Metodos matemáticos 2

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1. Función Gamma

L'imite al infinito(Euler)

$$\begin{aligned}
\Gamma(z) \equiv & \lim_{n\to\infty} \tfrac{1\cdot 2\cdot 3\cdots n}{z(z+1)(z+2)\cdots(z+n)} n^z, \quad z \in \mathbb{Z}^+ \ o \ z \in \mathbb{C} \\
\Gamma(z+1) = & \lim_{n\to\infty} \tfrac{1\cdot 2\cdot 3\cdots n}{(z+1)(z+2)\cdots(z+n+1)} n^{z+1} \\
= & z\cdot\Gamma(z) = \lim_{n\to\infty} \tfrac{nz}{z+n+1}\Gamma(z) = \Gamma(z+1) = z\cdot\Gamma(z)
\end{aligned}$$

Aplicando lo anterior a z=1,2,3...n

$$\begin{aligned}
\Gamma(1) = & \lim_{n\to\infty} \tfrac{1\cdot 2\cdot 3\cdots n}{(1)(2)\cdots(n+1)} n^z = \lim_{n\to\infty} \tfrac{n}{n+1} = 1 \\
\Gamma(2) = & z\cdot\Gamma(z) = \Gamma(z+1) = \Gamma(1+1) = 1 \\
\Gamma(3) = & 2\cdot\Gamma(1) = \Gamma(2+1) = 2\cdot 1 \\
\Gamma(4) = & 3\cdot\Gamma(3) = \Gamma(3+1) = 3\cdot 2\cdot 1 \\
& . \\
& . \\
& . \\
\Gamma(n) = & (n-1)!
\end{aligned}$$

Itegral Definida(Integral de Euler)

$$\Gamma(z) = \int_0^\infty e^{-t} t^{z-1} dt, \quad Re(z) > 0$$

Ej. como aparecen en f'isica

$$\Gamma(z) = 2\int_0^\infty e^{-t^2} t^{z+1} dt \quad o \quad \Gamma(z) = \int_0^1 \left[ln(\frac{1}{t}) \right]^{z-1} dt$$

Si $z = \frac{1}{2} \Rightarrow \Gamma(\frac{1}{2}) = \sqrt{\pi}$ es integral error de Gauss

$$F(z,n) = \int_0^n \left(1 - \frac{t}{n}\right)^n t^{z-1} dt, \quad Re(z) > 0 \ con \ n \ entero \ positivo \ni$$

$$e^{-t} = \lim_{n\to\infty} \left(1 - \frac{t}{n}\right)^n$$

$$\Rightarrow F(z,n) = \int_0^n \left(1 - \frac{t}{n}\right)^n t^{z-1} dt = \int_0^1 (1-u)^n (un)^{z-1} n du \ni u = \frac{t}{n}$$

$$\frac{F(z,n)}{n^z} = \int_0^1 (1-u)^n (u)^{z-1} du$$

Ahora por integraci'on por partes usando $u=(1-u)^n$, $du=n(1-u)^{n-1}du$, $v=\frac{u^z}{z}$, $dv=u^{z-1}du$ tenemos

$$\begin{aligned}
\frac{F(z,n)}{n^z} =&\quad (1-u)^n\ \frac{u^z}{z}\Big|_0^1 - \int_0^1 \frac{u^z}{z} n(1-u)^{n-1}du \\
F(z,n) =&\quad n^z\cdot\frac{n(n-1)\cdots 1}{z(z+1)\cdots(z+n-1)} = \int_0^1 u^{z+n-1}du \\
=&\quad \frac{1\cdot 2\cdots n}{z(z+1)\cdots(z+n)}n^z \Rightarrow \lim_{n\to\infty} F(z,n)=\Gamma(z)
\end{aligned}$$

Producto infinito(Weierstrass)

$$\frac{1}{\Gamma(z)} = ze^{\delta z}\prod_{n=1}^{\infty}\left(1+\frac{z}{n}\right)e^{-\frac{z}{n}}$$

$\delta :=$ constante de Euler-Mascheron

$\delta := 0.5772156619$

Ecuaci'on estad'istica de Maxwell-Boltzmann

$e^{-E/KT}$

K es la constante de Boltzmann

T es la temperatura absoluta

E(energ'ia): estado de energ'ia ocupada

probabilidad de estar en estado de energ'ia es $Y_{kt}=\beta$

$P(E)=Ce^{-\beta E}$	*Para un gas idel sin estructura*
$n(E)dE$	$n(E)^{1/2}$
$1=C\int n(E)e^{-\beta E}dE$	*E=energ'ia cin'etica*

$$\begin{aligned}
1 =&\quad c\int_0^\infty E^{1/2}e^{-\beta E}dE = \frac{C\Gamma(3/2)}{\beta^{3/2}}, \quad \beta E = T \Rightarrow dE = \frac{dt}{\beta} \\
1 =&\quad C\int_0^\infty e^{-t}\left(\frac{t}{\beta}\right)^{1/2}\frac{1}{\beta}dt \\
&= C\int_0^\infty e^{-t}t^{3/2-1}dt\cdot\frac{1}{\beta^{3/2}} \\
=&\quad \frac{C\Gamma(3/2)}{\beta^{3/2}} \\
=&\quad c\cdot\frac{\sqrt{\pi}}{2\cdot\beta^{3/2}}
\end{aligned}$$

$$\therefore C = \frac{2\cdot\beta^{3/2}}{\sqrt{\pi}}$$

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Relaciones de funciones

$$\Gamma(z+1) = z\Gamma(z)$$

Fomula de reflexi'on

$$\Gamma(z)\Gamma(1-z) = \frac{\pi}{sen(z\pi)}$$

F'ormuláde duplicaci'on de Legendre

$$\Gamma(1+z)\Gamma(z+\frac{1}{2}) = 2^{-2z}\sqrt{\pi}\Gamma(2z+1)$$

Notac'i'on doble factorial:

$$\begin{aligned} (2n+1)!! =&\quad 1\cdot 3\cdot 5\cdots(2n+1) = \tfrac{(2n+1)!}{2^n n} \\ (2n)!! =&\quad 2\cdot 4\cdot 6\cdots(2n) = 2^n n! \\ (-1)!! =&\quad 1 \end{aligned}$$

PROPIEDADES ANAL'ITICAS

$[\Gamma(z)]^{-1}$ tiene singularidades en $z = 0, -1, -2...$ y no tiene cuando $z = -1, 2, 3...$ y no tiene ceros en el plano complejo finito positivo.

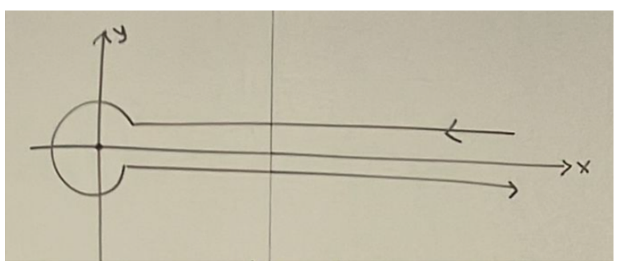
El residuo R_n cuando $z = -n$ donde n es un entero mayor o igual a cero.

$$\begin{aligned} R_n =&\quad \lim_{n\to 0}(\epsilon\Gamma(-n+\epsilon)) \\ =&\quad \lim_{n\to 0}\tfrac{\epsilon\Gamma(-n+1+\epsilon)}{-n+\epsilon} \\ =&\quad \lim_{n\to 0}\tfrac{\epsilon\Gamma(-n+2+\epsilon)}{(+\epsilon)(+1+\epsilon)} \\ =&\quad \lim_{n\to 0}\tfrac{\epsilon\Gamma(1+\epsilon)}{(-n+\epsilon)\cdots(\epsilon))} \\ =&\quad \tfrac{(-1)^n}{n!}, \textit{el residuo alterna signos en } z=-n \end{aligned}$$

Integral de Schaefli

$$\int_C e^{-t}t^v dt = (e^{2\pi i v})\Gamma(v+1)$$

donde C es el contorno



Esta integral es 'util cuando V no es entero.

Notaci'on factorial:

$$\prod(z) = z! = \Gamma(z+1)$$

Funci'on Digamma

$$\begin{aligned} \ln(\Gamma(z+1)) =&\quad z\cdot\Gamma(z) \\ =&\quad \lim_{n\to\infty}\tfrac{z}{z}\ln\left[\tfrac{n!\cdot n^z}{(z+1)(z_2)\cdots(z+n)}\right] \\ \tfrac{d}{dz}\ln(\Gamma(z+1)) =&\quad \tfrac{d}{dz}\lim_{n\to\infty}[\ln(n!) + z\cdot\ln(n) - \ln(z+1) - \ln(z+2) - ...\ln(z+n)] \\ \tfrac{d}{dz}\ln(\Gamma(z+1)) \equiv \underbrace{\varphi(z+1)}_{Funci'on\ digamma} =&\quad \lim_{n\to\infty}\left(\ln(n) - \tfrac{1}{z+1} - \tfrac{1}{z+2}\cdots - \tfrac{1}{z+n}\right) \end{aligned}$$

$$\varphi(z+1) = \quad \tfrac{[\Gamma(z+1)]'}{\Gamma(z+1)}$$

Si sumamos y restamos un n'umero arm'onico

$$H_n = \sum_{m=1}^{n} \frac{1}{m}$$

$$\varphi(z+1) = \lim_{n\to\infty} \underbrace{[(\ln(n) - H_n)}_{=-\delta \ cuando \ n\to\infty} - \sum\left(\frac{1}{z+m} - \frac{1}{m}\right)$$

$$\Rightarrow \varphi(z+1) = -\delta + \sum_{m=1}^{\infty} \frac{z}{m(m+z)}, \quad si \ z{=}0 \Rightarrow \varphi(1) = -\delta$$

para $n > 0$, n entero.

$$\begin{aligned} \varphi(n+1) = & \ -\delta + \textstyle\sum_{m=1}^{n} \frac{1}{m} \\ = & \ -\delta + H_n \end{aligned}$$

Funci'on Poligamma

Cuando se deriva muchas veces la funci'on digamma

$$\begin{aligned} \varphi^{(m)}(z+1) \equiv & \ \tfrac{d^{m+1}}{dz^{m+1}} \ln\Gamma(z+1) \\ \varphi^{(m)}(z+1) = & \ (-1)^{m+1} m! \textstyle\sum_{n=1}^{\infty} \frac{1}{(z+n)^{m+1}}, \quad m = 1,2,3... \end{aligned}$$

Si z=0 la funci'on se conoce como zeta de Riemann

$$\zeta(m) = \sum_{n-1}^{\infty} \frac{1}{n^m}$$

$$\Rightarrow \varphi_{(1)}^{(m)} = (-1)^{m+1} \cdot m! \cdot \zeta(m+1), \quad m = 1,2,3...$$

Expansi'on de Maclaurin

$$\begin{aligned} \ln\Gamma(z+1) = & \ \textstyle\sum_{n=1}^{\infty} \frac{z^n}{n!} \varphi_{(1)}^{(n-1)} \\ = & \ -\delta \cdot z + \textstyle\sum_{n=2}^{\infty} (-1)^n \frac{z^n}{n} \zeta(n) \end{aligned}$$

converge en $|z| < 1$ para $z = x \Rightarrow$ se puede calcular $\Gamma(z+1)$ para n'umeros reales o complejos.

Suma de series

Transformar la serie por medio de fracciones parciales y expresar la serie infinita como sumas finitas de funciones gamma y poligamma.

EJEMPLO:

$$\begin{aligned} \varphi^{(m)}(z+2) = & \ \varphi^{(m)}(z+1) + (-1)^m \tfrac{m!}{(z+1)^{m+1}} \\ \varphi^{(m)}(z+2) - \varphi^{(m)}(z+1) = & \ \tfrac{(-1)^m m!}{(z+1)^{m+1}} \end{aligned}$$

$$\frac{d^{(m)}}{dz^{(m)}} [\varphi(z+2) - \varphi(z+1)]$$

$$\begin{aligned} \varphi(z+2) - \varphi(z+1) = & \ \textstyle\sum_{m=1}^{\infty} \frac{z+1}{m(m+z+1)} - \sum_{m=1}^{\infty} \frac{z}{m(m+z)} \\ = & \ \textstyle\sum_{m=1}^{\infty} \left[\frac{1}{m} - \frac{1}{m+z+1} - \frac{1}{m} + \frac{1}{m+z}\right] \\ = & \ \tfrac{1}{z+1} \end{aligned}$$

$$\frac{d^m}{dz^m} [\varphi(z+2) - \varphi(z+1)] = \frac{d^m}{dz^m}\left[\frac{1}{z+1}\right] = \frac{(-1)^m m!}{(z+1)^{m+1}}$$

Demuestre que:

$$\frac{1}{2} \ln\left[\frac{\pi \cdot z}{sen(\pi \cdot z)}\right] = \sum_{n=1}^{\infty} \frac{\zeta(2n)}{2n} z^{2n}, \quad |z| < 1$$

- $\frac{1}{2} \ln\left[\underbrace{\Gamma(z)}_{z\Gamma(z)=\Gamma(z+1)} \quad \Gamma(z-1)\right] = \frac{1}{2} \ln\left[\frac{\pi}{sen(z\pi)}\right]$
- $\ln\Gamma(z+1) = \lim_{n\to\infty} [\ln(n!) + z\ln(n) - \ln(z+1) - ... - \ln(z+n)]$

$$\Gamma(z)\Gamma(1-z) \rightsquigarrow \Gamma(1+z)\Gamma(1-z)$$

$$\begin{aligned}
\ln\left[\Gamma(1+z)\Gamma(1-z)\right] = &\quad -\delta \cdot z + \textstyle\sum_{n=2}^{\infty}(-1)^n \frac{z^n}{n}\zeta(n) \\
&\quad +(-\delta(-z)) + \textstyle\sum_{n=2}^{\infty}(-1)^n \frac{(-z)^n}{n}\zeta(n) \\
= &\quad \textstyle\sum_{n=2}^{\infty}\zeta(n) \underbrace{\left[\dfrac{z^n}{n} + \dfrac{(-z)^n}{n}\right]}_{z^{2n}} \\
= &\quad 2\textstyle\sum_{n=1}^{\infty} \cancelto{1}{(-1)^{2n}} \zeta(2n)\frac{z^{2n}}{2n}
\end{aligned}$$

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La Funci'on Beta

$$\Gamma(p+q) \propto \Gamma(p)\Gamma(q)$$

$$n! = \Gamma(n+1) \& \int_0^\infty e^{-t} t^2 dt, \quad \textit{Variante de gamma para definir la funi'on beta}$$

$$\Gamma(m)\Gamma(n) = \int_0^\infty e^{-u} u^m fu \int_0^\infty e^{-v} v^n dv$$

$$\begin{array}{cc} u = x^2 & v = y^2 \\ du = 2xdx & dv = 2ydy \end{array}$$

$$\begin{aligned} \Gamma(m)\Gamma(n) = & \int_0^\infty \int_0^\infty e^{-x^2} e^{-y^2} x^{2m}(2x) y^{2m}(2y) dxdy \\ = & 4\int_0^\infty e^{-x^2} e^{-x^2} e^{-y^2} x^{2m+1} y^{2m+1} dxdy \end{aligned}$$

otro cambio de variables

$$x = rcos(\theta) \quad y = rsen(\theta)$$

Usando el jacobiano

$$\begin{vmatrix} \frac{\partial x}{\partial r} = cos\theta dr & \frac{\partial x}{\partial \theta} = -rsen\theta d\theta \\ \\ \frac{\partial y}{\partial r} = sen\theta dr & \frac{\partial y}{\partial \theta} rcos\theta d\theta \end{vmatrix} = rdrd\theta$$

$$\begin{aligned} = & \int_0^\infty \int_0^{\pi/2} \cancel{e^{-r^2}} \xrightarrow{} e^{-r^2 cos^2\theta} \; r^{2m+2n+3} cos^{2m+1}\theta sen^{2n+1} drd\theta \\ \Gamma(m)\Gamma(n) = & 2\Gamma(m+n+1) \int_0^{\pi/2} cos^{2m+1}\theta sen^{2n+1}\theta d\theta \end{aligned}$$

$$\begin{aligned} \Gamma(z) = & 2\int_0^\infty e^{-t^2} t^{2z-1} dt; \; Re(z > 0) \\ \Gamma(q)\Gamma(p) = & 4\int_0^\infty e^{-s^2} s^{2q-1} ds \int_0^\infty e^{-t^2} t^{2p-1} dt \end{aligned}$$

$$\begin{array}{cc} s = rcos\theta & t = rsen\theta \\ r^2 = s^2 + t^2 & dsdt = rdrd\theta \end{array}$$

$$\begin{aligned} \Gamma(q)\Gamma(p) = & 4\int\int e^{-r^2 cos^2\theta} e^{-r^2 sen^2\theta} r^{2q-1} cos^{2q-1} r^{2p-1} sen^{2p-1} rdrd\theta \\ = & 2 \cdot 2\int_0^\infty e^{-r^2} r^{2q+2p-1} cos^{2q-1} sen^{2p-1} drd\theta \\ \Gamma(p)\Gamma(q) = & 2\int cos^{2q-1} sen^{2p-1} d\theta \end{aligned}$$

$$\frac{m!n!}{(m+n+1)!} \star = 2\int_0^{pi/2} cos^{2m+1}\theta sen^{2n+1}\theta d\theta$$

$$\boxed{b(p,q) = \frac{\Gamma(q)\Gamma(p)}{\Gamma(q+p)}} \text{ definici'on de funci'on beta}$$

$$b(p,q) = b(q,p)$$

Si se sustituye $t = cos^2\theta$ y $dt = -2cos\theta sen\theta d\theta$

$$\begin{aligned} \cancel{cos^2\theta} \xrightarrow{} t + sen^2\theta = & \; 1 \\ sen^2\theta = & \; 1 - t \\ sen^{2n+1}\theta = & \; sen^{2n}\theta sen\theta \\ = & \; (1-t)^n sen\theta \end{aligned} \quad \begin{aligned} \star = & \; 2\int_0^1 cos^{2m}\theta \cancel{cos\theta} (1-t)^n \cancel{sen\theta} dt \\ = & \int_0^1 t^m (1-t)^n dt \Rightarrow \textit{ otra forma de escribir} B(m+1, n+1) \end{aligned}$$

lunes 17 de julio Ejemplo

Una particula de masa m se mueve en un potencial sim'etrico esta descrito por $v(x) = A|x|^n$ y energ'ia total como $\frac{1}{2}m\left(\frac{dx}{dt}\right)^2 + v(x) = E$

Resolviendo para $\frac{dx}{dt}$ e intgrando encontramos que el movimiento per'iodico es:

$$\zeta = \sqrt{2m}\int_0^{X_{max}} \frac{dx}{(E - Ax)^1}$$

donde X_{max} es el punto de innplexion clasico dado por $Ax^n_{max} = E$

Muestre que:

$$\tau = \frac{2}{n}\sqrt{\frac{e\pi m}{E}}\left(\frac{E}{A}\right)^{1/2}\left[\frac{\Gamma(1/n)}{\Gamma(\frac{1}{n}+\frac{1}{2})}\right]$$

$$\boxed{\begin{gathered} B(p+1, q+1) = B(1/n, 1/2) = \int_0^1 t^p(1 - t^q dt) \\ \frac{1}{n} = p + 1 \Rightarrow p = \frac{1}{n} - 1 \\ \frac{1}{2} = q + 1 = \frac{1}{2} - 2 \\ \int_0^1 t^{\frac{1}{n}-1}(1-t)^{-1} \end{gathered}}$$

$$t = \frac{Ax^n}{E} \Rightarrow \left(\frac{Et}{A}\right)^{1/n}$$

$$dt = \frac{A}{E}n(x^{n-1})dx$$

$$dx = \underbrace{\frac{E}{A}}_{\frac{1}{n}\frac{x}{t}dt}\frac{1}{n}\frac{x}{x^n}dt$$

$$\begin{aligned} \tau = \quad & 2\frac{\sqrt{2m}}{\sqrt{E}\int_0^1 \frac{\frac{x}{nt}dt}{(1-t)^{1/2}}} \\ = \quad & \left(\frac{E}{A}\right)^{1/n}\frac{2}{n}\sqrt{\frac{2m}{E}}\int_0^1 \frac{\frac{\frac{(Et)^{1/n}}{A^{1/n}}}{nt}dt}{(1-t)^{1/2}} \\ = \quad & \frac{2}{n}\sqrt{\frac{2m}{E}}\left(\frac{E}{A}\right)^{1/n}\underbrace{\int_0^1 t^{\frac{1}{n}-1}(1-t)^{-1/2}dt}_{\frac{\Gamma(\frac{1}{n})\Gamma\,\cancel{\Gamma(1/2)}\ \sqrt{\pi}}{\Gamma(\frac{1}{n}+\frac{1}{2})}} \\ \tau = \quad & \frac{2}{n}\sqrt{\frac{2\pi m}{E}}\left(\frac{E}{A}\right)^{1/n}\frac{\Gamma}{(}\frac{1}{n})\Gamma(\frac{1}{n}+\frac{1}{2}) \end{aligned}$$

2. SERIES DE STIRLING

$$\ln(n!) \rightsquigarrow \ln\Gamma(z)$$

para z no entero, $|z|$ muy grande, no Series de Maclaurin

$$\Gamma(z+1) = z\Gamma(z)$$

Se usa una expansi'on asint'otica para $\ln(\Gamma(z))$ esta expansi'on es la que se conoce como Serie de Stirling o Formula de Stirling

Series asint'oticas→ Series semi-convergentes→Series de Poincar'e

Se consideran 2 tipos de integrales

$$\begin{aligned} I_1(x) = \quad & \int_x^\infty e^{-u}f(u)du \\ I_2(x) = \quad & \int_0^\infty e^{-u}f\left(\frac{u}{x}\right)du \end{aligned}$$

Qu'e podemos hacer con las series asint'oticas:

1. Si multiplicamos 2 series asint'oticas vamos a tener otra serie asint'otica.
2. Se pueden integrar t'ermino a t'ermino u el resultado ser'a, otra serie asintotica de la forma

$$\int_x^\infty f(x)dx$$

3. La diferenciaci'on es v'alido 'unicamente bajo ciertas condiciones.

M'etodo para generar una serie asint'otica:

$$\left.\begin{array}{l}\textbf{M'etodo de Steepest Descents}\end{array}\right\} \begin{array}{l}\textit{a lo largo del camino}\\ \textit{m'aximo decrecimiento}\end{array}$$

$\rightarrow$ *Saddle point* $\rightarrow$ *punto silla*
$\rightarrow$ *M'etodo de integrandos Holomorfos*

Derivaci'on de la f'ormula de integraci'on de Euler-Maclaurin

$\longrightarrow$ *uno de los usos de los polinomios de Bernoulli* $\longleftarrow$

N'umeros de Bernoulli $\longrightarrow$ Polinomios de Bernoulli

$$B_n$$

$$\frac{t}{e^t - 1} = \sum_{n=0}^{\infty} \frac{B_n t^n}{n!}$$

Como $\frac{t}{e^t-1} = \sum_{n=0}^{\infty} \frac{B_n t^n}{n!}$ es una serie de Taylor se pueden identificar a B_n como una sucesi'on de derivadas de la funci'on generadora valuadas en cero.

$$\Rightarrow B_n = \frac{d^n}{dt^n}\left(\frac{t}{e^t-1}\right)\bigg|_{t=0}$$

$$\begin{aligned}
B_0 &= \lim_{n\to 0} \tfrac{t}{e^t-1} = 1 \\
B_1 &= \lim_{n\to 0} \tfrac{d}{dt}\left[\tfrac{t}{e^t-1}\right]\Big|_{t=0} = -\tfrac{1}{2} \\
B_1 &= \lim_{n\to 0} \left(\tfrac{1}{e^t-1} - \tfrac{te^t}{(e^t-1)^2}\right) \\
B_1 &= \lim_{n\to 0} \left(\tfrac{e^t-1-te^t}{(e^t-1)^2}\right) = \lim_{n\to 0} \tfrac{\cancel{e^t}-t\cancel{e^t}-\cancel{e^t}}{2(e^t-1)\cancel{e^t}} \\
&= -\tfrac{1}{2}
\end{aligned}$$

La f'ormula Euler-Maclaurin par aevaluar integrales definidas en $(0,\infty)$

$$\begin{aligned}
\textstyle\int_0^\infty = \;& \tfrac{1}{2} f(0) + f(1) + f(2) + f(3...) \\
& + \tfrac{B_2}{2!} f'(0) + \tfrac{B_4}{4!} f^{(3)}(0) + \tfrac{B_6}{6!} f^{(5)}(0)
\end{aligned}$$

Donde B_n son los n'umeros de Bernoulli

$$B_2 = \frac{1}{6},\ B_4 = -\frac{1}{30},\ B_6 = \frac{1}{42},\ B_8 = -\frac{1}{30}, ...$$

$$\int_0^\infty \frac{dx}{(z+x)^2} = \frac{1}{z}$$

z no est'a en el eje real negativo

$$f(1) + f(2) + f(3) + ... = \sum_{n=1}^{\infty} \frac{1}{(z+n)^2} =$$

$$\varphi^{(m)}(z+1) = \frac{d^{m+1}}{dz^{m+1}} \ln\Gamma(z+1)$$

$$\underset{m=1,2,3...}{=} (-1)^{m+1} m! \sum_{n=2}^{\infty} \frac{1}{(z+n)^{m+1}}$$

m=1

$$(-1)^2 \cdot 1! \sum_{n=1}^{\infty} \frac{1}{(z+n)^2} = \varphi^{(1)}(z+1)$$

$$f_{(0)}^{(2n+1)} = \left(\frac{d}{dx}\right)^{2n-1} \frac{1}{(z+x)^2}\Bigg|_{x=0} = -\frac{(2n)!}{z^{2n+1}}$$

$$\int_0^{\infty} f(x)dx = \int_0^{\infty} \frac{dx}{(z+x)^2} = \frac{1}{z} = \frac{1}{2}\frac{1}{z^2} + \underbrace{\varphi_{(z+1)}^{(1)}}_{\textit{poligamma con m=1}} - \frac{B_2}{z^3} - \frac{B_4}{z^5} - ...$$

$$\varphi_{(z+1)}^{(1)} = \frac{d}{dz} \underbrace{\varphi(z+1)}_{digamma}$$

$$\varphi_{(z+1)}^{(1)} = \frac{1}{z} - \frac{1}{2z^2} + \frac{B_2}{z^3} + \frac{B_4}{z^5} + ...$$

$$= \frac{1}{z} - \frac{1}{2z^2} + \underbrace{\sum_{n=1}^{\infty} \frac{B_{2n}}{z^{2n+1}}}_{\substack{\textit{Es una serie asintotica} \\ \textit{que es util si se} \\ \textit{utilizan pocos t'erminos}}}$$

A; integrar $\varphi_{(z+1)}^{(1)}$ obtenemos $\varphi_{(z+1)}$[Digamma]