



浙江大学 海洋学院  
OCEAN COLLEGE  
ZHEJIANG UNIVERSITY

# 声纳与雷达技术（物理篇）

## 3/8

浙江大学海洋学院

2023.11.29



图书馆整理了本课程的7本教参电子资源：《水声学》、《水声传播原理》、《数字式声纳设计原理》、《声呐信号处理引论》、《水下矢量声场理论与应用》、《水声学原理》、《声呐技术》，除Cadal平台需要简单注册后使用外，其他平台均可校网IP直接访问或校外RVPN访问。资源链接如下：

1. <http://book.sciencereading.cn/shop/book/Booksimple/show.do?id=B21515D73A2B14446AADA83450C7DB837000>
2. （1994版） <http://cadal.edu.cn/cardpage/bookCardPage?ssno=06509418>
3. <http://cadal.edu.cn/cardpage/bookCardPage?ssno=33041255>
4. <http://cadal.edu.cn/cardpage/bookCardPage?ssno=58003540>
5. <http://book.sciencereading.cn/shop/book/Booksimple/show.do?id=B1E02CA02B1C841859C72C461E7E34690000>
6. <http://book.sciencereading.cn/shop/book/Booksimple/show.do?id=B979E853CCAAE846FE053010B0A0AF15E000>
7. （2000版） <http://cadal.edu.cn/cardpage/bookCardPage?ssno=32033055>



# 内容提纲

## □ 声学基础 Acoustic Fundamental

- ▶ 波动方程 Wave Equation
- ▶ 简正模态 Normal Mode
- ▶ 射线理论 Ray Theory

## □ 声学工具箱 Acoustic Toolbox

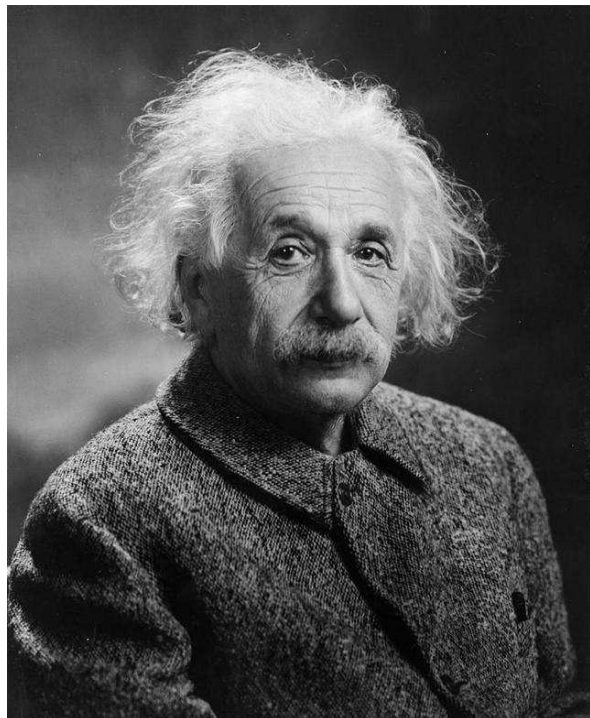
- ▶ 计算声学 Computational Acoustics
- ▶ Kraken Model
- ▶ Bellhop Model

## □ 水声传感器 Underwater Sensor

- ▶ 水听器 Hydrophone
- ▶ 换能器 Transducer
- ▶ 传感器阵列 Sensor Array



爱因斯坦



薛定谔



波尔




$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$$

$$E=mc^2$$

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi$$




# 麦克斯韦方程组

名称	微分形式	积分形式
高斯定律	$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$	$\oiint_S \mathbf{D} \cdot d\mathbf{s} = \frac{Q}{\epsilon_0}$
高斯磁定律	$\nabla \cdot \mathbf{B} = 0$	$\oiint_S \mathbf{B} \cdot d\mathbf{s} = 0$
法拉第感应定律	$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$	$\oint_L \mathbf{E} \cdot d\mathbf{l} = -\frac{d\Phi_B}{dt}$
麦克斯韦-安培定律	$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$	$\oint_L \mathbf{H} \cdot d\mathbf{l} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$

	曲面积分	曲线积分
积分形式	通量	环量
联系	高斯定理	斯托克斯定理
微分形式	散度	旋度

直角坐标系中梯度、散度、旋度和拉普拉斯展开式

- 方程1描述静电场 $\mathbf{E}$ 特征：静电场是有源场，来源于电荷 $\rho$ 。
- 方程2描述静磁场 $\mathbf{B}$ 特征：静磁场是无源场，散度为零 $\nabla \cdot \mathbf{B} = 0$ 。
- 方程3/4描述电场和磁场关联：变化的磁场 $\frac{\partial \mathbf{B}}{\partial t}$ 产生电场，变化的电场 $\frac{\partial \mathbf{E}}{\partial t}$ 产生磁场。
- 电场强度 $\mathbf{E}$ ，磁感应强度 $\mathbf{B}$ ，电流密度 $\mathbf{J}$ ，电位移矢量 $\mathbf{D}$ ，磁场强度 $\mathbf{H}$ ， $\epsilon_0$ 介电常数，磁导率 $\mu_0$ ，电导率 $\sigma_0$ 。
- $\mathbf{D} = \epsilon \mathbf{E}$ ， $\mathbf{B} = \mu \mathbf{H}$ ， $\mathbf{J} = \sigma \mathbf{E}$ ，对于各向同性线性介质。

$$\vec{\nabla} u = \vec{e}_x \frac{\partial u}{\partial x} + \vec{e}_y \frac{\partial u}{\partial y} + \vec{e}_z \frac{\partial u}{\partial z}$$

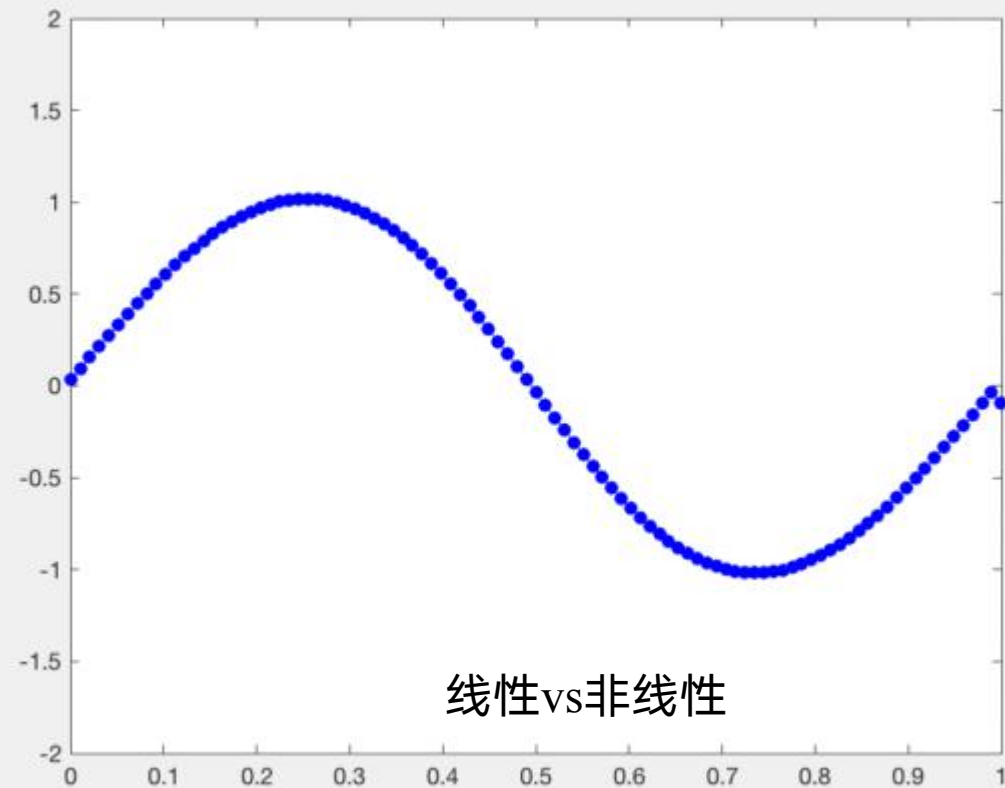
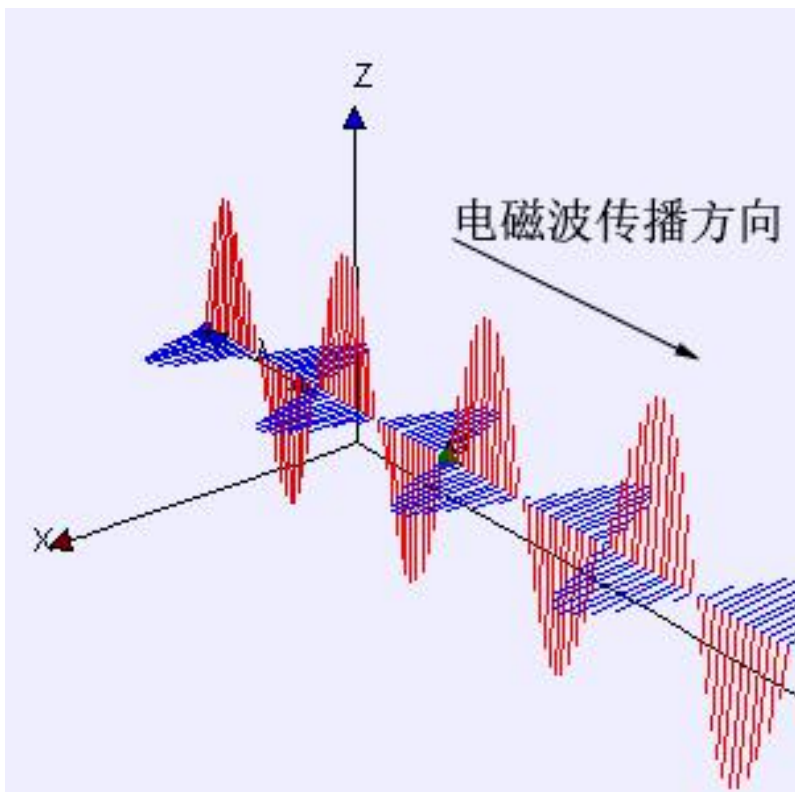
$$\vec{\nabla} \times \vec{A} = \begin{vmatrix} \vec{e}_x & \vec{e}_y & \vec{e}_z \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix}$$

$$\vec{\nabla} \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

$$\nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2}$$

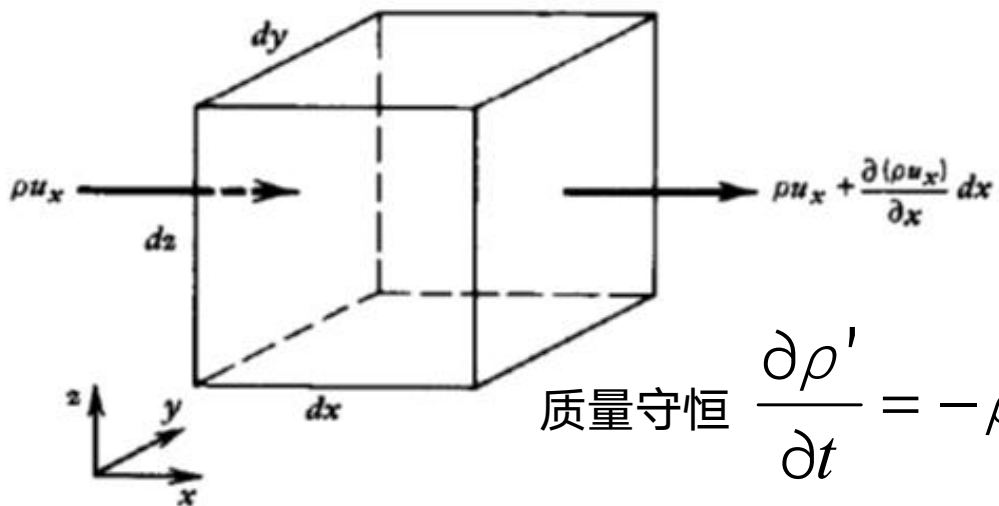
# 1、声学基础

- 波动方程
- 简正模态
- 射线理论



# 1.1、波动方程

$$\tilde{p}(x) - \tilde{p}(x + dx) \approx -\frac{\partial p}{\partial x} dx \quad -\frac{\partial p}{\partial x} = \rho \frac{\partial^2 u}{\partial t^2}$$



质量守恒  $\frac{\partial \rho'}{\partial t} = -\rho_0 \nabla \cdot \mathbf{v}$

动量守恒  $\frac{\partial \mathbf{v}}{\partial t} = -\frac{1}{\rho_0} \nabla p'(\rho)$

状态方程  $p' = \rho' c^2$

$\nabla$  哈密顿算子

$\Delta = \nabla^2$  拉普拉斯算子

$\rho \frac{\partial \mathbf{v}}{\partial t} = \nabla p$  欧拉方程

质点振速无旋  $\mathbf{v} = -\nabla \phi \quad \text{rot} \mathbf{v} = 0$

质点振速有旋  $\mathbf{v} = -\nabla \phi + \nabla \times \boldsymbol{\psi}$

$$\Rightarrow \left\{ \begin{array}{ll} \text{压缩} & \text{惯性力} \\ \rho \nabla \left( \frac{1}{\rho} \nabla p \right) - \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = 0 & \text{Pressure} \\ \frac{1}{\rho} \nabla (\rho c^2 \nabla \cdot \mathbf{v}) - \frac{\partial^2 \mathbf{v}}{\partial t^2} = 0 & \text{Particle Velocity} \\ \nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 0 & \text{Velocity Potential} \\ \nabla^2 \psi - \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2} = 0 & \text{Displacement Potential} \end{array} \right.$$



$$\left[ \nabla^2 + k^2(\mathbf{r}) \right] \psi(\mathbf{r}, \omega) = 0 \quad \text{Helmholtz Equation}$$

$$\mathbf{v} = -\nabla \phi, p = \rho \frac{\partial \phi}{\partial t}$$

$$\nabla^2 \Phi = \frac{1}{C^2} \frac{\partial^2 \Phi}{\partial t^2}$$

Wave equation

$$\Phi = \phi e^{-i\omega t}$$

Harmonic solution

$$\nabla^2 \phi + k^2 \phi = 0$$

Helmholtz equation 双曲型

$$\phi = F(x, y, z) e^{iG(x, y, z)}$$

Range-dependent  
(3D)

- Ray theory

$F \rightarrow$  Amplitude function

$G \rightarrow$  Phase function

Transport Eq 迁移方程

Eikonal Eq 程函方程

$$\phi = F(z) G(r)$$

Range-independent

- Normal mode  
– Extended to range dependence (2D/3D)
- Multipath expansion
- Fast field

$F \rightarrow$  Normal-mode equation

$\rightarrow$  Green's function

$G \rightarrow$  Bessel equation

$\rightarrow$  Hankel function

$$\phi = F(r, \theta, z) G(r)$$

Range-dependent  
(3D)

- Parabolic equation  
抛物型

$F \rightarrow$  Parabolic equation

$G \rightarrow$  Bessel equation

$\rightarrow$  Hankel function



# 海洋声传播模型

- 海洋声学环境通常非常复杂，具有距离和深度有关的特性。这种环境一般不适于对声传播作简单的分析预测。即使在距离无关的环境中，也存在许多路径(多路径)，这些路径组合起来形成复杂的干涉图样。例如，收敛区是一个无法用单调几何扩展定律描述的更复杂结构的例子。声学模型在声传播预测中起着重要作用；这些模型的输入是海洋学量，最终转化为与声学有关的参数，如声速、密度和衰减。
- 海洋中的声传播用波动方程进行数学描述，波动方程的系数和边界条件由海洋环境导出。基本上有四种类型的模型(波动方程的计算机解)来描述海洋中的声传播：射线、频谱或快速场程序(FFP)、简正模态(NM)和抛物方程(PE)。射线理论是波动方程的一种渐近高频近似，而后三种模型或多或少是波动方程在各种温和约束下的直接解。高频极限不包括衍射现象。这些模型都能很好地处理海洋声环境的深度变化。同时考虑到环境水平变化(例如，倾斜的海底或空间可变的海洋学)的模型被称为距离有关模型。对于高频率(几千赫或以上)，射线理论最实用。其他三种类型在较低频率(1kHz以下)下更适用和可用。水声模型的层次结构如图1所示。这些模型的输出通常是传播损失，即在单位距离上相对于单位源强度，以分贝表示。传输损耗Transmission loss是传播损失propagation loss的负值，因此是一个正值。
- 传播模型的输出示例如图2所示，表明模型之间的一致性。然而，我们也看到不同模型之间的差异，射线理论预测的阴影区比波动理论模型更尖锐(即图8B中的10-30 km区域)；这是无限频率射线近似的预期结果。

W.A. Kuperman, Acoustics, Deep Ocean, [Encyclopedia of Ocean Sciences \(Third Edition\) Volume 5](#), 2019, Pages 296-307

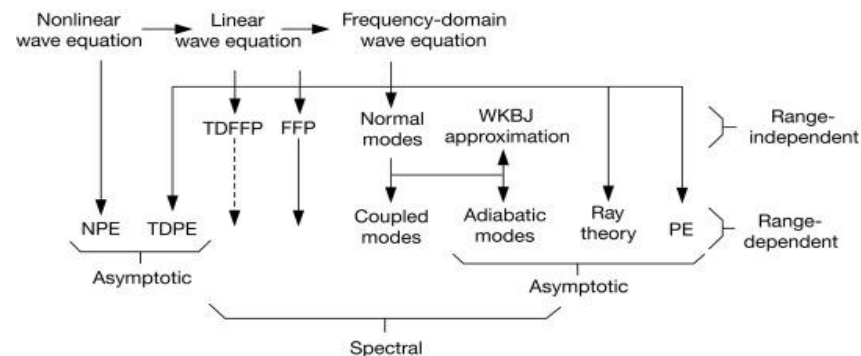


图1 水声模型层次结构。TD表示时域。NPE是描述高振幅(如冲击波)传播的非线性抛物方程。箭头指向模型的派生流程。

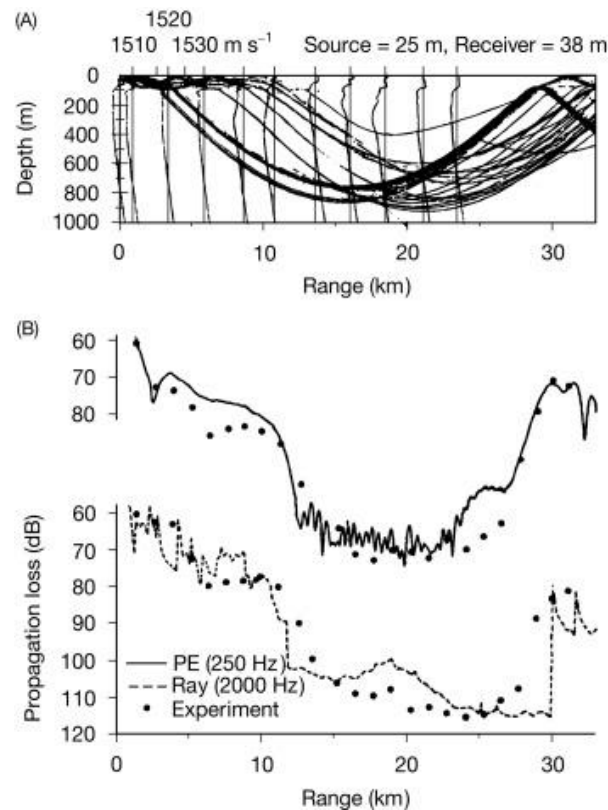


图2 一个距离有关的深水案例的模型和数据比较。(A)作为距离函数的声速剖面，以及显示海面波导传播分解的射线轨迹。(B) 250Hz数据的抛物线方程与2000Hz数据的射线理论比较。



Model type	Applications							
	Shallow water				Deep water			
	Low frequency		High frequency		Low frequency		High frequency	
	RI	RD	RI	RD	RI	RD	RI	RD
Ray theory	○	○	◐	●	◐	◐	●	●
Normal mode	●	◐	●	◐	●	◐	◐	○
Multipath expansion	○	○	◐	◐	◐	◐	●	◐
Fast field	●	◐	●	◐	●	◐	◐	◐
Parabolic equation	◐	●	○	○	◐	●	◐	◐

Low frequency (<500 Hz)

High frequency (>500 Hz)

RI: Range-independent environment

RD: Range-dependent environment

- Modeling approach is both applicable (physically) and practical (computationally)
- ◐ Limitations in accuracy or in speed of execution
- Neither applicable nor practical



# 1.2、简正模态

简正截止  
模态色散  
模态干涉  
模态转换

$$\phi = F(z) \cdot S(r)$$

$$\frac{d^2 F}{dz^2} + (k^2 - \xi^2) F = 0$$

$$\frac{d^2 S}{dr^2} + \frac{1}{r} \frac{dS}{dr} + \xi^2 S = 0$$

$$\phi = \int_{-\infty}^{\infty} G(z, z_0; \xi) \cdot H_0^{(1)}(\xi r) \cdot \xi d\xi$$

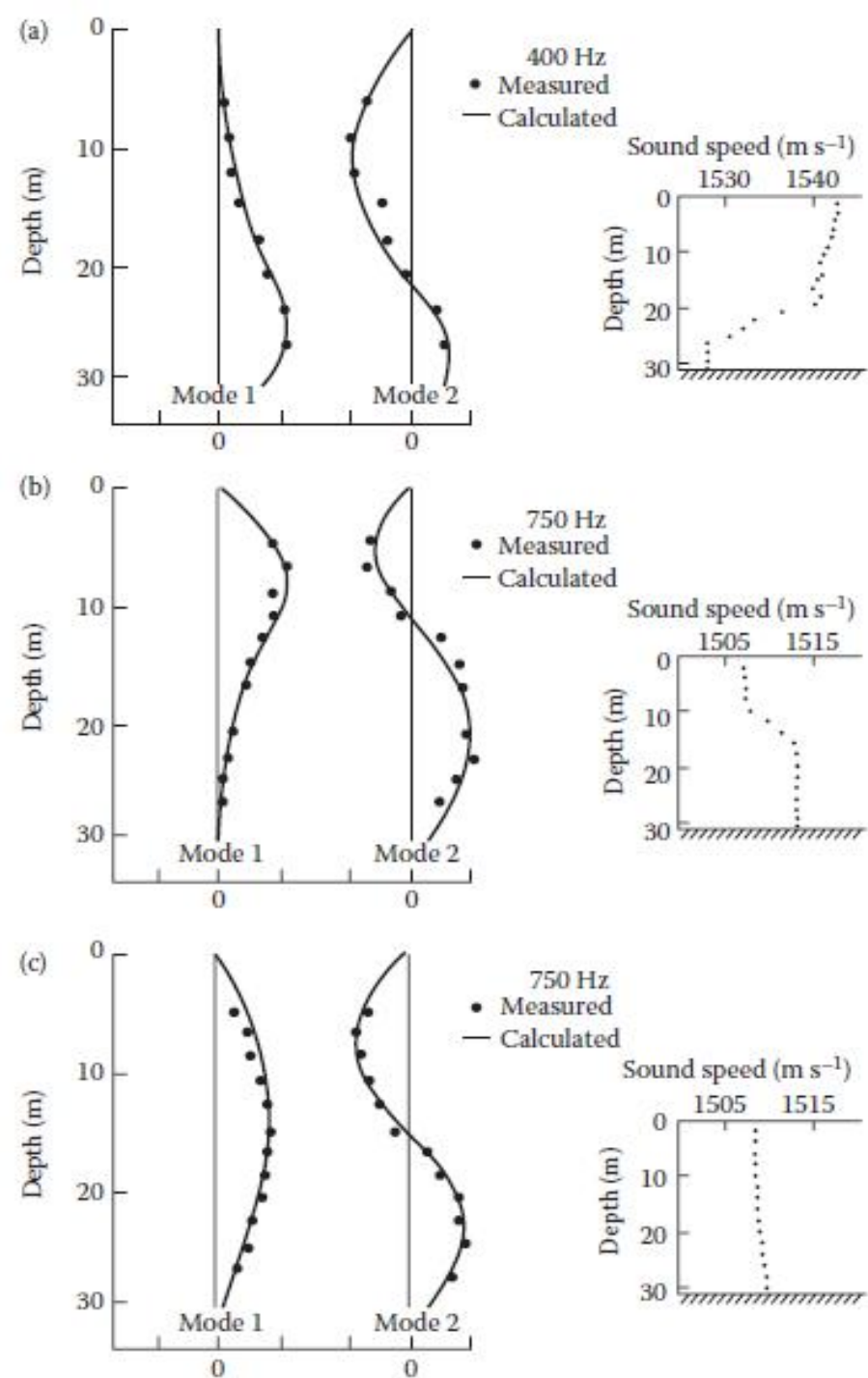
$G$  is Green's function,  $H_0^{(1)}$  is a zero-order Hankel function of the first kind

$$\phi = \oint \sum \frac{u_n(z) \cdot u_n(z_0)}{\xi^2 - \xi_n^2} H_0^{(1)}(\xi r) \cdot \xi d\xi + \text{branch-cut integral}$$

$$H_0^{(1)}(\xi r) \approx \sqrt{\frac{2}{\pi \xi r}} e^{i(\xi r - \pi/4)} \quad \text{for } \xi r \gg 1$$

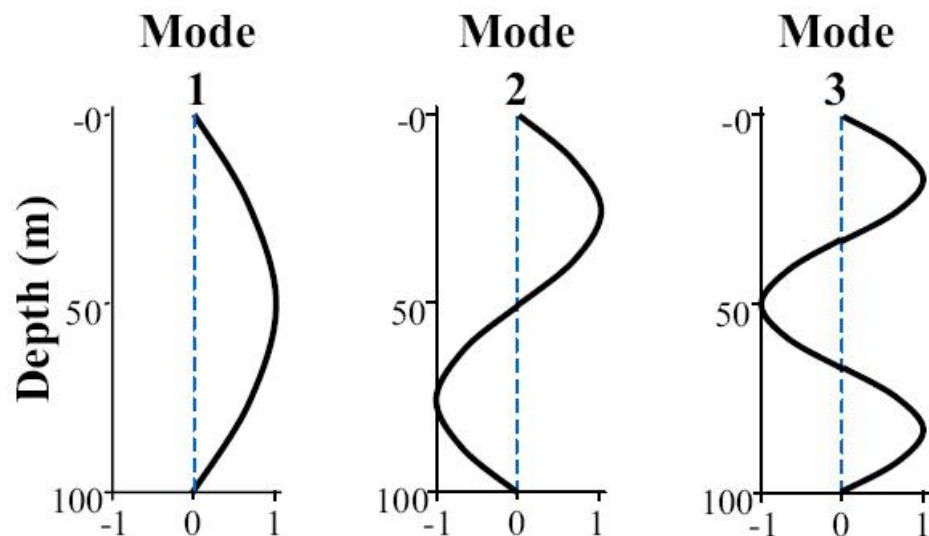
$$\phi = g(r, \rho) \sum \frac{u_n(z) \cdot u_n(z_0)}{\sqrt{\xi_n}} \exp[i(\xi_n r - \pi/4)] \exp(-\delta_n r)$$

$g(r, \rho)$  is a general function of range ( $r$ ) and water density ( $\rho$ ).





# 理想波导Ideal Waveguide



Depth dependence of the first 3 normal modes in ideal waveguide at 20Hz.

First Three Modal Shapes

Modal Expansion

$$\psi(r, z) = -\frac{iS_\omega}{2D} \sum_{m=1}^{\infty} \sin(k_{zm}z) \sin(k_{zm}z_s) H_0^{(1)}(k_{rm}r)$$

$$k_{zm} = \frac{m\pi}{D} \text{ Vertical Wavenumber}$$

$$f_{0m} = \frac{mc}{2D} \text{ Modal Cut-off Frequency}$$

Modal Phase Velocity

$$v_m = \frac{\omega}{k_{rm}}$$

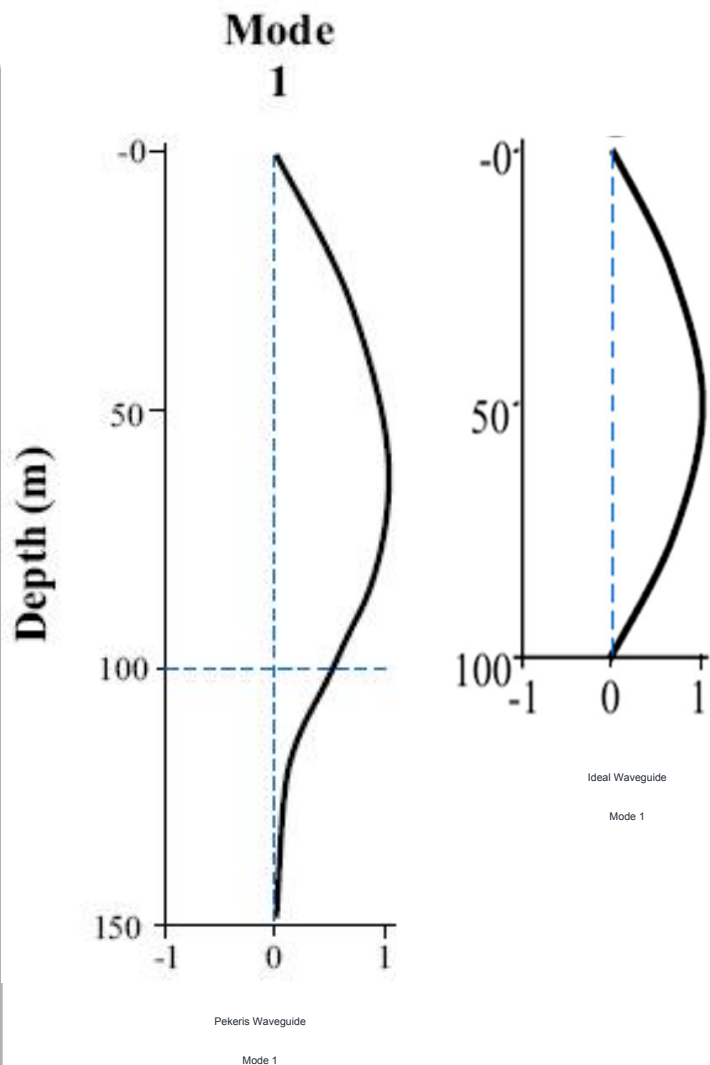
$$k_{rm} = \sqrt{k^2 - k_{zm}^2} = \sqrt{\left(\frac{\omega}{c}\right)^2 - \left(\frac{m\pi}{D}\right)^2} \text{ Horizontal Wavenumber}$$

Modal Group Velocity

$$u_m = \frac{d\omega}{dk_{rm}}$$



# 匹科瑞斯波导Pekeris Waveguide



Modal Expansion

$$\psi(r, z) \approx -\frac{iS_\omega}{2D} \sum_{m=1}^{\infty} a_m(k_{rm}) \sin(k_{zm}z) \sin(k_{zm}z_s) H_0^{(1)}(k_{rm}r)$$

Modal Cut-off Frequency

$$f_{0m} = \frac{(m - 0.5)c_1c_2}{2D\sqrt{c_2^2 - c_1^2}}$$





# 1.3、射线理论

$$\phi = Ae^{iP}$$

程函方程

$$\frac{1}{A} \nabla^2 A - [\nabla P]^2 + k^2 = 0$$

$$[\nabla P]^2 = k^2 \quad \frac{1}{A} \nabla^2 A \ll k^2$$

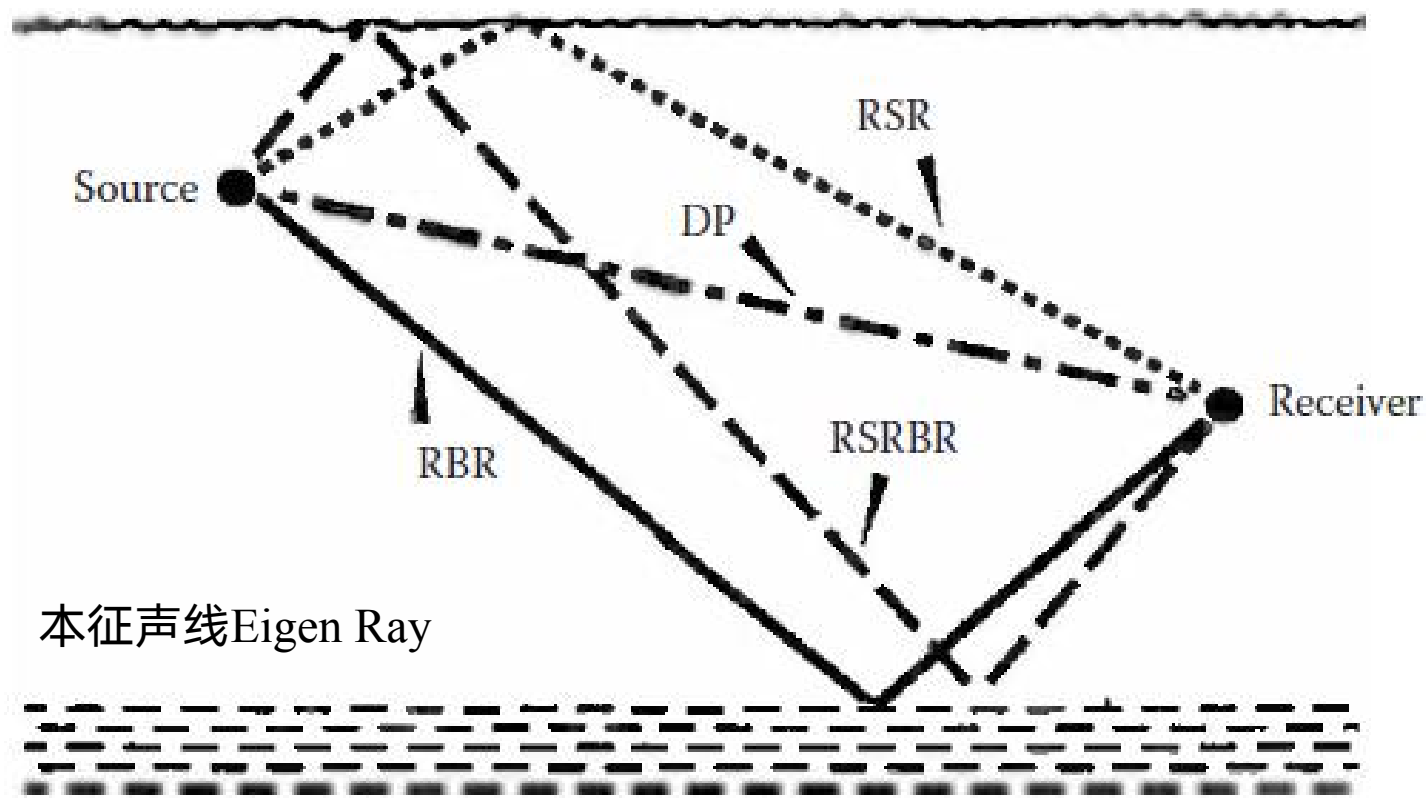
迁移方程

$$2[\nabla A \cdot \nabla P] + A \nabla^2 P = 0$$

$$f > 10 \frac{c}{H} \quad \text{高斯射线}$$

距离有关

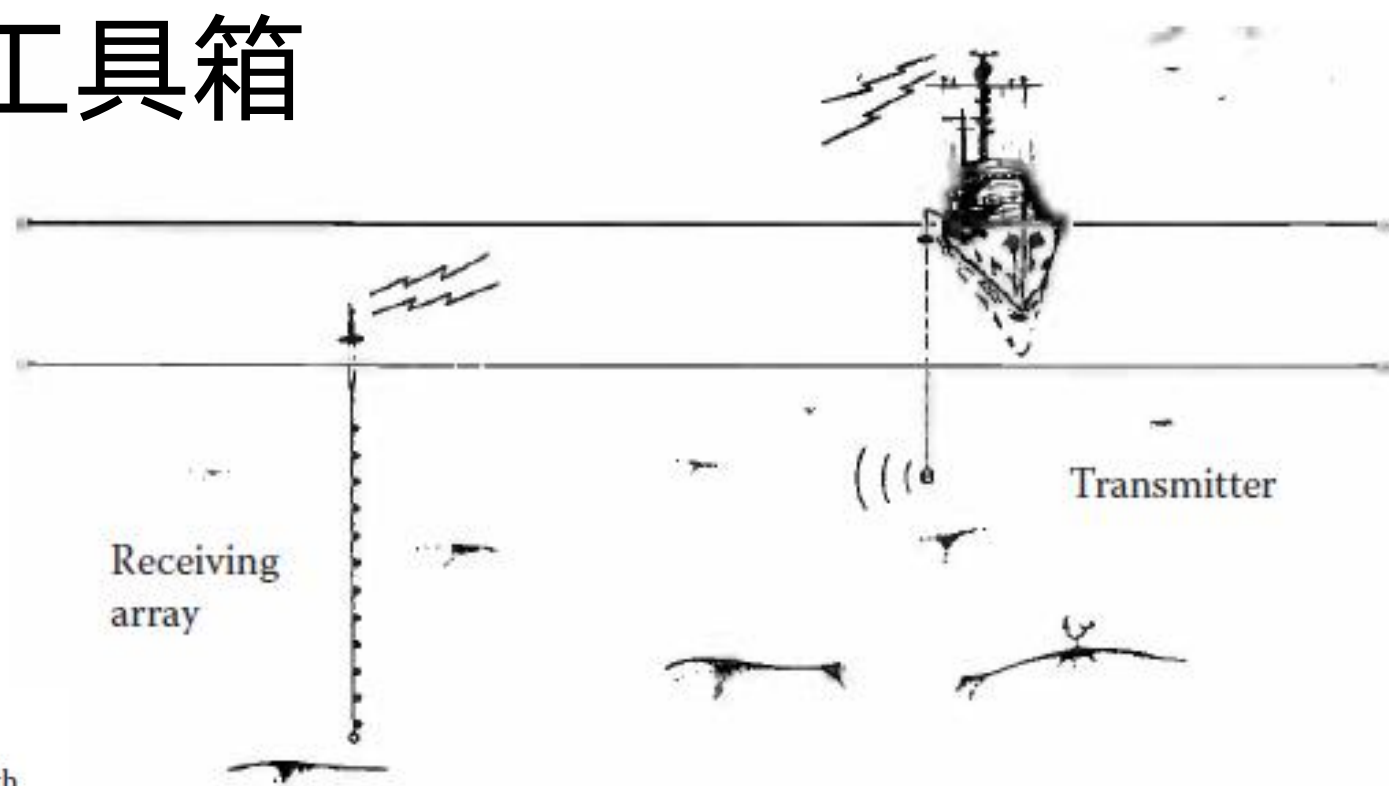
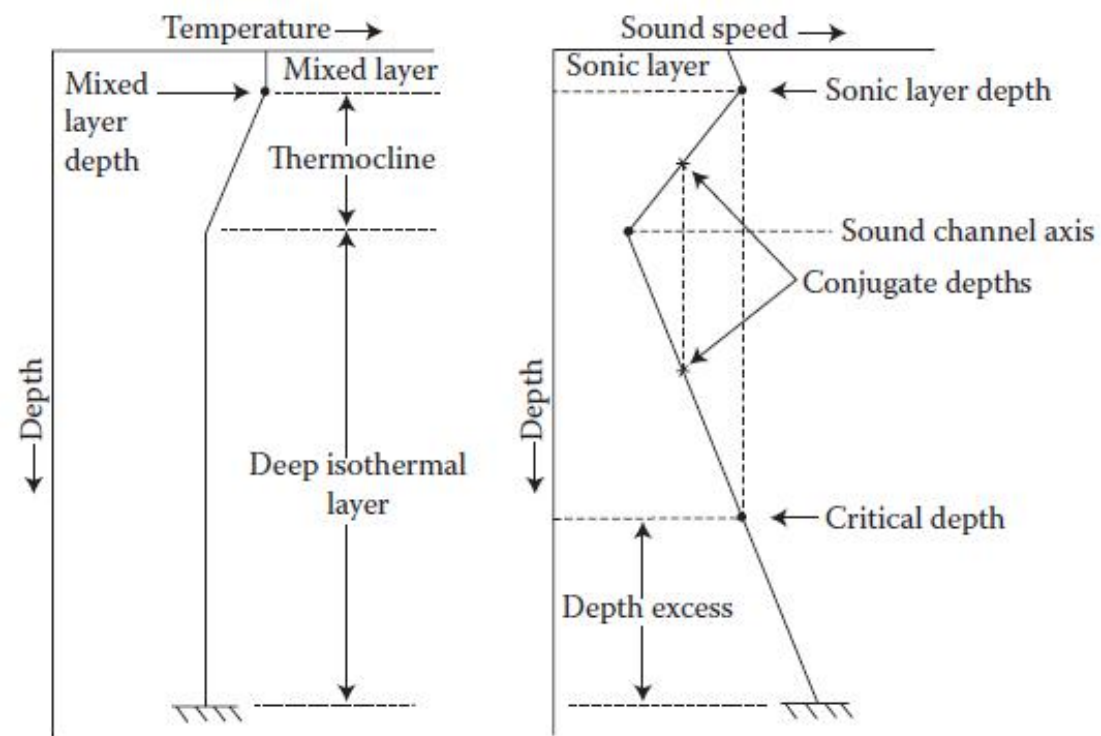
$f$  is the frequency,  $H$  is the duct depth, and  $c$  is the speed of sound



波粒二象性 vs 量子理论  
wave-particle dualism  
quantum theory

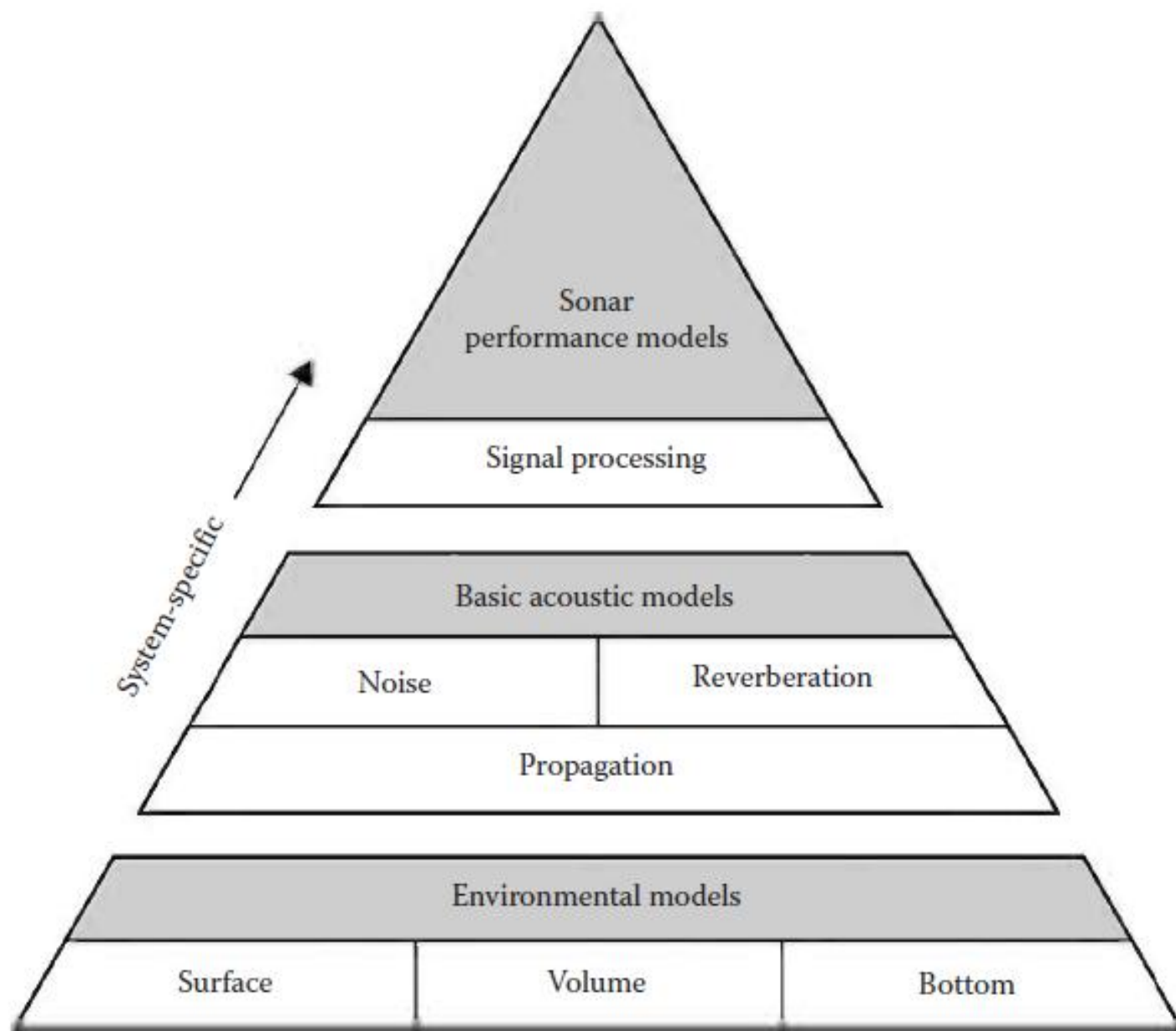
## 2、声学工具箱

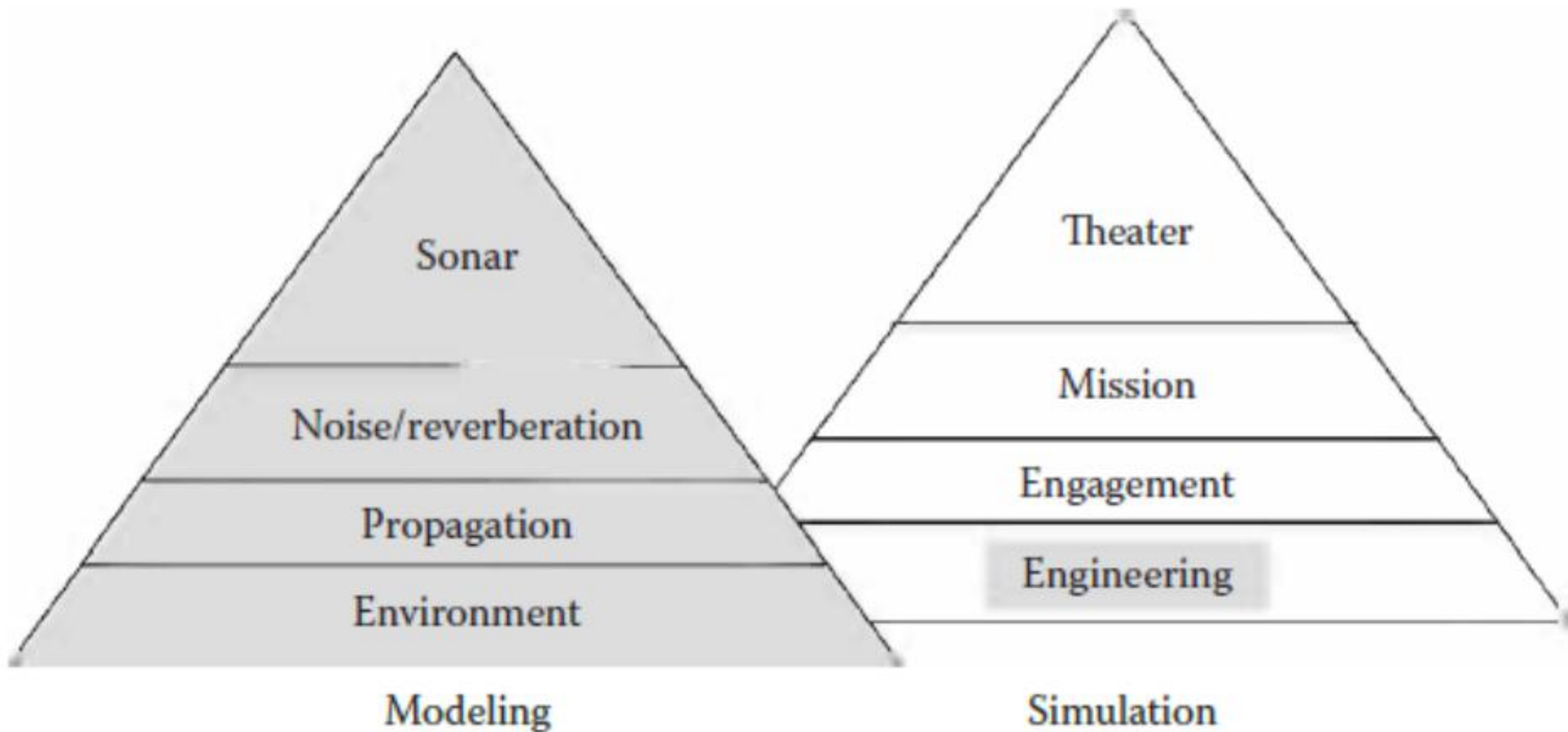
- 声场建模
- Kraken模型
- Bellhop模型

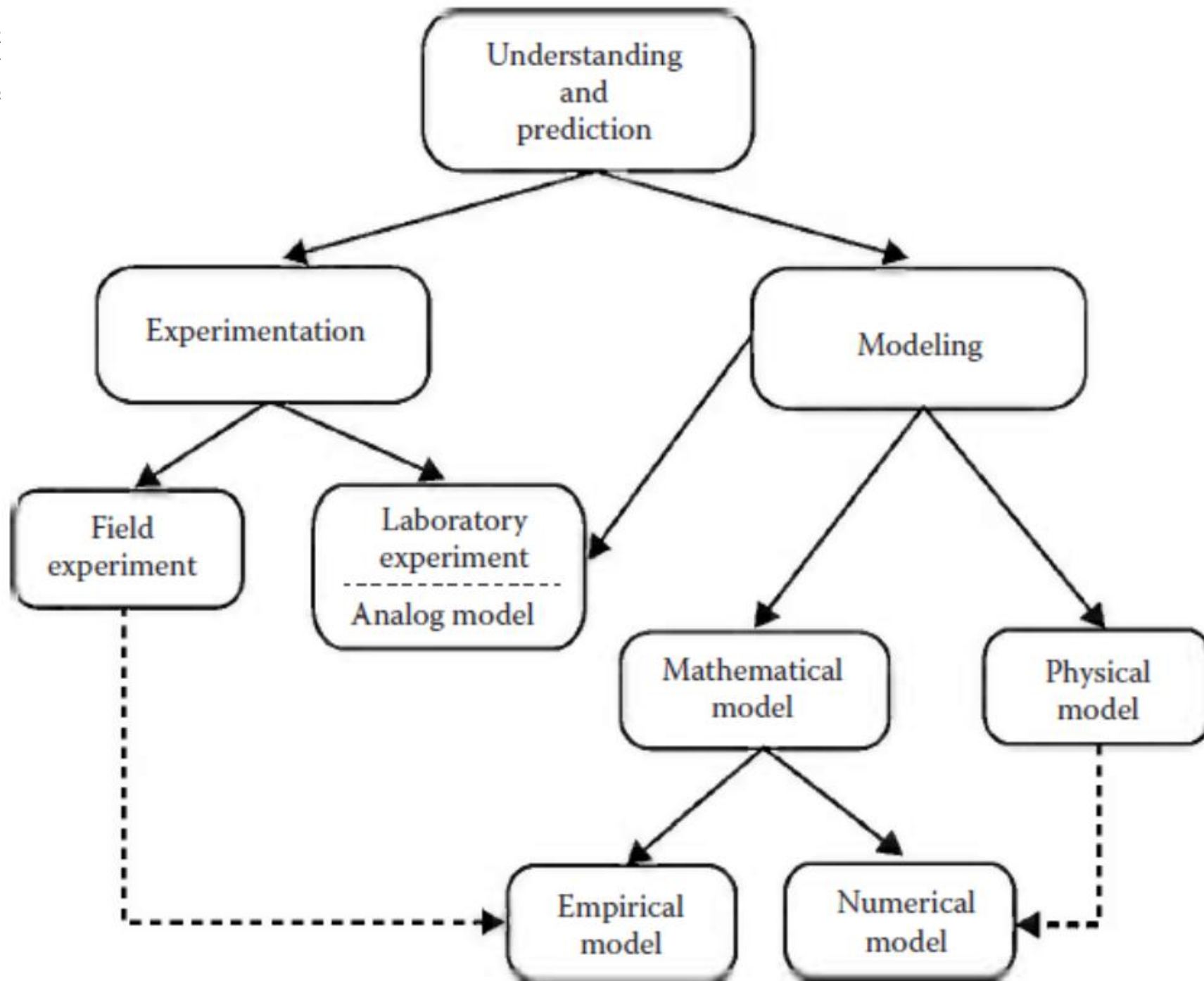




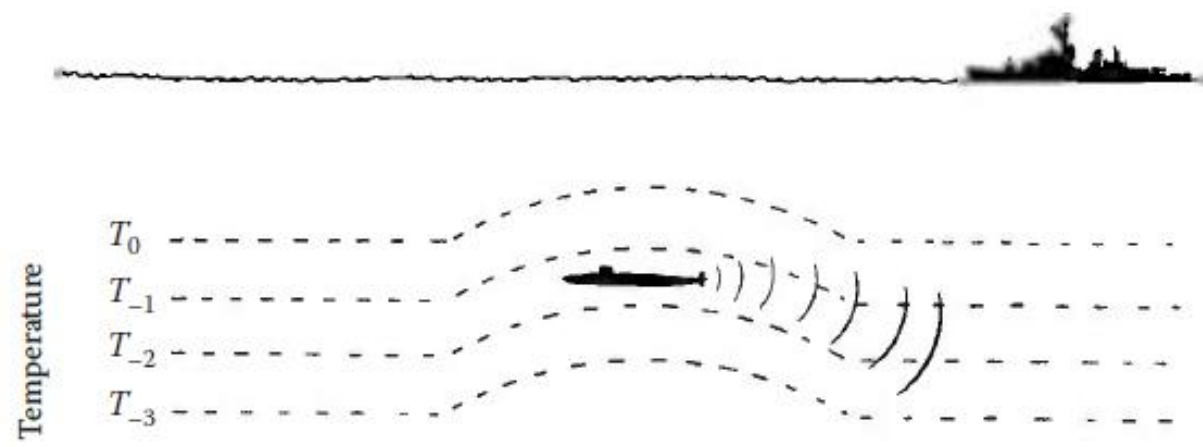
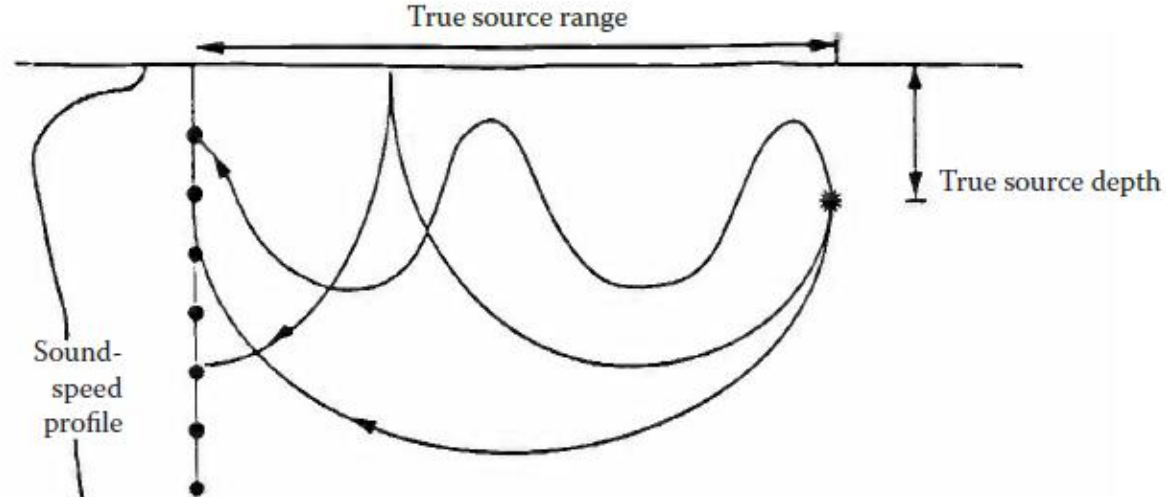
## 2.1、声场建模







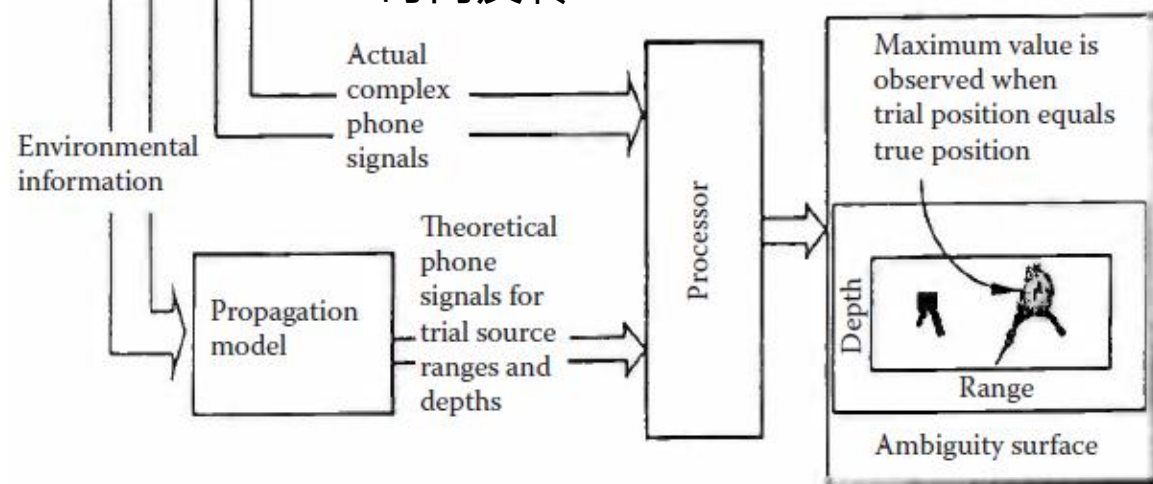




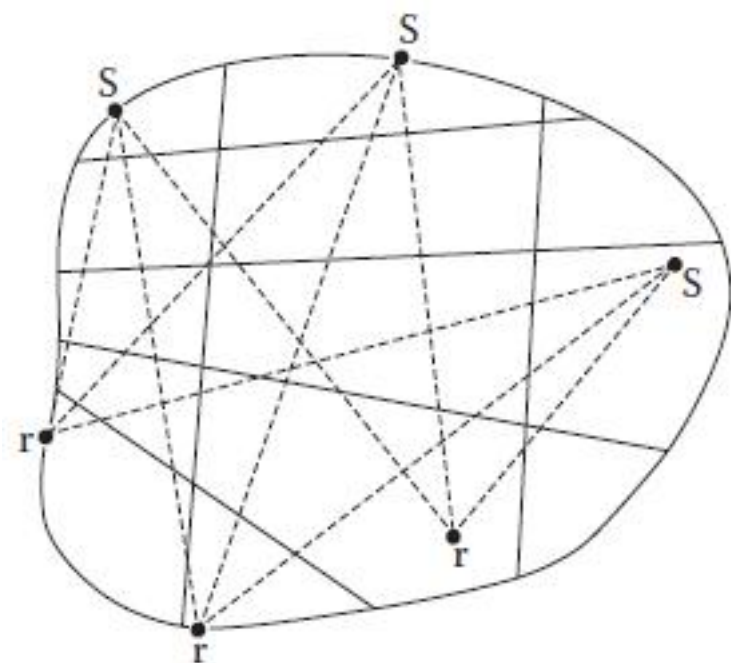
潜艇隐身



模基处理 Model Based Processing  
 匹配场 Matched Field  
 匹配模 Matched Mode  
 时间反转 Time Reversal



海洋参数声学反演

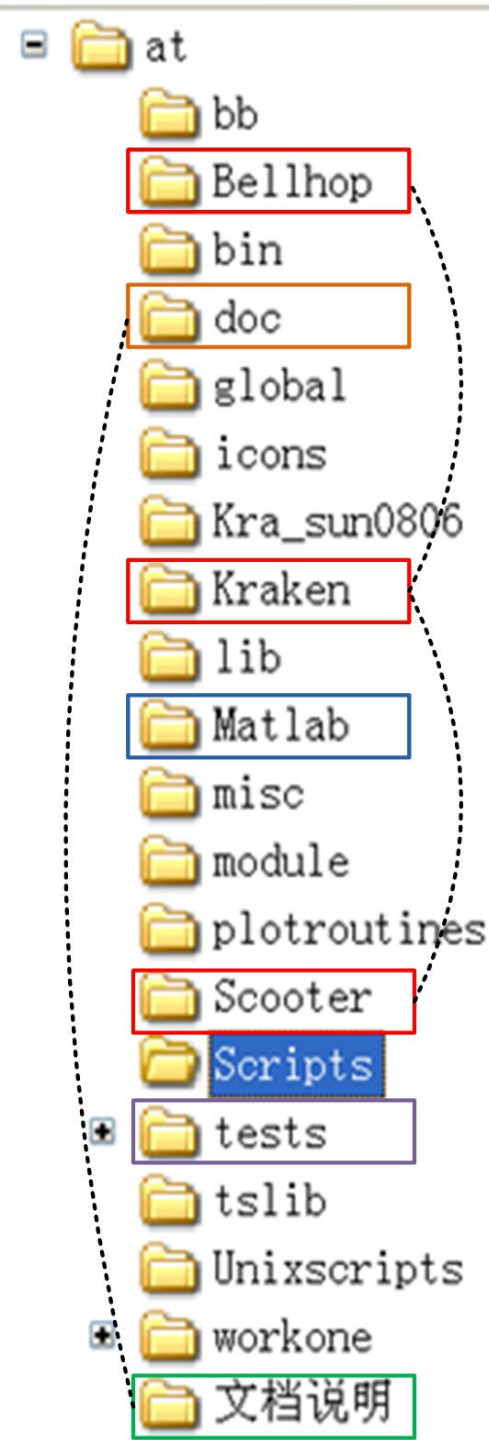


# What's **AT**

AT是声学工具箱AcousticToolbox的英文字头缩写，是由Michael Porter等开发的用于声场建模计算的开源水声场计算软件，主要程序用Fortran90语言编写，后置处理和显示部分则用Matlab、Python编写。AT代表了已被广泛认可的现代海洋计算声学的高水准。

简单实用、精度较高、速度较快。已被广泛用于海洋声学计算。

除AT之外，还有ComSol多物理场有限元软件等。



## 2.2、Kraken简正模型

语法：

- 1) ! c:\at\bin\kraken.exe < EnvFile.env ( 命令窗口调用kraken.exe )
- 2) ! c:\at\bin\kraken EnvFile ( 命令窗口调用kraken.bat )
- 3) eval ( [ ' ! c:\at\scripts\kraken ' EnvFile ] ); ( 用于程序调用kraken.bat )
- 4) 批处理文件 ( bat ) 比执行文件 ( exe ) 多一个计算声场的函数field.bat , 需要输入文件field.flp
- 5) Krakenc用于计算复数本征值问题, 后缀c表示complex。
- 6) 两个重要的输入配置文件: EnvFile.env和field.flp, EnvFile可以任意取名, field.flp则不能更名。

## 2.2.1、Kraken输入配置文件env

Kraken.exe需要\*.env文件计算声场模态\*.mod

'Pekeris problem'

100.0

1

'NVF'

500 0.0 100.0

0.0 1500.0 /

100.0 1500.0 /

'A' 0.0

100.0 2000.0 0.0 2.0 /

1400.0 2000.0 ! CMIN CMAX (m/s)

100.0 ! RMAX (km)

1 ! NSD

50.0 / ! SD(1:NSD) (m)

501 ! NRD

0.0 100.0 / ! RD(1:NRD) (m)

## 2.2.2、Kraken输入配置文件flp

Field.exe需要field.flp文件计算整个声场\*.shd。

```
/,                ! TITLE takes from *.mod file
'RA'              ! OPT 'X/R' (coords), 'C/A' (couple/adiab)
9999,             ! M (number of modes to include)
1                 ! NPROF number of SSP
0.0               ! RPROF(1:NPROF) (km)
1001              ! NR number of range
0.01  1.01/       ! R(1:NR ) (km)
1                 ! NSD number of source depth
10.0 /            ! SD(1:NSD) (m)
1                 ! NRD number of receiver depth
40.0 /            ! RD(1:NRD) (m)
1                 ! NRR = NRD (must be) number of receiver range
0.0 /             ! RR(1:NRR) (m) means perfect vertical array
```



## 2.2.3、Kraken's Example

Kraken.bat需要两个输入文件:

pekeris\_kraken.env和field.flp

主要程序代码:

%==running kraken and producing .mod & .shd files ==

! c:\at\scripts\kraken pekeris\_kraken;% (调用kraken.bat)

%==plot modes ==

plotmode('pekeris\_kraken.mod');

%== reading from the shade file ==

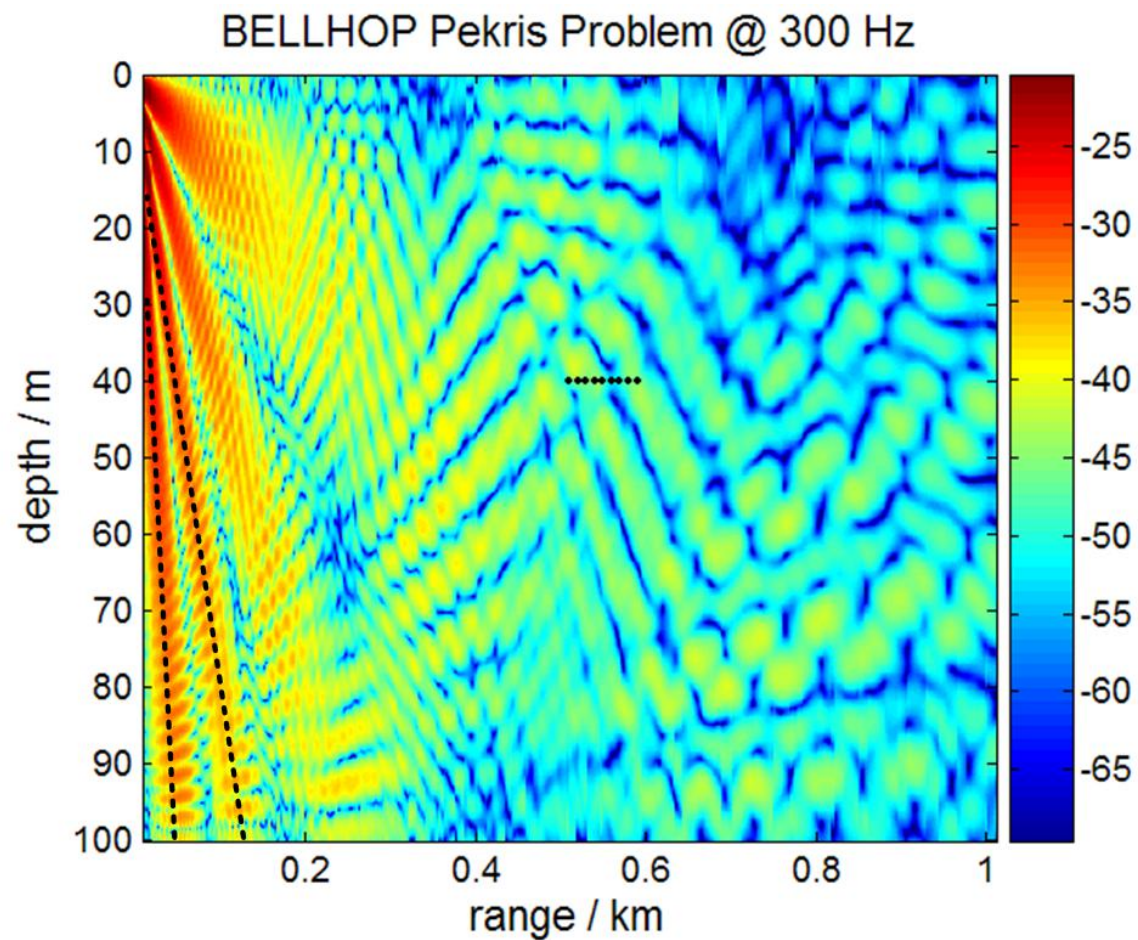
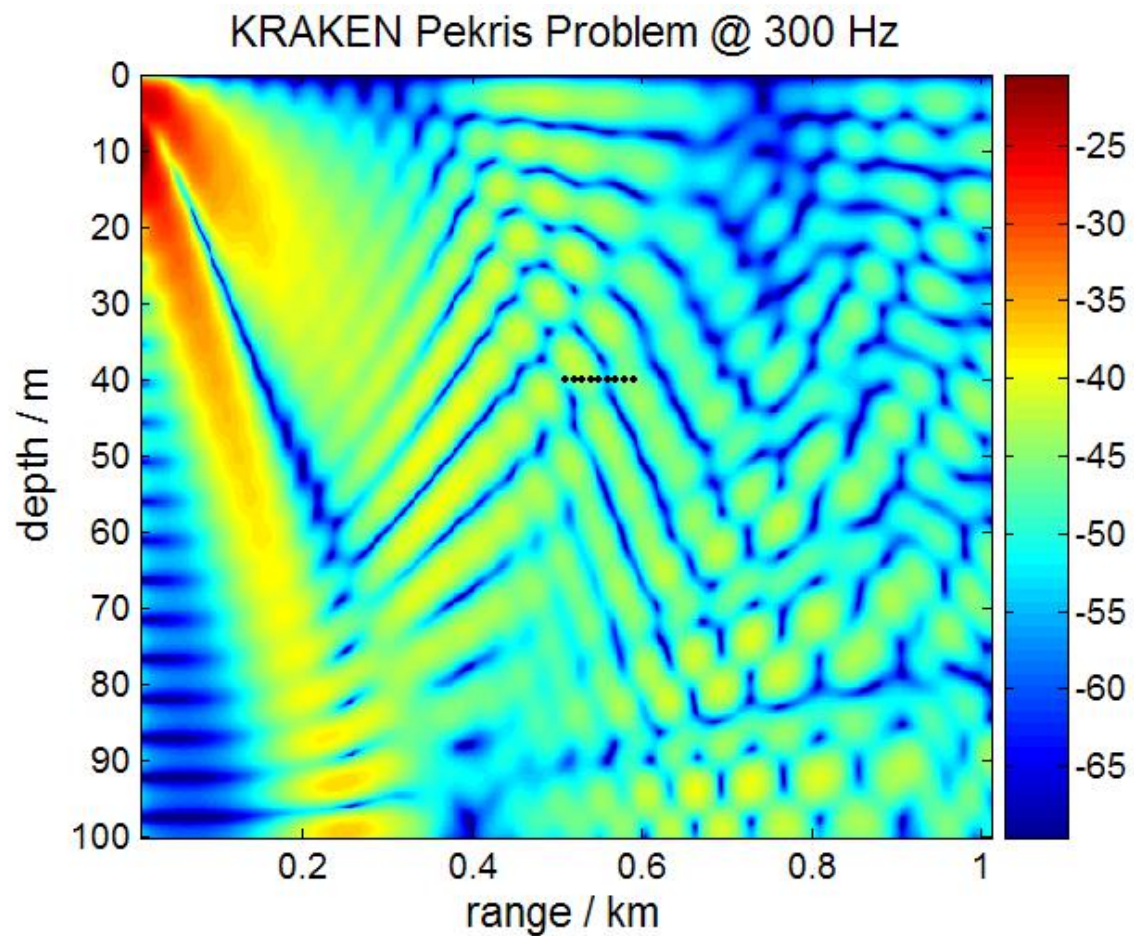
[ pltitl,freq,nsd,nrd,nrr,sd,rd,rr,p1] = read\_shd(('pekeris\_kraken.shd' );

%==transmission loss ==

tl( 1, : ) = 20.0 \* log10( abs( p1( ird, : ) ) ); % ird – receiver depth



# Kraken vs Bellhop



距离 - 深度平面上的声场分布

# 练习：利用AT进行声场建模仿真

发射信号 $x(p, t, f)$ 相当于水声信道的输入信号，若视水声信道为线性时不变系统，其冲激响应函数为 $h(p, t, f)$ ，则接收信号 $y(p, t, f)$ 相当于信道的输出信号。符号 $p$ 代表空间位置， $t$ 代表时间， $f$ 代表频率。



接收声场仿真的原理框图

$$y(p, t, f) = \int_{-\infty}^T h(p, t - \tau, f) x(p, \tau, f) d\tau$$





# 3、水声传感器

- 水听器Hydrophone for Receiving Sound
- 换能器Transducer for Transmitting Sound
- 传感器阵列 Sensor Array

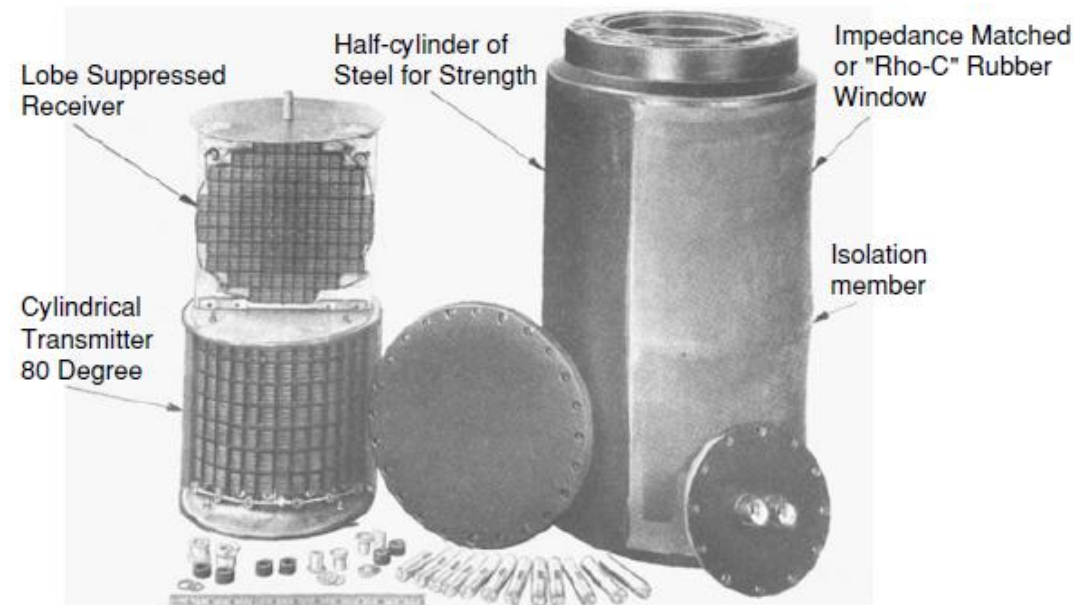


FIGURE 1.7. Early sonar array system with cylindrical transmitter for wide beam pattern and planar array for narrow beam receiving. A baffle between the two minimizes the sound fed from transmitter to receiver [7].

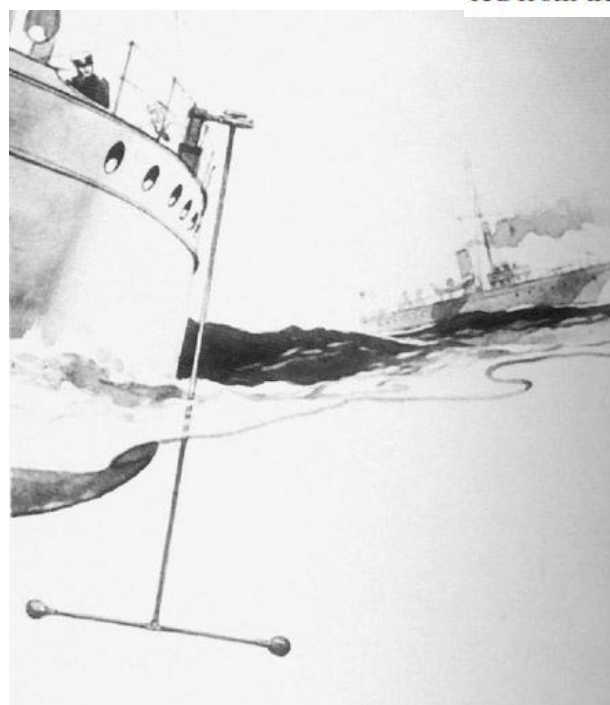
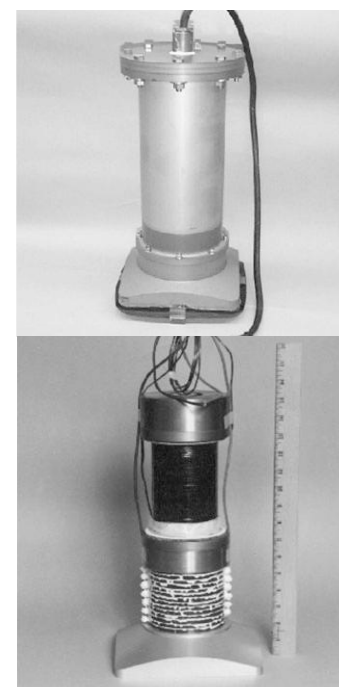


FIGURE 1.15. Submarine sonar spherical array undergoing tests.



# 3.1、水听器

- 声压灵敏度
- 接收指向性

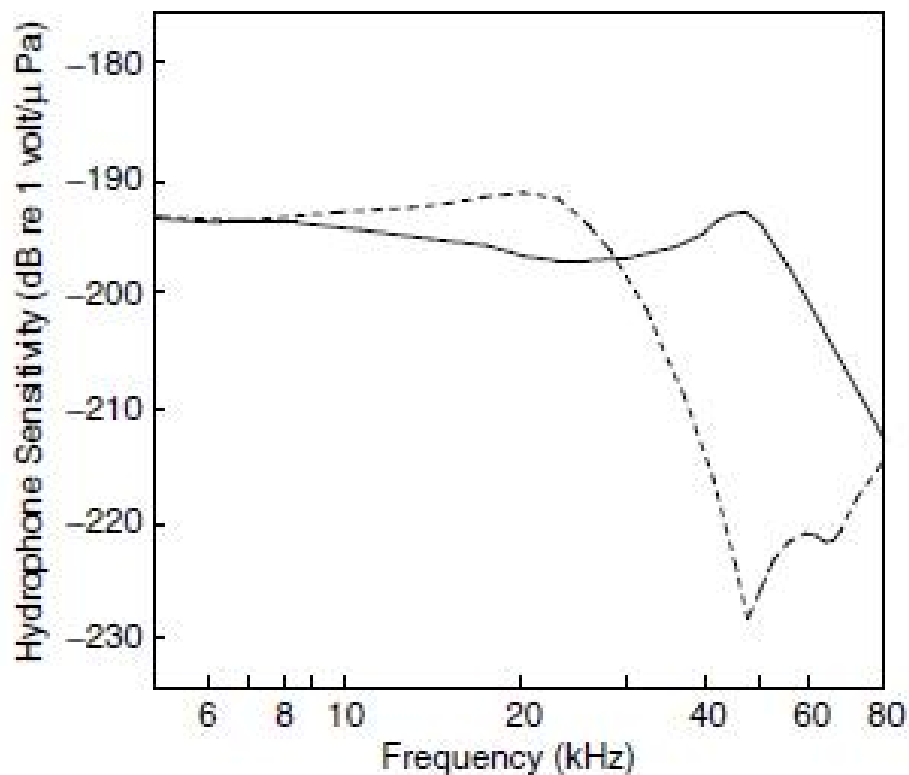


FIGURE 4.11. Comparison of theoretical RVS results for a sphere (—) and a cylinder (---) of diameter and height equal to the diameter of the sphere with radius  $a = 0.0222\text{m}$  [3].

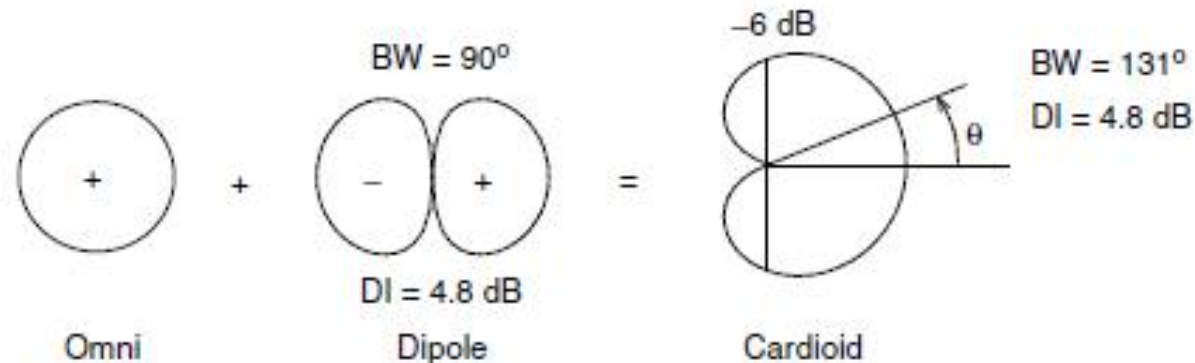


FIGURE 4.30. Synthesis of a cardioid pattern.

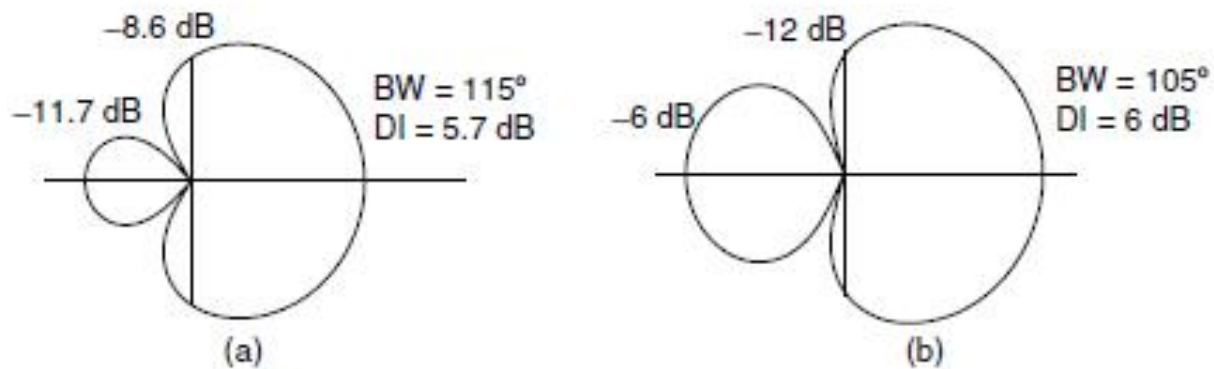


FIGURE 4.31. Super cardioid (a) and hyper cardioid (b) beam patterns.



## 3.2、换能器

- 发射电压响应
- 发射指向性

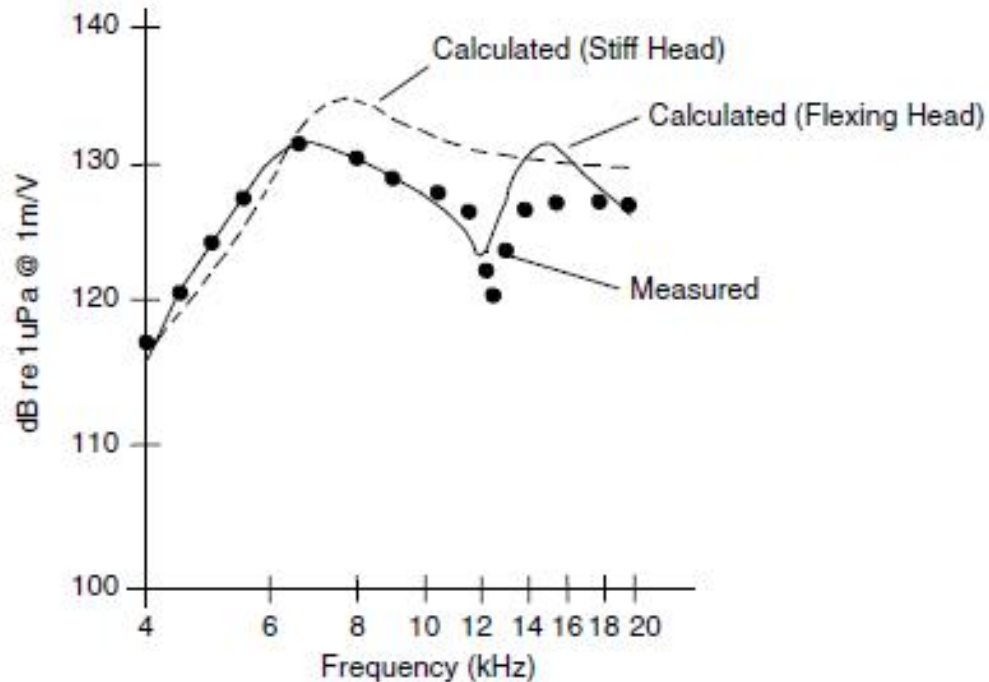


FIGURE 3.21. Comparison of the stiff (---) and flexing (—) head calculated responses with the measured (●●●) constant voltage transmitting response of flexing head Tonpilz transducer [23].

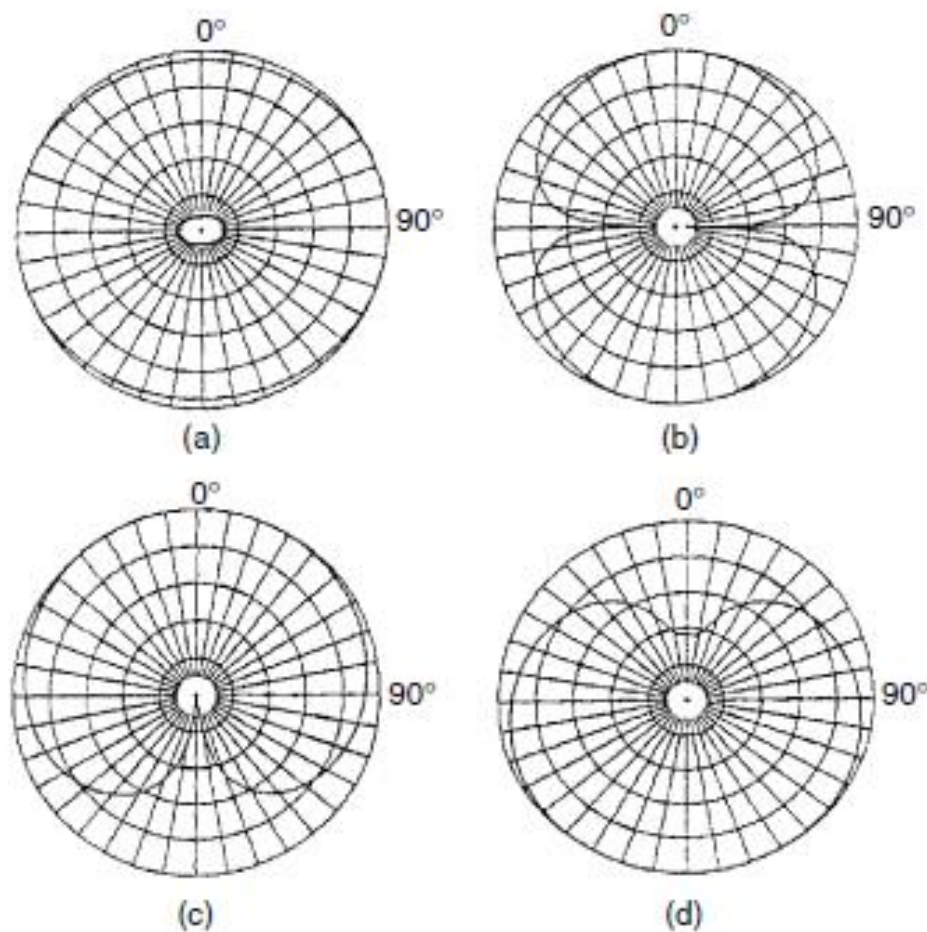


FIGURE 3.42. Measured single element 900 Hz radiation patterns operating in the (a) quadrupole mode, (b) dipole mode, (c) directional mode, and (d) directional mode drive leads reversed (10 dB/division) [56b].

## 3.3、传感器阵列

矩形面阵

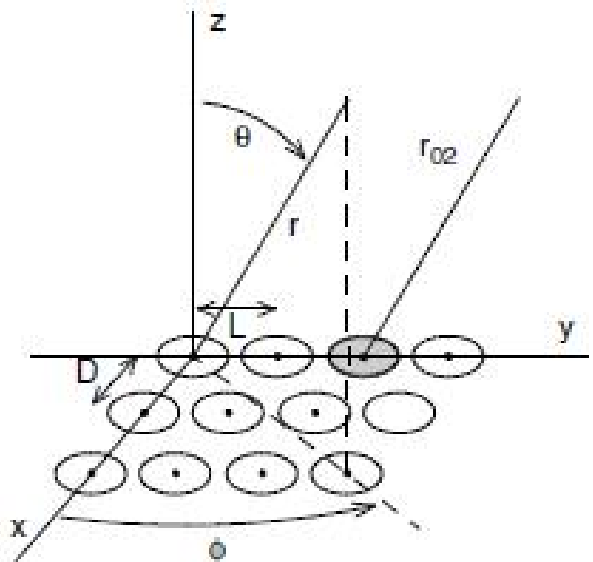


FIGURE 5.4. Coordinates for calculating the far-field of a rectangular array of identical piston transducers. The shaded transducer is located at  $n = 0$ ,  $m = 2$ .

$$p(r, \theta, \phi) = \left[ \frac{jNM\rho ckuAf(\theta, \phi)e^{-jkR}}{2\pi r} \right] \\
 \times \frac{\sin\left(\frac{1}{2}NkD\sin\theta\cos\phi\right)}{N\sin\left(\frac{1}{2}kD\sin\theta\cos\phi\right)} \frac{\sin\left(\frac{1}{2}MkL\sin\theta\sin\phi\right)}{M\sin\left(\frac{1}{2}kL\sin\theta\sin\phi\right)}$$

圆形阵列

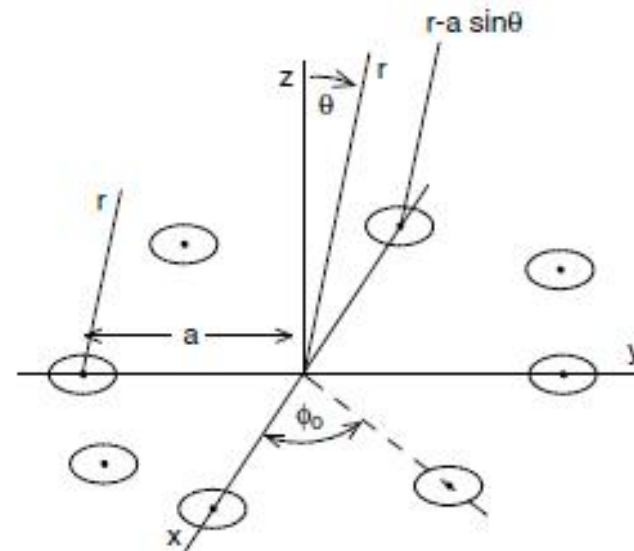


FIGURE 5.5. Coordinates for calculating the far-field in the  $xz$  plane of a circular array of identical transducers. The far-field is the same in the  $xz$ ,  $yz$  and other planes passing through two transducers and the center of the array.

$$p(r, \theta) = \frac{j\rho ckuA}{2\pi} f(\theta) \frac{e^{-jkR}}{r} \left[ 2\cos(ka\sin\theta) + 2 + 4\sum_{n=1}^{\frac{N}{4}-1} \cos(ka\sin\theta\cos n\phi_0) \right]$$

波束形成Beamforming



# 波束图：主瓣、旁瓣、栅瓣

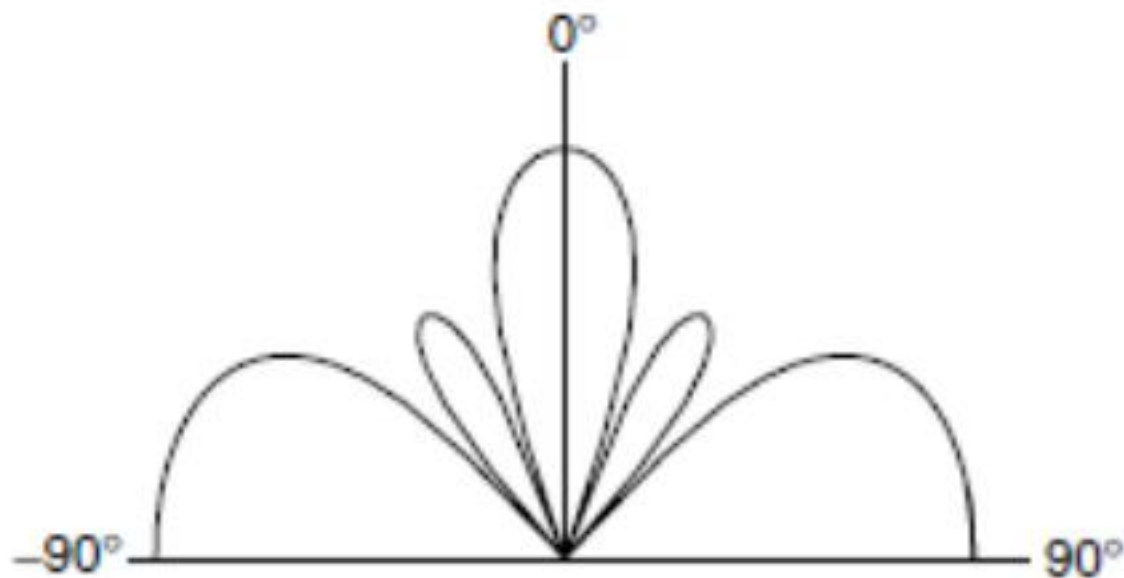


FIGURE 5.6. Example of grating lobes at  $+/- 90^\circ$  for a line array of three transducers with  $kD = 2\pi$ .