## ACSE-2

# Lecture 8 Conservation Equations & Rheology

## Outline

- Conservation equations
- Energy equation
- Rheology
- Elasticity and Wave Equation
- Newtonian Viscosity and Navier Stokes

## Learning Objectives

- Learn main conservation equations used in continuum mechanics modelling and understand what different terms in these equations represent
- Be able to solve conservation equations for basic analytical solutions given boundary/initial conditions.
- Understand basic properties of elastic and viscous rheology and understand how the choice of rheology leads to different forms of the momentum conservation equation
- Using tensor analysis to obtain relations between the main isotropic elastic parameters

## Continuum Mechanics Equations

#### General:

- 1. <u>Kinematics</u> describing deformation and velocity without considering forces
- 2. <u>Dynamics</u> equations that describe force balance, conservation of linear and angular momentum
- 3. <u>Thermodynamics</u> relations temperature, heatflux, stress, entropy

#### **Material-specific**

4. <u>Constitutive equations</u> – relations describing how material properties vary as a function of T,P, stress,.... Such material properties govern dynamics (e.g., *density*), response to stress (*viscosity*, *elastic parameters*), heat transport (*thermal conductivity*, *heat capacity*)

## Conservation equations

- Conservation of mass
  - Kinematics

$$\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0$$

- Conservation of linear momentum
  - Dynamics
  - Newton's second law

$$\rho \mathbf{a} = \nabla \cdot \underline{\underline{\sigma}} + \mathbf{f} = \rho \frac{D\mathbf{v}}{Dt}$$

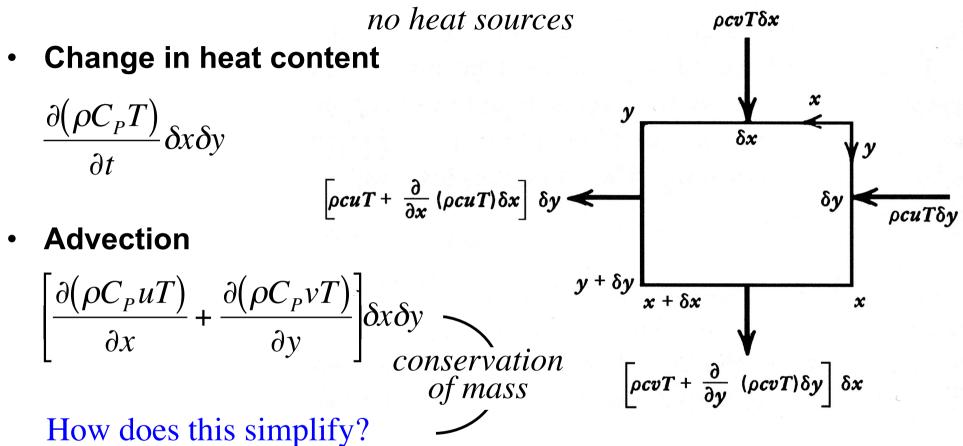
- Conservation of energy

- First law of thermodynamics 
$$\frac{D(K+U)}{Dt} = W + Q$$

K- kinetic energy, U- internal energy, W – power input, Q – heat input

## 2-D energy equation

Spatial, constant  $\rho$ ,  $C_P$ , k, incompressible

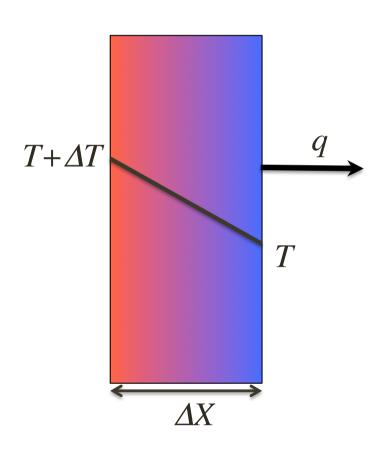


Conduction

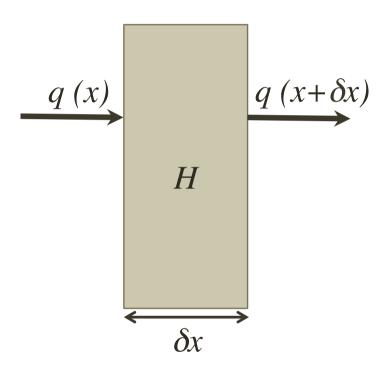
## Fourier's Law for conduction

- *Heat flux*, q, = heat/area = energy/time/area, unit: J/s/m<sup>2</sup> = W/m<sup>2</sup>
- Heat flux proportional to temperature gradient
- Minus sign because heat flows from hot to cold
- Constant of proportionality: *thermal conductivity*, *k*, unit: W/m/K

$$q = -k \frac{dT}{dx}$$



## 1-D Steady State Conduction



$$-k\frac{d^2T}{dx^2} = \rho H = A$$

net heat flow/unit area/unit time =

$$q(x+\delta x)-q(x)$$

$$q(x+\delta x) = q(x) + \delta x \frac{dq}{dx} + \dots$$

$$q(x+\delta x) - q(x) \approx \delta x \frac{dq}{dx}$$

$$q(x + \delta x) - q(x) \approx \delta x \frac{dq}{dx}$$
$$\delta x \frac{dq}{dx} = \delta x \left[ \frac{d}{dx} \left( -k \frac{dT}{dx} \right) \right]$$

$$\delta x \frac{dq}{dx} = \delta x \left[ -k \frac{d^2 T}{dx^2} \right] \qquad for \ constant \ k$$

• heat produced =  $\rho H \delta x = A \delta x$ 

H - heat production rate/unit mass (W/kg)

A – heat production/unit volume  $(W/m^3)$ 

## 2-D energy equation

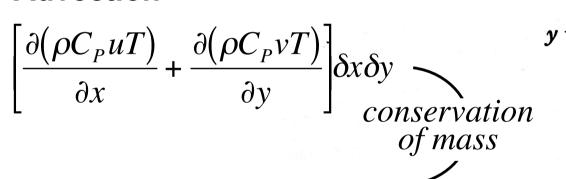
Spatial, constant  $\rho$ ,  $C_P$ , k, incompressible

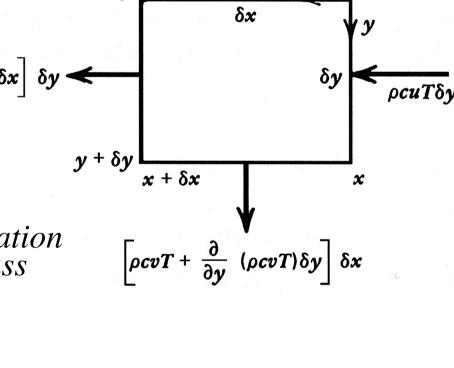


$$\frac{\partial(\rho C_P T)}{\partial t} \delta x \delta y$$

$$\left[\rho cuT + \frac{\partial}{\partial x} \left(\rho cuT\right) \delta x\right] \delta y \blacktriangleleft$$

Advection





 $\rho cvT\delta x$ 

#### Conduction

$$-k\left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2}\right) \delta x \delta y$$

$$\rho C_P \left[ \frac{\partial T}{\partial t} + u \cdot \nabla T \right] = k \nabla^2 T$$

## **Energy equation**

$$\frac{D(K+U)}{Dt} = W + Q$$

Material derivative internal heat

$$\rho C_P \left[ \frac{\partial T}{\partial t} + u \cdot \nabla T \right] = \rho C_P \frac{DT}{Dt} \Rightarrow \frac{D(\rho C_P T)}{Dt}$$

Heat input

$$k\nabla^2 T \Rightarrow \nabla \cdot k\nabla T$$

Conduction

+*A* 

Allowing for spatial variations of material parameters

Internal heat production

- Work done
  - ⇒ Changes in *motion* (kinetic energy) and *internal deformation*

Net effect of 
$$W - \frac{DK}{Dt}$$
 becomes  $\sigma : \mathbf{D}$ 

$$\mathbf{D} - \text{strain rate}$$

### 1-D advection-diffusion solution

$$-v_z \frac{\partial T}{\partial z} = \kappa \frac{\partial^2 T}{\partial z^2}$$

Take 
$$f(z) = \frac{\partial T}{\partial z}$$
 and  $c = \frac{-v_z}{\kappa}$ 

Then 
$$\frac{\partial f}{\partial z} = -cf(z)$$

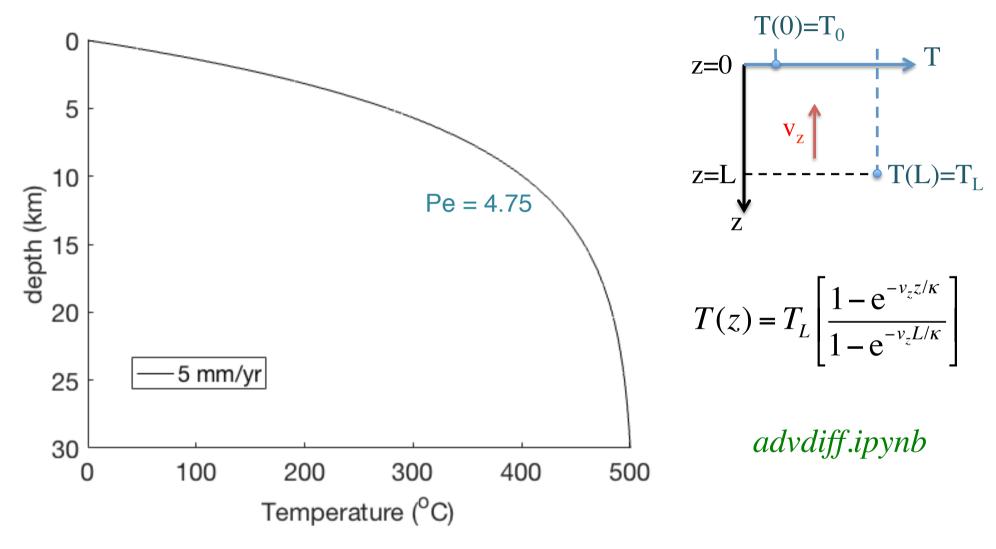
$$T(0)=T_0$$
 temperature  $T$ 
 $z=0$ 
 $z=L$ 
 $z=L$ 
 $z=depth$  upward (downward) advection

This yields 
$$f(z) = f(0)e^{-cz}$$
, i.e.  $\frac{\partial T}{\partial z}(z) = Ae^{-v_z z/\kappa}$  where A, B are  $T(z) = B - \frac{A}{v_z/\kappa}e^{-v_z z/\kappa}$  integration constants

For constant temperature boundary conditions at z=0 and z=L

$$\Rightarrow \text{Integration gives:} \qquad T(z) = T_L \left[ \frac{1 - e^{-v_z z/\kappa}}{1 - e^{-v_z L/\kappa}} \right]$$

#### 1-D advection-diffusion solution



Peclet number, measure of relative importance advection/diffusion

$$Pe = \frac{v_z L}{\kappa} = \frac{[(m/s)m]}{[m^2/s]}$$

#### **Energy equation**

conservation of heat

I III III IV V VI 
$$D(\rho C_p T)/Dt = \nabla \cdot k \nabla T + A + \sigma : D + \alpha T v \cdot \nabla P \dots)$$

- I change in temperature with time
- **II** heat transfer by conduction (and radiation)
- **III** heat production (including latent heat)
- IV heat generated by internal deformation
- V heat generated by adiabatic compression
- VI other heat sources, e.g. latent heat

## Conservation equations

- Conservation of mass  $\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0$
- Conservation of linear momentum  $\rho \frac{D\mathbf{v}}{Dt} = \nabla \cdot \underline{\boldsymbol{\varphi}} + \mathbf{f}$
- Conservation of angular momentum:  $\sigma = \sigma^T$
- Conservation of energy  $\frac{D(\rho C_P T)}{Dt} = \nabla \cdot k \nabla T + A + \sigma : \mathbf{D}$
- Entropy inequality Rate of entropy increase of a particle always  $\geq$  entropy supply

## Continuum Mechanics Equations

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#### **Material-specific**

4. <u>Constitutive equations</u> – relations describing how material properties vary as a function of T,P, stress,.... Such material properties govern dynamics (e.g., *density*), response to stress (*viscosity*, *elastic parameters*), heat transport (*thermal conductivity*, *heat capacity*)

## Thermal parameters

Can you name 4 material parameters that affect temperatures or how material responds to changes in temperature

Each of these may depend on T, P, phase, composition,...

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## Rheology

deformation (**E**) =  $rheology \cdot stress (\sigma)$ 

material <u>response to stress</u>, depends on material, P,T, time, deformation history, environment (volatiles, water)

- elastic
- viscous
- brittle
- plastic

- experiments under simple stress conditions
- ⇒ strain evolution under constant stress, stress-strain rate diagrams
- thermodynamics + experimental parameters
- ab-initio calculations

## Recap Fluid - Solid

What is a solid?
 A solid acquires <u>finite deformation under stress</u>
 stress σ ~ strain ε

What is a fluid?

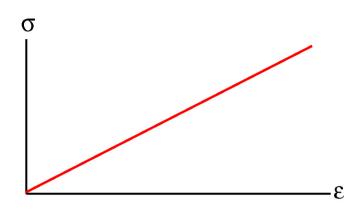
A material that flows in response to applied stress stress  $\sigma \sim strain\ rate\ D\epsilon/Dt$ 

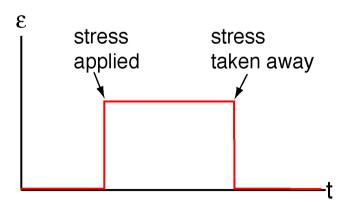


Figures from Funiciello et al. (2003a)

### **Elasticity**

- linear response to load applied
- instantaneous
- completely recoverable
- below threshold (yielding) stress
- dominates behaviour of coldest part of the lithosphere on time scales of up to 100 m.y. -> fault loading
- on time scale of seismic waves the whole Earth is elastic
- $\sigma_{ij} = C_{ijkl} \varepsilon_{kl}$  Hooke's law  $C_{ijkl}$  rank 4 elasticity tensor  $3^4$  elements, up to 21 independent





ε=E, to avoid confusion with Young's modulus E

#### **Elasticity tensor**

$$C_{ijkl}$$
 34=81 elements (for n=3)

- symmetry of  $\sigma_{ij}$  and  $\varepsilon_{kl}$ 
  - ⇒ only 36 independent elements

Why 36?

- **■** conservation of elastic energy  $U=\sigma:\epsilon=C:\epsilon:\epsilon \ge 0$ 
  - $\Rightarrow C_{ijkl} = C_{klij}$
  - ⇒ only 21 independent elements most general form of C

• other symmetries further reduce the number of independent elements

#### **Elasticity tensor**

• for example for *isotropic* media Only 2 independent elements  $(\lambda, \mu)$ :

$$\begin{split} \sigma_{ij} &= \lambda \delta_{ij} \delta_{kl} \epsilon_{kl} + \alpha \delta_{ik} \delta_{jl} \epsilon_{kl} + \beta \delta_{il} \delta_{jk} \epsilon_{kl} \\ &= \lambda \delta_{ij} \epsilon_{kk} + \alpha \epsilon_{ij} + \beta \epsilon_{ji} \\ &= \lambda \delta_{ij} \theta + (\alpha + \beta) \epsilon_{ij} \end{split}$$

3 isotropic rank  
4 tensors:  
$$\delta_{ii}\delta_{kl}$$
,  $\delta_{ik}\delta_{il}$ ,  $\delta_{il}\delta_{ik}$ 

$$\Rightarrow \sigma_{ij} = \lambda \theta \delta_{ij} + 2\mu \epsilon_{ij}$$

#### Hooke's law for isotropic material: 2 independent coefficients

Lamé constants

$$\lambda$$
 and  $\mu$ :  $\sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij}$ 

Bulk and shear modulus

K and 
$$\mu$$
=G:  $\sigma_{ij} = -p\delta_{ij} + \sigma'_{ij}$   
where:  $p = \frac{\sigma_{kk}}{3} = \left(\lambda + \frac{2}{3}\mu\right)\theta = K\theta$  hydrostatic  $\sigma'_{ij} = \lambda\theta\delta_{ij} + 2\mu\epsilon_{ij} - p\delta_{ij} = 2\mu\epsilon'_{ij}$  deviatoric

Young's modulus and Poisson's ratio

**E** and 
$$\mathbf{v}$$
:  $\mathbf{E} = \sigma_{11}/\epsilon_{11}$ ,  $\mathbf{v} = -\epsilon_{33}/\epsilon_{11}$  (uniaxial stress)

Determine in problem set

#### Wave equation

For infinitesimal deformation: spatial coordinates ≈ material coordinates

$$v_i \text{ (spatial)} \approx \partial u_i / \partial t$$
  
 $a_i \text{ (spatial)} \approx \partial v_i / \partial t = \partial^2 u_i / \partial t^2$ 

Equation of motion: 
$$f_i + \partial \sigma_{ii}/\partial x_i = \rho \partial^2 u_i/\partial t^2$$
 (1)

Elastic rheology: 
$$\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2\mu \varepsilon_{ij}$$
 (2)

Substitute (2) in (1) if (infinitesimal) deformation is consequence of force balance

#### Wave equation

Equation of motion: 
$$f_i + \partial \sigma_{ji}/\partial x_j = \rho \partial^2 u_i/\partial t^2$$

Elastic rheology: 
$$\sigma_{ij} = \lambda \epsilon_{kk} \delta_{ij} + 2\mu \epsilon_{ij}$$

$$\begin{split} \partial \sigma_{ji}/\partial x_j &= \lambda \partial \epsilon_{kk}/\partial x_i + \mu \partial (\partial u_i/\partial x_j + \partial u_j/\partial x_i)/\partial x_j \\ &= \lambda \partial (\partial u_k/\partial x_k)/\partial x_i + \mu \partial^2 u_i/\partial^2 x_j + \mu \partial (\partial u_j/\partial x_j)/\partial x_i \end{split}$$

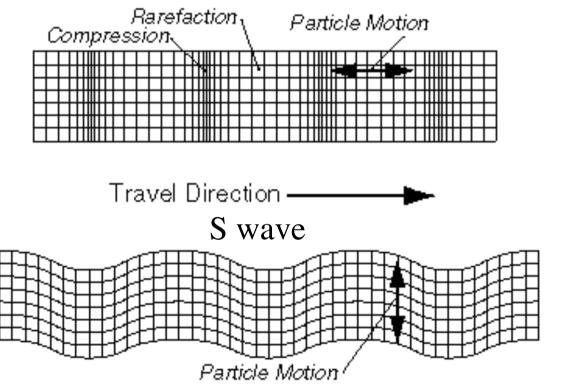
$$\nabla \cdot \sigma$$
 = Write vector equation

Using: 
$$\nabla^2 \mathbf{u} = \nabla(\nabla \cdot \mathbf{u}) - \nabla \times \nabla \times \mathbf{u}$$

$$=> \left|\rho \partial^2 \mathbf{u}/\partial t^2 = \mathbf{f} + (\lambda + 2\mu)\nabla(\nabla \cdot \mathbf{u}) - \mu \nabla \times \nabla \times \mathbf{u}\right|$$

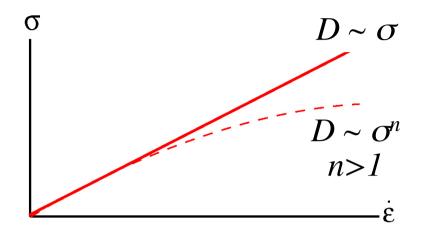
what type of deformation do the two terms represent?

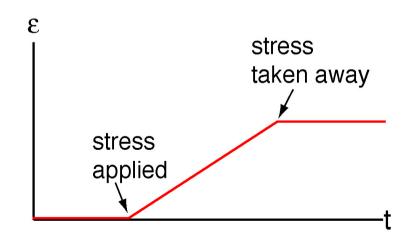
#### P wave



#### **Viscous Flow**

- steady state flow at constant stress
- permanent deformation
- linear (Newtonian) or nonlinear (e.g., Powerlaw)
   relation between strain rate and stress
- isotropic stress does not cause flow
- on timescales > years lower lithosphere and mantle deform predominantly viscously -> plate motions, postseismic deformation, but also glaciers, magmas





## Hydrostatics

Fluids can not support shear stresses

i.e. if in rest/rigid body motion:  $\mathbf{\sigma} \cdot \hat{\mathbf{n}} = \lambda \hat{\mathbf{n}}$  and this normal stress is the same on any plane:  $\mathbf{\sigma} = -p\mathbf{I}$ 

p is hydrostatic pressure

In force balance: 
$$\nabla \cdot \mathbf{\sigma} + \mathbf{f} = 0$$
  
 $-\nabla p = -\mathbf{f}$ 

In gravity field 
$$\frac{\partial p}{\partial z} = \rho g$$
  $\Rightarrow p_2 - p_1 = \rho g h$ 

### Newtonian Fluids

In general motion:

$$\sigma = -p\mathbf{I} + \sigma'$$

In Newtonian fluids, deviatoric stress varies *linearly* with *strain rate*, **D**  $D_{ij} = (\partial v_i/\partial x_j + \partial v_j/\partial x_i)/2$ 

For isotropic, Newtonian fluids, 2 material parameters:

Viscous stress tensor  $\sigma'_{ij} = -\xi D_{kk} \delta_{ij} + 2\eta D_{ij}$ 

where  $\zeta$  is *bulk viscosity* and  $\eta$  (*shear*) *viscosity*,  $\Delta = D_{kk}$ 

$$\mathbf{\sigma} = (-p + \varsigma \Delta)\mathbf{I} + 2\eta \mathbf{D}$$

p not always mean normal stress:  $\sigma_{kk} = -3p + (3\xi + 2\mu)D_{kk}$ 

## Consider a Newtonian shear flow with velocity field $v_1(x_2)$ , $v_2=v_3=0$

What is **D**? What is  $\sigma$ ?

## Navier-Stokes for incompressible Newtonian Flow

For incompressible fluids  $\Delta = 0$ , so that:  $\sigma = -p\mathbf{I} + 2\eta\mathbf{D}$ 

Force balance: 
$$\nabla \cdot \underline{\underline{\sigma}} + \mathbf{f} = \rho \frac{D\mathbf{v}}{Dt}$$

Show that: 
$$\frac{\partial \sigma_{ij}}{\partial x_i} = -\frac{\partial p}{\partial x_i} + \eta \frac{\partial^2 v_i}{\partial x_j \partial x_j}$$
 Assuming constant  $\eta$ 

$$\nabla \cdot \underline{\underline{\sigma}} = -\nabla p + \eta \nabla^2 \mathbf{v}$$

## Navier-Stokes for incompressible Newtonian Flow

For incompressible fluids  $\Delta = 0$ , so that:  $\sigma = -p\mathbf{I} + 2\eta\mathbf{D}$ 

Force balance: 
$$\nabla \cdot \underline{\underline{\sigma}} + \mathbf{f} = \rho \frac{D\mathbf{v}}{Dt}$$

Navier Stokes equation of motion:

$$-\nabla p + \eta \nabla^2 \mathbf{v} + \mathbf{f} = \rho \left[ \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right]$$
 Assuming constant  $\eta$ 

Together with continuity, 4 equations, 4 unknowns  $(p, v_x, v_y, v_z)$ 

$$\nabla \cdot \mathbf{v} = 0$$

## Navier-Stokes for compressible Newtonian Flow

$$\mathbf{\sigma} = (-p + \varsigma \Delta)\mathbf{I} + 2\eta \mathbf{D} \qquad \qquad \nabla \cdot \underline{\underline{\sigma}} + \mathbf{f} = \rho \frac{D\mathbf{v}}{Dt}$$

Navier Stokes equation of motion:

$$-\nabla p + (\zeta + \eta)\nabla \Delta + \eta \nabla^2 \mathbf{v} + \mathbf{f} = \rho \left[ \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right]$$
 Assuming constant  $\zeta, \eta$ 

Assuming

+ Conservation of mass:

$$\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0$$

6 equations 6 unknowns

$$(p, v_{\rm x}, v_{\rm y}, v_{\rm z}, \rho, T)$$

+ Energy equation

+ Equation of state for  $\rho(T,p)$ 

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- Energy equation
- Rheology
- Elasticity and Wave Equation
- Newtonian Viscosity and Navier Stokes

More reading on the topics covered in this lecture can be found in, for example: Lai et al. Ch 4.14-4-16, 6.18, Ch 5.1-5.6, Ch 6.1-6.7; Reddy parts of Ch 5 & Ch 6

## Outline of course

- ➤ 1. Mathematical essentials *Matthew Piggott*
- **≥2.** Linear Algebra I *Matthew Piggott*
- ▶3. Linear Algebra II, ODEs– Matthew Piggott
- **▶4.** Verifying models– *Matthew Piggott*
- ▶ 5. Vector and Tensor Calculus Saskia Goes
- **>6.** Stress principles Saskia Goes
- >7. Kinematics and strain Saskia Goes
- ▶8. Rheology and conservation equations Saskia Goes
- **>9.** Potential flow Stephen Neethling
- **▶10.** Fluid flow I Stephen Neethling
- ➤ 11. Fluid flow II Stephen Neethling
- ➤ 12. Wave propagation Adrian Umpleby

## Exam

- 1.5 hours
- Analytical/essay-style questions
- From first 8 lectures
- Questions as shorter <u>non-numerical</u> ones <u>from</u> the <u>problem sets/lecture examples</u>
- An example exam will be released before Christmas break
- Only most basic equations expected to remember. Others will be given.