

EXAMINATION SOLUTION FORM, Semester 1, 2012-13

Examiner (initials) JMA

MAIN

Module Code PHY2063

Question Number 1.

Page 1 (please number each page)

of

| <ul style="list-style-type: none"> • Use BLACK ink only • Writing must be LEGIBLE • Indicate whether "bookwork" or not | MARKS (to total 20 for each question) |
|--|--|
| <p>a) i) For a thermally isolated system the change in entropy for any process $dS \geq 0$.</p> <p>ii) W — the number of microstates associated with a particular set of constraints (e.g. volume, energy, pressure etc.)</p> <p>b) i) For N objects $(N-m)$ type A m type B</p> $\Rightarrow W = \frac{N!}{m!(N-m)!}$ $\Rightarrow \ln W = N \ln N - N - m \ln m + m - (N-m) \ln (N-m) + N-m$ $= N \ln N - m \ln m - (N-m) \ln (N-m)$ $\therefore S = k_B \ln W$ $= k_B [N \ln N - m \ln m - (N-m) \ln (N-m)]$ | <p style="text-align: center;">3</p> <p style="text-align: center;">3</p> <p style="text-align: center;">2</p> <p style="text-align: center;">3</p> <p style="text-align: center;">1</p> |

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EXAMINATION SOLUTION FORM, Semester 1, 2012-13

Examiner (initials) SMA

Module Code

Question Number 1

(please number each page) Page 2

of

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|---|--|
| <p>ii) In equilibrium $\frac{\partial f}{\partial x} = 0$</p> <p>$\Rightarrow 0 = \epsilon + k_B T (\ln x + 1 - \ln(1-x) - 1)$</p> <p>$\therefore \frac{x}{1-x} = e^{-\epsilon/k_B T}$</p> <p>c) Set $x = 0.03$</p> <p>$\Rightarrow \epsilon = -k_B T \ln \left(\frac{x}{1-x} \right)$</p> <p>$\Rightarrow \epsilon = -k_B T (-3.476) = +1.44 \times 10^{-20} \text{ J}$ $\approx 90 \text{ meV}$</p> <p>\Rightarrow at 100 K</p> <p>$\frac{x}{1-x} = 2.92 \times 10^{-5} \Rightarrow x \approx \frac{2.92 \times 10^{-5}}{1 + 2.92 \times 10^{-5}}$ $\approx 2.9 \times 10^{-5}$ $\approx 3 \times 10^{-3} \%$</p> | <p style="text-align: center;">2</p> <p style="text-align: center;">1</p> <p style="text-align: center;">1</p> <p style="text-align: center;">2</p> <p style="text-align: center;">2</p> |

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EXAMINATION SOLUTION FORM, Semester 1, 2012-13

Examiner (initials)

Module Code

Question Number 2

Page 3 of

(please number each page)

| <ul style="list-style-type: none"> • Use BLACK ink only • Writing must be LEGIBLE • Indicate whether "bookwork" or not | MARKS (to total 20 for each question) |
|---|---|
| <p>a) $Z = \sum_i e^{-\epsilon_i/k_B T}$</p> <p style="margin-left: 100px;">↑ Sum over all states</p> <p>The probability of state with energy ϵ_i is</p> $p_i = \frac{e^{-\epsilon_i/k_B T}}{Z}$ <p>The Helmholtz free energy is given</p> $F = -k_B T \ln Z$ <p>b) i) $Z = \sum_n e^{-E_n/k_B T} = \sum_n e^{-\hbar\omega/k_B T (n + \frac{1}{2})}$</p> $= e^{-\hbar\omega/2k_B T} \sum_{n=0}^{\infty} e^{-\hbar\omega/k_B T n}$ $= \frac{e^{-\hbar\omega\beta/2}}{1 - e^{-\beta\hbar\omega}}$ <p>$\langle E \rangle = - \frac{\partial}{\partial \beta} \ln Z$</p> | <p style="text-align: center;">1</p> <p style="text-align: center;">2</p> <p style="text-align: center;">1</p> <p style="text-align: center;">2</p> <p style="text-align: center;">2</p> <p style="text-align: center;">1</p> |

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EXAMINATION SOLUTION FORM, Semester 1, 2012-13

Examiner (initials)

Module Code

Question Number 2

Page (please number each page) 4

of

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|---|--|
| $\langle E \rangle = - \frac{\partial}{\partial \beta} \left(- \frac{h\nu \beta}{2} - \ln(1 - e^{-\beta h\nu}) \right)$ $= \frac{h\nu}{2} + \frac{1}{1 - e^{-\beta h\nu}} (+ h\nu e^{-\beta h\nu})$ $= \frac{h\nu}{2} + \frac{h\nu}{e^{\beta h\nu} - 1}$ <p>c) i) Each degree of freedom associated with a quadratic energy has $\frac{1}{2} k_B T$ of energy associated with it.</p> <p>ii) As $T = 1000K \gg T_{rot}$ the rotational degree of freedom has $\frac{1}{2} k_B T \times 2$</p> <p>$T \ll T_{vib} \Rightarrow$ hence this degree of freedom is quenched & doesn't contribute to heat capacity.</p> <p>$T \gg T_{tr} \Rightarrow \frac{3}{2} k_B T$ for translational capacity</p> <p>$\Rightarrow \langle E \rangle = \frac{5}{2} k_B T \Rightarrow C_v = \frac{5}{2} k_B$</p> | <p style="text-align: center;">2</p> <p style="text-align: center;">1</p> <p style="text-align: center;">2</p> <p style="text-align: center;">2</p> <p style="text-align: center;">2</p> <p style="text-align: center;">1</p> <p style="text-align: center;">1</p> |

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Question Number 3.

(please number each page) Page 6

of

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|---|--|
| <p> $\therefore p \propto \exp \left\{ \frac{S(u)}{k_B} + \frac{\partial^2 S}{\partial u^2} \frac{\Delta u^2}{2 k_B} \right\}$ </p> <p> Since $\frac{\partial S}{\partial u} = \frac{1}{T} \Rightarrow \frac{\partial^2 S}{\partial u^2} = -\frac{1}{T^2} \frac{\partial T}{\partial u}$ $= -\frac{1}{T^2} \frac{1}{C_V}$ </p> <p> $\Rightarrow p \propto \exp \left\{ -\frac{\Delta u^2}{2 C_V k_B T^2} \right\}$ </p> <p> Comparing with $p(x)$ $\Rightarrow \langle \Delta u^2 \rangle = C_V k_B T^2$ </p> <p> b) i) $\overline{\epsilon} = \sum_i e^{-\frac{(\epsilon_i \mu N_i)}{k_B T}}$ <div style="display: flex; justify-content: space-around; margin-top: 10px;"> <div style="text-align: center;"> \uparrow Sum over all states </div> <div style="text-align: center;"> \nwarrow # particles </div> <div style="text-align: center;"> \nearrow energy of state i </div> </div> </p> | <div style="margin-bottom: 20px;">2</div> <div style="margin-bottom: 20px;">2</div> <div>2</div> |

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Examiner (initials)

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Question Number 3

(please number each page)

Page 7

of

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|--|--|
| <p>ii) Fermions obey the Pauli exclusion principle \Rightarrow only 1 particle can occupy any state particular quantum state.</p> $\boxed{n_i} = \frac{1}{1 + e^{-(E_i - \mu)/k_B T}}$ <div style="display: flex; justify-content: space-around; align-items: center; margin-top: 10px;"> <div style="text-align: center;"> \uparrow $N_i = 0$ $E_i = 0$ </div> <div style="text-align: center;"> \uparrow $N_i = 1$ $E_i = E$ </div> </div> $\langle n \rangle = \frac{1}{\boxed{n_i}} \cdot 0 + \frac{e^{-(E - \mu)/k_B T}}{\boxed{n_i}}$ $= \frac{1}{1 + e^{(E - \mu)/k_B T}}$ <p>iv) $\langle n \rangle$ vs E graph:</p> <p>Low T \Rightarrow only $E < \mu$ occupied. (Quantum \Rightarrow Pauli Exclusion)</p> | <p style="text-align: center; font-size: 2em;">3</p> <p style="text-align: center; font-size: 2em;">3</p> <p style="text-align: center; font-size: 2em;">2</p> |

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