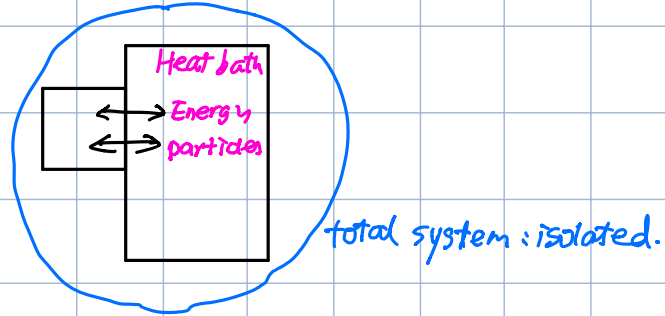
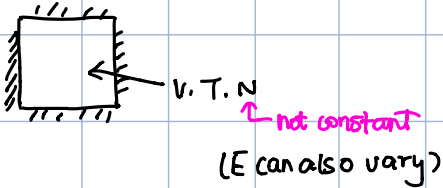


Ensembles	variables	system
Microcanonical	E, N, V	isolated system
Canonical	T, N, V	heat bath (for maintaining the temperature a constant)
Grand canonical	T, μ, V	heat bath & particle bath (maintain T & μ)

Chapter 10 Grand Canonical Ensemble

1.



Assumptions: 1° $E_{\text{system}} + E_{\text{bath}} = E_{\text{total}}$, $E_{\text{sys}} \ll E_{\text{total}} \sim E_{\text{bath}}$

2° $N_{\text{system}} + N_{\text{bath}} = N_{\text{total}}$. $N_{\text{sys}} \ll N_{\text{total}} \sim N_{\text{bath}}$

3° The total system follows microcanonical distribution.

2. Derivation

$$\rightarrow P_{\text{total}}^{\text{"i"}} = \begin{cases} a, & \text{if } E^{\text{"i"}}_{\text{tot}} = E_L \\ 0, & \text{otherwise} \end{cases}$$

$$^{\text{"i"}}_{\text{tot}} \rightarrow ^{\text{"i"}}_{\text{bath}} + ^{\text{"i"}}_{\text{sys}}$$

$$P_{\text{"i"}}^{\text{sys}} = \sum_{\text{"i"}_{\text{bath}}} P_{\text{"i"}_{\text{bath}}} = \sum_{\text{"i"}_{\text{bath}}} \begin{cases} a, & \text{if } E^{\text{"i"}}_{\text{total}} = E_L = E^{\text{"i"}}_{\text{sys}} + E^{\text{"i"}}_{\text{bath}} \\ 0, & \text{otherwise} \end{cases}$$

$$= a \cdot \# \text{ of "i" bath having } E_{\text{bath}} = E_{\text{total}} - E^{\text{"i"}}_{\text{sys}}$$

$$= a \Omega(E_{\text{total}} - E_{\text{sys}}, N_{\text{total}} - N_{\text{sys}})$$

$k_B \ln \Omega \rightarrow \text{entropy}$

$$\ln \Omega = \ln \Omega(E_{\text{tot}} - E_{\text{sys}}, N_{\text{tot}} - N_{\text{sys}}) \approx \ln \Omega(E_{\text{tot}}, N_{\text{tot}}) - E_{\text{sys}} \frac{\partial \ln \Omega}{\partial E} - N_{\text{sys}} \frac{\partial \ln \Omega}{\partial N} - \dots$$

1st + 2nd laws of thermodynamics

$$T^{-1} = \frac{\partial S}{\partial E} \quad \frac{1}{k_B T} = \frac{1}{k_B} \frac{\partial S}{\partial E} = \frac{\partial \ln \Omega_{tot}}{\partial E}$$

$$-\mu T^{-1} = \frac{\partial S}{\partial N} \quad \frac{-\mu}{k_B T} = \frac{\partial \ln \Omega_{tot}}{\partial N}$$

$$\frac{\partial \ln \Omega}{\partial E} = \beta \quad \frac{\partial \ln \Omega}{\partial N} = \alpha = \mu \beta$$

$$\ln P_{i,N}^{sys} = \ln a + \ln \Omega = \ln a + \ln \Omega(E_t, N_t) - E_t \frac{1}{k_B T} + N_t \frac{\mu}{k_B T} + \dots$$

$$\Rightarrow P_{i,N}^{sys} \propto \exp(-\beta E_{i,N}^{sys} + \mu \beta N_{i,N}^{sys})$$

3. Normalization

$$1 = \sum_{i,N} P_{i,N} = A \sum_{i,N} \exp(-\beta E_{i,N}^{sys} + \alpha N)$$

$$A = \left(\sum \exp(-\beta E_{i,N}^{sys} + \alpha N) \right)^{-1}$$

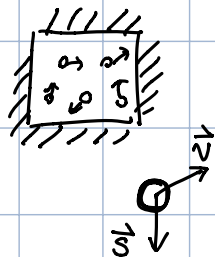
$$P_{i,N} = \frac{\exp(-\beta E_{i,N}^{sys} + \alpha N)}{\sum_{i,N} \exp(-\beta E_{i,N}^{sys} + \alpha N)}$$

$$Z \equiv \sum_{i,N} \exp(-\beta E_{i,N}^{sys} + \mu N)$$

Grand Canonical
Partition Function.

Chapter 11 Quantum Ideal gas

1.



Ideal gas \rightarrow No interaction between particles.
weak

particles \rightarrow Quantum particles: Fermions: spin half integer $\cdot \hbar$ e^-
Bosons: spin integer $\cdot \hbar$ γ
 \uparrow
carry spins.
intrinsic

- Pauli Exclusion Principle: No two fermions can occupy same Quantum State.
- Indistinguishability: cannot label the particles.

2. What do we determine?

Quantum state energy levels $\left\{ \begin{array}{l} \dots \\ \epsilon_4 \\ \epsilon_3 \\ \epsilon_2 \\ \epsilon_1 \end{array} \right.$ $\epsilon_{i,j}$ \rightarrow energy states of a single particle \uparrow ideal gas
 i_1, i_2, i_3

On average $\langle n_i \rangle \rightarrow$ average occupation number (can be non-integer)
as a function of ϵ_i

3. Derivation

consider a distribution of occupation number

$$\begin{array}{ccccccc} \epsilon_1 & , & \epsilon_2 & , & \epsilon_3 & , & \epsilon_4 \\ \downarrow & & \downarrow & & \downarrow & & \downarrow \\ n_1 & & n_2 & & n_3 & & n_4 \end{array} \text{ integers}$$

Note: No labels.

We know: $N = n_1 + n_2 + n_3 + \dots$ # of particles in the system.
label of quantum states

system energy $E_i = n_1 \epsilon_1 + n_2 \epsilon_2 + \dots$
of quantum state i

Gibbs Distribution of Grand Canonical Ensemble.

$$P_{i,N} = P(n_1, n_2, n_3, \dots)$$

$$\propto \exp(-\beta E_{i,\text{sys}} + \alpha N) = \exp(-\beta(n_1 \epsilon_1 + n_2 \epsilon_2 + \dots) + \alpha(n_1 + n_2 + \dots))$$

ideal gas

Grand canonical partition function

$$\begin{aligned} Z(\alpha, \beta) &= \sum_{\{n_i\}} \exp(-\beta(n_1 \epsilon_1 + n_2 \epsilon_2 + \dots) + \alpha(n_1 + n_2 + \dots)) \\ &= \prod_i \sum_{n_i=0}^{\infty} \exp(-\beta n_i \epsilon_i + \alpha n_i) = \prod_i Z_i(\alpha, \beta, \epsilon_i) \end{aligned}$$

system

Fermions (Fermi-Dirac Theory) $n=0,1$. No double/triple... occupation.

$$Z_{i,\text{FD}} = 1 + \exp(-\beta \epsilon_i + \alpha) \text{ Fermi-Dirac partial partition function.}$$

Bosons (Bose-Einstein Theory) $n=0,1,2,\dots$

$$Z_{i,\text{BE}} = \sum_{n=0}^{\infty} \exp(-\beta n \epsilon_i + \alpha n) = \sum_{n=0}^{\infty} \exp[(-\beta \epsilon_i + \alpha) \cdot n] = \frac{1}{1 - \exp(-\beta \epsilon_i + \alpha)}$$

$$\langle n_k \rangle \equiv \sum n_k P(n_k)$$

$$P(n_k) = \sum_{\{n_i\}} P(n_1, n_2, \dots, n_k, \dots)$$

↑
keep n_k unsummed
sum over all other n_i 's
Probability Reduction.