
QUANTUM OPERATIONS

REPORT

Ameya Deshmukh

ABSTRACT

The formalism offered by quantum operations allows easy handling of *open* quantum systems, which are those that interact with the environment. This formalism takes care of a variety of cases: coupled unitary evolution of the environment + system, interactions in which measurements are involved.

1 Basic cases

We denote the transformation of our open quantum system in terms of the density operator representation:

$$\rho' = \mathcal{E}(\rho)$$

For the simple case of unitary evolution, this map is simply: $\rho' = U\rho U^\dagger$.

Now, consider, $\{M_m\}$, a set of measurement operators ($\implies \sum_m M_m^\dagger M_m = I$). If a measurement of our system $\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$ results in the outcome m , the new state will be:

$$\begin{aligned} \rho' &= \sum_i \frac{p(i|m) M_m |\psi_i\rangle \langle \psi_i| M_m^\dagger}{\left(\sqrt{\langle \psi_i | M_m^\dagger M_m | \psi_i \rangle} \right)^2} \\ &= \sum_i \frac{p(m|i) p(i) M_m |\psi_i\rangle \langle \psi_i| M_m^\dagger}{p(m) \langle \psi_i | M_m^\dagger M_m | \psi_i \rangle} \\ &\quad (\text{since: } p(m|i) = \langle \psi_i | M_m^\dagger M_m | \psi_i \rangle) \\ &= \frac{M_m \rho M_m^\dagger}{p(m)} \end{aligned}$$

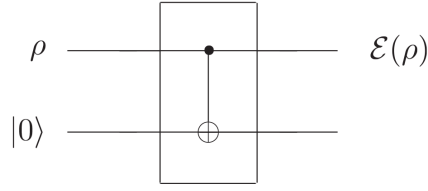
Moreover,

$$\begin{aligned} p(m) &= \sum_i p(m|i) p(i) = \sum_i \langle \psi_i | M_m^\dagger M_m | \psi_i \rangle p_i \\ &= \sum_i \text{tr}(M_m |\psi_i\rangle \langle \psi_i| M_m^\dagger) p_i = \text{tr}(M_m \rho M_m^\dagger) \\ \implies \mathcal{E}(\rho) &= \frac{M_m \rho M_m^\dagger}{\text{tr}(M_m \rho M_m^\dagger)} \end{aligned}$$

Moving ahead, we introduce an environment, along with which the system forms a *closed* joint system. **An assumption** for now: Before the coupled evolution, the joint system is in a product state: $\rho \otimes \rho_{\text{env}}$. Hence:

$$\mathcal{E}(\rho) = \text{tr}_{\text{env}} (U(\rho \otimes \rho_{\text{env}}) U^\dagger)$$

Consider the example in the figure below. In this case, we begin with $\rho \otimes |0\rangle \langle 0|$. The controlled-NOT gate, serves as the unitary $U = |00\rangle \langle 00| + |01\rangle \langle 01| + |10\rangle \langle 11| + |11\rangle \langle 10| = P_0 \otimes I + P_1 \otimes X$,



where P_0, P_1 are the projection operators for the system qubit.

$$\begin{aligned}
 \implies U(\rho \otimes |0\rangle \langle 0|)U^\dagger &= P_0 \rho P_0 \otimes |0\rangle \langle 0| + P_0 \rho P_0 \otimes |1\rangle \langle 1| \\
 \implies \rho' &= \text{tr}_{\text{env}}(P_0 \rho P_0 \otimes |0\rangle \langle 0| + P_0 \rho P_0 \otimes |1\rangle \langle 1|) \\
 &= P_0 \rho P_0 + P_1 \rho P_1
 \end{aligned}$$

2 Operator-sum Representation

The transformations we have seen till now all have a common structure which arises because:

$$\begin{aligned}
 &\text{tr}_{\text{env}}(U(\rho \otimes |e_0\rangle \langle e_0|)U^\dagger) \\
 &= \sum_k (I \otimes \langle e_k|)(U(\rho \otimes |e_0\rangle \langle e_0|)U^\dagger)(I \otimes |e_k\rangle)
 \end{aligned}$$

Writing

$$(I \otimes \langle e_k|)U|\psi e_0\rangle = E_k|\psi\rangle$$

we end up getting:

$$\begin{aligned}
 \mathcal{E}(\rho) &= \sum_i p_i \text{tr}_{\text{env}}(U(|\psi_i e_0\rangle \langle \psi_i e_0|)U^\dagger) \\
 &= \sum_i p_i \sum_k (I \otimes \langle e_k|)U(|\psi_i e_0\rangle \langle \psi_i e_0|)U^\dagger(I \otimes |e_k\rangle) \\
 &= \sum_i p_i \sum_k E_k |\psi_i\rangle \langle \psi_i| E_k^\dagger = \sum_k E_k \sum_i (p_i |\psi_i\rangle \langle \psi_i|) E_k^\dagger \\
 \mathcal{E}(\rho) &= \sum_k E_k \rho E_k^\dagger
 \end{aligned}$$

These operators $\{E_k\}$ which define \mathcal{E} are known as its operation elements. Requiring $\mathcal{E}(\rho)$ to be a valid density operator for all $\rho \implies \text{tr}(\sum_k E_k^\dagger E_k \rho) = 1 \forall \rho$

Hence,

$$\sum_k E_k^\dagger E_k = I$$

Introducing measurements