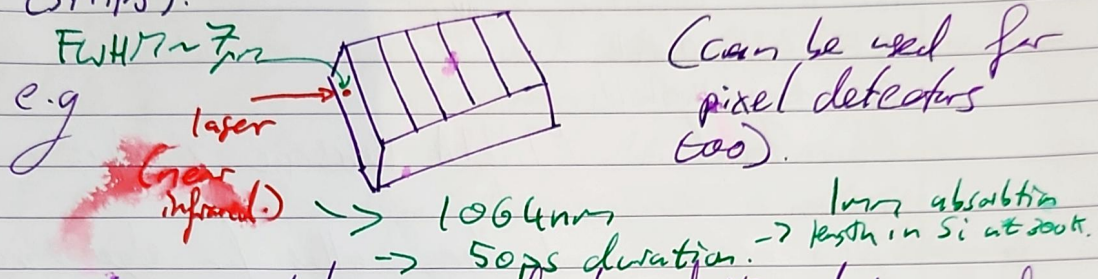


## Edge-TCT measurements.

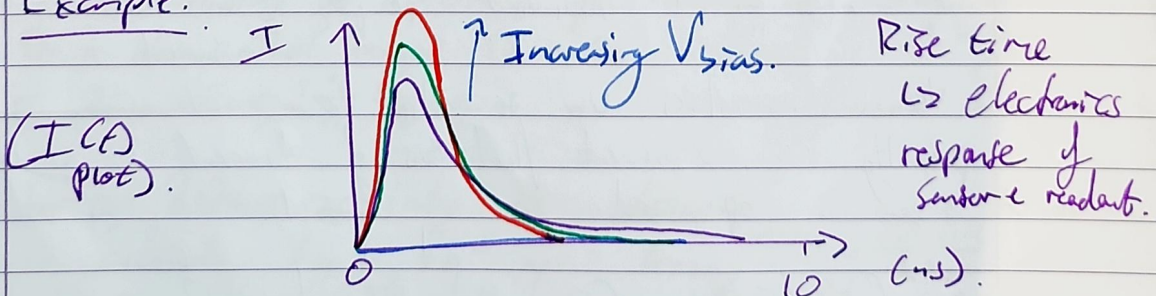
- Edge - Transient Current Technique.
- Sensor irradiated by near infrared laser (focused)
  - ↳ Parallel to sensor surface and normal to silicon (strips).



Current recorded as a function of the distance of the laser beam to the readout plane.  
Initially induced current pulse  $\propto E_{VF}/V_h \rightarrow$  electric field can be determined.

- One strip connected via current amplifier to a 1.5 GHz Scope. (amplifier 6 kHz - 1 GHz)
- Remaining strips connected to a fixed potential via 50  $\Omega$  resistors.
- Average of 400 transients recorded.
- Carefully polished edge for laser injection.

Example.



Peak  $\rightarrow$  dominated by electron current.  
Tail  $\rightarrow$  holes (field decreases towards back of sensor).



# MALTA 2

- 180 nm process.
- 100-300 nm process.
- $\sim 1 \mu W$  / pixel power dissipation.
- $36.4 \times 36.4 \mu m^2$  pixel pitch.
- low  $f_{cut}$  capacitance  $< 5$  aF
- $512 \times 224$  pixels.
- Operating thresholds  $\sim 100 e^-$ .
- Characterised in MALTA telescope  $\rightarrow$  before & after neutron & x-ray irradiation.
- New camera board design allows for edge measurements.
- Polish edge as such there is uniform surface.

Spot size & focus position.

$$f(z, A, \mu, \sigma) = A \left( \operatorname{erf} \left( \frac{z - \mu}{\sigma} \right) + 1 \right)$$

Fitting Parameter

$$\operatorname{erf} \left( \frac{z - \mu}{\sigma} \right) = \frac{2}{\sqrt{\pi}} \int_0^{\frac{z - \mu}{\sigma}} e^{-t^2} dt$$

Width parameter

$$w(z) = w_0 \sqrt{1 + \left( \frac{z - z_0}{z_R} \right)^2}$$

Rayleigh length

$$z_R = \frac{\pi w_0^2}{\lambda} = \frac{1}{2} k w_0^2$$

distance for double the cross section of the beam.

What is a weighting field?

chapter in Speiser.

$$I(t) = Q \cdot v(E(t)) \cdot E_w(r(t))$$

$$\chi^2 = \frac{1}{\sigma^2_{tot}} \sum_{k=1}^{n_k} \left( 1 - \frac{v_{thr}}{v_{scale} \cdot v_k} \right)^2 + w_{pen} \sum_{i=2}^{n_E-1} \left( \frac{0.5(E_{i+1} + E_{i-1}) - E_i}{E_i} \right)^2$$

assumed

$v_{scale}$  normalises  $v_k$  to  $v_{thr}$ .  $\sigma^2 \sim 2\%$ .  $w_{pen}$  is a penalty term.  $\rightarrow$  prevents fluctuations of adjacent  $E_i$  values.  $w_{pen}$  is optimised.

$$v_k = (j_{ph}(E_k) + j_{pe}(E_k)) \int_0^d E(y) dy = V.$$

E-field. bias voltage.

edge TCT setup.

- $\rightarrow$  Junctions Thesis
- $\rightarrow$  reproducible results in LABS.
- $\rightarrow$  Beam spot too large?
- $\rightarrow$  weighting field.



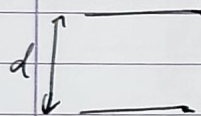
Weighting potential,  $\Phi \rightarrow$  Coupling of a charge at a position to electrode A.

If a charge moves along a path  $s$  from 1  $\rightarrow$  2, induced charge is:  $\Delta Q_A = q(V_{q,1}(z) - V_{q,2}(z)) \equiv q(\Phi_A(1) - \Phi_A(2))$

Instantaneous current expressed in terms of the weighting field

$$i_c = -q \frac{V}{\tau_{\text{velocity}}} \cdot E_{\text{w}} \text{ --- weighting field.}$$

- $E_q$  determined by applying unit potential to the measurement electrode & zero to the others.
  - Electric field & the weighting field are distinct.
    - $\hookrightarrow E$ -field determines the charge trajectory & velocity.
    - $\hookrightarrow$  Weighting field depends only on geometry & determines how charge motion couples to a specific electrode.
- e.g. Parallel Plate geometry.



Very large overbias, can be approximated by a uniform field.  $V_b$  voltage across electrode spacing  $d$ .  $\Rightarrow E = \frac{V_b}{d}$  determines the motion of a charge carrier in the detector.  $v = \mu E = \mu \frac{V_b}{d}$

Weighting field determined by applying unit potential to collection electrode & grounding the other.  $E_q = 1/d$

$\therefore$  induced current  $i = qvE_q = q\mu \frac{V_b}{d} \cdot \frac{1}{d} = q\mu \frac{V_b}{d^2}$

In this geometry,  $E$  &  $E_q$  are uniform across the detector.

$\therefore i$  is constant as a charge flows towards the collector.

For a charge that traverses the entire thickness:  $t_c = \frac{d}{v} = \frac{d^2}{\mu V_b}$

&  $Q = i t_c = q\mu \frac{V_b}{d^2} \frac{d^2}{\mu V_b} = q$

$t_c$  collection time.

For a distance  $x$  into the detector,  $\frac{x}{d}$  a e-h pair

is created.  $t_{ec} = \frac{x}{v_c} = \frac{x d}{\mu V_b}$   $t_{en} = \frac{d-x}{v_n} = \frac{(d-x)d}{\mu V_b}$

$\therefore Q_e = e\mu \frac{V_b}{d^2} \frac{x d}{\mu V_b} = e \frac{x}{d}$   $Q_n = q\mu \frac{V_b}{d^2} \frac{(d-x)d}{\mu V_b} = e(1 - \frac{x}{d})$