Instruction 1: warming up

Course: Partial differential equation (3341), 2020/21

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A. Ordinary differential equations

- 1. Linear, first order equations. Use integrating factors for finding the general solution to the equations below $(a \in \mathbf{R} \text{ is a constant})$. If applicable, find the solution that satisfies the initial condition:
 - a) u'(t) + a t u(t) = 0;

Solution. Standard Form (S.F.):

$$u'(t) + a t u(t) = 0. (1)$$

Integrating Factor (I.F.):

$$I.F. = e^{\int a \, t \, dt} = e^{\frac{a \, t^2}{2}}.$$
 (2)

Remark: Using product rule, chain rule and fundamental theorem of calculus, we can write

$$\frac{d}{dt} \left(e^{\frac{a t^2}{2}} u(t) \right) = u(t) \frac{d}{dt} \left(e^{\frac{a t^2}{2}} \right) + e^{\frac{a t^2}{2}} \frac{d}{dt} u(t),$$

$$= u(t) e^{\frac{a t^2}{2}} \frac{d}{dt} \left(\frac{a t^2}{2} \right) + e^{\frac{a t^2}{2}} u'(t),$$

$$= u(t) e^{\frac{a t^2}{2}} \left(\frac{2 a t}{2} \right) + e^{\frac{a t^2}{2}} u'(t),$$

$$= a t u(t) e^{\frac{a t^2}{2}} + e^{\frac{a t^2}{2}} u'(t),$$

$$= e^{\frac{a t^2}{2}} \left(u'(t) + a t u(t) \right).$$
(3)

Multiplying S.F. by I.F.:

$$e^{\frac{at^2}{2}} \left(u'(t) + a t u(t) \right) = 0,$$

$$\Longrightarrow e^{\frac{at^2}{2}} \frac{d}{dt} u(t) + u(t) e^{\frac{at^2}{2}} \frac{d}{dt} \left(\frac{at^2}{2} \right) = 0,$$

$$\Longrightarrow \frac{d}{dt} \left(e^{\frac{at^2}{2}} u(t) \right) = 0.$$
(4)

Integration with respect to (w.r.t.) t:

$$\int \frac{d}{dt} \left(e^{\frac{at^2}{2}} u(t) \right) dt = \int 0 dt,$$

$$\implies e^{\frac{at^2}{2}} u(t) = C,$$
where C is a integrating constant,
$$\implies u(t) = C e^{-\frac{at^2}{2}}.$$
(5)

b) u'(t) + a t u(t) = t;

Solution. Standard Form (S.F.):

$$u'(t) + a t u(t) = t. (6)$$

Integrating Factor (I.F.):

$$I.F. = e^{\int a t dt} = e^{\frac{a t^2}{2}}.$$

Multiplying S.F. by I.F.:

$$e^{\frac{at^2}{2}} \left(u'(t) + a t u(t) \right) = t e^{\frac{at^2}{2}},$$

$$\Longrightarrow e^{\frac{at^2}{2}} \frac{d}{dt} u(t) + u(t) e^{\frac{at^2}{2}} \frac{d}{dt} \left(\frac{at^2}{2} \right) = t e^{\frac{at^2}{2}},$$

$$\Longrightarrow \frac{d}{dt} \left(e^{\frac{at^2}{2}} u(t) \right) = t e^{\frac{at^2}{2}}.$$
(7)

Integration with respect to (w.r.t.) t:

$$\int \frac{d}{dt} \left(e^{\frac{at^2}{2}} u(t) \right) dt = \int t e^{\frac{at^2}{2}} dt,$$

$$\implies e^{\frac{at^2}{2}} u(t) = \int \frac{1}{a} \frac{d}{dt} \left(e^{\frac{at^2}{2}} \right) dt + C,$$
where C is a integrating constant,
$$\implies e^{\frac{at^2}{2}} u(t) = \frac{1}{a} e^{\frac{at^2}{2}} + C,$$

$$\implies u(t) = \frac{1}{a} + C e^{-\frac{at^2}{2}}.$$
(8)

c) $t^2 u'(t) + 2 t u(t) = t^2 \text{ with } u(1) = 1;$

Solution. Standard Form (S.F.):

$$u'(t) + \frac{2}{t}u(t) = 1. (9)$$

Integrating Factor (I.F.):

$$I.F. = e^{\int \frac{2}{t} dt} = e^{2 \ln(t)} = e^{\ln(t^2)} = t^2. \tag{10}$$

Multiplying S.F. by I.F.:

$$t^{2}\left(u'(t) + \frac{2}{t}u(t)\right) = t^{2},$$

$$\Rightarrow t^{2}\frac{d}{dt}u(t) + u(t) \ 2 \ t = t^{2},$$

$$\Rightarrow t^{2}\frac{d}{dt}u(t) + u(t) \ \frac{d}{dt}\left(t^{2}\right) = t^{2},$$

$$\Rightarrow \frac{d}{dt}\left(t^{2}u(t)\right) = t^{2}.$$
(11)

Integration with respect to (w.r.t.) t:

$$\int \frac{d}{dt} (t^2 u(t)) dt = \int t^2 dt,$$

$$\implies t^2 u(t) = \frac{t^3}{3} + C,$$
where C is a integrating constant,
$$\implies u(t) = \frac{t}{3} + \frac{C}{t^2},$$
(12)

which is the general solution.

Particular solution: At t=1, we can write,

$$u(t=1) = \frac{1}{3} + C,$$
Given, initial condition is $u(1) = 1$,
$$\implies 1 = \frac{1}{3} + C,$$

$$\implies C = \frac{2}{3}.$$
(13)

Hence the particular solution is $u(t) = \frac{t}{3} + \frac{2}{3t^2}$.

d) $t^2 u'(t) + a t u(t) = t^2 \text{ with } u(0) = 1.$

Solution. General solution:

$$u(t) = \frac{t}{3} + \frac{C}{t^2},$$

$$\Longrightarrow C = t^2 u(t) - \frac{t^3}{3}.$$
(14)

Particular solution: At t=0,

$$C = 0^2 u(t)|_{t=0} - \frac{0^3}{3} = 0.$$
 (15)

Hence the particular solution is $u_p(t) = \frac{t}{3}$.

Since given initial condition is u(0) = 1 and at t = 0, $u_p(0) = 0$, which is a contradiction. Hence we can not find a particular solution with the given initial condition u(0) = 1.

Remark:

$$\frac{d}{dt}u(t) = 1/3 - \frac{C}{t} \tag{16}$$

which is not defined for t = 0. This is why we can not find a particular solution at t = 0, with u(0) = 1.

2. Linear, second order equations with constant coefficients. Find the solution to the equations below:

a)
$$u''(t) - 3 u'(t) + 2 u(t) = 0$$
 with $u(0) = 1$ and $u'(1) = 1$;

Solution. Characteristic equation: We have

$$u''(t) - 3 u'(t) + 2 u(t) = 0. (17)$$

Let us consider that

$$u(t) = e^{r t},$$

is a solution of (17). Then using

$$u'(t) = r e^{r t}, \ u''(t) = r^2 e^{r t},$$
 (18)

in (17), we can write

$$r^{2} e^{r t} - 3 r e^{r t} + 2 e^{r t} = 0,$$

$$\implies e^{r t} (r^{2} - 3 r + 2) = 0.$$
(19)

Hence the characteristic equation can be written as

$$r^2 - 3 r + 2 = 0, (20)$$

with a=-3, b=2.

Here $\Delta = a^2 - 4 \ b = 1 > 0$.

General Solution: The roots of (20) are

$$r_1 = \frac{-a + \sqrt{\Delta}}{2} = 2,$$

$$r_2 = \frac{-a - \sqrt{\Delta}}{2} = 1.$$
(21)

which are real and distinct.

Then we can write the general solution as

$$u(t) = c_1 e^{r_1 t} + c_2 e^{r_2 t},$$

$$\implies u(t) = c_1 e^{2t} + c_2 e^{t},$$
(22)

where c_1, c_2 are constant.

Here

$$u'(t) = 2 c_1 e^{2t} + c_2 e^t. (23)$$

Particular solution: Given initial conditions are u(0) = 1 and u'(1) = 1.

Using u(t=0)=1 in (22), we can write

$$u(0) = c_1 e^0 + c_2 e^0,$$

$$\Longrightarrow 1 = c_1 + c_2,$$

$$\Longrightarrow c_1 = 1 - c_2.$$
(24)

Using u'(t=1) = 1 in (23), we obtain

$$u'(1) = 2 c_1 e^2 + c_2 e,$$

$$\implies 1 = 2 (1 - c_2) e^2 + c_2 e,$$

$$\implies 1 - 2 e^2 = -2 c_2 e^2 + c_2 e,$$

$$\implies 1 - 2 e^2 = c_2 (-2 e^2 + e),$$

$$\implies c_2 = \frac{1 - 2 e^2}{e (1 - 2 e)}.$$
(25)

Using 25 in 24, one gets

$$c_{1} = 1 - \frac{1 - 2 e^{2}}{e (1 - 2 e)},$$

$$\implies c_{1} = \frac{e - 2 e^{2} - 1 + 2 e^{2}}{e (1 - 2 e)},$$

$$\implies c_{1} = \frac{(e - 1)}{e (1 - 2 e)}.$$
(26)

Hence, the particular solution is

$$u(t) = \frac{(e-1)}{e(1-2e)} e^{2t} + \frac{1-2e^2}{e(1-2e)} e^t.$$
 (27)

b) u''(t) - 2 u'(t) + 2 u(t) = 0 with u(0) = 1 and u'(0) = 0;

Solution. Characteristic equation: We have

$$u''(t) - 2 u'(t) + 2 u(t) = 0. (28)$$

Let us consider that

$$u(t) = e^{r t}$$

is a solution of (28). Then using

$$u'(t) = r e^{r t}, \ u''(t) = r^2 e^{r t},$$
 (29)

in (28), we can write

$$r^{2} e^{r t} - 2 r e^{r t} + 2 e^{r t} = 0,$$

$$\implies e^{r t} (r^{2} - 2 r + 2) = 0,$$
(30)

Hence the characteristic equation can be written as

$$r^2 - 2r + 2 = 0, (31)$$

with a= -2, b=2. Here $\Delta = a^2 - 4 b = -4 < 0$.

General Solution: The roots of (31) are

$$r_{1} = -\frac{a}{2} + \frac{i}{2}\sqrt{-\Delta} = 1 + i,$$

$$r_{2} = -\frac{a}{2} - \frac{i}{2}\sqrt{-\Delta} = 1 - i,$$
(32)

which are imaginary and distinct.

Then we can write the general solution as

$$u(t) = e^{-a t/2} \left(c_1 \cos(\frac{\sqrt{-\Delta}}{2} t) + c_2 \sin(\frac{\sqrt{-\Delta}}{2} t) \right),$$

$$\implies u(t) = e^t \left(c_1 \cos(t) + c_2 \sin(t) \right),$$
(33)

where c_1, c_2 are constant.

Here,

$$u'(t) = e^{t} (c_1 \cos(t) + c_2 \sin(t)) + e^{t} (-c_1 \sin(t) + c_2 \cos(t)).$$
(34)

Particular solution: Initial conditions are u(0) = 1 and u'(0) = 0. At t = 0, u(0) = 1 and from (33),

$$u(0) = e^{0} (c_{1} \cos(0) + c_{2} \sin(0)),$$

 $\implies 1 = c_{1} 1 + c_{2} 0,$
 $\implies c_{1} = 1.$ (35)

At t = 0, u'(0) = 0 and from (34),

$$u'(0) = e^{0} (c_{1} \cos(0) + c_{2} \sin(0)) + e^{0} (-c_{1} \sin(0) + c_{2} \cos(0)),$$

$$\implies 0 = c_{1} 1 + c_{2} 0 - c_{1} 0 + c_{2} 1,$$

$$\implies 0 = c_{1} + c_{2}$$

$$\implies c_{2} = -c_{1} = -1.$$
(36)

Hence, the particular solution is

$$u(t) = e^t \left(\cos(t) - \sin(t)\right). \tag{37}$$

c)
$$u''(t) + 4 u'(t) + 4 u(t) = 0$$
 with $u'(0) = 0$ and $u(1) = 1$;

Solution. Characteristic equation: We have

$$u''(t) + 4 u'(t) + 4 u(t) = 0. (38)$$

Let us consider that

$$u(t) = e^{r t},$$

is a solution of (38). Then using

$$u'(t) = r e^{r t}, \ u''(t) = r^2 e^{r t},$$
 (39)

in (38), we can write

$$r^{2} e^{rt} + 4 r e^{rt} + 4 e^{rt} = 0,$$

$$\Rightarrow e^{rt} (r^{2} + 4 r + 4) = 0.$$
(40)

Hence the characteristic equation can be written as

$$r^2 + 4 r + 4 = 0, (41)$$

with a=4, b=4.

Here $\Delta = a^2 - 4 \ b = 0$.

General Solution: The roots of (41) are

$$r_1 = r_2 = -\frac{a}{2} = -2, (42)$$

which are real and equal.

Then we can write the general solution as

$$u(t) = c_1 e^{-\frac{a}{2}} + c_2 t e^{-\frac{a}{2}},$$

$$\implies u(t) = e^{-2t} (c_1 + t c_2),$$
(43)

where c_1, c_2 are constant.

Here,

$$u'(t) = -2 e^{-2t} (c_1 + t c_2) + e^{-2t} (0 + c_2), \qquad (44)$$

$$\implies u'(t) = -2 c_1 e^{-2t} + c_2 (1 - 2t) e^{-2t}. \tag{45}$$

Particular solution: Given initial conditions are u'(0) = 0 and u(1) = 1.

At t = 1, u(1) = 1 and from (43),

$$u(1) = e^{-2} (c_1 + c_2),$$

$$\implies 1 = e^{-2} (c_1 + c_2),$$

$$\implies c_1 + c_2 = 1/e^{-2} = e^2,$$

$$\implies c_1 = e^2 - c_2.$$
(46)

At t = 0, u'(0) = 0 and from (44),

$$u'(0) = -2 c_1 e^0 + c_2 (1 - 0) e^0,$$

$$\implies 0 = -2 (e^2 - c_2) + c_2,$$

$$\implies c_2 = \frac{2 e^2}{3}.$$
(47)

Therefore,

$$c_1 = e^2 - \frac{2 e^2}{3} = \frac{e^2}{3}. (48)$$

Hence, the particular solution is

$$u(t) = e^{-2t} \left(\frac{e^2}{3} + t \frac{2e^2}{3} \right),$$

$$\implies u(t) = e^{-2t} (1 + 2t) \frac{e^2}{3}.$$
(49)

- B. Partial integration, Gauss/Divergece Theorem Let $\Omega \subset \mathbb{R}^d$ (d=2 of d=3) be a bounded domain with smooth boundary $\partial\Omega$ and $\vec{\nu}=(\nu_1,\ldots,\nu_d)^T\in\mathbb{R}^d$ the unit normal to $\partial\Omega$ pointing outwards Ω .
 - 1. Prove that for any $u \in C^1(\Omega)$ and $v \in C^2(\Omega)$ one has

$$\nabla \cdot (u \nabla v) = \nabla u \cdot \nabla v + u \Delta v.$$

Solution. Let us consider, d=3, then we can write

$$\vec{\nabla} = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right). \tag{50}$$

Hence,

$$\vec{\nabla}u = \left(\frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z}\right),$$

$$\vec{\nabla}v = \left(\frac{\partial v}{\partial x}, \frac{\partial v}{\partial y}, \frac{\partial v}{\partial z}\right),$$

$$\nabla u \cdot \nabla v = \left(\frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z}\right) \cdot \left(\frac{\partial v}{\partial x}, \frac{\partial v}{\partial y}, \frac{\partial v}{\partial z}\right) = \left(\frac{\partial u}{\partial x}, \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y}, \frac{\partial v}{\partial y} + \frac{\partial u}{\partial z}, \frac{\partial v}{\partial z}\right),$$

$$\Delta v = \nabla \cdot \nabla v = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right) \cdot \left(\frac{\partial v}{\partial x}, \frac{\partial v}{\partial y}, \frac{\partial v}{\partial z}\right) = \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2}\right),$$

$$u \vec{\nabla}v = \left(u \frac{\partial v}{\partial x}, u \frac{\partial v}{\partial y}, u \frac{\partial v}{\partial z}\right).$$
(51)

Therefore,

$$\vec{\nabla} \cdot \left(u \, \vec{\nabla} v \right) = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \cdot \left(u \, \frac{\partial v}{\partial x}, u \, \frac{\partial v}{\partial y}, u \, \frac{\partial v}{\partial z} \right),$$

$$= \frac{\partial}{\partial x} \left(u \, \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial y} \left(u \, \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial z} \left(u \, \frac{\partial v}{\partial z} \right),$$

$$= \left(\frac{\partial u}{\partial x} \frac{\partial v}{\partial x} + u \frac{\partial^2 v}{\partial x^2} \right) + \left(\frac{\partial u}{\partial y} \frac{\partial v}{\partial y} + u \frac{\partial^2 v}{\partial y^2} \right) + \left(\frac{\partial u}{\partial z} \frac{\partial v}{\partial z} + u \frac{\partial^2 v}{\partial z^2} \right) \text{ (product rule)},$$

$$= \left(\frac{\partial u}{\partial x} \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \frac{\partial v}{\partial y} + \frac{\partial u}{\partial z} \frac{\partial v}{\partial z} \right) + u \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right),$$

$$= \nabla u \cdot \nabla v + u \, \Delta v.$$

$$(52)$$

2. Prove that for any $u \in C^2(\Omega) \cap C^1(\bar{\Omega})$ one has

$$\int_{\Omega} \Delta u(\vec{x}) d\vec{x} = \int_{\partial \Omega} \frac{\partial u}{\partial \vec{\nu}} dS.$$

Solution. We know that for a scalar function $u:\Omega\to\mathbb{R}$,

$$\Delta u := \nabla \cdot (\nabla u) \text{ and } \frac{\partial u}{\partial \vec{\nu}} := \vec{\nu} \cdot \nabla u.$$
 (53)

Therefore,

$$\int_{\Omega} \Delta u(\vec{x}) d\vec{x} = \int_{\Omega} \nabla \cdot (\nabla u)(\vec{x}) d\vec{x},$$

$$= \int_{\partial \Omega} \vec{\nu} \cdot \nabla u dS \text{ (By the Gauss divergence theorem)},$$

$$= \int_{\partial \Omega} \frac{\partial u}{\partial \vec{\nu}} dS.$$
(54)

3. Prove that for any $u, v \in C^1(\Omega) \cap C(\bar{\Omega})$ one has

$$\int_{\Omega} \frac{\partial u}{\partial x_k} v d\vec{x} = \int_{\partial \Omega} u v \nu_k dS - \int_{\Omega} u \frac{\partial v}{\partial x_k} d\vec{x}, \quad k = 1, \dots, d.$$

Hint: Consider the vector-valued function $\vec{F}: \bar{\Omega} \to \mathbb{R}^d, \vec{F} = (f_1, \dots, f_d)^T$ s.t. $f_i \equiv 0$ (i.e. $f_i(x) = 0$ for all $x \in \bar{\Omega}$) for all $i \neq k$ and take $f_k = uv$.

Solution. Let us consider

$$\vec{F} = (f_1, \dots, f_k, \dots, f_d)^T = (0, \dots, u \ v, \dots, 0)^T,$$
 (55)

where $f_k := u v$, is the k-th component of \vec{F} . Then,

$$\vec{\nabla} \cdot \vec{F} = \frac{\partial f_1}{\partial x_1} + \dots + \frac{\partial f_k}{\partial x_k} + \dots + \frac{\partial f_d}{\partial x_d},$$

$$= 0 + \dots + \frac{\partial (u \ v)}{\partial x_k} + \dots + 0,$$

$$= \frac{\partial (u \ v)}{\partial x_k},$$
(56)

and

$$\vec{\nu} \cdot \vec{F} = \nu_1 f_1 + \dots + \nu_k f_k + \dots + \nu_d f_d,$$

$$= 0 + \dots + \nu_k (u \ v) + \dots + 0,$$

$$= \nu_k (u \ v).$$
(57)

By the Gauss divergence theorem, we can write,

$$\int_{\Omega} \nabla \cdot \vec{F} \, d\vec{x} = \int_{\partial \Omega} \vec{\nu} \cdot \vec{F} \, dS,$$

$$\implies \int_{\Omega} \frac{\partial (u \, v)}{\partial x_k} = \int_{\partial \Omega} \nu_k \, (u \, v) \, dS,$$

$$\implies \int_{\Omega} \left(u \frac{\partial v}{\partial x_k} + \frac{\partial u}{\partial x_k} v \right) \, d\vec{x} = \int_{\partial \Omega} (u \, v) \, \nu_k \, dS,$$

$$\implies \int_{\Omega} \frac{\partial u}{\partial x_k} v \, d\vec{x} = \int_{\partial \Omega} (u \, v) \, \nu_k \, dS - \int_{\Omega} u \, \frac{\partial v}{\partial x_k} \, d\vec{x}.$$
(58)

4. Let $\Omega = (a, b)$ be a finite, open interval. Prove that for all $u, v \in$ $C^1([a,b])$ one has

$$\int_a^b u'vdx = u(b)v(b) - u(a)v(a) - \int_a^b uv'dx.$$

Solution. Clearly,

$$\int_{a}^{b} \frac{d}{dx}(u \ v) \ dx = \left[(u \ v) \ (x) \right]_{x=a}^{x=b},$$

$$\implies \int_{a}^{b} \left(u \ \frac{dv}{dx} + v \ \frac{du}{dx} \right) dx = u(b)v(b) - u(a)v(a),$$

$$\implies \int_{a}^{b} u'v dx = u(b)v(b) - u(a)v(a) - \int_{a}^{b} uv' dx.$$
(59)

5. Explain why this situation is similar to the one in \mathbb{R}^d .

Solution. Let k = 1, $\Omega = (a, b)$, $\partial \Omega = \{a\} \cup \{b\}$.

At x = a, $\nu_a = -1$ and x = b, $\nu_b = 1$.

We showed that for \mathbb{R}^d ,

$$\int_{\Omega} \frac{\partial u}{\partial x_k} v \, d\vec{x} = \int_{\partial \Omega} (u \, v) \, \nu_k \, dS - \int_{\Omega} u \, \frac{\partial v}{\partial x_k} \, d\vec{x}$$
 (60)

Then for k = 1,

$$\int_{a}^{b} \frac{\partial u}{\partial x} v \, d\vec{x} = u(b) \, v(b) \, \nu_{b} + u(a) \, v(a) \nu_{a} - \int_{a}^{b} u \, \frac{\partial v}{\partial x} \, d\vec{x},$$

$$\implies \int_{a}^{b} \frac{\partial u}{\partial x} v \, d\vec{x} = u(b) \, v(b) - u(a) \, v(a) - \int_{a}^{b} u \, \frac{\partial v}{\partial x} \, d\vec{x}.$$
(61)

6. Prove that for all $u \in C^1(\Omega) \cap C(\bar{\Omega})$ and $v \in C^2(\Omega) \cap C^1(\bar{\Omega})$ one has

$$\int_{\Omega} \nabla u \nabla v d\vec{x} = \int_{\partial \Omega} u \partial_{\vec{v}} v dS - \int_{\Omega} u \Delta v d\vec{x}.$$

Hint: Use the Divergence Theorem for de vector-valued function $\vec{F} = u\nabla v$.

Solution. From problem 1, we can write,

$$\int_{\Omega} \nabla u \, \nabla v \, d\vec{x} + \int_{\Omega} u \, \Delta v \, d\vec{x} = \int_{\Omega} \nabla \cdot (u \, \nabla v) \, d\vec{x},$$

$$= \int_{\partial \Omega} \vec{v} \cdot (u \, \nabla v) \, dS \text{ (Gauss divergence)},$$

$$= \int_{\partial \Omega} u (\vec{v} \cdot \nabla v) \, dS.$$
(62)

We know that for a scalar function $u:\Omega\to\mathbb{R}$,

$$\frac{\partial u}{\partial \vec{\nu}} := \vec{\nu} \cdot \nabla u. \tag{63}$$

Hence,

$$\int_{\Omega} \nabla u \, \nabla v \, d\vec{x} + \int_{\Omega} u \, \Delta v \, d\vec{x} = \int_{\partial \Omega} u \, \partial_{\vec{v}} v \, dS,$$

$$\implies \int_{\Omega} \nabla u \, \nabla v \, d\vec{x} = \int_{\partial \Omega} u \, \partial_{\vec{v}} v \, dS - \int_{\Omega} u \, \Delta v \, d\vec{x}.$$
(64)

C. First order PDEs Below $f: \mathbb{R} \times (0, \infty) \to \mathbb{R}$ is a given function. Solve the first order equations below with the given initial conditions (do not forget to check your answers!):

$$\begin{cases} \frac{\partial u}{\partial t} + 3x^{\frac{4}{3}} \frac{\partial u}{\partial x} = f, & \text{for } x > 0, t > 0, \\ u(x, 0) = x^{\frac{1}{3}}, & \text{for } x > 0. \end{cases}$$

Consider three situations:

1. f(x,t,u) = 0,

Solution. **PDE:** Given PDE,

$$\frac{\partial u}{\partial t} + 3x^{\frac{4}{3}} \frac{\partial u}{\partial x} = 0, \text{for } x > 0, t > 0,$$

$$u(x,0) = x^{\frac{1}{3}} = u_0(x), \text{for } x > 0 \text{ (Initial condition)}.$$
(65)

Step1: ODE (non-linear): Let x(t) be the solution of the following ODE s.t.

$$\frac{dx(t)}{dt} = 3x^{\frac{4}{3}}, \quad x(0) = x_0 \text{ (characteristic curve)}.$$
 (66)

Step2: Solving ODE: Integrating,

$$\frac{1}{3} \int x^{-\frac{4}{3}} dx = \int 1 dt$$

$$\implies \frac{1}{3} \frac{x^{-\frac{1}{3}}}{(-1/3)} = t + C,$$

$$\implies -x^{-\frac{1}{3}} - t = C.$$
(67)

At t = 0,

$$-x_0^{-\frac{1}{3}} - 0 = C,$$

$$\implies C = -x_0^{-\frac{1}{3}}.$$
(68)

Hence,

$$-x^{-\frac{1}{3}} - t = -x_0^{-\frac{1}{3}},$$

$$\implies x^{-\frac{1}{3}} = x_0^{-\frac{1}{3}} - t,$$

$$\implies \frac{1}{x^{\frac{1}{3}}} = \left(x_0^{-\frac{1}{3}} - t\right),$$

$$\implies x = x(t) = \frac{1}{\left(x_0^{-\frac{1}{3}} - t\right)^3},$$
(69)

which is the solution of the non-linear ODE. It is the characteristic curves of the original PDE. Then we get,

$$x_0 = x(0) = \frac{1}{\left(x^{-\frac{1}{3}} + t\right)^3}. (70)$$

Step3: Solution of the original PDE, We define,

$$v(t) := u(x(t), t),$$

$$\implies v'(t) = \frac{d}{dt}u(x(t), t),$$

$$= u_t(x(t), t) + \frac{dx(t)}{dt}u_x(x(t), t),$$

$$= u_t(x(t), t) + 3x^{\frac{4}{3}}u_x(x(t), t),$$

$$= 0 \quad \text{(given from PDE)}.$$

$$(71)$$

Integrating both sides w.r.t. t,

$$v(t) = c$$
, where c is an integrating constant,
 $\Rightarrow u(x(t), t) = c$. (72)

At t = 0,

$$u(x(0), 0) = c = c,$$

 $\implies x_0^{1/3} = c \text{ (initial condition)}.$
(73)

Then the solution becomes,

$$u(x(t),t) = x_0^{1/3} = \frac{1}{\left(x^{-\frac{1}{3}} + t\right)}. (74)$$

It is the general solution of the PDE, it is constant on each characteristic curve.

Step4:Validation

Note: Every solution of the PDE is constant on the solution curves of the ODE.

Moral: Solutions of PDEs generally depend on arbitrary functions (instead of arbitrary constants). You need an auxiliary condition if you want to determine a unique solution. Such conditions are usually called initial or boundary conditions.

2. f(x,t,u) = 1,

Solution. **PDE:** Given PDE,

$$\frac{\partial u}{\partial t} + 3x^{\frac{4}{3}} \frac{\partial u}{\partial x} = 1, \quad \text{for } x > 0, t > 0,$$

$$u(x,0) = x^{\frac{1}{3}}, \quad \text{for } x > 0.$$
(75)

Step1: ODE (non-linear)

$$\frac{dx(t)}{dt} = 3x^{\frac{4}{3}}, \ x(0) = x_0 \text{ (characteristic curve)}. \tag{76}$$

Step2: Solving ODE In (1.), we show that the solution of the ODE is

$$x = x(t) = \frac{1}{\left(x_0^{-\frac{1}{3}} - t\right)^3}$$
 and $x_0 = x(0) = \frac{1}{\left(x^{-\frac{1}{3}} + t\right)^3}$. (77)

Step3: We define,

$$v(t) := u(x(t), t),$$

$$\implies v'(t) = \frac{d}{dt}u(x(t), t),$$

$$= u_t(x(t), t) + \frac{dx(t)}{dt}u_x(x(t), t),$$

$$= u_t + 3 x^{\frac{4}{3}} u_x,$$

$$= 1 \quad \text{(given PDE)}.$$

$$(78)$$

Integrating both sides w.r.t. t,

$$v(t) = t + c$$
, where c is an integrating constant,
 $\Rightarrow u(x(t), t) = t + c$. (79)

At
$$t = 0$$
,

$$u(x(0), 0) = c + 0 = c,$$

 $\implies x_0^{1/3} = c \text{ (initial condition)}.$
(80)

Then the solution becomes,

$$u(x(t),t) = x_0^{1/3} + t,$$

$$\implies u(x(t),t) = \frac{1}{\left(x^{-\frac{1}{3}} + t\right)} + t.$$
(81)

Step4:Validation

3. f(x, t, u) = u.

Solution. PDE(non-homogeneous): Given PDE,

$$\frac{\partial u}{\partial t} + 3x^{\frac{4}{3}} \frac{\partial u}{\partial x} = u, \quad \text{for } x > 0, t > 0,$$

$$u(x, 0) = x^{\frac{1}{3}} = \phi(x), \quad \text{for } x > 0.$$
(82)

Step1: ODE (non-linear)

$$\frac{dx(t)}{dt} = 3x^{\frac{4}{3}}, \ x(0) = x_0 \text{ (characteristic curve)}. \tag{83}$$

Step2: Solving ODE In (1.), we show that the solution of the ODE is

$$x = x(t) = \frac{1}{\left(x_0^{-\frac{1}{3}} - t\right)^3}$$
, and $x_0 = x(0) = \frac{1}{\left(x^{-\frac{1}{3}} + t\right)^3}$. (84)

Step3: We define,

$$v(t) := u(x(t), t),$$

$$\Rightarrow v'(t) = \frac{d}{dt}u(x(t), t),$$

$$= u_t(x(t), t) + \frac{dx(t)}{dt}u_x(x(t), t),$$

$$= u(x(t), t), \quad \text{given from PDE}$$

$$= v(t),$$

$$\Rightarrow v'(t) - v(t) = 0, \quad \text{Integrating factor is } e^{-t},$$

$$\Rightarrow v(t) = v(0)e^t,$$

$$= u(x(0), 0)e^t,$$

$$= x_0^{1/3}e^t,$$

$$\Rightarrow u(x(t), t) = \frac{e^t}{\left(x^{-\frac{1}{3}} + t\right)}.$$
(85)

Step4:Validation

D. A bit of Analysis:

Prove the 2^{nd} "Vanishing lemma":

Lemma 1. Let $\Omega \subset \mathbb{R}^d$ be a domain and $f: \Omega \to \mathbb{R}$ a continuous, positive function $(f(\vec{x}) \geq 0 \text{ for all } \vec{x} \in \Omega)$. If $\int_{\Omega} f(\vec{x}) d\vec{x} = 0$ then one has $f \equiv 0$ (f is 0 overall in Ω , or $f(\vec{x}) = 0$ for all $\vec{x} \in \Omega$).

Solution. We assume that $\exists \vec{x_0} \in \Omega : f(\vec{x_0}) = f_0 > 0$. Since f is continuous, by the definition of continuity, \exists a ball $B_{\vec{x_0}} \subseteq \Omega$ s.t. $\forall \vec{x} \in B_{\vec{x_0}}$,

$$|f(\vec{x}) - f(\vec{x_0})| < \frac{1}{2}f(\vec{x_0}), \text{ (here } \frac{1}{2}f(\vec{x_0}) = \epsilon \text{ in the def. of cont.)},$$

$$-\frac{1}{2}f(\vec{x_0}) < f(\vec{x}) - f(\vec{x_0}) < \frac{1}{2}f(\vec{x_0}),$$

$$f(\vec{x_0}) - \frac{1}{2}f(\vec{x_0}) < f(\vec{x}) < \frac{1}{2}f(\vec{x_0}) + f(\vec{x_0}),$$

$$\frac{1}{2}f(\vec{x_0}) < f(\vec{x}) < \frac{3}{2}f(\vec{x_0}),$$
(86)

which implying that

$$f(\vec{x}) > \frac{1}{2}f(\vec{x_0}) = \frac{1}{2}f_0 > 0, \ \forall \vec{x} \in B_{\vec{x_0}}.$$
 (87)

Let us consider that $\chi_{B_{\vec{x_0}}}$ be the characteristic function of $B_{\vec{x_0}}$, i.e. the function

$$\chi_{B_{\vec{x_0}}}: \Omega \longrightarrow \mathbb{R},$$

$$\vec{x} \longmapsto \chi_{B_{\vec{x_0}}}(\vec{x}) = \begin{cases} 1 & \text{if } \vec{x} \in B_{\vec{x_0}}, \\ 0, & \text{if } \vec{x} \notin B_{\vec{x_0}}. \end{cases}$$
(88)

Clearly, from (87) and (88), we can write

$$f(\vec{x}) \ge \chi_{B_{\vec{x}\vec{0}}}(\vec{x}) \ f(\vec{x}) \ge 0, \ \forall \vec{x} \in \Omega$$
 (89)

If

$$0 = \int_{\Omega} f(\vec{x}) d\vec{x},$$

$$\geq \int_{\Omega} \chi_{B_{\vec{x_0}}}(\vec{x}) f(\vec{x}) d\vec{x},$$

$$= \int_{B_{\vec{x_0}}} f(\vec{x}) d\vec{x}, \text{ (Using (88))},$$

$$\geq \int_{B_{\vec{x_0}}} \frac{1}{2} f(\vec{x_0}) d\vec{x}, \text{ (Using (87))},$$

$$= \frac{1}{2} f(\vec{x_0}) \int_{B_{\vec{x_0}}} 1 d\vec{x},$$

$$= \frac{1}{2} f(\vec{x_0}) |B_{\vec{x_0}}|,$$

$$> 0, \text{ (Since } \frac{1}{2} f(\vec{x_0}) > 0, \text{ the volume of } B_{\vec{x_0}} = |B_{\vec{x_0}}| > 0),$$

which is a contradiction.

Hence, if $\int_{\Omega} f(\vec{x}) \ d\vec{x} = 0$ then one has $f \equiv 0$ (f is 0 overall in Ω .