

WARSAW UNIVERSITY OF TECHNOLOGY
FACULTY OF POWER AND AERONAUTICAL ENGINEERING

COMPUTER METHODS IN COMBUSTION

Calculation of the detonation velocity of the hydrogen-air mixture in a constant volume for variable initial parameters

Author:

Wiktoria MIECZKOWSKA (313825)

Supervisor:

dr inż. Mateusz ŻBIKOWSKI

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1 Introduction

Detonation is a compression wave propagating supersonically, which can be produced by driving a piston at an appropriate speed into the mixture. It compresses the burned gas, slowing it down. The downstream flow is:

- sonic ($M_2 = 1$) for Chapman-Jouguet detonation
- supersonic ($M_2 > 1$) for weak detonations
- subsonic ($M_2 < 1$) for strong detonations

Weak detonations can propagate only at one speed.

2 Theoretical background

2.1 Rankine-Hugoniot equations

Consider a mixture of combustible gases in a straight pipe of constant cross section in which we have a plane combustion wave propagating along the pipe axes. The wave is described in coordinate system stationary to the wave.

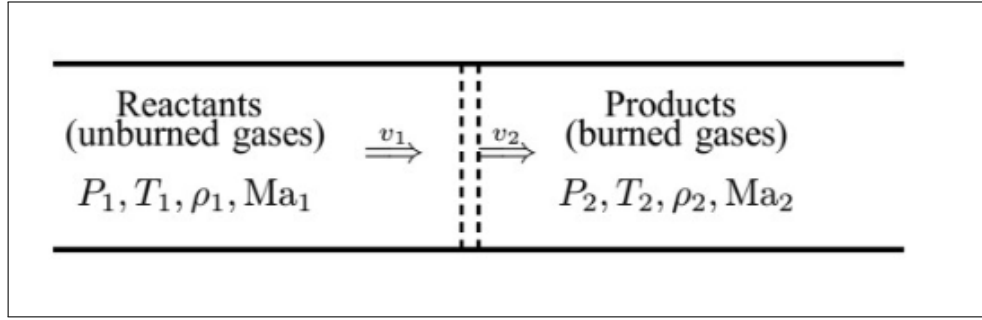


Figure 1: Detonation wave visualization

Then the mass, momentum and the energy preservation equations are:

$$\rho_1 v_1 = \rho_2 v_2 = m \quad (1)$$

$$P_1 + \rho_1 v_1^2 = P_2 + \rho_2 v_2^2 \quad (2)$$

$$h_1 + \frac{v_1^2}{2} = h_2 + \frac{v_2^2}{2} \quad (3)$$

Where h - enthalpy equals to:

$$h(T) = \sum Y_i h_{fi}^o + \sum Y_i \int_{T_o}^T c_{pi} dT \quad (4)$$

Taking into consideration the fact that:

$$p = \rho R T \quad (5)$$

Assume $c_{p1} = c_{p2} = c_p$ and q - heat combustion per mass unit of combustible mixture

$$\rho_1 v_1 = \rho_2 v_2 \quad (6)$$

$$P_1 + \rho_1 v_1^2 = P_2 + \rho_2 v_2^2 \quad (7)$$

$$c_p T_1 + \frac{v_1^2}{2} = c_p T_2 + \frac{v_2^2}{2} - q \quad (8)$$

$$\frac{p_1}{\rho_1 T_1} = \frac{p_2}{\rho_2 T_2} = \frac{\kappa - 1}{\kappa} c_p \quad (9)$$

After some manipulations and use of the equation of state to eliminate T:

$$\frac{p_2 - p_1}{\frac{1}{\rho_2} - \frac{1}{\rho_1}} = -m^2 - \text{Rayleigh Line} \quad (10)$$

$$\frac{\kappa}{\kappa - 1} \left(\frac{p_2}{\rho_2} - \frac{p_1}{\rho_1} \right) - 0.5 \left(\frac{1}{\rho_2} + \frac{1}{\rho_1} \right) (p_2 - p_1) = q - \text{Hugoniot curve} \quad (11)$$

By drawing the Rayleigh Line and the Hugoniot curve on the same plot allows to visualize solutions of above equations:

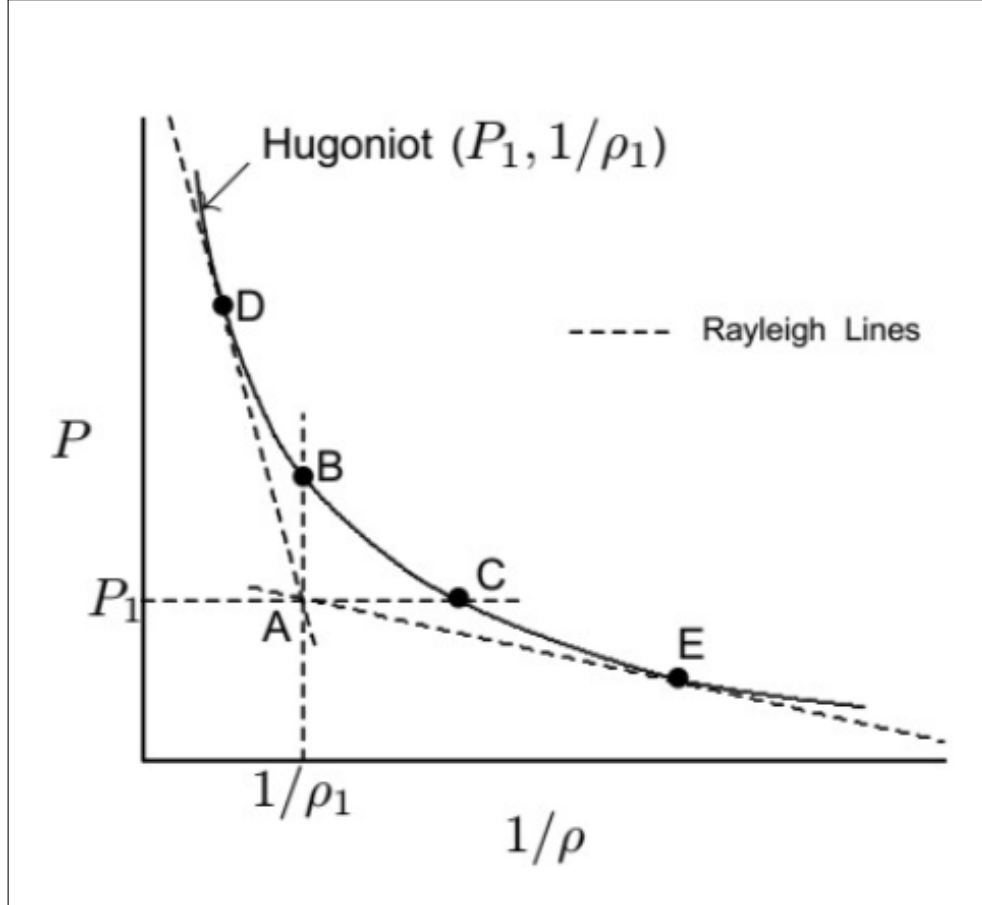


Figure 2: Rankine-Hugoniot curve

As we can see, four Rayleigh Lines divide a Hugoniot curve into five regions:

- above point D - strong detonation
- B-D - weak detonation
- C-E - weak deflagration
- below point E - strong deflagration

2.2 Chapman-Jouguet detonation

The Chapman-Jouguet (C-J) detonation is a detonation, where the outlet gas velocity is sonic. Assuming constant heat capacities of the reactants and transforming mass, momentum and energy conservation equations the detonation velocity may be calculated as:

$$v_{C-J} = \sqrt{2(\kappa_2^2 - 1)(q + c_{p1}T_1)} \quad (12)$$

In the Figure 2 Chapman - Jouguet detonation is represented by the point D.

2.3 ZND detonation model

ZND detonation model is a developed C-J model that takes into consideration the finite time of the reaction. Due to that theory the structure of the detonation looks as below:

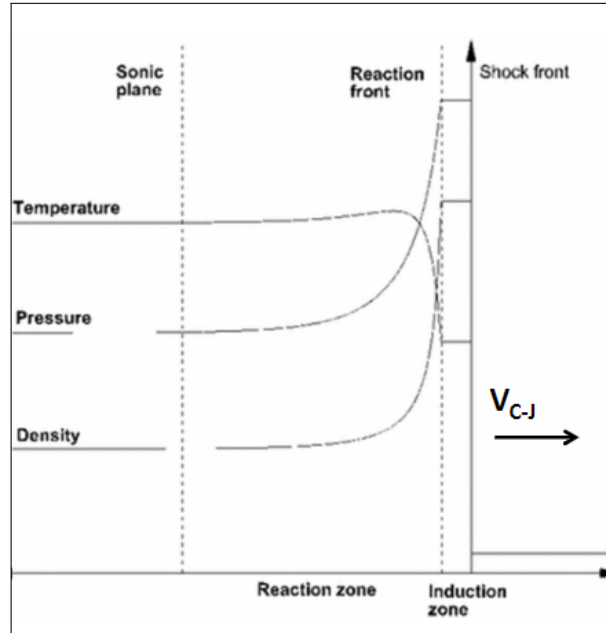


Figure 3: ZND detonation

2.4 Real detonation structure

The real detonation structure due to experiments may be:

2.4.1 Cellular detonation structure

That is the most common detonation structure

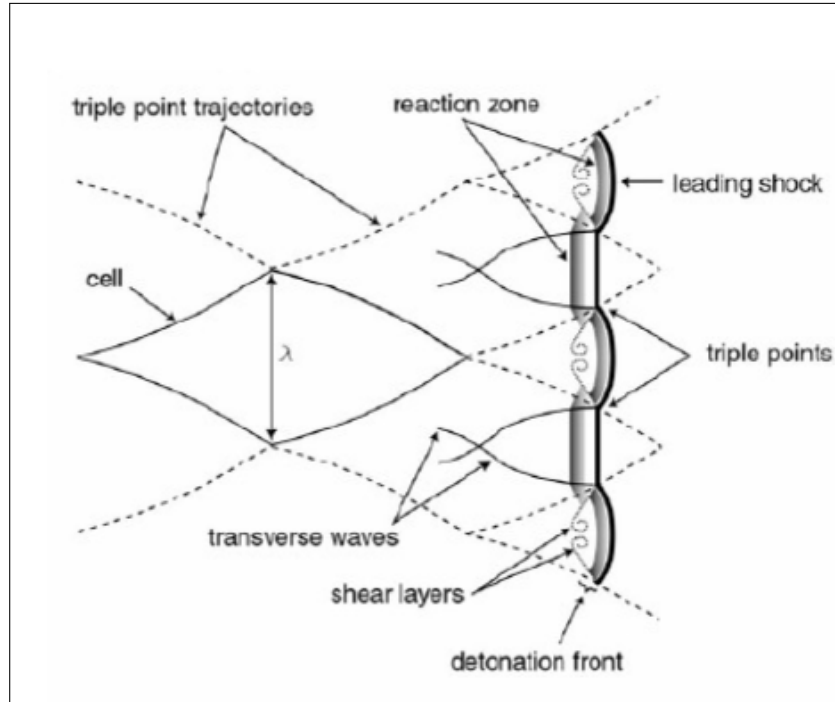


Figure 4: Cellular detonation

2.4.2 Rotating detonation structure

This structure is characteristic for parameters next to the detonation limits

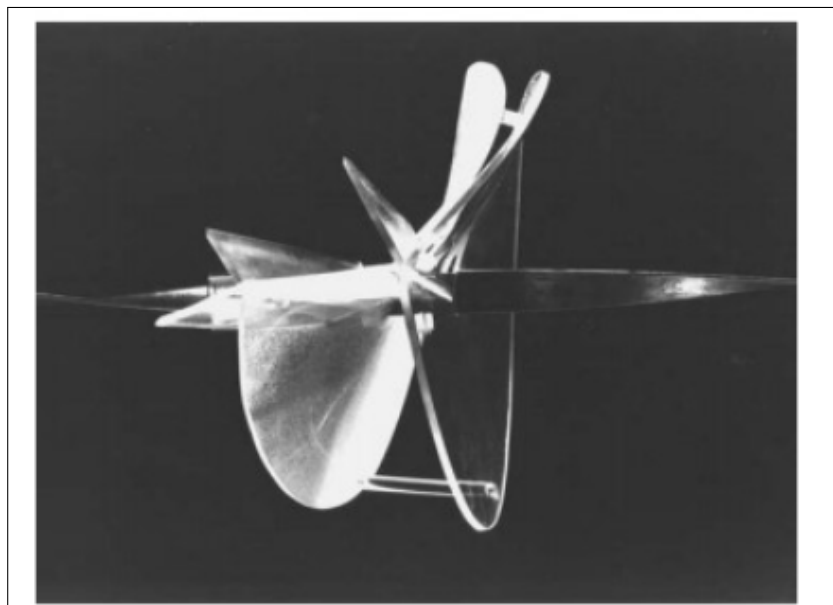


Figure 5: Rotating detonation

2.4.3 Zig-zag detonation structure

Zig-zag structure appears in rectangular pipes:

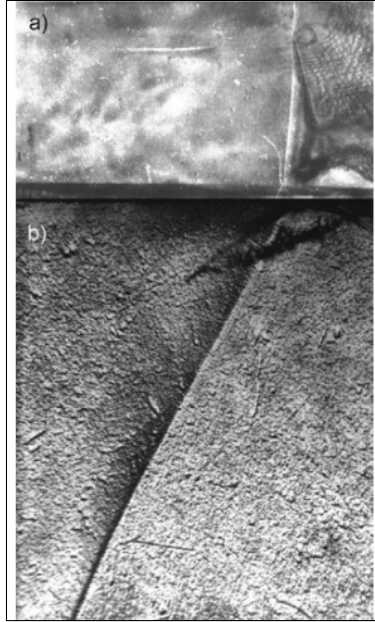


Figure 6: Zig-zag detonation

3 Model description

Detonation models were implemented using Cantera software. All of the assumptions are stated according to these models as well.

Initial parameters were changed and the C-J wave velocity was calculated for each combination. For thermodynamic and transport properties or the reaction rates GRI - 3.0 mechanism was used (a chemical mechanism developed by the Gas Research Institute to model natural gas combustion, including full NO_x chemistry, a compilation of 325 reactions involving 53 species).

4 Results

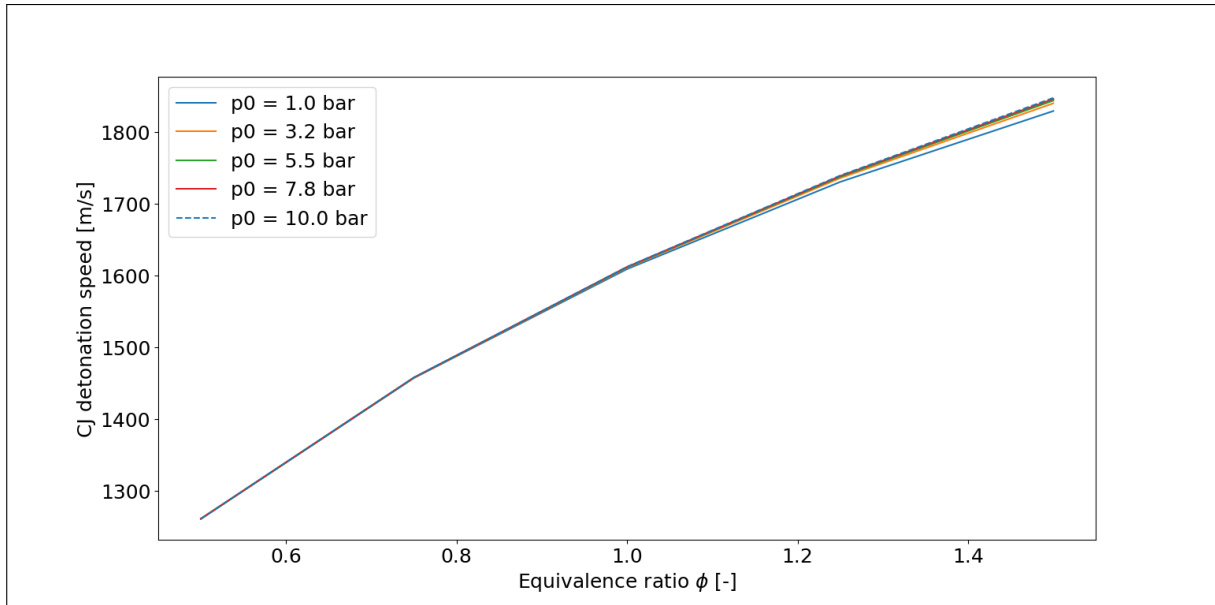


Figure 7: Relation of C-J speed and equivalence ratio ϕ for different initial pressure, const temperature

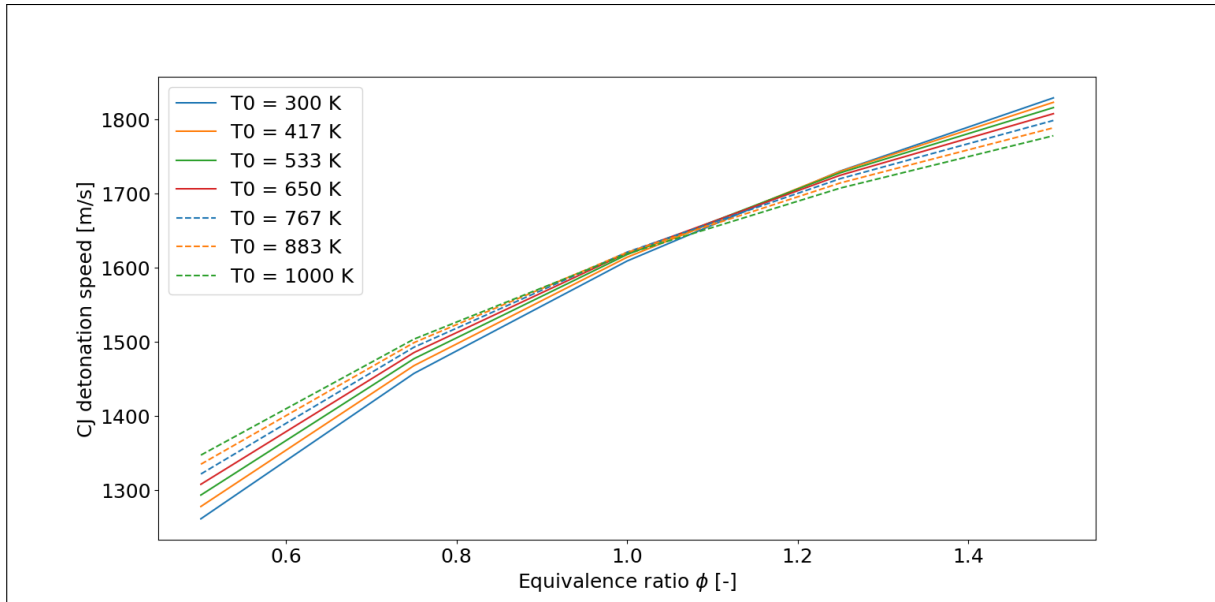


Figure 8: Relation of C-J speed and equivalence ratio ϕ for different initial temperature, const pressure

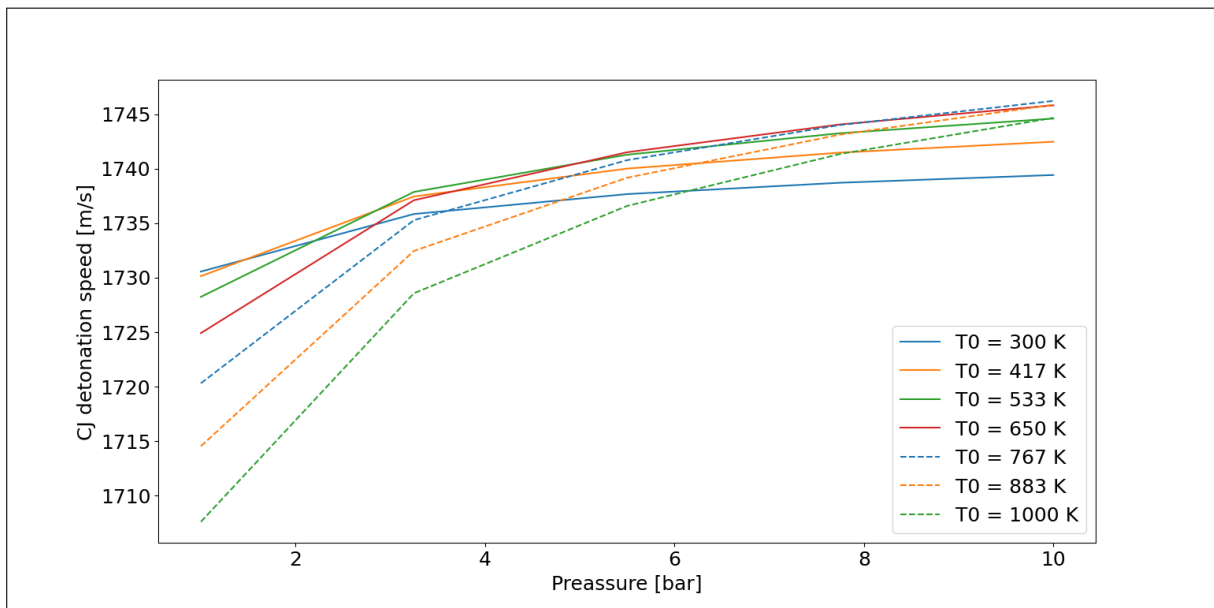


Figure 9: Relation of C-J speed and pressure for different initial temperature, const ϕ

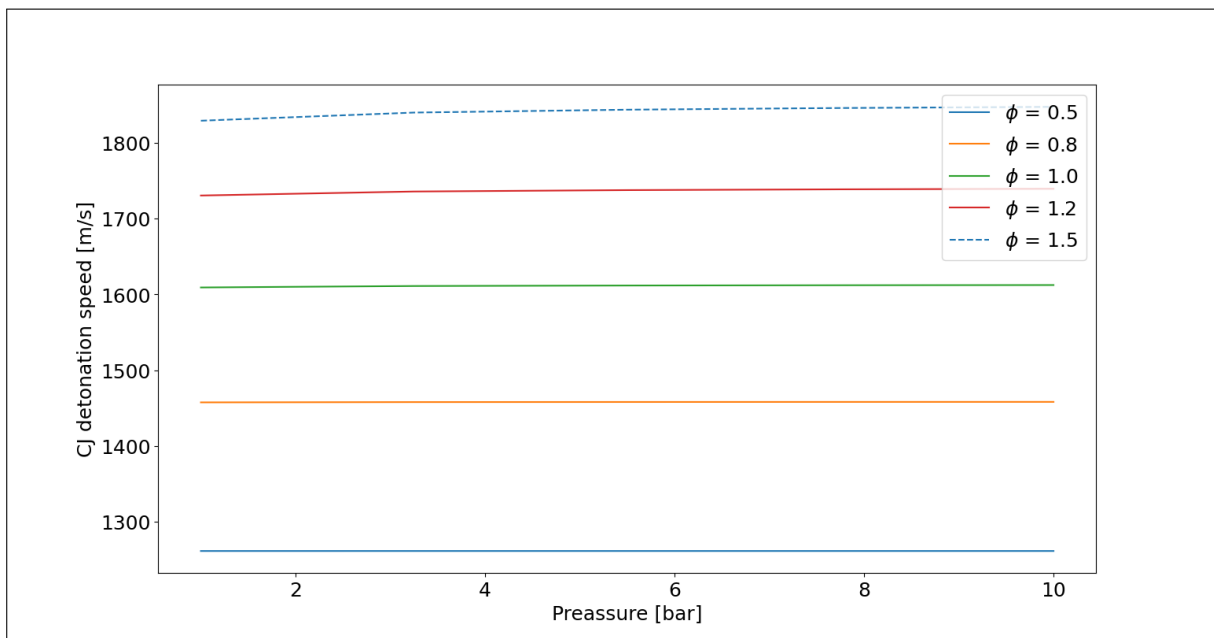


Figure 10: Relation of C-J speed and pressure for different initial ϕ , const temperature

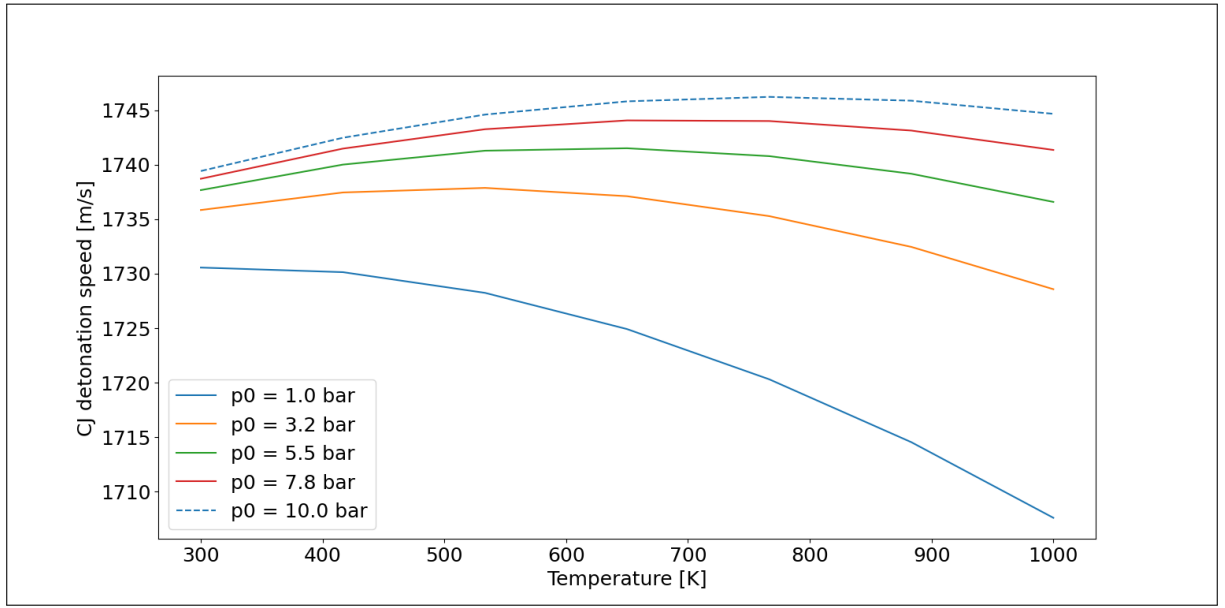


Figure 11: Relation of C-J speed and temperature for different initial pressure, const ϕ

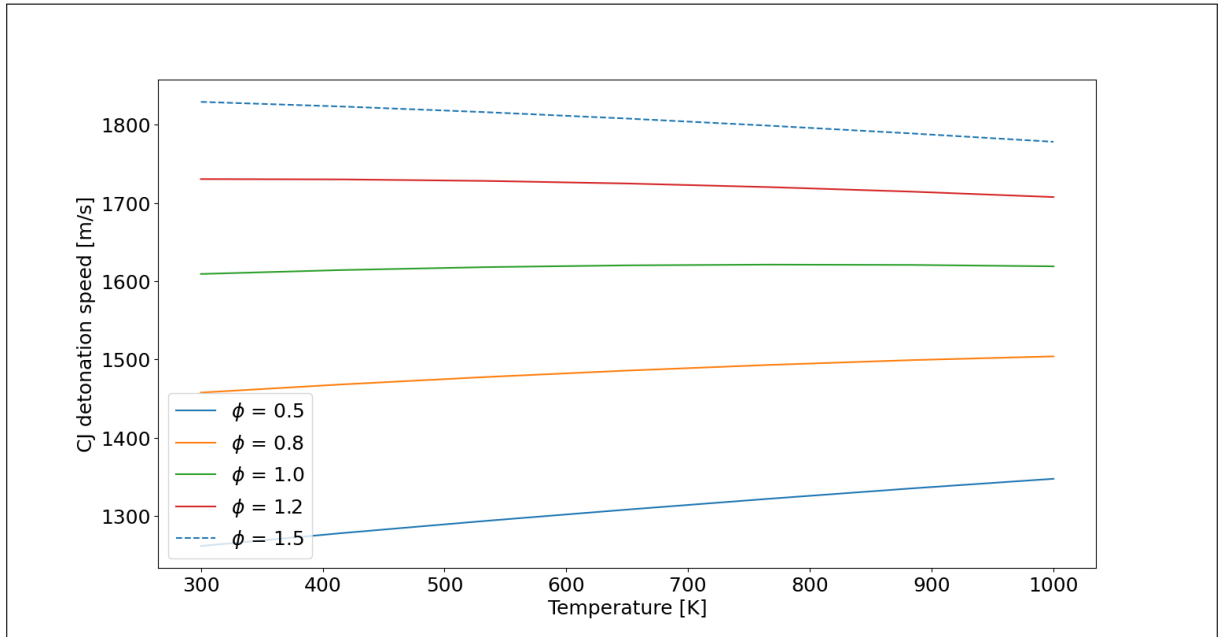


Figure 12: Relation of C-J speed and temperature for different initial ϕ , const pressure

5 Conclusions

- CJ speed (ϕ) for different initial pressure almost doesn't differ, for rising temperature, the curve is more horizontal. (figures 7,8)
- In Figure 9 it is seen that the CJ speed increases with rising pressure and has lower values for higher initial temperatures. CJ speed (pressure) graph is almost flat, detonation speed increases with ϕ . (Fig 10)
- According to Figure 11 the highest detonation speed is for initial temperature 800K, for the highest initial pressure. In the Figure 12 lines are almost flat - detonation speed almost does not depend on initial temperature for established ϕ .

6 Appendix

1. mkws.py

7 Bibliography

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3. 'Towards Detonation Theory' A. N. Dremin
4. *[https : //shepherd.caltech.edu/EDL/PublicResources/sdt/cti_mech.html](https://shepherd.caltech.edu/EDL/PublicResources/sdt/cti_mech.html)*
5. <https://github.com/Bazyl29/MKWS>