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Constitutive model

Von Mises Model

Plane strain or Axisymmetric

PEEQ

$\bar{\epsilon}^p$ is the effective/accumulated/equivalent plastic strain(PEEQ). It is used as a **scalar** measure of the tensor plastic strain ϵ^p .

$$\bar{\epsilon}^p \equiv \int_0^t \sqrt{\frac{2}{3} \dot{\epsilon}^p : \dot{\epsilon}^p} dt = \int_0^t \sqrt{\frac{2}{3}} \|\dot{\epsilon}^p\| dt$$

$$\dot{\bar{\epsilon}}^p = \sqrt{\frac{2}{3} \dot{\epsilon}^p : \dot{\epsilon}^p} = \sqrt{\frac{2}{3}} \|\dot{\epsilon}^p\| = \dot{\gamma}$$

In return mapping, we have

$$\epsilon_{n+1}^p = \epsilon_n^p + \Delta\gamma \sqrt{\frac{3}{2}} \frac{\mathbf{s}_{n+1}}{\|\mathbf{s}_{n+1}\|}$$

which denotes the relationship between $\bar{\epsilon}^p$ and ϵ^p .

Material parameters

Shear modulus:

$$G = \frac{E}{2(1 + \nu)}$$

Bulk modulus:

$$K = \frac{E}{3(1 - 2\nu)}$$

Tensor operator

[Tensor double contraction]

: denotes double contraction of tensors (internal product of second-order tensors).

For 2 second-order tensors A and B:

$$A : B = \text{tr}(A \cdot B^T) = A_{ij} B_{ij} \text{ (scalar)}$$

For A with order m and B with order n ($m > n$), their double contraction is a tensor C with order (m-n):

eg. for A with order 5 and B with order 2:

$$C_{ijk} = (A : B)_{ijk} = A_{ijkpq} B_{pq}$$

[Tensor product]

" \otimes " denotes tensor product of tensors.

For tensor A with order m and tensor B with order n, their tensor product is a tensor with order (m+n).

Thus each item is:

$$(A \otimes B)_{i_1 i_2 \dots i_{m+n}} = A_{i_1 i_2 \dots i_m} \cdot B_{i_{m+1} \dots i_{m+n}}$$

$$A \otimes B \neq B \otimes A$$

Second-order tensor analysis

Symmetric tensor: $\mathbf{T} = \mathbf{T}^T$

Skew symmetric tensor: $\mathbf{T} = -\mathbf{T}^T$

Any tensor can be decomposed as

$$\mathbf{T} = \text{sym}(\mathbf{T}) + \text{skew}(\mathbf{T})$$

The symmetric part: $\text{sym}(\mathbf{T}) \equiv \frac{1}{2}(\mathbf{T} + \mathbf{T}^T)$

The skew part: $\text{skew}(\mathbf{T}) \equiv \frac{1}{2}(\mathbf{T} - \mathbf{T}^T)$

Common tool tensors

\mathbf{I} is a second-order identity tensor:

$$I_{ij} = \mathbf{e}_i \cdot \mathbf{e}_j = \delta_{ij}$$

\mathbf{I} is a 4th-order identity tensor:

$$I_{ijkl} = \delta_{ik} \delta_{jl}$$

\mathbf{I}_S is a 4th-order symmetric identity tensor:

$$I_{ijkl} = \frac{1}{2}(\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk})$$

\mathbf{I}_d is a 4th-order deviatoric projection tensor:

$$I_d \equiv \mathbf{I}_S - \frac{1}{3} \mathbf{I} \otimes \mathbf{I}$$

Array notation of some tensors

[Second-order identity tensor]

For plane strain and axisymmetric problems

$$\mathbf{I} \rightarrow \mathbf{i} = [1, 1, 0, 1]^T$$

[4th-order symmetric identity tensor]

For plane problems

$$\mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{1}{2} \end{bmatrix}$$

For axisymmetric problems

$$\mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 & 0 \\ & 1 & 0 & 0 \\ \text{sym} & & \frac{1}{2} & 0 \\ & & & 1 \end{bmatrix}$$

For 3D cases

$$\mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ & 1 & 0 & 0 & 0 & 0 \\ & & 1 & 0 & 0 & 0 \\ \text{sym} & & & \frac{1}{2} & 0 & 0 \\ & & & & \frac{1}{2} & 0 \\ & & & & & \frac{1}{2} \end{bmatrix}$$

[4th-order isotropic elasticity tensor]

$$\mathbf{D}^e = 2G\mathbf{I}_S + A(K - \frac{2}{3}G)\mathbf{ii}^T$$

where \mathbf{I}_S is the 2D array form.

In plane strain, axisymmetric and 3D analyses: $A = 1$

In plane stress: $A = \frac{2G}{K + \frac{4}{3}G}$

For 3D problems:

$$\begin{pmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{12} \\ \sigma_{23} \\ \sigma_{13} \end{pmatrix} = \begin{bmatrix} K + \frac{4}{3}G & K - \frac{2}{3}G & K - \frac{2}{3}G & 0 & 0 & 0 \\ K - \frac{2}{3}G & K + \frac{4}{3}G & K - \frac{2}{3}G & 0 & 0 & 0 \\ K - \frac{2}{3}G & K - \frac{2}{3}G & K + \frac{4}{3}G & 0 & 0 & 0 \\ 0 & 0 & 0 & G & 0 & 0 \\ 0 & 0 & 0 & 0 & G & 0 \\ 0 & 0 & 0 & 0 & 0 & G \end{bmatrix} \begin{pmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{33} \\ 2\varepsilon_{12} \\ 2\varepsilon_{23} \\ 2\varepsilon_{13} \end{pmatrix}$$

Split of tensors

The strain tensor can be written as:

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}_d + \boldsymbol{\varepsilon}_v$$

volumetric strain tensor:

$$\boldsymbol{\epsilon}_v \equiv \frac{1}{3}\epsilon_v \mathbf{I}$$
$$\epsilon_v = \frac{1}{3}(\mathbf{I} \otimes \mathbf{I}) : \boldsymbol{\epsilon}$$

volumetric strain:

$$\epsilon_v \equiv I_1(\boldsymbol{\epsilon}) = \text{tr} \boldsymbol{\epsilon} = \text{tr} \nabla^s \mathbf{u} = \text{tr} \nabla \mathbf{u}$$

deviatoric strain tensor:

$$\boldsymbol{\epsilon}_d = \boldsymbol{\epsilon} - \boldsymbol{\epsilon}_v$$
$$\boldsymbol{\epsilon}_d = [\mathbf{I}_S - \frac{1}{3}\mathbf{I} \otimes \mathbf{I}] : \boldsymbol{\epsilon} = \mathbf{I}_d : \boldsymbol{\epsilon}$$

where $\text{tr} \boldsymbol{\epsilon}_d = 0$

The stress tensor can be written as

$$\boldsymbol{\sigma} = \mathbf{s} + p\mathbf{I}$$

hydrostatic stress/hydrostatic pressure/mean stress/mean normal pressure:

$$p \equiv \frac{1}{3}\text{tr} \boldsymbol{\sigma}$$

deviatoric stress tensor/stress deviator:

$$\mathbf{s} \equiv \boldsymbol{\sigma} - p\mathbf{I} = \mathbf{I}_d : \boldsymbol{\sigma}$$

\mathbf{I}_d is a 4th-order tensor named deviatoric projection tensor.

$$\mathbf{I}_d \equiv \mathbf{I}_S - \frac{1}{3}\mathbf{I} \otimes \mathbf{I}$$

\mathbf{I}_S is a 4th-order tensor named symmetric identity tensor:

$$\mathbf{I}_{ijkl} = \frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})$$

This deviatoric projection tensor can project a second-order symmetric tensor \mathbf{S} into the deviatoric subspace, i.e. into the space of traceless tensors.

$$\text{dev}(\mathbf{S}) \equiv \mathbf{I}_d : \mathbf{S}$$

$p\mathbf{I}$ is called spherical stress tensor:

$$p\mathbf{I} = \frac{1}{3}(\mathbf{I} \otimes \mathbf{I}) : \boldsymbol{\sigma}$$

Invariant (scalar)

[First principal invariant I_1]

$$\begin{aligned} I_1(\mathbf{S}) &\equiv \text{tr} \mathbf{S} = S_{ii} \\ &= s_{11} + s_{22} + s_{33} \end{aligned}$$

[Second principal invariant I_2]

$$\begin{aligned} I_2(\mathbf{S}) &\equiv \frac{1}{2}[(\text{tr} \mathbf{S})^2 - \text{tr}(\mathbf{S}^2)] = \frac{1}{2}(S_{ii}S_{jj} - S_{ij}S_{ji}) \\ &= s_{23}^2 + s_{31}^2 + s_{12}^2 - s_{22}s_{33} - s_{33}s_{11} - s_{11}s_{22} \end{aligned}$$

[Third principal invariant I_3]

$$\begin{aligned} I_3(\mathbf{S}) &\equiv \det \mathbf{S} = \frac{1}{6}\epsilon_{ijk}\epsilon_{pqr}S_{ip}S_{jq}S_{kr} \\ &= 2s_{23}s_{31}s_{12} + s_{11}s_{22}s_{33} - s_{11}s_{23}^2 - s_{22}s_{31}^2 - s_{33}s_{12}^2 \end{aligned}$$

Here ϵ_{ijk} is the **permutation tensor**

$$\epsilon_{ijk} = \frac{1}{2}(i-j)(j-k)(k-i)$$

Thus

ϵ_{ijk}	Condition
$\epsilon_{ijk} = 0$	if any two of i,j,k are equal
$\epsilon_{ijk} = 1$	for an even permutation (eg. 123,231,312)
$\epsilon_{ijk} = -1$	for an odd permutation (eg. 132,213,321)

[J_2 invariant of the stress deviator]

Note:

stress deviator \mathbf{s} is symmetric $\Rightarrow \mathbf{s} : \mathbf{s} = \text{tr}(\mathbf{s} \cdot \mathbf{s}^T) = \mathbf{s} : \mathbf{s} = \text{tr}(\mathbf{s} \cdot \mathbf{s})$

trace of this deviatoric stress tensor is 0 $\Rightarrow \text{tr} \mathbf{s} = 0$

$$J_2 \equiv -I_2(\mathbf{s}) = \frac{1}{2}\mathbf{s} : \mathbf{s} = \frac{1}{2}\text{tr}[\mathbf{s} \cdot \mathbf{s}] = \frac{1}{2}\|\mathbf{s}\|^2$$

[J_3 invariant of the stress deviator]

$$J_3 \equiv I_3(\mathbf{s}) \equiv \det \mathbf{s} = \frac{1}{3}\text{tr}(\mathbf{s})^3$$

Engineering strain/stress

Plane strain/stress:

$$\boldsymbol{\epsilon} = [\epsilon_{11}, \epsilon_{22}, 2\epsilon_{12}]^T$$

$$\boldsymbol{\sigma} = [\sigma_{11}, \sigma_{22}, \sigma_{12}]^T$$

Axisymmetric:

$$\boldsymbol{\epsilon} = [\epsilon_{11}, \epsilon_{22}, 2\epsilon_{12}, \epsilon_{33}]^T$$

$$\boldsymbol{\sigma} = [\sigma_{11}, \sigma_{22}, \sigma_{12}, \sigma_{33}]^T$$

3D analysis:

$$\boldsymbol{\epsilon} = [\epsilon_{11}, \epsilon_{22}, \epsilon_{33}, 2\epsilon_{12}, 2\epsilon_{23}, 2\epsilon_{13}]^T$$

$$\boldsymbol{\sigma} = [\sigma_{11}, \sigma_{22}, \sigma_{33}, \sigma_{12}, \sigma_{23}, \sigma_{13}]^T$$

Associate flow rule

Associative flow rule assumes plastic strain increment is normal to the yield at that point. Thus we can assume yield function Φ to be the same as the flow potential function Ψ .

$$\Phi \equiv \Psi$$

In associative models, the evolution equations are given by

$$\dot{\boldsymbol{\epsilon}}^p = \dot{\gamma} \frac{\partial \Phi}{\partial \boldsymbol{\sigma}} = \dot{\gamma} \mathbf{N}$$

\mathbf{N} is called plastic flow vector, $\dot{\gamma}$ is the plastic multiplier.

Infinitesimal Isotropic elasticity tensor (fourth-order tensor)

$$\mathbf{D}^e = 2G\mathbf{I}_S + A(K - \frac{2}{3}G)\mathbf{I} \otimes \mathbf{I}$$

where G and K are shear and bulk moduli respectively, \mathbf{I} is the second-order identity and \mathbf{I}_S is the fourth-order symmetric identity tensor.

4th-order identity tensor \mathbf{I} : $\mathbf{I}_{ijkl} = \delta_{ik}\delta_{jl}$

4th-order symmetric identity tensor \mathbf{I}_S : $\mathbf{I}_{ijkl} = \frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})$

In plane strain, axisymmetric and 3D analyses: $A = 1$

In plane stress: $A = \frac{2G}{K + \frac{4}{3}G}$

DDSDDE

DDSDDE in Abaqus:

$$\mathbf{D} = \frac{\partial \Delta \boldsymbol{\sigma}}{\partial \Delta \boldsymbol{\epsilon}}$$

In other books this is expressed as:

Consistent tangent modulus(change within each NR iteration!!):

$$\mathbf{D} = \frac{\partial \boldsymbol{\sigma}_{n+1}}{\partial \boldsymbol{\epsilon}_{n+1}}$$

Continuum tangent modulus(Also change within each NR iteration):

not known yet

Continuum tangent operator has nothing to do with $\Delta \gamma$ thus is not suitable for return mapping method considering convergence rate.

Anyway, consistent tangent operator rather than continuum is recommended.

UMAT is to give out stress tensor(SUVM), consistent tangent operator (DDSDDE in abaqus)(CTVM) and updated state variables given current state variables, stress tensor, strain tensor and trial strain increment.

Full Newton-Raphson Scheme for FEM

(i) $k := 0$. Set initial guess and residual

$$\mathbf{u}_{n+1}^{(0)} := \mathbf{u}_n; \quad \mathbf{r} := \mathbf{f}^{\text{int}}(\mathbf{u}_n) - \lambda_{n+1} \bar{\mathbf{f}}^{\text{ext}}$$

(ii) Compute consistent tangent matrices [MATICT]

$$\mathbf{D} := \partial \hat{\boldsymbol{\sigma}} / \partial \boldsymbol{\varepsilon}_{n+1}$$

(iii) Assemble element tangent stiffness matrices [ELEIST, STSTD2]

$$\mathbf{K}_T^{(e)} := \sum_{i=1}^{n_{\text{gausp}}} w_i j_i \mathbf{B}_i^T \mathbf{D}_i \mathbf{B}_i$$

(iv) $k := k + 1$. Assemble global stiffness and solve for $\delta \mathbf{u}^{(k)}$ [FRONT]

$$\mathbf{K}_T \delta \mathbf{u}^{(k)} = -\mathbf{r}^{(k-1)}$$

(v) Apply Newton correction to displacements [UPCONF]

$$\mathbf{u}_{n+1}^{(k)} := \mathbf{u}_{n+1}^{(k-1)} + \delta \mathbf{u}^{(k)}$$

(vi) Update strains [IFSTD2]

$$\boldsymbol{\varepsilon}_{n+1}^{(k)} := \mathbf{B} \mathbf{u}_{n+1}^{(k)}$$

(vii) Use constitutive integration algorithm to update stresses and other state variables [MATISU]

$$\boldsymbol{\sigma}_{n+1}^{(k)} := \hat{\boldsymbol{\sigma}}(\boldsymbol{\alpha}_n, \boldsymbol{\varepsilon}_{n+1}^{(k)}); \quad \boldsymbol{\alpha}_{n+1}^{(k)} := \hat{\boldsymbol{\alpha}}(\boldsymbol{\alpha}_n, \boldsymbol{\varepsilon}_{n+1}^{(k)})$$

(viii) Compute element internal force vectors [INTFOR, IFSTD2]

$$\mathbf{f}_{(e)}^{\text{int}} := \sum_{i=1}^{n_{\text{gausp}}} w_i j_i \mathbf{B}_i^T \boldsymbol{\sigma}_{n+1}^{(k)} \Big|_i$$

(ix) Assemble global internal force vector and update residual [CONVER]

$$\mathbf{r} := \mathbf{f}^{\text{int}} - \lambda_{n+1} \bar{\mathbf{f}}^{\text{ext}}$$

(x) Check for convergence [CONVER]

IF $\|\mathbf{r}\| / \|\mathbf{f}^{\text{ext}}\| \leq \epsilon_{\text{tol}}$ THEN set $(\cdot)_{n+1} := (\cdot)_{n+1}^{(k)}$ and EXIT
ELSE GOTO (ii)

In this scheme, K_T is updated in each iteration during the increment.

A variant of this is the modified NR scheme where K_T remains unchanged during each increment with the initial value.

In my previous multiscale framework, the strain increment $\delta \mathbf{u}$ is checked as the threshold, while in this scheme, the residual \mathbf{r} is checked. This results in the difference that in this scheme the

consistent tangent operator (k) is calculated based on the state of the previous iteration (k-1). In essence they are the same.

State Update for Von Mises Model (Fortran program SUVVM)

Given state variables and trial strain, give out updated state variables and stress.

Input:

- **ntype** : analysis type (2→plane strain;3→axisymmetric;others→ not covered for SUVVM)
- **iprops** : array of integer materials properties (the only one used in SUVVM is nhard: number of sampling points for the piecewise hardening curve)
- **rprops** : array of real material properties. For SUVVM, the needs are $[E, \nu, \bar{\epsilon}_0^p, \sigma_{y0}, \bar{\epsilon}_1^p, \sigma_{y1}, \dots, \bar{\epsilon}_{\text{nhard}}^p, \sigma_{y \text{nhard}}]$
- **rstava** : array of state variables other than stress components $[\epsilon^e * 4, \bar{\epsilon}^p * 1]$, 4 elements for elastic strain, 1 for equivalent plastic strain. For SUVVM input, only $\bar{\epsilon}^p$ of the previous iteration is required.
- **strat** : $\epsilon_{n+1}^{e \text{ trial}}$, array of the elastic trial strains in engineering form. $[\epsilon_{11}, \epsilon_{22}, 2\epsilon_{12}, \epsilon_{33}]$

Return:

- $\Delta\gamma$: plastic multiplier from return mapping (0 if purely elastic)
- **lalgva**: array of logical algorithmic flags or variables. (For SUVVM, this contains [ifplas(whether plastic corrector involved), sufail(True: the NR iteration of return mapping fails to converge and the state variables are not updated)])
- **stres** : array of updated stress components $[\sigma_{11}, \sigma_{22}, \sigma_{12}, \sigma_{33}]$
- **rstava** : array of state variables other than stress components $[\epsilon^e * 4, \bar{\epsilon}^p * 1]$, 4 elements for elastic strain, 1 for equivalent plastic strain. For SUVVM output, all these elements are updated.

1. Check type

If $\text{ntype} \neq 2$ (plane strain) and $\text{ntype} \neq 3$ (axisymmetric), return error.

2. Assign and assemble basic parameters

$$\Delta\gamma = 0 \quad G = \frac{E}{2(1+\nu)} \quad K = \frac{E}{3(1-2\nu)}$$

3. Elastic predictor

$$\epsilon_{n+1}^{e \text{ trial}} := \epsilon_n^e + \Delta\epsilon$$

This step is implemented before calling SUVVM and is stored in form $[\epsilon_{11}, \epsilon_{22}, 2\epsilon_{12}, \epsilon_{33}]$.

$$\bar{\epsilon}_{n+1}^{p \text{ trial}} := \bar{\epsilon}_n^p$$

Equivalent plastic strain (scalar) is directly obtained from state variables.

$$\epsilon_{v n+1}^{e \text{ trial}} := \text{tr}(\epsilon_{n+1}^{e \text{ trial}}) = \epsilon_{11} + \epsilon_{22} + \epsilon_{33}$$

$$\epsilon_{d n+1}^{e \text{ trial}} = \epsilon_{n+1}^{e \text{ trial}} - \frac{1}{3}\epsilon_{v n+1}^{e \text{ trial}} \mathbf{I}$$

Note: for strain in engineering form $\epsilon_{d n+1 \ 12}^{e \text{ trial}} = \frac{1}{2}\epsilon_{n+1 \ 12}^{e \text{ trial}}$

$$p_{n+1}^{\text{trial}} := K \epsilon_{v n+1}^{e \text{ trial}} \quad \mathbf{s}_{n+1}^{\text{trial}} := 2G \epsilon_{d n+1}^{e \text{ trial}}$$

$$q_{n+1}^{\text{trial}} := \sqrt{3 \mathbf{J}_2(\mathbf{s}_{n+1}^{\text{trial}})} = \sqrt{\frac{3}{2} \mathbf{s}_{n+1}^{\text{trial}} : \mathbf{s}_{n+1}^{\text{trial}}} \quad (\text{Trial effective stress})$$

4. Check plastic admissibility

The yield function for Von Mises model ($J - 2$ flow theory) is

$$\Phi(\boldsymbol{\sigma}, \sigma_y) = \sqrt{3 \mathbf{J}_2(\mathbf{s}(\boldsymbol{\sigma}))} - \sigma_y$$

$$\text{where } \mathbf{J}_2(\mathbf{s}(\boldsymbol{\sigma})) = \frac{1}{2} \mathbf{s}(\boldsymbol{\sigma}) : \mathbf{s}(\boldsymbol{\sigma}) \quad \sigma_y = \sigma_y(\bar{\epsilon}^p)$$

Calculate $\sigma_y(\bar{\epsilon}_{n+1}^{p \text{ trial}})$

This yield stress is obtained from interpolation of the hardening curve based on current trial equivalent plastic strain.

If $q_{n+1}^{\text{trial}} - \sigma_y(\bar{\epsilon}_{n+1}^{p \text{ trial}}) \leq 0$

In fact $\frac{q_{n+1}^{\text{trial}} - \sigma_y(\bar{\epsilon}_{n+1}^{p \text{ trial}})}{\sigma_y(\bar{\epsilon}_{n+1}^{p \text{ trial}})} \leq 10^{-6}$ is checked in codes.

Then purely elastic:

$$\boldsymbol{\epsilon}_{n+1}^e := \boldsymbol{\epsilon}_{n+1}^{e \text{ trial}}$$

$$\bar{\epsilon}_{n+1}^p := \bar{\epsilon}_{n+1}^{p \text{ trial}} = \bar{\epsilon}_n^p \quad (\text{not changed})$$

$$\boldsymbol{\sigma}_{n+1} := \mathbf{s}_{n+1}^{\text{trial}} + p_{n+1}^{\text{trial}} \mathbf{I}$$

Goto 7.

Else:

Goto 5.

5. Return mapping

Solve the equation about the yield function with respect to $\Delta\gamma$.

$$\tilde{\Phi}(\Delta\gamma) = q_{n+1}^{\text{trial}} - 3G\Delta\gamma - \sigma_y(\bar{\epsilon}_n^p + \Delta\gamma) = 0$$

Newton-Raphson for solving $\Delta\gamma$:

1. For initial iteration $k := 0$,
set $\Delta\gamma^{(0)} := 0$ and $\tilde{\Phi} = q_{n+1}^{\text{trial}} - \sigma_y(\bar{\epsilon}_n^p)$
2. Perform N-R iteration
 $H := \left. \frac{d\sigma_y}{d\bar{\epsilon}^p} \right|_{\bar{\epsilon}_n^p + \Delta\gamma}$ (Hardening slope)
 $d := \frac{d\tilde{\Phi}}{d\Delta\gamma} = -3G - H$ (Residual derivative)
 $\Delta\gamma := \Delta\gamma - \frac{\tilde{\Phi}}{d}$ (update $\Delta\gamma$)
3. Check for convergence
 $\tilde{\Phi}(\Delta\gamma) := q_{n+1}^{\text{trial}} - 3G\Delta\gamma - \sigma_y(\bar{\epsilon}_n^p + \Delta\gamma)$
 If $|\tilde{\Phi}| \leq \epsilon_{tol}$
 In fact $\frac{|\tilde{\Phi}|}{\sigma_y(\bar{\epsilon}_n^p + \Delta\gamma)} \leq \epsilon_{tol}$ is checked in codes.
 Then Goto 6.
 Else Goto local 2.

6. Update state variables and stress

$$p_{n+1} := p_{n+1}^{\text{trial}}; \quad \mathbf{s}_{n+1} := \left(1 - \frac{\Delta\gamma 3G}{q_{n+1}^{\text{trial}}}\right) \mathbf{s}_{n+1}^{\text{trial}}$$

$$\begin{aligned}\boldsymbol{\sigma}_{n+1} &:= \boldsymbol{s}_{n+1} + p_{n+1} \boldsymbol{I} \\ \boldsymbol{\epsilon}_{n+1}^e &:= \frac{1}{2G} \boldsymbol{s}_{n+1} + \frac{1}{3} \epsilon_{v\,n+1}^{e\,\text{trial}} \boldsymbol{I} \\ \bar{\epsilon}_{n+1}^p &:= \bar{\epsilon}_n^p + \Delta\gamma\end{aligned}$$

7. Return other parameters

$\Delta\gamma$, ifplas, sufail

Consistent tangent operator for Von Mises Model (CTVM)

Given state variables, give out consistent tangent operator (DDSDDE).

Input:

- $\Delta\gamma$: plastic multiplier from SUVM
- epflag : whether elastic or elastoplastic (from SUVM)
- ntype : analysis type (plane strain(2) or axisymmetric(3))
- iprops : array of integer material parameters (num of points of the hardening curve...)
- rprops : array of real material parameters ($E, \nu, \bar{\epsilon}_0^p, \sigma_{y0}, \bar{\epsilon}_1^p, \sigma_{y1} \dots$)
- rstava : array of updated state parameters ($\boldsymbol{\epsilon}^e * 4, \bar{\epsilon}^p * 1$) from SUVM
- stres : array of updated stress tensor $\boldsymbol{\sigma}_{n+1}^{k-1}$ from SUVM

Return:

\mathbf{D}^e or \mathbf{D}^{ep}

1. Initialization

4th-order symmetric identity tensor:

$$\mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 & 0 \\ & 1 & 0 & 0 \\ \text{sym} & & \frac{1}{2} & 0 \\ & & & 1 \end{bmatrix}$$

2nd-order identity tensor:

$$\boldsymbol{I} \rightarrow \mathbf{i} = [1, 1, 0, 1]^T$$

2. Check ntype

if ntype $\neq 2$ and ntype $\neq 3$, return Error.

3. Get and assemble material properties and projection tensor

Shear modulus:

$$G = \frac{E}{2(1 + \nu)}$$

Bulk modulus:

$$K = \frac{E}{3(1 - 2\nu)}$$

Deviatoric projection tensor:

$$\begin{aligned} \mathbf{I}_d &\equiv \mathbf{I}_S - \frac{1}{3} \mathbf{I} \otimes \mathbf{I} \\ &= \mathbf{I}_S - \frac{1}{3} \mathbf{i} \cdot \mathbf{i}^T \end{aligned}$$

4. If epflag = True, goto 5(elsatoplastic); else goto 6(elastic).

5. Compute elastoplastic consistent tangent operator

$$\begin{aligned} \mathbf{D}^{ep} &= \mathbf{D}^e - \frac{\Delta\gamma 6G^2}{q_{n+1}^{\text{trial}}} \mathbf{I}_d + 6G^2 \left(\frac{\Delta\gamma}{q_{n+1}^{\text{trial}}} - \frac{1}{3G + H} \right) \bar{\mathbf{N}}_{n+1} \otimes \bar{\mathbf{N}}_{n+1} \\ &= 2G \left(1 - \frac{\Delta\gamma 3G}{q_{n+1}^{\text{trial}}} \right) \mathbf{I}_d + K \mathbf{I} \otimes \mathbf{I} \\ &\quad + 6G^2 \left(\frac{\Delta\gamma}{q_{n+1}^{\text{trial}}} - \frac{1}{3G + H} \right) \bar{\mathbf{N}}_{n+1} \otimes \bar{\mathbf{N}}_{n+1} \end{aligned}$$

The unknowns here are q_{n+1}^{trial} and $\bar{\mathbf{N}}_{n+1}$.

Since we have

$$\begin{aligned} q_{n+1}^{\text{trial}} &\equiv \sqrt{3J_2(\mathbf{s}_{n+1}^{\text{trial}})} = \sqrt{\frac{3}{2}} \|\mathbf{s}_{n+1}^{\text{trial}}\| \\ \bar{\mathbf{N}}_{n+1} &= \frac{\mathbf{s}_{n+1}^{\text{trial}}}{\|\mathbf{s}_{n+1}^{\text{trial}}\|} = \frac{\mathbf{s}_{n+1}}{\|\mathbf{s}_{n+1}\|} \end{aligned}$$

And also we have these knowns

- Hydrostatic pressure

$$p = \frac{1}{3} \sigma_{ii}$$

- Deviatoric stress components $\mathbf{s}_{n+1} = [s_{11}, s_{22}, s_{12}, s_{33}]$

$$s_i = \sigma_i - p \quad (i = 1, 2, 4)$$

$$s_i = \sigma_i \quad (i = 3)$$

Thus the only gap is the relationship between \mathbf{s}_{n+1} and $\mathbf{s}_{n+1}^{\text{trial}}$.

$$\begin{aligned} \mathbf{s}_{n+1} &= \left(1 - \frac{3\Delta\gamma G}{q_{n+1}^{\text{trial}}} \right) \mathbf{s}_{n+1}^{\text{trial}} \\ &= \left(1 - \sqrt{\frac{3}{2}} \frac{\Delta\gamma 2G}{\|\mathbf{s}_{n+1}^{\text{trial}}\|} \right) \mathbf{s}_{n+1}^{\text{trial}} \\ \Rightarrow \|\mathbf{s}_{n+1}^{\text{trial}}\| &= \|\mathbf{s}_{n+1}\| + \sqrt{6}\Delta\gamma G \end{aligned}$$

Therefore, q_{n+1}^{trial} can be expressed as

$$q_{n+1}^{\text{trial}} = \sqrt{\frac{3}{2}} \|\mathbf{s}_{n+1}^{\text{trial}}\| = \sqrt{\frac{3}{2}} \|\mathbf{s}_{n+1}\| + 3\Delta\gamma G$$

6. Compute elastic tangent operator

$$\mathbf{D}^e = 2G\mathbf{I}_S + A(K - \frac{2}{3}G)\mathbf{I} \otimes \mathbf{I}$$

In plane strain, axisymmetric and 3D analyses: $A = 1$

Thus

$$\begin{aligned}\mathbf{D}^e &= 2G\mathbf{I}_S + (K - \frac{2}{3}G)\mathbf{I} \otimes \mathbf{I} \\ &= 2G(\mathbf{I}_S - \frac{1}{3}\mathbf{I} \otimes \mathbf{I}) + K * \mathbf{I} \otimes \mathbf{I} \\ &= 2G\mathbf{I}_d + K * \mathbf{I} \otimes \mathbf{I}\end{aligned}$$

Plane Stress Plasticity

Why plane stress is special?

The generic algorithmic formulas for **3-dimensional** cases are based on the fact that stress is unknown and is updated based on strain.

Plane strain/axisymmetric conditions can be seen as degeneration of 3D cases where some components of strain are prescribed and the stress are updated in the same way based on input strain.

While for **plane stress** problems, stress components rather than strain are prescribed. Thus the normal strain \Rightarrow stress flow cannot be obtained and the 3D numerical integration algorithms should be modified.

Plane stress linear elasticity

For common 3D cases:

$$\begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{12} \\ \sigma_{23} \\ \sigma_{13} \end{Bmatrix} = \begin{bmatrix} K + \frac{4}{3}G & K - \frac{2}{3}G & K - \frac{2}{3}G & 0 & 0 & 0 \\ K - \frac{2}{3}G & K + \frac{4}{3}G & K - \frac{2}{3}G & 0 & 0 & 0 \\ K - \frac{2}{3}G & K - \frac{2}{3}G & K + \frac{4}{3}G & 0 & 0 & 0 \\ 0 & 0 & 0 & G & 0 & 0 \\ 0 & 0 & 0 & 0 & G & 0 \\ 0 & 0 & 0 & 0 & 0 & G \end{bmatrix} \begin{Bmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{33} \\ 2\varepsilon_{12} \\ 2\varepsilon_{23} \\ 2\varepsilon_{13} \end{Bmatrix}$$

Based on constraints $\sigma_{33} = \sigma_{23} = \sigma_{13} = 0$

↓

$$\varepsilon_{13} = 0, \quad \varepsilon_{23} = 0$$

$$\varepsilon_{33} = -\frac{3K - 2G}{3K + 4G}(\varepsilon_{11} + \varepsilon_{22}) = -\frac{\nu}{1 - \nu}(\varepsilon_{11} + \varepsilon_{22})$$

↓

Plane stress cases (Only pay attention to in-plane components):

$$\begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{Bmatrix} = \frac{E}{1 - \nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} \begin{Bmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ 2\varepsilon_{12} \end{Bmatrix}$$

Or equivalently

$$\begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{Bmatrix} = 2G \begin{bmatrix} 1 + \alpha & \alpha & 0 \\ \alpha & 1 + \alpha & 0 \\ 0 & 0 & \frac{1}{2} \end{bmatrix} \begin{Bmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ 2\varepsilon_{12} \end{Bmatrix}$$

where $\alpha = \frac{3K-2G}{3K+4G}$

The tangent operator here is equivalent to the one mentioned in previous section:

$$\mathbf{D}^e = 2G\mathbf{I}_S + A(K - \frac{2}{3}G)\mathbf{ii}^T$$

where for plane stress condition,

$$A = \frac{2G}{K + \frac{4}{3}G} \quad \mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 \\ & 1 & 0 \\ \text{sym} & & \frac{1}{2} \end{bmatrix} \quad \mathbf{i} = [1, 1, 0]^T$$

Plane stress constraints

For plane stress cases, the plane stress constraints are always considered while deriving formulas.

- Basics

The linear elastic tangent operator and the relationship between ε_{33} and $\varepsilon_{11} + \varepsilon_{22}$. (Refer to **Plane stress linear elasticity**)

- In **matrix** form:

$$\boldsymbol{\sigma} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & 0 \\ \sigma_{21} & \sigma_{22} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\mathbf{s} = \begin{bmatrix} s_{11} & s_{12} & 0 \\ s_{21} & s_{22} & 0 \\ 0 & 0 & s_{33} \end{bmatrix} = \begin{bmatrix} \frac{2}{3}\sigma_{11} - \frac{1}{3}\sigma_{22} & \sigma_{12} & 0 \\ \sigma_{21} & \frac{2}{3}\sigma_{22} - \frac{1}{3}\sigma_{11} & 0 \\ 0 & 0 & -\frac{1}{3}(\sigma_{11} + \sigma_{22}) \end{bmatrix}$$

In **array** form:

$$\boldsymbol{\sigma} = [\sigma_{11}, \sigma_{22}, \sigma_{12}]^T, \quad \mathbf{s} = [s_{11}, s_{22}, s_{12}]^T.$$

And we have $\mathbf{s} = \mathbf{P}\boldsymbol{\sigma}$.

$$\text{where } \mathbf{P} = \frac{1}{3} \begin{bmatrix} 2 & -1 & 0 \\ -1 & 2 & 0 \\ 0 & 0 & 6 \end{bmatrix}.$$

From this equation we can get

$$\begin{cases} s_{11} = \frac{2}{3}\sigma_{11} - \frac{1}{3}\sigma_{22} \\ s_{22} = \frac{2}{3}\sigma_{22} - \frac{1}{3}\sigma_{11} \\ s_{12} = 2\sigma_{12} \end{cases}$$

Note: In matrix form, $s_{12} = \sigma_{12}$. In array form, $s_{12} = 2\sigma_{12}$.

They are different!!

Matrix form is the exact general form for deriving formulas while array form is simply an approach of storing data whose formulas should be derived specifically. The formulas for matrix form can not automatically apply to array form in most cases!!

eg.

$$\begin{aligned} \|\mathbf{s}\| &= \sqrt{\mathbf{s} : \mathbf{s}} \\ &= \sqrt{s_{11}^2 + s_{22}^2 + 2s_{12}^2 + s_{33}^2} \quad (\text{Matrix form}) \\ &= \sqrt{\sigma_{\alpha\beta}\sigma_{\alpha\beta} - \frac{1}{3}(\sigma_{\alpha\alpha})^2} \quad \alpha, \beta \in \{1, 2\} \quad (\text{Matrix form}) \\ &= \sqrt{\frac{2}{3}(\sigma_{11}^2 + \sigma_{22}^2 - \sigma_{11}\sigma_{22}) + 2\sigma_{12}^2} \quad (\text{Matrix form}) \\ &= \sqrt{\boldsymbol{\sigma}^T \mathbf{P} \boldsymbol{\sigma}} \quad (\text{Array form}) \\ &\neq \sqrt{s_{11}^2 + s_{22}^2 + s_{12}^2} \quad (\text{Array form}) \end{aligned}$$

Three methods

Plane stress constraint at the Gauss point level

Plane stress constraint at the structural level

Plane stress-projected plasticity models

This method only applies to sufficiently simple models which allow closed-form relationship between in-plane and out-of-plane variables to be derived from the plane stress constraint (such as the linear elastic model),

$$\sigma_{13} = \sigma_{23} = \sigma_{33} = 0$$

For other complex models, the other two methods are recommended.

Plane stress-projected Von Mises model

The general 3D von Mises model with isotropic strain hardening:

$$\begin{aligned}\dot{\boldsymbol{\epsilon}} &= \dot{\boldsymbol{\epsilon}}^e + \dot{\boldsymbol{\epsilon}}^p \\ \boldsymbol{\sigma} &= \mathbf{D}^e : \boldsymbol{\epsilon}^e \\ \Phi &= \sqrt{3J_2(\mathbf{s})} - \sigma_y(\bar{\epsilon}^p) \\ \dot{\boldsymbol{\epsilon}} &= \dot{\gamma} \frac{\partial \Phi}{\partial \boldsymbol{\sigma}} = \dot{\gamma} \sqrt{\frac{3}{2}} \frac{\mathbf{s}}{\|\mathbf{s}\|} \\ \dot{\bar{\epsilon}}^p &= \dot{\gamma} \\ \dot{\gamma} &\geq 0, \quad \Phi \leq 0, \quad \dot{\gamma} \Phi = 0.\end{aligned}$$

Elastic predictor

Plastic corrector

The derivative of Φ with respect to $\Delta\gamma$ is difficult to obtain.

Orthogonal matrix A (正交矩阵) :

$$AA^T = E$$

The column/row vectors of A are orthogonal to each other and are all unit vectors. And we have $A^T = A^{-1}$.

This matrix can be obtained from the eigenvectors as basis.

Similarity matrix (相似矩阵) :

n 阶方阵 A 相似于 B 定义: 存在可逆矩阵 P , 使得 $P^{-1}AP = B$
 P 成为相似变换矩阵

Diagonal representation (矩阵对角化) :

n 阶方阵 A 可对角化 $\Leftrightarrow A$ 有 n 个线性无关特征向量 \Leftrightarrow

A 相似于对角矩阵 Λ , $P^{-1}AP = \Lambda$

此处相似变换矩阵 P 各列对应 A 的特征向量, Λ 对角线元素为对应特征值,
若 P 为正交矩阵, 有 $P^TAP = \Lambda$

SUVMPS(State update procedur for the von Mises model in plane stress)

Pure input:

- **ntype** : analysis type (1→plane stress;2→plane strain;3→axisymmetric)
- **iprops** : array of integer materials properties (the only one used in SUVM is nhard: number of sampling points for the piecewise hardening curve)
- **rprops** : array of real material properties. For SUVMPS, the needs are
 $[E, \nu, \bar{\epsilon}_0^p, \sigma_{y0}, \bar{\epsilon}_1^p, \sigma_{y1}, \dots, \bar{\epsilon}_{\text{nhard}}^p, \sigma_{y \text{nhard}}]$
- **strat** : $\boldsymbol{\epsilon}_{n+1}^{e \text{ trial}}$, array of the elastic trial strains in engineering form. $[\epsilon_{11}, \epsilon_{22}, 2\epsilon_{12}, \epsilon_{33}]$

Input and updated:

- **rstava** : array of state variables other than stress components $[\varepsilon_{11}^e, \varepsilon_{22}^e, 2\varepsilon_{12}^e, \varepsilon_{33}^e, \bar{\varepsilon}^p]$, 4 elements for elastic strain in engineering form, 1 for equivalent plastic strain. For SUVMPs input, only $\bar{\varepsilon}^p$ of the previous iteration is required. These elements should be updated through this program.

Pure output:

- $\Delta\gamma$: plastic multiplier from return mapping (0 if purely elastic)
- **lalgva**: array of logical algorithmic flags or variables. (For SUVMPs, this contains [ifplas(whether plastic corrector involved), sufail(True: the NR iteration of return mapping fails to converge and the state variables are not updated)])
- **stres** : array of updated stress components $[\sigma_{11}, \sigma_{22}, \sigma_{12}]$

Matrix notation

In the program, all the matrices are in array form.

Array and matrix forms are different. Formulas under matrix form cannot directly apply to array form.

Refer

$$\begin{aligned}\boldsymbol{\sigma} &= [\sigma_{11}, \sigma_{22}, \sigma_{12}]^T, & \mathbf{s} &= [s_{11}, s_{22}, s_{12}]^T \\ \boldsymbol{\varepsilon} &= [\varepsilon_{11}, \varepsilon_{22}, 2\varepsilon_{12}]^T, & \boldsymbol{\varepsilon}^e &= [\varepsilon_{11}^e, \varepsilon_{22}^e, 2\varepsilon_{12}^e]^T \\ \boldsymbol{\varepsilon} &= [\varepsilon_{11}^p, \varepsilon_{22}^p, 2\varepsilon_{12}^p]^T\end{aligned}$$

1. Check type

If $\text{ntype} \neq 1$ (plane stress), return error.

2. Assign and assemble basic parameters

$$\Delta\gamma = 0 \quad G = \frac{E}{2(1+\nu)} \quad K = \frac{E}{3(1-2\nu)}$$

3. Elastic predictor

$$\boldsymbol{\varepsilon}_{n+1}^{e \text{ trial}} := \boldsymbol{\varepsilon}_n^e + \Delta\boldsymbol{\varepsilon}$$

This step is implemented before calling SUVMPs and is stored in form $[\varepsilon_{11}, \varepsilon_{22}, 2\varepsilon_{12}]$.

$$\bar{\varepsilon}_{n+1}^{p \text{ trial}} := \bar{\varepsilon}_n^p$$

Equivalent plastic strain (scalar) is directly obtained from state variables.

To obtain $\boldsymbol{\sigma}_{n+1}^{\text{trial}}$:

- One way:

$$\boldsymbol{\sigma}_{n+1}^{\text{trial}} := \mathbf{D}^e \boldsymbol{\varepsilon}_{n+1}^{e \text{ trial}} \quad (\text{Elastic law})$$

$$\text{where } \mathbf{D}^e \equiv \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} \equiv 2G \begin{bmatrix} 1+\alpha & \alpha & 0 \\ \alpha & 1+\alpha & 0 \\ 0 & 0 & \frac{1}{2} \end{bmatrix}$$

$$\text{and } \alpha = \frac{3K-2G}{3K+4G}.$$

- Another way (implemented in codes):

$$\begin{aligned}\varepsilon_{v n+1}^{e \text{ trial}} &:= \text{tr}(\boldsymbol{\varepsilon}_{n+1}^{e \text{ trial}}) = \varepsilon_{11} + \varepsilon_{22} + \varepsilon_{33} \\ &= \varepsilon_{11} + \varepsilon_{22} + \left(-\frac{3K-2G}{3K+4G}(\varepsilon_{11} + \varepsilon_{22})\right) = \frac{6G}{3K+4G}(\varepsilon_{11} + \varepsilon_{22})\end{aligned}$$

$$\boldsymbol{\epsilon}_{dn+1}^{e \text{ trial}} = \boldsymbol{\epsilon}_{n+1}^{e \text{ trial}} - \frac{1}{3} \epsilon_{vn+1}^{e \text{ trial}} \mathbf{I}$$

[Note: for strain in engineering form $\epsilon_{dn+1 \ 12}^{e \text{ trial}} = \frac{1}{2} \epsilon_{n+1 \ 12}^{e \text{ trial}}$]

$$p_{n+1}^{\text{trial}} := K \epsilon_{vn+1}^{e \text{ trial}} \quad \mathbf{s}_{n+1}^{\text{trial}} := 2G \boldsymbol{\epsilon}_{dn+1}^{e \text{ trial}}$$

Therefore,

$$\boldsymbol{\sigma}_{n+1}^{\text{trial}} = \mathbf{s}_{n+1}^{\text{trial}} + p_{n+1}^{\text{trial}} \mathbf{I}$$

4. Check plastic admissibility

The yield function for von Mises model is

$$\Phi(\boldsymbol{\sigma}, \sigma_y) = \sqrt{3J_2(\mathbf{s}(\boldsymbol{\sigma}))} - \sigma_y$$

$$\text{where } J_2(\mathbf{s}(\boldsymbol{\sigma})) = \frac{1}{2} \mathbf{s}(\boldsymbol{\sigma}) : \mathbf{s}(\boldsymbol{\sigma}) \quad \sigma_y = \sigma_y(\bar{\epsilon}^p)$$

In plane stress, this can be written as:

$$\Phi^{\text{trial}} = \frac{1}{2} \xi^{\text{trial}} - \frac{1}{3} \sigma_y^2(\bar{\epsilon}_{n+1}^{p \text{ trial}})$$

$$\text{where } \xi^{\text{trial}} := \frac{1}{6} a_1 + \frac{1}{2} a_2 + 2a_3$$

$$\text{and } a_1 := (\sigma_{11}^{\text{trial}} + \sigma_{22}^{\text{trial}})^2; a_2 := (\sigma_{22}^{\text{trial}} - \sigma_{11}^{\text{trial}})^2; a_3 := (\sigma_{12}^{\text{trial}})^2.$$

$\sigma_y(\bar{\epsilon}_{n+1}^{p \text{ trial}})$ is obtained from interpolation of the hardening curve based on current trial equivalent plastic strain $\bar{\epsilon}_{n+1}^{p \text{ trial}}$.

$$\text{If } \Phi^{\text{trial}} \leq 0$$

$$\text{In fact } \frac{\Phi^{\text{trial}}}{\sigma_y(\bar{\epsilon}_{n+1}^{p \text{ trial}})} \leq 10^{-8} \text{ is checked in codes.}$$

Then purely elastic:

$$\boldsymbol{\epsilon}_{n+1}^e := \boldsymbol{\epsilon}_{n+1}^{e \text{ trial}}$$

$$\bar{\epsilon}_{n+1}^p := \bar{\epsilon}_{n+1}^{p \text{ trial}} = \bar{\epsilon}_n^p \text{ (not changed)}$$

$$\boldsymbol{\sigma}_{n+1} := \mathbf{s}_{n+1}^{\text{trial}} + p_{n+1}^{\text{trial}} \mathbf{I} = \boldsymbol{\sigma}_{n+1}^{\text{trial}}$$

Goto 7.

Else:

Goto 5.

5. Return mapping

Solve the equation about the yield function with respect to $\Delta\gamma$.

$$\tilde{\Phi}(\Delta\gamma) \equiv \frac{1}{2} \xi(\Delta\gamma) - \frac{1}{3} \sigma_y^2(\bar{\epsilon}_n^p + \Delta\gamma \sqrt{\frac{2}{3} \xi(\Delta\gamma)}) = 0$$

Newton-Raphson for solving $\Delta\gamma$:

1. For initial iteration $k := 0$,
set $\Delta\gamma := 0$ and $\tilde{\Phi} = \Phi^{\text{trial}}, \xi = \xi^{\text{trial}}$
2. Perform N-R iteration

$$H := \frac{d\sigma_y}{d\bar{\epsilon}^p} \big|_{\bar{\epsilon}_n^p + \Delta\gamma} \quad (\text{hardening slope})$$

$$\xi' := -\frac{(\sigma_{11}^{\text{trial}} + \sigma_{22}^{\text{trial}})^2}{9[1 + \frac{E\Delta\gamma}{3(1-\nu)}]^3} \frac{E}{1-\nu} - 2G \frac{(\sigma_{22}^{\text{trial}} - \sigma_{11}^{\text{trial}})^2 + 4(\sigma_{12}^{\text{trial}})^2}{(1 + 2G\Delta\gamma)^3}$$

$$\bar{H} := 2[\sigma_y(\bar{\epsilon}_n^p + \Delta\gamma\sqrt{\frac{2}{3}\xi})]H\sqrt{\frac{2}{3}}(\sqrt{\xi} + \frac{\Delta\gamma\xi'}{2\sqrt{\xi}})$$

$$\tilde{\Phi}' := \frac{1}{2}\xi' - \frac{1}{3}\bar{H} \quad (\text{residual derivative})$$

$$\Delta\gamma := \Delta\gamma - \frac{\tilde{\Phi}}{\tilde{\Phi}'} \quad (\text{new guess for } \Delta\gamma)$$

3. Check for convergence

$$\xi := \frac{(\sigma_{11}^{\text{trial}} + \sigma_{22}^{\text{trial}})^2}{6[1 + \frac{E\Delta\gamma}{3(1-\nu)}]^2} + \frac{\frac{1}{2}(\sigma_{22}^{\text{trial}} - \sigma_{11}^{\text{trial}})^2 + 2(\sigma_{12}^{\text{trial}})^2}{(1 + 2G\Delta\gamma)^2}$$

$$\tilde{\Phi} = \frac{1}{2}\xi - \frac{1}{3}\sigma_y^2(\bar{\epsilon}_n^p + \Delta\gamma\sqrt{\frac{2}{3}\xi})$$

If $|\tilde{\Phi}| \leq \epsilon_{tol}$ (In fact $|\frac{\tilde{\Phi}}{\sigma_y(\bar{\epsilon}_n^p + \Delta\gamma)}| \leq \epsilon_{tol}$ is checked in codes.)

Then Goto 6.

Else Goto local 2.

6. Update state variables and stress

$$\bar{\epsilon}_{n+1}^p := \bar{\epsilon}_n^p + \Delta\gamma\sqrt{\frac{2}{3}\xi(\Delta\gamma)}$$

$$\boldsymbol{\sigma}_{n+1} := \mathbf{A}(\Delta\gamma)\boldsymbol{\sigma}_{n+1}^{\text{trial}}$$

$$\text{where } \mathbf{A}(\Delta\gamma) = \begin{bmatrix} \frac{1}{2}(A_{11}^* + A_{22}^*) & \frac{1}{2}(A_{11}^* - A_{22}^*) & 0 \\ \frac{1}{2}(A_{11}^* - A_{22}^*) & \frac{1}{2}(A_{11}^* + A_{22}^*) & 0 \\ 0 & 0 & A_{33}^* \end{bmatrix}$$

$$\text{with } A_{11}^* = \frac{3(1-\nu)}{3(1-\nu) + E\Delta\gamma}, \quad A_{22}^* = \frac{1}{1 + 2G\Delta\gamma}, \quad A_{33}^* = A_{22}^*$$

For $\boldsymbol{\epsilon}_{n+1}^e$, there are two approaches:

- $\boldsymbol{\epsilon}_{n+1}^e = [\mathbf{D}^e]^{-1}\boldsymbol{\sigma}_{n+1}$

- (Implemented in codes)

$$p_{n+1} = \frac{1}{3}(\sigma_{11} + \sigma_{22}) = K\epsilon_{vn+1}^e \Rightarrow \epsilon_{vn+1}^e$$

$$\mathbf{s}_{n+1} = 2G\boldsymbol{\epsilon}_{dn+1}^e \Rightarrow \boldsymbol{\epsilon}_{dn+1}^e$$

$$\text{where } \begin{cases} s_{n+1\ 11} = \sigma_{11} - \frac{1}{3}(\sigma_{11} + \sigma_{22}) = \frac{2}{3}\sigma_{11} - \frac{1}{3}\sigma_{22} \\ s_{n+1\ 22} = \sigma_{22} - \frac{1}{3}(\sigma_{11} + \sigma_{22}) = \frac{2}{3}\sigma_{22} - \frac{1}{3}\sigma_{11} \\ s_{n+1\ 12} = \sigma_{12} \end{cases}$$

Here \mathbf{s} is in matrix form and the formulas can be applied.

Thus $\boldsymbol{\epsilon}_{n+1}^e$ can be obtained.

$$\begin{cases} \varepsilon_{n+1\ 11}^e = \left(\frac{2}{3}\sigma_{11} - \frac{1}{3}\sigma_{22}\right)\frac{1}{2G} + \frac{1}{3}\varepsilon_{v\ n+1}^e \\ \varepsilon_{n+1\ 22}^e = \left(\frac{2}{3}\sigma_{22} - \frac{1}{3}\sigma_{11}\right)\frac{1}{2G} + \frac{1}{3}\varepsilon_{v\ n+1}^e \\ \varepsilon_{n+1\ 12}^e = \frac{\sigma_{12}}{2G} * 2 \quad (\text{Engineering form}) \\ \varepsilon_{n+1\ 33}^e = -\frac{\nu}{1-\nu}(\varepsilon_{n+1\ 11}^e + \varepsilon_{n+1\ 22}^e) \quad (\text{Constraint}) \end{cases}$$

7. Return other parameters

$\Delta\gamma$, ifplas, sufail

Elastoplastic consistent tangent operator

$$\mathbf{D}^{ep} \equiv \frac{d\boldsymbol{\sigma}_{n+1}}{d\boldsymbol{\varepsilon}_{n+1}} = \frac{d\boldsymbol{\sigma}_{n+1}}{d\boldsymbol{\varepsilon}_{n+1}^{\text{trial}}}$$

CTVMPS(Consistent tangent operator for the von Mises model in plane stress)

Pure input:

- $\Delta\gamma$: plastic multiplier from SUVMPs (0 for purely elastic)
- stres : array of updated stress tensor $\boldsymbol{\sigma}_{n+1} = [\sigma_{11}, \sigma_{22}, \sigma_{12}]$ from SUVMPs
- rstava : array of updated state parameters $[\varepsilon_{11}^e, \varepsilon_{22}^e, 2\varepsilon_{12}^e, \varepsilon_{33}^e, \bar{\varepsilon}^p]$ (4+1) from SUVMPs. Only $\bar{\varepsilon}_{n+1}^p$ is used in CTVMPS.
- iprops : array of integer material parameters (num of points of the hardening curve...)
- rprops : array of real material parameters ($E, \nu, \bar{\varepsilon}_0^p, \sigma_{y0}, \bar{\varepsilon}_1^p, \sigma_{y1} \dots$)
- epflag : whether elastic or elastoplastic (from SUVMPs)
- ntype : analysis type (plane stress(1); plane strain(2); axisymmetric(3))

Pure output:

- \mathbf{D}^e or \mathbf{D}^{ep}

1. Initialization

4th-order symmetric identity tensor:

$$\mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 & 0 \\ & 1 & 0 & 0 \\ \text{sym} & & \frac{1}{2} & 0 \\ & & & 1 \end{bmatrix}$$

2nd-order identity tensor:

$$\mathbf{I} \rightarrow \mathbf{i} = [1, 1, 0, 1]^T$$

2. Check ntype

if ntype \neq 1, return Error.

3. Get and assemble material properties

Shear modulus:

$$G = \frac{E}{2(1 + \nu)}$$

Bulk modulus:

$$K = \frac{E}{3(1 - 2\nu)}$$

4. If epflag = True, goto 5(elsatoplastic); else goto 6(elastic).

5. Compute elastoplastic consistent tangent operator

◦ Compute ξ

$$\begin{aligned}\xi &:= \boldsymbol{\sigma}_{n+1}^T \mathbf{P} \boldsymbol{\sigma}_{n+1} \\ &= [\sigma_{11}, \sigma_{22}, \sigma_{12}] \cdot \frac{1}{3} \begin{bmatrix} 2 & -1 & 0 \\ -1 & 2 & 0 \\ 0 & 0 & 6 \end{bmatrix} \cdot \begin{bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{bmatrix} \\ &= \frac{2}{3}(\sigma_{11}^2 + \sigma_{22}^2 - \sigma_{11}\sigma_{22}) + 2\sigma_{12}^2\end{aligned}$$

◦ Compute H

$$H := \left. \frac{d\sigma_y}{d\bar{\epsilon}^p} \right|_{\bar{\epsilon}_{n+1}^p} \text{ (Hardening slope from the hardening curve)}$$

◦ Compute matrix \mathbf{E}

$$\begin{aligned}\mathbf{E} &:= [\mathbf{C} + \Delta\gamma \mathbf{P}]^{-1} \\ &= \begin{bmatrix} \frac{1}{2}(E_{11}^* + E_{22}^*) & \frac{1}{2}(E_{11}^* - E_{22}^*) & 0 \\ \frac{1}{2}(E_{11}^* - E_{22}^*) & \frac{1}{2}(E_{11}^* + E_{22}^*) & 0 \\ 0 & 0 & E_{33}^* \end{bmatrix}\end{aligned}$$

$$\text{with } E_{11}^* = \frac{3E}{3(1-\nu)+E\Delta\gamma}, E_{22}^* = \frac{2G}{1+2G\Delta\gamma}, E_{33}^* = \frac{E_{22}^*}{2}$$

◦ Compute matrix \mathbf{EP}

$$\begin{aligned}\mathbf{EP} &= \begin{bmatrix} \frac{1}{2}(E_{11}^* + E_{22}^*) & \frac{1}{2}(E_{11}^* - E_{22}^*) & 0 \\ \frac{1}{2}(E_{11}^* - E_{22}^*) & \frac{1}{2}(E_{11}^* + E_{22}^*) & 0 \\ 0 & 0 & E_{33}^* \end{bmatrix} \cdot \frac{1}{3} \begin{bmatrix} 2 & -1 & 0 \\ -1 & 2 & 0 \\ 0 & 0 & 6 \end{bmatrix} \\ &= \begin{bmatrix} \frac{1}{6}E_{11}^* + \frac{1}{2}E_{22}^* & \frac{1}{6}E_{11}^* - \frac{1}{2}E_{22}^* & 0 \\ \frac{1}{6}E_{11}^* - \frac{1}{2}E_{22}^* & \frac{1}{6}E_{11}^* + \frac{1}{2}E_{22}^* & 0 \\ 0 & 0 & E_{22}^* \end{bmatrix}\end{aligned}$$

◦ Compute vector \mathbf{n}

$$\begin{aligned}
\mathbf{n} &:= \mathbf{E} \mathbf{P} \boldsymbol{\sigma}_{n+1} \\
&= \begin{bmatrix} \frac{1}{6} E_{11}^* + \frac{1}{2} E_{22}^* & \frac{1}{6} E_{11}^* - \frac{1}{2} E_{22}^* & 0 \\ \frac{1}{6} E_{11}^* - \frac{1}{2} E_{22}^* & \frac{1}{6} E_{11}^* + \frac{1}{2} E_{22}^* & 0 \\ 0 & 0 & E_{22}^* \end{bmatrix} \cdot \begin{bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{bmatrix} \\
&= \begin{bmatrix} (EP)_{11} \sigma_{11} + (EP)_{12} \sigma_{22} \\ (EP)_{21} \sigma_{11} + (EP)_{22} \sigma_{22} \\ (EP)_{33} \sigma_{12} \end{bmatrix}
\end{aligned}$$

◦ Compute scalar α

$$\begin{aligned}
\alpha &:= \frac{1}{\boldsymbol{\sigma}_{n+1}^T \mathbf{P} \mathbf{n} + \frac{2\xi H}{3-2H\Delta\gamma}} \\
&= \frac{1}{[\sigma_{11}, \sigma_{22}, \sigma_{12}] \cdot \frac{1}{3} \begin{bmatrix} 2 & -1 & 0 \\ -1 & 2 & 0 \\ 0 & 0 & 6 \end{bmatrix} \cdot \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix} + \frac{2\xi H}{3-2H\Delta\gamma}} \\
&= \frac{1}{\left[\frac{\sigma_{11}}{3} (2n_1 - n_2) + \frac{\sigma_{22}}{3} (2n_2 - n_1) + 2\sigma_{12} n_3 \right] + \frac{2\xi H}{3-2H\Delta\gamma}}
\end{aligned}$$

◦ Assemble elastoplastic tangent:

$$\begin{aligned}
\mathbf{D}^{ep} &:= \mathbf{E} - \alpha \mathbf{n} \otimes \mathbf{n} \\
&= \begin{bmatrix} E_{11} & E_{12} & 0 \\ E_{12} & E_{22} & 0 \\ 0 & 0 & E_{33} \end{bmatrix} - \alpha \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix} \begin{bmatrix} n_1 & n_2 & n_3 \end{bmatrix} \\
&= \begin{bmatrix} E_{11} - \alpha n_1^2 & E_{12} - \alpha n_1 n_2 & -\alpha n_1 n_3 \\ E_{12} - \alpha n_1 n_2 & E_{22} - \alpha n_2^2 & -\alpha n_2 n_3 \\ -\alpha n_1 n_3 & -\alpha n_2 n_3 & E_{33} - \alpha n_3^2 \end{bmatrix}
\end{aligned}$$

6. Compute elastic tangent operator

$$\mathbf{D}^e = 2G \mathbf{I}_S + A \left(K - \frac{2}{3} G \right) \mathbf{I} \otimes \mathbf{I}$$

In plane strain, axisymmetric and 3D analyses: $A = 1$

In plane stress analysis: $A = \frac{2G}{K + \frac{4}{3}G}$

$$\text{And also } \mathbf{I}_S = \begin{bmatrix} 1 & 0 & 0 \\ & 1 & 0 \\ \text{sym} & & \frac{1}{2} \end{bmatrix} \quad \mathbf{i} = [1, 1, 0]^T$$

Thus

$$\begin{aligned}\mathbf{D}^e &= 2G\mathbf{I}_S + \left(\frac{2G}{K + \frac{4}{3}G}\right)\left(K - \frac{2}{3}G\right)\mathbf{I} \otimes \mathbf{I} \\ &= 2G\mathbf{I}_S + \left(\frac{2G}{K + \frac{4}{3}G}\right)\left(K - \frac{2}{3}G\right)\mathbf{ii}^T\end{aligned}$$

Abaqus Subroutine

User subroutine interface

```
SUBROUTINE UMAT(STRESS,STATEV,DDSDDE,SSE,SPD,SCD,
  1 RPL,DDSDDT,DRPLDE,DRPLDT,
  2 STRAN,DSTRAN,TIME,DTIME,TEMP,DTEMP,PREDEF,DPRED,CMNAME,
  3 NDI,NSHR,NTENS,NSTATV,PROPS,NPROPS,COORDS,DROT,PNEWDT,
  4 CELENT,DFGRD0,DFGRD1,NOEL,NPT,LAYER,KSPT,JSTEP,KINC)
C
  INCLUDE 'ABA_PARAM.INC'
C
  CHARACTER*80 CMNAME
  DIMENSION STRESS(NTENS),STATEV(NSTATV),
  1 DDSDDE(NTENS,NTENS),DDSDDT(NTENS),DRPLDE(NTENS),
  2 STRAN(NTENS),DSTRAN(NTENS),TIME(2),PREDEF(1),DPRED(1),
  3 PROPS(NPROPS),COORDS(3),DROT(3,3),DFGRD0(3,3),DFGRD1(3,3),
  4 JSTEP(4)
```

Variables to be defined (Inputs and outputs)

In all situations

- DDSDDE(NTENS,NTENS)

Jacobian matrix of the constitutive model to be calculated by UMAT.

$$\mathbf{D} = \frac{\partial \Delta \boldsymbol{\sigma}}{\partial \Delta \boldsymbol{\epsilon}}$$

Output.

- STRESS(NTENS)

Array of the stress tensor.

Passed in and updated.

- STATEV(NSTAVA)

Array containing the solution-dependent state variables.

The size of it is specified in *DEPVAR in Abaqus.

Passed in and updated.

- SSE, SPD, SCD

Specific elastic strain energy(SSE), specific plastic dissipation(SPD), specific creep dissipation(SCD).

Passed in and updated.

Only in a fully coupled thermal-stress or a coupled thermal-electrical-structural analysis

- RPL

Volumetric heat generation per unit time at the end of the increment caused by mechanical working of the material.

Output

- DDSDDT(NTENS)

Variation of the stress increments with respect to temperature

Output

- DRPLDE(NTENS)

Variation of RPL with respect to strain increments.

Output

- DRPLDT

Variation of RPL with respect to temperature

Output

Only in a geostatic stress procedure or a coupled pore fluid diffusion/stress analysis for pore pressure cohesive elements

- RPL

RPL is used to indicate whether or not a cohesive element is open to the tangential flow of pore fluid. Set RPL equal to 0 if there is no tangential flow; otherwise, assign a nonzero value to RPL if an element is open. Once opened, a cohesive element will remain open to the fluid flow.

Variables that can be updated

- PNEWDT

Ratio of suggested new time increment to the time increment being used.

Variables passed in for information

All the variables are passed in if not specifically pointed out.

- STRAN(NTENS)

Array of total strains at the beginning of the increment.

For cases with thermal expansion, this strain is the one with thermal strains subtracted.

- DSTRAN(NTENS)
Array of strain increments.
For cases with thermal expansion, this is the one with thermal strain increment subtracted.
- TIME(1)
Value of time step at the beginning of the current increment.
- TIME(2)
Value of total time at the beginning of the current increment.
- DTIME
Time increment.
- TEMP
Temperature at the start of the increment.
- DTEMP
Increment of temperature.
- PREDEF
Array of interpolated values of predefined field variables based on values read at the nodes at this point at the start of this increment.
- DPRED
Array of increments of predefined field variables.
- CMNAME
User-defined material name. Do not start with "ABQ_".
- NDI
Number of direct stress components at this point.
- NSHR
Number of engineering shear stress components at this point.
- NTENS
Size of the stress or strain component array(NDI + NSHR).
- NSTATV
Number of solution-dependent state variables associated with this material>(*DEPVAR)
- NPROPS
Number of material constants associated with this user material.
- PROPS(NPROPS)
Array of material constants associated with this user material.(Defined in material constants in Abaqus/CAE)
- COORDS
Array of the coordinates of this point. (For cases with geometric nonlinearity, it contains the current coordinates otherwise the original ones.)
- DROT(3,3)
Rotation increment matrix.

This matrix represents the increment of rigid body rotation of the basis system in which the components of stress (STRESS) and strain (STRAN) are stored. It is provided so that vector- or tensor-valued state variables can be rotated appropriately in this subroutine: stress and strain

components are already rotated by this amount **before** UMAT is called. This matrix is passed in as a unit matrix for small-displacement analysis and for large-displacement analysis if the basis system for the material point rotates with the material (as in a shell element or when a local orientation is used).

- CELENT

Characteristic element length.

This is the typical length of a line across an element for a first-order element; it is half of the same typical length for a second-order element. For beams and trusses it is a characteristic length along the element axis. For membranes and shells it is a characteristic length in the reference surface. For axisymmetric elements it is a characteristic length in the plane only. For cohesive elements it is equal to the constitutive thickness.

- DFGRD0(3,3)

Array containing the deformation gradient at the beginning of the increment.

If a local orientation is defined at the material point, the deformation gradient components are expressed in the local coordinate system defined by the orientation at the beginning of the increment.

- DFGRD1(3,3)

Array containing the deformation gradient at the end of the increment.

If a local orientation is defined at the material point, the deformation gradient components are expressed in the local coordinate system defined by the orientation. This array is set to the identity matrix if nonlinear geometric effects are not included in the step definition associated with this increment.

Output

- NOEL

Element number.

- NPT

Integration point number.

- LAYER

Layer number.(for composite shells and layered solids.)

- KSPT

Section point number within the current layer.

- JSTEP(1)

Step number.

- JSTEP(2)

Procedure type key.

- JSTEP(3)

Flag of geometric nonlinearity.

If NLGEOM = YES: =1, else: =0

- KINC

Increment number.

Fortran

Fortran子程序或函数都是传址调用。

方法说明	颜色名称	颜色
此处实现方法利用 CSDN-markdown 内嵌 html 语言的优势	Hotpink	rgb(240, 248, 255)
借助 table, tr, td 等表格标签的 bgcolor 属性实现背景色设置	AntiqueWhite	rgb(255, 192, 203)

