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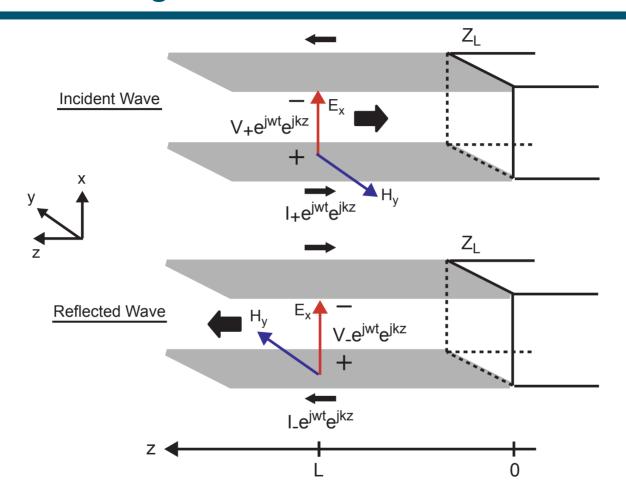
6.976
High Speed Communication Circuits and Systems
Lecture 4
Generalized Reflection Coefficient, Smith Chart,
Integrated Passive Components

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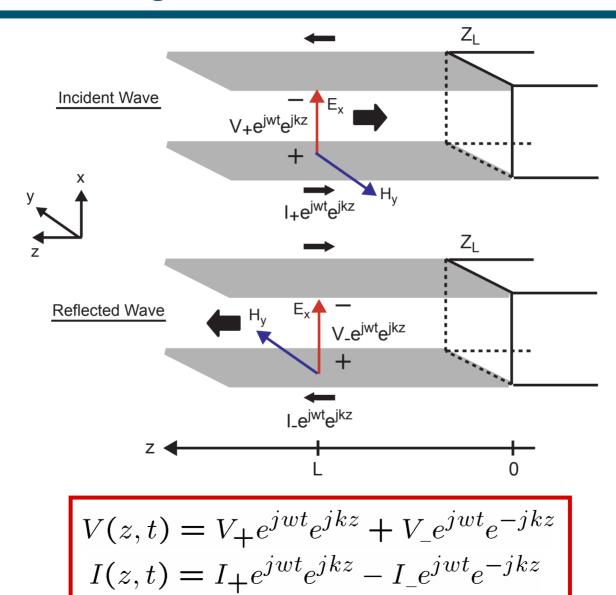
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Determine Voltage and Current At Different Positions



Incident and reflected waves must be added together

Determine Voltage and Current At Different Positions



Define Generalized Reflection Coefficient

$$V(z,t) = V_{+}e^{jwt}e^{jkz} + V_{-}e^{jwt}e^{-jkz}$$
$$I(z,t) = I_{+}e^{jwt}e^{jkz} - I_{-}e^{jwt}e^{-jkz}$$

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \frac{V_{-}}{V_{+}}e^{-2jkz}\right)$$

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \Gamma_{L}e^{-2jkz}\right)$$

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \Gamma(z)\right)$$

Similarly:
$$I(z,t) = I_{+}e^{jwt}e^{jkz}(1-\Gamma(z))$$

$$\Rightarrow \Gamma(z) = \Gamma_L e^{-2jkz}$$

A Closer Look at $\Gamma(z)$

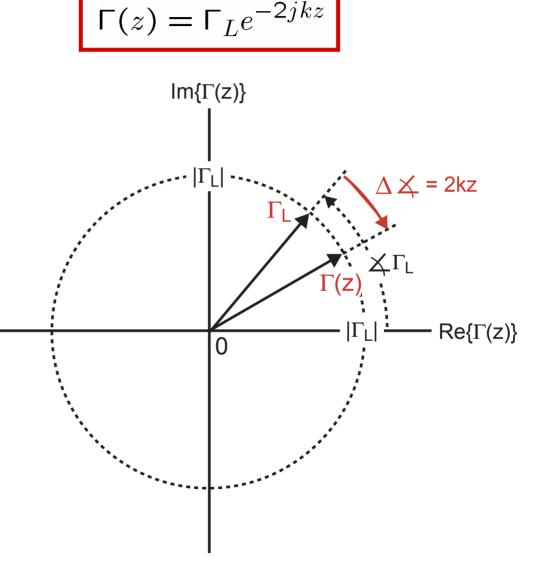
• Recall Γ_{L} is

$$\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o}$$

Note: $|\Gamma_L| \leq 1$

for
$$Re\{Z_L/Z_o\} \ge 0$$

We can view Γ(z) as a complex number that rotates clockwise as z (distance from the load) increases



Calculate $|V_{max}|$ and $|V_{min}|$ Across The Transmission Line

We found that

$$V(z,t) = V_{+}e^{jwt}e^{jkz}\left(1 + \Gamma(z)\right)$$

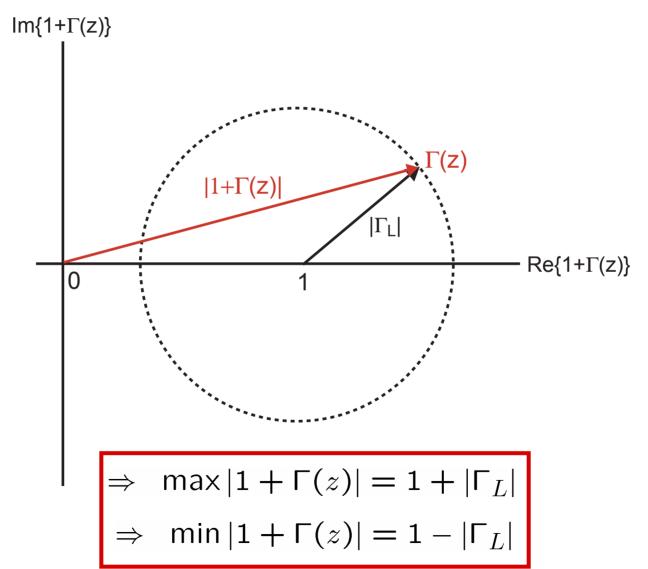
So that the max and min of V(z,t) are calculated as

$$\Rightarrow V_{max} = \max |V(z,t)| = |V_{+}| \max |1 + \Gamma(z)|$$

$$\Rightarrow V_{min} = \min |V(z,t)| = |V_{+}| \min |1 + \Gamma(z)|$$

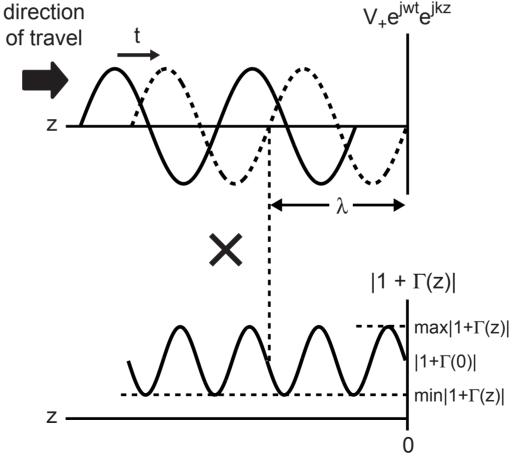
We can calculate this geometrically!

A Geometric View of $|1 + \Gamma(z)|$



Reflections Cause Amplitude to Vary Across Line

- Equation: $V(z,t) = V_{+}e^{jwt}e^{jkz}|1 + \Gamma(z)|e^{j\angle(1+\Gamma(z))}|$
- Graphical representation:



Voltage Standing Wave Ratio (VSWR)

Definition

VSWR =
$$\frac{V_{max}}{V_{min}} = \frac{|V_{+}|(1+|\Gamma_{L}|)}{|V_{+}|(1-|\Gamma_{L}|)} = \frac{1+|\Gamma_{L}|}{1-|\Gamma_{L}|}$$

For passive load (and line)

$$|\Gamma_L| \leq 1 \;\; \Rightarrow \;\; 1 \leq \mathsf{VSWR} \leq \infty$$
 $|\Gamma_L| = 0 \;\; |\Gamma_L| = 1$

We can infer the magnitude of the reflection coefficient based on VSWR

$$|\Gamma_L| = rac{\mathsf{VSWR} - 1}{\mathsf{VSWR} + 1}$$

Reflections Influence Impedance Across The Line

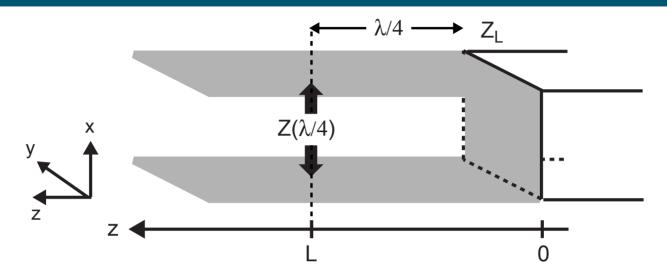
• From Slide 4 $V(z,t) = V_{+}e^{jwt}e^{jkz}(1+\Gamma(z))$ $I(z,t) = I_{+}e^{jwt}e^{jkz}(1-\Gamma(z))$

$$\Rightarrow Z(z,t) = \frac{V_{+}(1+\Gamma(z))}{I_{+}(1-\Gamma(z))} = Z_{o}\frac{1+\Gamma(z)}{1-\Gamma(z)}$$

- Note: not a function of time! (only of distance from load)
- Alternatively $Z(z) = Z_o \frac{1 + \Gamma_L e^{-2jkz}}{1 \Gamma_L e^{-2jkz}}$
 - From Lecture 2: $\lambda = \frac{T}{\sqrt{\mu\epsilon}} = \frac{wT}{w\sqrt{\mu\epsilon}} = \frac{2\pi fT}{k} = \frac{2\pi}{k}$

$$Z(z) = Z_o \frac{1 + \Gamma_L e^{-j(4\pi/\lambda)z}}{1 - \Gamma_L e^{-j(4\pi/\lambda)z}}$$

Example: $Z(\lambda/4)$ with Shorted Load



Calculate reflection coefficient

$$\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o} = \frac{0 - Z_o}{0 + Z_o} = -1$$

Calculate generalized reflection coefficient

$$\Gamma(\lambda/4) = \Gamma_L e^{-j(4\pi/\lambda)(\lambda/4)} = \Gamma_L e^{-j\pi} = -\Gamma_L = 1$$

■ Calculate impedance $Z(\lambda/4) = Z_0 \frac{1 + \Gamma(z)}{1 - \Gamma(z)} = \infty$!

Generalize Relationship Between $Z(\lambda/4)$ and Z(0)

General formulation

$$Z(z) = Z_o \frac{1 + \Gamma_L e^{-j(4\pi/\lambda)z}}{1 - \Gamma_L e^{-j(4\pi/\lambda)z}}$$

At load (z=0)

$$Z_L = Z(0) = Z_0 \frac{1 + \Gamma_L}{1 - \Gamma_L}$$

• At quarter wavelength away (z = $\lambda/4$)

$$Z(\lambda/4) = Z_o \frac{1 - \Gamma_L}{1 + \Gamma_L} = \frac{Z_o^2}{Z_L}$$

- Impedance is inverted!
 - Shorts turn into opens
 - Capacitors turn into inductors

Now Look At Z(△) (Impedance Close to Load)

Impedance formula (∆ very small)

$$Z(\Delta) = Z_o \frac{1 + \Gamma_L e^{-2jk\Delta}}{1 - \Gamma_L e^{-2jk\Delta}}$$

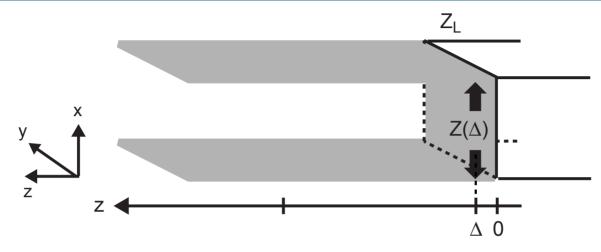
A useful approximation: $e^{-jx} \approx 1 - jx$ for $x \ll 1$

$$\Rightarrow e^{-2jk\Delta} \approx 1 - 2jk\Delta$$

- Recall from Lecture 2: $k = w\sqrt{LC}, \quad Z_o = \sqrt{\frac{L}{C}}$
- Overall approximation:

$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 + \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}{1 - \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}$$

Example: Look At Z(△) With Load Shorted

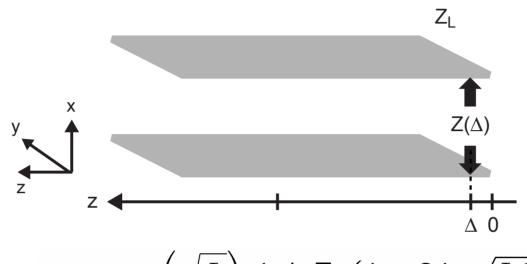


$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 + \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}{1 - \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}$$

- Reflection coefficient: $\Gamma_L = \frac{Z_L Z_o}{Z_L + Z_o} = \frac{0 Z_o}{0 + Z_o} = -1$
- Resulting impedance looks inductive!

$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 - (1 - 2jw\sqrt{LC}\Delta)}{1 + (1 - 2jw\sqrt{LC}\Delta)} pprox jwL\Delta$$

Example: Look At Z(△) With Load Open



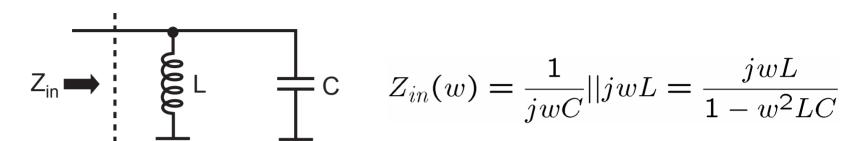
$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 + \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}{1 - \Gamma_L(1 - 2jw\sqrt{LC}\Delta)}$$

Reflection coefficient: $\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o} = \frac{\infty - Z_o}{\infty + Z_o} = 1$

Resulting impedance looks capacitive!

$$Z(\Delta) pprox \left(\sqrt{\frac{L}{C}}\right) rac{1 + (1 - 2jw\sqrt{LC}\Delta)}{1 - (1 - 2jw\sqrt{LC}\Delta)} pprox rac{1}{jwC\Delta}$$

Consider an Ideal LC Tank Circuit



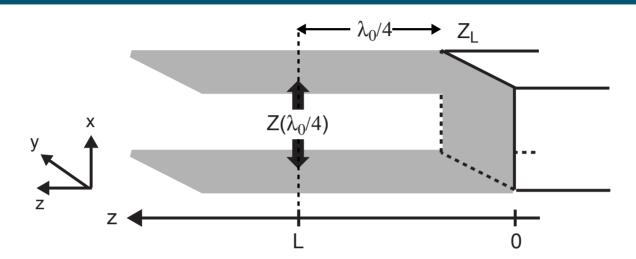
Calculate input impedance about resonance

Consider
$$w=w_o+\Delta w$$
, where $w_o=\frac{1}{\sqrt{LC}}$
$$Z_{in}(\Delta w)=\frac{j(w_o+\Delta w)L}{1-(w_o+\Delta w)^2LC} = \frac{j(w_o+\Delta w)L}{\frac{1-w_o^2LC}{2\Delta w(w_oLC)-\Delta w^2LC}} = \mathbf{0}$$
 negligible

$$\Rightarrow Z_{in}(\Delta w) \approx \frac{j(w_o + \Delta w)L}{-2\Delta w(w_o LC)} \approx \frac{jw_o L}{-2\Delta w(w_o LC)} = -\frac{j}{2}\sqrt{\frac{L}{C}}\left(\frac{w_o}{\Delta w}\right)$$

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Transmission Line Version: $Z(\lambda_0/4)$ with Shorted Load



As previously calculated

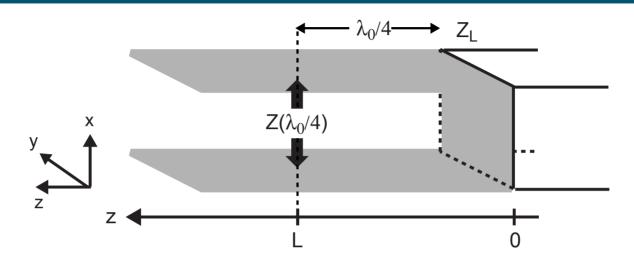
$$\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o} = \frac{0 - Z_o}{0 + Z_o} = -1$$

Impedance calculation

$$Z(z) = Z_o \frac{1 + \Gamma(z)}{1 - \Gamma(z)}$$
, where $\Gamma(z) = \Gamma_L e^{-j(4\pi/\lambda)z}$

Relate λ to frequency $\lambda = \frac{1}{f\sqrt{\mu\epsilon}} = \frac{1}{(f_o + \Delta f)\sqrt{\mu\epsilon}}$

Calculate $Z(\Delta f)$ – Step 1



Wavelength as a function of ∆ f

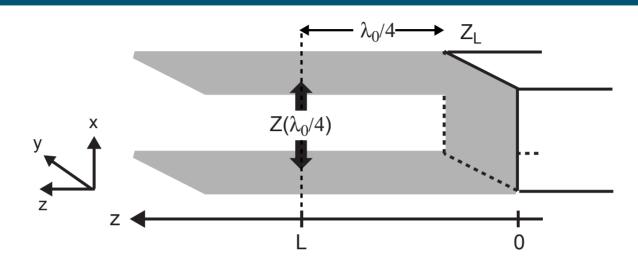
$$\lambda = \frac{1}{(f_o + \Delta f)\sqrt{\mu\epsilon}} = \frac{1}{f_o\sqrt{\mu\epsilon}(1 + \Delta f/f_o)} = \frac{\lambda_o}{1 + \Delta f/f_o}$$

Generalized reflection coefficient

$$\Gamma(\lambda_o/4) = \Gamma_L e^{-j(4\pi/\lambda)\lambda_o/4} = \Gamma_L e^{-j\pi\lambda_o/\lambda} = \Gamma_L e^{-j\pi\lambda_o/\lambda}$$

$$\Rightarrow \Gamma(\lambda_o/4) = \Gamma_L e^{-j\pi(1+\Delta f/f_o)} = -\Gamma_L e^{-j\pi\Delta f/f_o}$$

Calculate $Z(\Delta f)$ – Step 2



Impedance calculation

$$Z(\lambda_o/4) = Z_o \frac{1 - \Gamma_L e^{-j\pi\Delta f/f_o}}{1 + \Gamma_L e^{-j\pi\Delta f/f_o}} = Z_o \frac{1 + e^{-j\pi\Delta f/f_o}}{1 - e^{-j\pi\Delta f/f_o}}$$

• Recall $e^{-j\pi\Delta f/f_o} \approx 1 - j\pi\Delta f/f_o$

$$\Rightarrow Z(z) \approx Z_o \frac{1 + 1 - j\pi\Delta f/f_o}{1 - 1 + j\pi\Delta f/f_o} \approx Z_o \frac{2}{j\pi\Delta f/f_o} = -j\frac{2}{\pi}\sqrt{\frac{L}{C}} \left(\frac{w_o}{\Delta w}\right)$$

Looks like LC tank circuit about frequency w_o!

Smith Chart

Define normalized impedance

$$Z_n = \frac{Z_L}{Z_O}$$

Mapping from normalized impedance to Γ is one-to-one

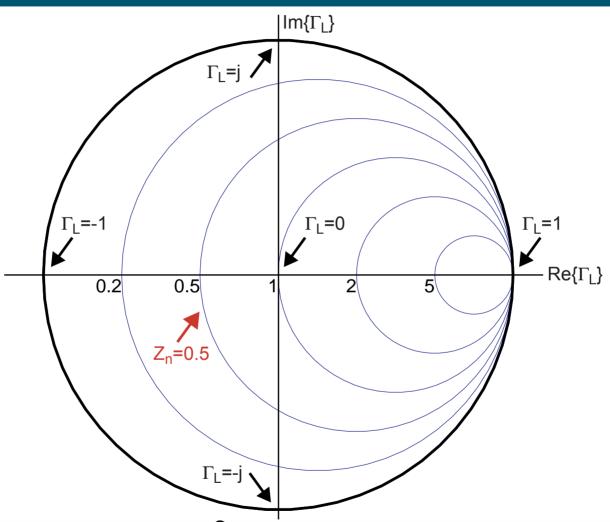
$$Z_n = \frac{1 + \Gamma_L}{1 - \Gamma_L}$$

- lacktriangle Consider working in coordinate system based on Γ
- Key relationship between Z_n and Γ

$$Re\{Z_n\} + jIm\{Z_n\} = \frac{1 + Re\{\Gamma_L\} + jIm\{\Gamma_L\}}{1 - (Re\{\Gamma_L\} + jIm\{\Gamma_L\})}$$

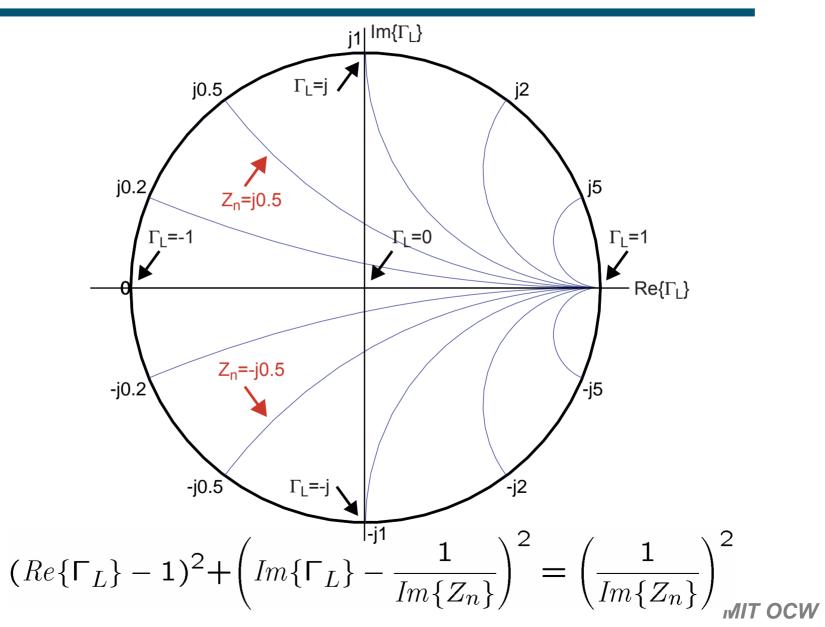
Equate real and imaginary parts to get Smith Chart

Real Impedance in Γ Coordinates (Equate Real Parts)



$$\left(Re\{\Gamma_L\} - \frac{Re\{Z_n\}}{1 + Re\{Z_n\}}\right)^2 + (Im\{\Gamma_L\})^2 = \left(\frac{1}{1 + Re\{Z_n\}}\right)^2$$
W.H. Periou

Imag. Impedance in Γ Coordinates (Equate Imag. Parts)



M.H. Perrott

What Happens When We Invert the Impedance?

Fundamental formulas

$$Z_n = \frac{1 + \Gamma_L}{1 - \Gamma_L} \Rightarrow \Gamma_L = \frac{Z_n - 1}{Z_n + 1}$$

Impact of inverting the impedance

$$Z_n \to 1/Z_n \Rightarrow \Gamma_L \to -\Gamma_L$$

Derivation:

$$\frac{1/Z_n - 1}{1/Z_n + 1} = \frac{1 - Z_n}{1 + Z_n} = -\left(\frac{Z_n - 1}{Z_n + 1}\right)$$

- We can invert complex impedances in Γ plane by simply changing the sign of Γ !
- How can we best exploit this?

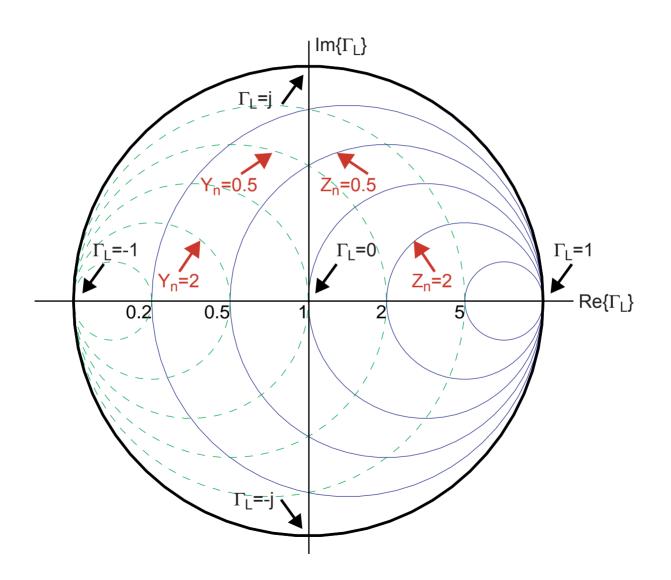
The Smith Chart as a Calculator for Matching Networks

 Consider constructing both impedance and admittance curves on Smith chart

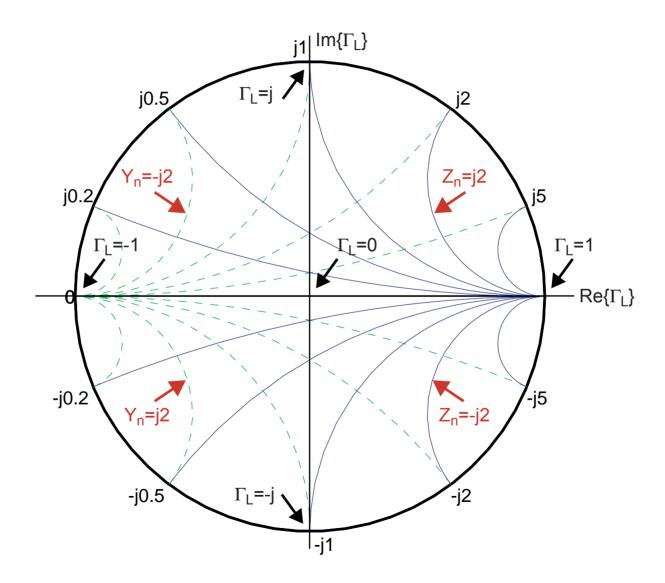
$$Z_n \to 1/Z_n \Rightarrow \Gamma_L \to -\Gamma_L$$

- Conductance curves derived from resistance curves
- Susceptance curves derived from reactance curves
- For series circuits, work with impedance
 - Impedances add for series circuits
- For parallel circuits, work with admittance
 - Admittances add for parallel circuits

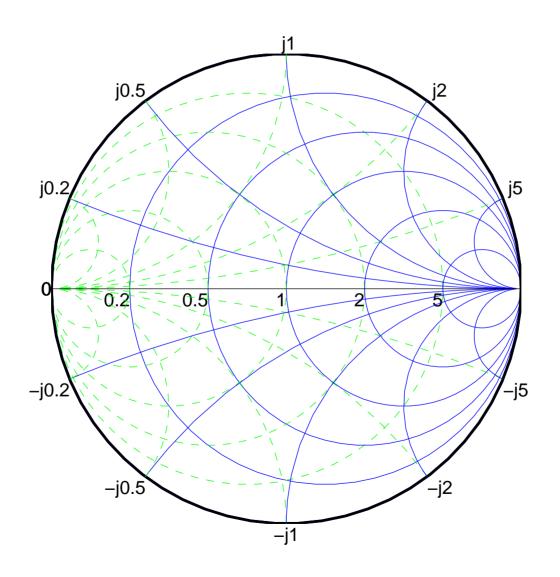
Resistance and Conductance on the Smith Chart



Reactance and Susceptance on the Smith Chart

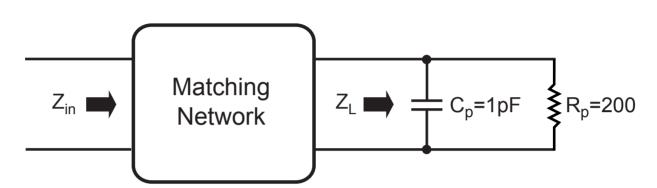


Overall Smith Chart



Example – Match RC Network to 50 Ohms at 2.5 GHz

Circuit

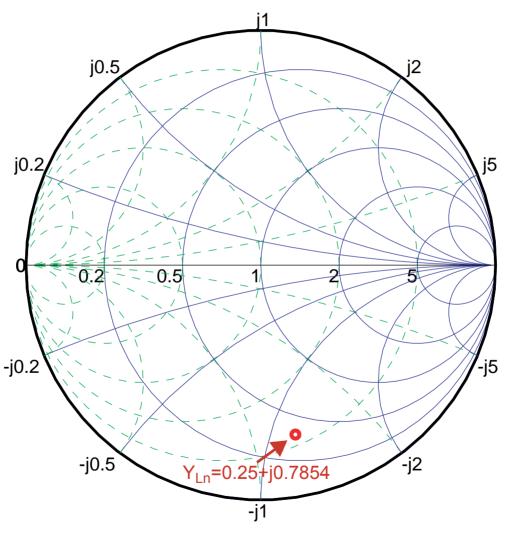


Step 1: Calculate Z_{Ln}

$$Z_{Ln} = \frac{Z_L}{Z_o} = \frac{R_L || (1/jwC)}{50} = \frac{1}{50(1/R_L + jwC)}$$
$$= \frac{1}{50(1/200 + j2\pi(2.5e9)10^{-12})} = \frac{1}{0.25 + j.7854}$$

Step 2: Plot Z_{Ln} on Smith Chart (use admittance, Y_{Ln})

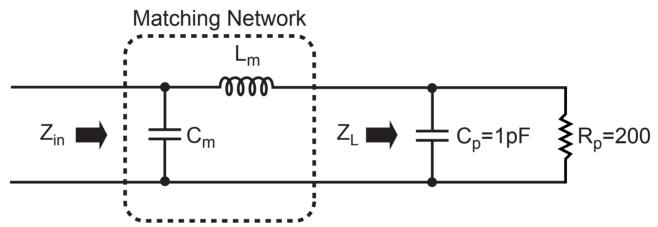
Plot Starting Impedance (Admittance) on Smith Chart



(Note: Z_{Ln} =0.37-j1.16)

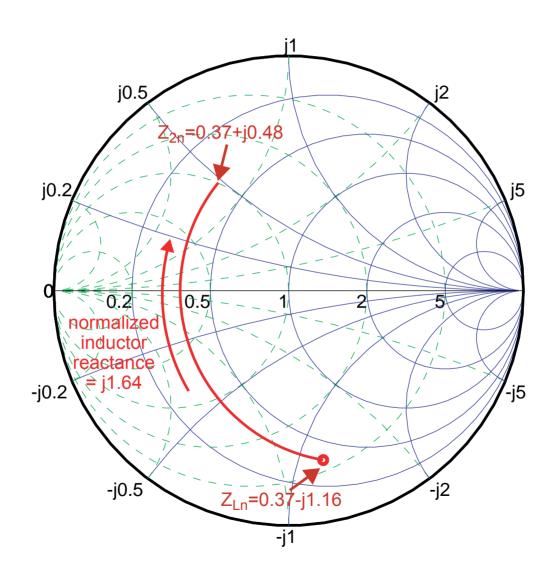
Develop Matching "Game Plan" Based on Smith Chart

By inspection, we see that the following matching network can bring us to Z_{in} = 50 Ohms (center of Smith chart)



- Use the Smith chart to come up with component values
 - Inductance L_m shifts impedance up along reactance curve
 - Capacitance C_m shifts impedance down along susceptance curve

Add Reactance of Inductor L_m



Inductor Value Calculation Using Smith Chart

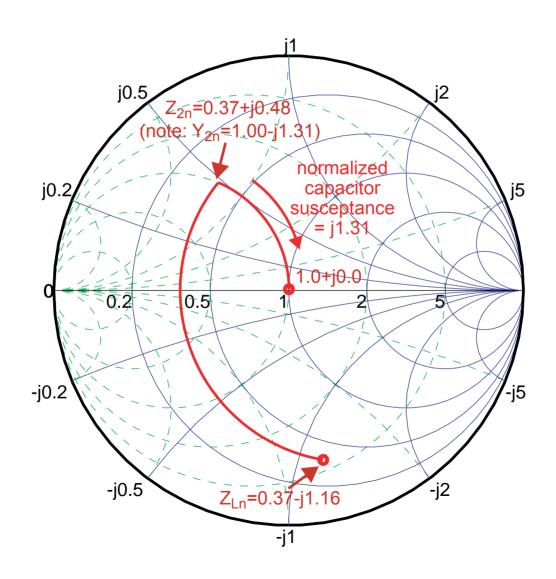
 From Smith chart, we found that the desired normalized inductor reactance is

$$\frac{jwL_m}{Z_0} = \frac{jwL_m}{50} = j1.64$$

Required inductor value is therefore

$$\Rightarrow L_m = \frac{50(1.64)}{2\pi 2.5e9} = 5.2nH$$

Add Susceptance of Capacitor C_m (Achieves Match!)



Capacitor Value Calculation Using Smith Chart

From Smith chart, we found that the desired normalized capacitor susceptance is

$$Z_{o}jwC_{m} = 50jwC_{m} = j1.31$$

Required capacitor value is therefore

$$\Rightarrow C_m = \frac{1.31}{50(2\pi 2.5e9)} = 1.67pF$$

Just For Fun

Play the "matching game" at

http://contact.tm.agilent.com/Agilent/tmo/an-95-1/classes/imatch.html

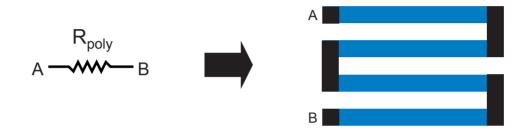
- Allows you to graphically tune several matching networks
- Note: game is set up to match source to load impedance rather than match the load to the source impedance

Same results, just different viewpoint



Polysilicon Resistors

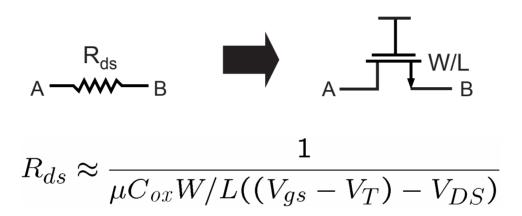
Use unsilicided polysilicon to create resistor



- Key parameters
 - Resistance (usually 100- 200 Ohms per square)
 - Parasitic capacitance (usually small)
 - Appropriate for high speed amplifiers
 - Linearity (quite linear compared to other options)
 - Accuracy (usually can be set within \pm 15%)

MOS Resistors

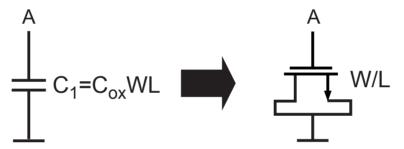
Bias a MOS device in its triode region



- High resistance values can be achieved in a small area (MegaOhms within tens of square microns)
- Resistance is quite nonlinear
 - Appropriate for small swing circuits

High Density Capacitors (Biasing, Decoupling)

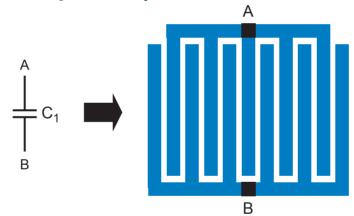
- MOS devices offer the highest capacitance per unit area
 - Limited to a one terminal device
 - Voltage must be high enough to invert the channel



- Key parameters
 - Capacitance value
 - Raw cap value from MOS device is 6.1 fF/μ m² for 0.24u
 CMOS
 - Q (i.e., amount of series resistance)
 - Maximized with minimum L (tradeoff with area efficiency)
- See pages 39-40 of Tom Lee's book

High Q Capacitors (Signal Path)

- Lateral metal capacitors offer high Q and reasonably large capacitance per unit area
 - Stack many levels of metal on top of each other (best layers are the top ones), via them at maximum density

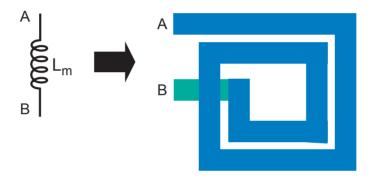


- lacktriangle Accuracy often better than $\pm 10\%$
- Parasitic side cap is symmetric, less than 10% of cap value
- **Example:** $C_T = 1.5$ fF/μm² for 0.24μm process with 7 metals, $L_{min} = W_{min} = 0.24$ μm, $t_{metal} = 0.53$ μm
 - See "Capacity Limits and Matching Properties of Integrated Capacitors", Aparicio et. al., JSSC, Mar 2002

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Spiral Inductors

 Create integrated inductor using spiral shape on top level metals (may also want a patterned ground shield)

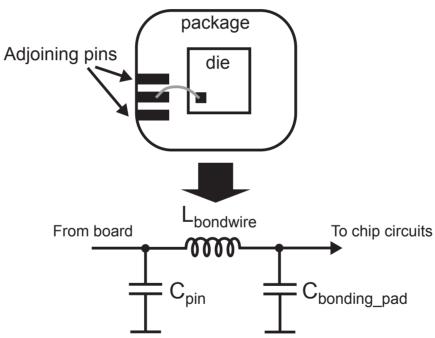


- Key parameters are Q (< 10), L (1-10 nH), self resonant freq.</p>
- Usually implemented in top metal layers to minimize series resistance, coupling to substrate
- Design using Mohan et. al, "Simple, Accurate Expressions for Planar Spiral Inductances, JSSC, Oct, 1999, pp 1419-1424

Verify inductor parameters (L, Q, etc.) using ASITIC http://formosa.eecs.berkeley.edu/~niknejad/asitic.html

Bondwire Inductors

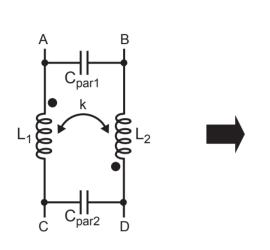
- Used to bond from the package to die
 - Can be used to advantage

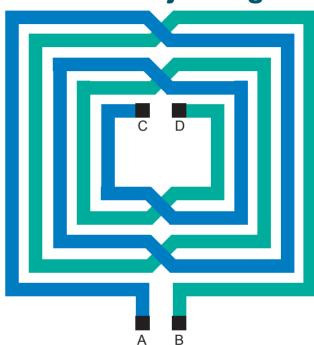


- Key parameters
 - Inductance (\approx 1 nH/mm usually achieve 1-5 nH)
 - Q (much higher than spiral inductors typically > 40)

Integrated Transformers

Utilize magnetic coupling between adjoining wires





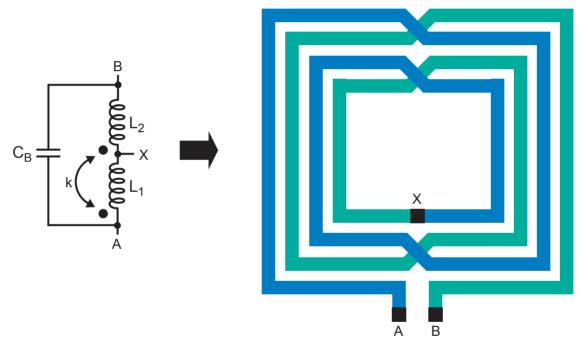
- Key parameters
 - L (self inductance for primary and secondary windings)
 - k (coupling coefficient between primary and secondary)

Note: $k = \frac{M}{\sqrt{L_1 L_2}}$ where M = mutual inductance

Design – ASITIC, other CAD packages

High Speed Transformer Example – A T-Coil Network

 A T-coil consists of a center-tapped inductor with mutual coupling between each inductor half



- Used for bandwidth enhancement
 - See S. Galal, B. Ravazi, "10 Gb/s Limiting Amplifier and Laser/Modulator Driver in 0.18u CMOS", ISSCC 2003, pp 188-189 and "Broadband ESD Protection ...", pp. 182-183