

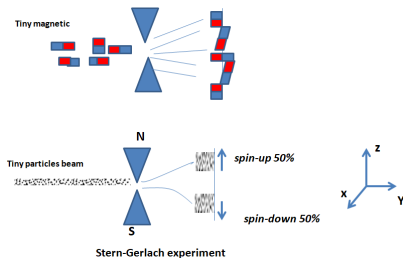
Quantum Computing (量子计算)

Get ready to grasp the leading edge technologies

Driss BOUTAT

Institut National des Sciences Appliquées Centre Val de Loire .

Spin operator¹: Stern–Gerlach experiment

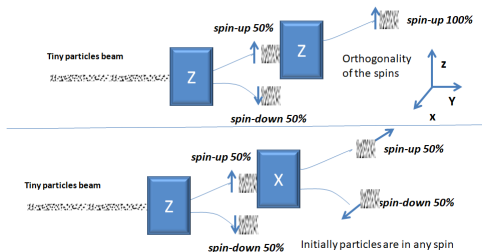


which corresponds to the following operator σ_z .

$\sigma_z = Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ has two eigenvalues 1 and -1 with eigenstates $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ spin-up state and $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ spin-down state.

Spin operator: Stern–Gerlach experiment

If we act a second time using Z then nothing changes because we act on its eigenstate (first figure below):



Let now acts on the spin-up eigenstate of Z using the following operator (second figure above)

$$\sigma_x = X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

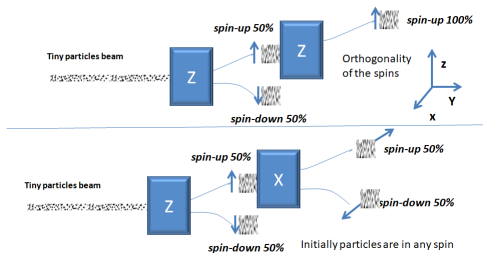
it splits this state on two new eigenstates of X

Spin operator: Stern–Gerlach experiment

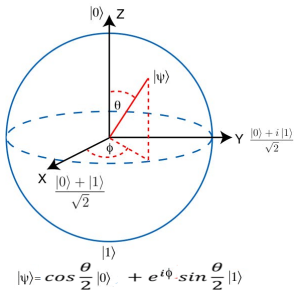
corresponds to eigenstates as linear combination of $|0\rangle$ and $|1\rangle$

$$\langle + | = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle), \quad \langle - | = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

with eigenvalues 1 and -1 .



Spin states: Bloch sphere



Each state in Bloch sphere is a linear combination of $|0\rangle$ and $|1\rangle$

$$\langle\psi| = \cos(\frac{\theta}{2})|0\rangle + e^{i\phi}\sin(\frac{\theta}{2})|1\rangle$$

and its orthogonal state

$$\langle\psi| = \sin(\frac{\theta}{2})|0\rangle - e^{i\phi}\cos(\frac{\theta}{2})|1\rangle$$

A class of transformations

Each state in Bloch sphere is a linear combination of $|0\rangle$ and $|1\rangle$

$$\langle\psi| = \cos\left(\frac{\theta}{2}\right) |0\rangle + e^{i\phi} \sin\left(\frac{\theta}{2}\right) |1\rangle$$

and we have the general unitary transformation (酉变换)

$$U(\theta, \phi, \lambda) = \begin{pmatrix} \cos \frac{\theta}{2} & -e^{i\lambda} \sin \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} & e^{i(\phi+\lambda)} \cos \frac{\theta}{2} \end{pmatrix}$$

Transpose Conjugate (共轭转置)

$$U^*(\theta, \phi, \lambda) = \begin{pmatrix} \cos \frac{\theta}{2} & e^{-i\phi} \sin \frac{\theta}{2} \\ -e^{-i\lambda} \sin \frac{\theta}{2} & e^{-i(\phi+\lambda)} \cos \frac{\theta}{2} \end{pmatrix} \text{ with } UU^* = I_d$$

thus U is hermitian.

Pauli operators²

Three important matrices (矩阵) in quantum computation

$$\sigma_x = X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix},$$
$$\sigma_z = Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

They fulfil the following relationships

$$\sigma_i^2 = I_d, \quad i = x, y, z$$
$$[\sigma_x, \sigma_y] = 2i\sigma_z, \quad [\sigma_y, \sigma_z] = 2i\sigma_x, \quad [\sigma_z, \sigma_x] = 2i\sigma_y.$$

⁵Pauli 算符

Rotation about axes : X, Y and Z

Let σ be an unitary hermitian operator : $\sigma^* = \sigma$ and $\sigma^2 = I_d$ then

$$e^{\frac{-i\theta\sigma}{2}} = \cos\left(\frac{\theta}{2}\right) - i \sin\left(\frac{\theta}{2}\right)\sigma$$

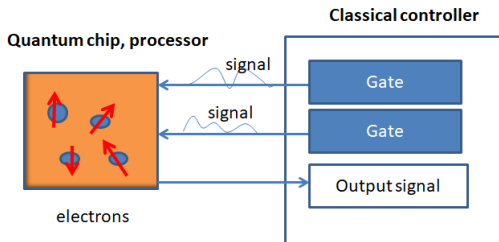
The rotation operator on the Bloch sphere axes are given

$$R_X = e^{\frac{-i\theta\sigma_x}{2}} = \begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ -i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}, \quad (1)$$

$$R_Y = e^{\frac{-i\theta\sigma_y}{2}} = \begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}, \quad (2)$$

$$R_Z = e^{\frac{-i\theta\sigma_z}{2}} = \begin{pmatrix} e^{\frac{-i\theta}{2}} & 0 \\ 0 & e^{\frac{i\theta}{2}} \end{pmatrix}. \quad (3)$$

Simplified diagram of quantum computer



In the controlled part each wire that support a gate is as follows



Qubits³



One wire support unitary complex numbers of the vector space \mathbb{C} endowed with the basis:

$$|0\rangle := \begin{pmatrix} 1 \\ 0 \end{pmatrix} : \text{spin up}$$

$$|1\rangle := \begin{pmatrix} 0 \\ 1 \end{pmatrix} : \text{spin down}$$

The canonical basis is formed by $|0\rangle$ and $|1\rangle$. Thus the state of the qubit are linear combination of them, for example

$$|\psi\rangle = a|0\rangle + b|0\rangle,$$

where a and b are complex numbers such that

$$|a|^2 + |b|^2 = 1$$

Not X gate⁴

To change the initial state $|0\rangle$ and $|1\rangle$ into the state $|1\rangle$ and $|0\rangle$ respectively, we use the NotX

For example the Not-gate (Pauli X gate) that swaps the two vectors of the basis

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \oplus$$



$$X |0\rangle = |1\rangle, \quad X |1\rangle = |0\rangle$$

Two wires

The two states on two wires (电路线) give rise to the tensor product (张量积).

Two qubits : two wires = tensor product. $|\alpha\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}, |\beta\rangle = \begin{pmatrix} \gamma \\ \delta \end{pmatrix} \quad |\alpha\rangle|\beta\rangle = |\alpha\rangle \otimes |\beta\rangle = |\alpha\beta\rangle = \begin{pmatrix} \alpha\gamma \\ \alpha\delta \\ \beta\gamma \\ \beta\delta \end{pmatrix}$
tensor product

