

# Exercises from Chapter 1

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**Problem 1.1.** Beginning with the formula for the sum of a geometric series, use differentiation to obtain the identity.

$$\sum_{n=0}^{\infty} ne^{-An} = \frac{e^{-A}}{(1 - e^{-A})^2}$$

*Solution.* First, we integrate the left side to get the expression into geometric series form.

$$\int \sum_{n=0}^{\infty} ne^{-An} dA = \sum_{n=0}^{\infty} -e^{-An}$$

Next, recall that the sum of a geometric series is  $\sum_{n=0}^{\infty} ar^n = \frac{a}{1-r}$ . Using this fact, the previous result can be rewritten,

$$\frac{-1}{1 - e^A}$$

Finally, we take the derivative with respect to A, to undo our previous integration.

$$\frac{d}{dA} \frac{-1}{1 - e^A} = \frac{e^{-A}}{(1 - e^{-A})^2}$$

□

**Problem 1.2.** In Planck's model of blackbody radiation, the energy in a given frequency  $\omega$  of electromagnetic radiation is distributed randomly over all numbers of the form  $n\hbar\omega$ , where  $n = 0, 1, 2, \dots$ . Specifically, the likelihood of finding energy  $n\hbar\omega$  is postulated to be

$$p(E = n\hbar\omega) = \frac{1}{Z} e^{-\beta n\hbar\omega},$$
$$Z = \frac{1}{1 - e^{-\beta\hbar\omega}}$$

Where  $Z$  is a normalization constant, which is chosen so that the sum over  $n$  of the probabilities is 1. Here  $\beta = \frac{1}{k_B T}$ , where  $T$  is the temperature and  $k_B$  is Boltzman's constant. The expected value of the energy, denoted  $\langle E \rangle$ , is defined to be

$$\langle E \rangle = \frac{1}{Z} \sum_{n=0}^{\infty} (n\hbar\omega) e^{-\beta n\hbar\omega}.$$

(a) using exercise 1, show that

$$\langle E \rangle = \frac{\hbar\omega}{e^{\beta\hbar\omega} - 1}.$$

(b) Show that  $\langle E \rangle$  behaves like  $\frac{1}{\beta} = k_B T$  for small  $\omega$ , but that  $\langle E \rangle$  decays exponentially as  $\omega$  tends to infinity.

*Proof.* (a) From exercise 1, we can rewrite the sum term as

$$\frac{\hbar\omega e^{-\hbar\omega\beta}}{(1 - e^{-\hbar\omega\beta})^2}$$

Multiply this by  $\frac{1}{Z} = 1 - e^{-\hbar\omega\beta}$ .

$$\begin{aligned} 1 - e^{-\hbar\omega\beta} & \frac{\hbar\omega e^{-\hbar\omega\beta}}{(1 - e^{-\hbar\omega\beta})^2} \\ &= \frac{\hbar\omega e^{-\hbar\omega\beta}}{1 - e^{-\hbar\omega\beta}} \\ &= \frac{\hbar\omega}{e^{\hbar\omega\beta} - 1} \end{aligned}$$

(b) Using the Taylor series expansion of  $e^x$ , we can rewrite  $\langle E \rangle$

$$\begin{aligned} & \frac{\hbar\omega}{e^{\hbar\omega\beta} - 1} \\ &= \frac{\hbar\omega}{(1 + \hbar\omega\beta + \sum_{n=2}^{\infty} \frac{(\hbar\omega\beta)^n}{n}) - 1} \\ &= \frac{1}{\beta + \sum_{n=2}^{\infty} \frac{(\hbar\omega)^{n-1}\beta^n}{n}} \end{aligned}$$

as  $\omega$  approaches 0, the sum disappears and the fraction approaches  $\frac{1}{\beta}$ .  
It is easy to find the limit of  $\langle E \rangle$  as  $\omega$  approaches  $\infty$  using L'Hospital's rule

$$\frac{\hbar\omega}{e^{\hbar\omega\beta} - 1} \stackrel{\text{H}}{=} \frac{\hbar}{\hbar\beta e^{\hbar\omega\beta}}$$

□