1\,\gamma_{kk}}}\left[-\overline{S}_1 \pm \sqrt{\gamma_{11}\left(\overline{\tau}^2 + 2\,\overline{D}\,\overline{\tau}\right)}\right]. \tag{73}$$

We see that we must have that $\overline{\tau}^2 + 2\,\overline{D}\,\overline{\tau} > 0 \implies \overline{\tau} > -2\,\overline{D}$. But, since we know that from the second condition of $\boldsymbol{H}_2$ that $\overline{\tau} > -\overline{D}$, this requirement is met.

3. Recovery of Primitive Variables

In order to recover the primitive from the conserved variables we need to solve the nonlinear equation:

$$f\left(p\right)=p-\overline{p}\left(p\right)=0, \tag{74}$$

where $\overline{p}\left(p\right)$ is the pressure as obtained via the ideal gas equation of state with an initial guess, p:

$$\overline{p}=\left(\Gamma-1\right)\rho\,\epsilon, \tag{75}$$

where

$$\rho=\rho\left(\boldsymbol{U},p\right),\quad \epsilon=\epsilon\left(\boldsymbol{U},p\right). \tag{76}$$

In order to solve this equation we make use of the bisection method, and therefore need bounds on our initial guess for the pressure.

3.1. Upper and Lower Bounds for Pressure

We obtain a lower bound for the pressure with:

$$\tau=D\left(h\,W-1\right)-p \implies p=-\left(\tau+D\right)+D\,h\,W\geq-\left(\tau+D\right)+D\,h\,W\,\sqrt{v^i\,v_i}=-\left(\tau+D\right)+\sqrt{S^i\,S_i}. \tag{77}$$

So, since the pressure must be non-negative, we have:

$$p\geq \text{MAX}\left[-\left(\tau+D\right)+\sqrt{S^i\,S_i},\text{SqrtTiny}\right]. \tag{78}$$

For an upper bound, we first note that:

$$h=1+\frac{e+p}{\rho}=1+\frac{\Gamma}{\Gamma-1}\frac{p}{\rho}=1+\frac{\Gamma}{\Gamma-1}\frac{p\,W}{D}, \tag{79}$$

so,

$$\tau=D\left(W+\frac{\Gamma}{\Gamma-1}\frac{p\,W^2}{D}-1\right)-p=D\left(W-1\right)+p\left(\frac{\Gamma}{\Gamma-1}W^2-1\right). \tag{80}$$

So,

$$p=\frac{\tau-D\left(W-1\right)}{\frac{\Gamma}{\Gamma-1}W^2-1}. \tag{81}$$

We also have:

$$W=\left(1-v^i\,v_i\right)^{-1/2}=\left(1-\frac{S^i\,S_i}{\left(\tau+D+p\right)^2}\right)^{-1/2}. \tag{82}$$

Treating p as an independent variable (SD: is this valid?), we have:

$$W\Big|_{p\to\infty}=1, \tag{83}$$

which gives us an upper limit:

$$p\leq\frac{\Gamma-1}{\Gamma}\,\tau. \tag{84}$$

Just to be safe, in the code we multiply this by two, so that:

$$p\leq 2\,\frac{\Gamma-1}{\Gamma}\,\tau. \tag{85}$$

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This preprint was prepared with the AAS LATEX macros v5.0.