# A constrained Fermi-Pasta-Ulam-Tsingou problem

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## Introduction

The goal is to establish a non linear stiff DAE toy model to test the behaviour of the corrected discrete gradient on a DAE. We take the well-known Fermi-Pasta-Ulam-Tsingou (FPUT) problem and propose a simple DAE version with kinematic and holonomic constraints. We then implement a simulation without simplifying the constraints.

#### 1 The model

For symbolic computations, see https://github.com/WetzelVictor/FPUT\_DAE\_PHS

We consider a 1-dimensional chain of :  $N_m$  mass points and  $N_k$  springs;  $N_m = 2 * N_k$ . First, we describe a cell, comprised of 2 masses and 1 spring. We then connect all the subsystems using kinematic constraints between the masses.

#### Sub-system 1.1

You may change the springs law in order to make the systelm non linear:

$$\mathbf{X} = \begin{bmatrix} \Pi_{iL} & \Pi_{iR} & x_{ki} \end{bmatrix}^{\mathsf{T}} \tag{1}$$

$$\mathbf{X} = [\Pi_{iL} \quad \Pi_{iR} \quad x_{ki}]^{\mathsf{T}}$$

$$\mathbf{H}(\mathbf{X}) = \frac{\Pi_{iL}^{2}}{2m_{iL}} + \frac{\Pi_{iR}^{2}}{2m_{iR}} + \frac{1}{2}x_{ki}^{2}k_{i}$$
(2)

$$\theta = \begin{bmatrix} m_{iL} & m_{iR} & k_i \end{bmatrix} \tag{3}$$

**Hypothesis**: here,  $m_{iL} \ll m_{iR}$ .

$$\begin{pmatrix}
\dot{\Pi}_{iL} \\
\dot{\Pi}_{iR} \\
\dot{x}_{ki}
\end{pmatrix} = \begin{bmatrix}
\cdot & \cdot & 1 \\
\cdot & \cdot & -1 \\
-1 & 1 & \cdot
\end{bmatrix} \cdot \begin{pmatrix}
v_{m_{i,L}} \\
v_{m_{i,R}} \\
F_{ki}(x_{ki})
\end{pmatrix}$$
(4)

#### DAE-pHs 1.2

See [2] for more details:

Dynamic Outputs 
$$\begin{cases} \dot{\mathbf{X}} &= [(\mathbf{J}(\mathbf{X}) - \mathbf{R}(\mathbf{X})] \nabla_{\mathbf{X}} \mathbf{H}(\mathbf{X}) + \mathbf{G} \mathbf{u} + \mathbf{B} \lambda \\ \mathbf{y} &= \mathbf{G}^{\mathsf{T}} \nabla_{\mathbf{X}} \mathbf{H}(\mathbf{X}) + \mathbf{G}_{yy} \mathbf{u} \\ 0 &= \mathbf{B}^{\mathsf{T}} \nabla_{\mathbf{X}} \mathbf{H}(\mathbf{X}) \end{cases}$$
(5)

where:

— **X** and  $\nabla_{\mathbf{X}}\mathbf{H}(\mathbf{X})$  are vectors of  $N_x$  elements

—  $\mathbf{J}(\mathbf{X})$  and  $\mathbf{R}(\mathbf{X})$  are  $N_x \times N_x$  matrices

— **u** and **y** are vector of  $N_p$  elements

— **G** is a  $N_x \times N_p$  matrix

—  $\lambda$  is the vector of the  $N_{\lambda}$  Lagrange multipliers

— **B** is a  $N_x \times N_\lambda$  matrix

—  $\mathbf{G}_{yy}$  is a  $N_p \times N_p$  matrix

#### 1.3 Assembly

The kinematic constraints between system i and i + 1 is:

$$0 = v_{m_{i,R}} - v_{m_{(i+1),L}} \tag{6}$$

We assume the state vector of the full system follows the following order:

$$\mathbf{X} = \begin{bmatrix} \Pi_{1L} & \Pi_{1R} & \dots & \Pi_{N_sL} & \Pi_{N_sR} & ; & x_{k1} & \dots & x_{kN_s} \end{bmatrix}^\mathsf{T}$$

The transposed constraint matrix  $\mathbf{B}^\intercal$  reads :

$$\mathbf{B}_{ij}^{\mathsf{T}} = \begin{cases} 1 & \text{if } i = 2j \\ -1 & \text{if } i = 2j+1 \\ 0 & \text{otherwise} \end{cases}$$
 (7)

### 1.4 Variable change

We use the method also described in [1]. The left annihilator  $\mathbf{B}^{\perp}$  of  $\mathbf{B}$  is:

The variable change is:

$$\mathbf{M} = \begin{bmatrix} \mathbf{B}^{\perp} \\ \mathbf{B}^{\star} = (\mathbf{B}^{\mathsf{T}}\mathbf{B})^{-1}\mathbf{B}^{\mathsf{T}} \end{bmatrix}$$
 (9)

where  $\mathbf{B}^{\star}$  is the Moore-Penrose inverse of  $\mathbf{B}$ .

#### 1.5 New variables

Denoting  $\mathbf{Z} = \mathbf{M}\mathbf{X}$  the new states, we have :

$$\mathbf{Z} = \begin{bmatrix} \Pi_{L1} \\ \Pi_{12} \\ \vdots \\ \Pi_{(N_s-1)N_s} \\ \Pi_{RN_s} \\ x_{k1} \\ \vdots \\ x_{kN_s} \end{bmatrix} = \begin{bmatrix} \Pi_{L1} \\ \Pi_{R1} + \Pi_{L2} \\ \vdots \\ \Pi_{R(N_s-1)} + \Pi_{LN_s} \\ \Pi_{RN_s} \\ x_{k1} \\ \vdots \\ x_{kN_s} \end{bmatrix}$$

$$\frac{\Delta \Pi_{12}}{\vdots}$$

$$\vdots$$

$$\Delta \Pi_{(N_s-1)N_s} \end{bmatrix}$$

$$\frac{\Pi_{L1}}{\Pi_{R1} + \Pi_{L2}}$$

$$\vdots$$

$$\frac{\pi_{R(N_s-1)} - \Pi_{LN_s}}{2}$$

### 1.6 Hamiltonian

We assume springs are linear for now, to simplify notations. Therefore the Hamiltonian can be expressed in its quadratic form:

$$\mathbf{H}(\mathbf{X}) = \frac{1}{2} \mathbf{X}^{\mathsf{T}} \mathbf{Q} \mathbf{X}$$
 with  $\mathbb{Q} = \operatorname{diag} \begin{pmatrix} m_{1L}^{-1} & m_{1R}^{-1} & \dots & m_{N_sL}^{-1} & m_{N_sR}^{-1} & ; x_{k1} & \dots & x_{kN_s} \end{pmatrix}$ 

With the variable change, the Hamiltonian reads :

$$H(\mathbf{X} = \mathbf{M}^{-1}\mathbf{Z}) = \frac{1}{2}\mathbf{Z}^{\mathsf{T}}\underbrace{\mathbf{M}_{i}^{-\mathsf{T}}\mathbb{Q}^{(1)}\mathbf{M}_{i}^{-1}}_{\mathbb{Q}^{(2)}}\mathbf{Z}$$

#### 1.7 Constraints

Let us split **Z** into:

- $\mathbf{Z}_1$  the unconstrained states, which dynamic is determined by on ODE
- $\mathbf{Z}_2$  the constrained states, which value is computed by solving an algebraic system (variables  $\Delta_{i(i+1)} \forall i \in [1, N_s 1]$ )

After applying the variable change to the system, the  $N_{\lambda}$  constraints are given by :

$$0 = \frac{\partial \mathbf{H}(\mathbf{Z})}{\partial (\mathbf{Z}_2) i}$$

For each  $\mathbf{Z}_{2i}$ , the constraint is:

$$0 = \Delta \Pi_{i(i+1)} \left( \frac{1}{m_{Ri}} + \frac{1}{m_{L(i+1)}} \right) + \Pi_{i(i+1)} \left( \frac{1}{2m_{Ri}} - \frac{1}{m_{2L(i+1)}} \right)$$
(10)

### 1.8 Matrix formulation

Let us divide  $\mathbb{Q}^{(2)}$  accordingly to **Z**, the constraints system is:

find 
$$\mathbf{Z}_2$$
 such as 
$$[\mathbb{Q}^{(2)}]_{12}\mathbf{Z}_1 + [\mathbb{Q}^{(2)}]_{22}\mathbf{Z}_2 = 0$$
 where  $\mathbb{Q} = \begin{bmatrix} [\mathbb{Q}^{(2)}]_{11} & [\mathbb{Q}^{(2)}]_{12} \\ [\mathbb{Q}^{(2)}]_{12} & [\mathbb{Q}^{(2)}]_{22} \end{bmatrix}$ 

### Références

- [1] Flavio Cardoso Ribeiro. "Port-hamiltonian Modeling and Control of a Fluid-structure System: Application to Sloshing Phenomena in a Moving Container Coupled to a Flexible Structure." Thèse de doct. Institut Supérieur de l'Aéronautique et de l'Espace, déc. 2016. URL: https://depozit.isae.fr/theses/2016/2016\_Cardoso\_Ribeiro\_Flavio\_Luiz\_D.pdf.
- [2] A. J. van der Schaft. "Port-Hamiltonian Differential-Algebraic Systems". In: Surveys in Differential-Algebraic Equations I. Sous la dir. d'Achim Ilchmann et Timo Reis. Berlin, Heidelberg: Springer Berlin Heidelberg, 2013, p. 173-226. ISBN: 978-3-642-34928-7. Doi: 10.1007/978-3-642-34928-7\_5. URL: https://doi.org/10.1007/978-3-642-34928-7\_5.