

# SIMULATION METHODS

FACULTY 2, M.SC. PROGRAM HIGH INTEGRITY SYSTEMS

*Winter-Semester 2020/2021, Prof. Dr.-Ing. Peter Thoma*

# CONTENTS OF THIS MODULE

1 INTRODUCTION & EXAMPLE

2 ORDINARY DIFFERENTIAL  
EQUATIONS

3 SYSTEM SIMULATION

4 FINITE DIFFERENCE METHOD

5 FINITE ELEMENT METHOD

6 OUTLOOK

# • SUMMARY WHERE WE ARE...

$$\frac{\partial^2 u(x, t)}{\partial x^2} - k \frac{\partial^2 u(x, t)}{\partial t^2} = 0$$

Wave Equation



Initial value problem



Eigenvalue problem

- Time Domain simulation
- Start from an initial state and simulate how the system behaves as function of time (iterative)
- No further excitation during the simulation

- Frequency Domain simulation
- Determine natural resonances of a system (solve eigenvalue problem)
- No excitation at all
- Eigenvectors (eigenmodes) describe the shapes of the resonances
- Eigenvalues describe the resonance frequencies





# THE FINITE DIFFERENCE METHOD

DRIVEN PROBLEMS



# FINITE DIFFERENCE METHOD AGENDA

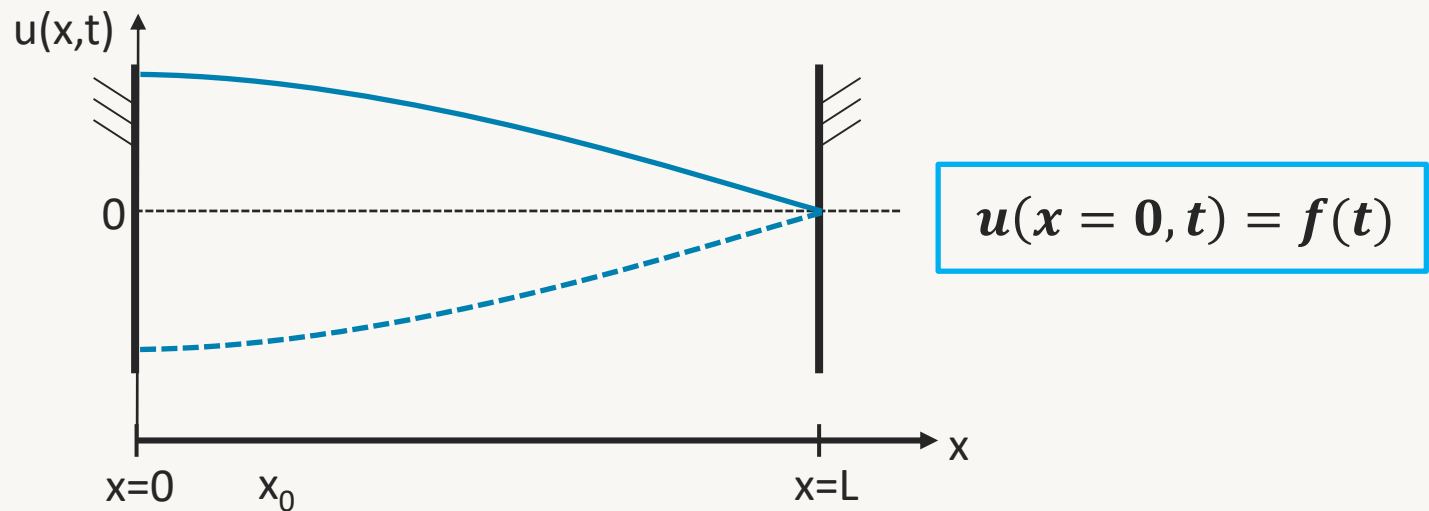
1. Wave equation with excitation
2. Excitation functions
3. Driven problems in Frequency Domain
4. Results in Time and Frequency Domain
5. Solutions of the wave equation
6. Varying material properties
7. Commercial software CST Studio Suite

# • WAVE EQUATION WITH EXCITATION

## SOLUTIONS OF THE WAVE EQUATION

$$\frac{\partial^2 u(x, t)}{\partial x^2} - k \frac{\partial^2 u(x, t)}{\partial t^2} = 0$$

Now we will consider a **driven** problem where we will have an **excitation** at the left end of the string.





# BOUNDARY CONDITIONS FOR DIFF. EQUATIONS

## TYPES OF BOUNDARY CONDITIONS FOR DIFFERENTIAL EQUATIONS

### 1. Dirichlet (first-type) boundary condition

Describe the **value** of the solution **at** the **domain boundary**, e.g.

$$u(x = 0, t) = f(t)$$

### 2. Neumann (second-type) boundary condition

Describe the **derivative** of the solution **normal** to the **domain boundary**, e.g.

$$\left. \frac{\partial u(x, t)}{\partial x} \right|_{x=0} = f(t)$$

### 3. Robin (third-type) boundary condition

**Combination** of **Dirichlet** and **Neumann** boundary conditions, e.g.:

$$\alpha \left. \frac{\partial u(x, t)}{\partial x} \right|_{x=0} + \beta u(x = 0, t) = f(t)$$

# • BOUNDARY CONDITIONS FOR DIFF. EQUATIONS

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# • WAVE EQUATION WITH EXCITATION

## TRANSIENT SOLUTION OF THE DRIVEN PROBLEM

$$\frac{\partial^2 u(x, t)}{\partial x^2} - k \frac{\partial^2 u(x, t)}{\partial t^2} = 0$$

Previously we derived the following **transient iteration scheme**:

$$u_{i,j} = cu_{i+1,j-1} + cu_{i-1,j-1} + 2(1-c)u_{i,j-1} - u_{i,j-2} \quad \text{with} \quad c = \frac{\Delta t^2}{k\Delta x^2}$$

Now we assume that the string is **straight** and **at rest** at the **beginning**:

$$\left. \begin{aligned} u(x, t = 0) = 0 &\Rightarrow u_{i,0} = 0 \\ \frac{\partial u(x, t)}{\partial t} \Big|_{t=0} = 0 &\Rightarrow u_{i,-1} = u_{i,0} = 0 \end{aligned} \right\} \text{Initial conditions}$$

The **end** of the string is still **fixed**, but we apply an **excitation** at the **beginning**:

$$\left. \begin{aligned} u(0, t) = f(t) &\Rightarrow u_{0,j} = f_j \\ u(L, t) = 0 &\Rightarrow u_{N,j} = 0 \end{aligned} \right\} \text{Boundary conditions (Dirichlet type)}$$

# • WAVE EQUATION WITH EXCITATION

## TRANSIENT SOLUTION OF THE DRIVEN PROBLEM

$$c = \frac{\Delta t^2}{k \Delta x^2}$$

Based on this information, we obtain the following **iteration scheme** in the **same way as before**:

$$i = 1, \dots, N - 1 \quad : \quad u_{i,0} = 0 \quad \text{Initialization}$$

---

$$\begin{aligned} & u_{0,1} = f_1 \quad \text{First step} \\ i = 1, \dots, N - 1 \quad : \quad & u_{i,1} = cu_{i+1,0} + cu_{i-1,0} + 2(1 - c)u_{i,0} - u_{i,0} \\ & u_{N,1} = 0 \end{aligned}$$

---

$$\begin{aligned} & j = 2, \dots : u_{0,j} = f_j \quad \text{Iteration} \\ i = 1, \dots, N - 1, j = 2, \dots : & u_{i,j} = cu_{i+1,j-1} + cu_{i-1,j-1} + 2(1 - c)u_{i,j-1} - u_{i,j-2} \\ & j = 2, \dots : u_{N,j} = 0 \end{aligned}$$

# • WAVE EQUATION WITH EXCITATION

## TRANSIENT SOLUTION OF THE DRIVEN PROBLEM

$$c = \frac{\Delta t^2}{k \Delta x^2}$$

Based on this information, we obtain the following **iteration scheme** in the **same way as before**:

---

$$i = 1, \dots, N - 1 \quad : \quad u_{i,0} = \boxed{0} \quad \text{Straight line} \quad \text{Initialization}$$

---

$$u_{0,1} = \boxed{f_1} \quad \text{Excitation} \quad \text{First step}$$

$$i = 1, \dots, N - 1 \quad : \quad u_{i,1} = cu_{i+1,0} + cu_{i-1,0} + 2(1 - c)u_{i,0} - u_{i,0}$$
$$u_{N,1} = 0$$

---

$$j = 2, \dots : u_{0,j} = \boxed{f_j} \quad \text{Excitation} \quad \text{Iteration}$$

$$i = 1, \dots, N - 1, j = 2, \dots : u_{i,j} = cu_{i+1,j-1} + cu_{i-1,j-1} + 2(1 - c)u_{i,j-1} - u_{i,j-2}$$

$$j = 2, \dots : u_{N,j} = 0$$

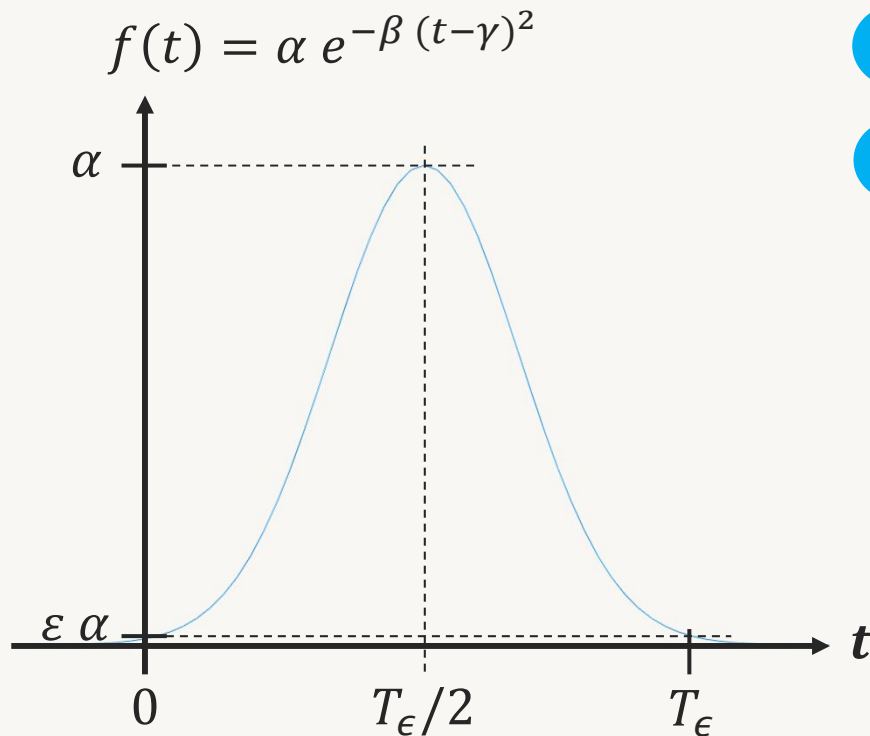
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# EXCITATION FUNCTIONS

## EXCITATION FUNCTIONS

The **excitation function** should be zero at the beginning of the simulation.  
An often-used smooth excitation function is a shifted **gaussian pulse**:



1  $f(T_\epsilon/2) = \alpha \Rightarrow \gamma = T_\epsilon/2$

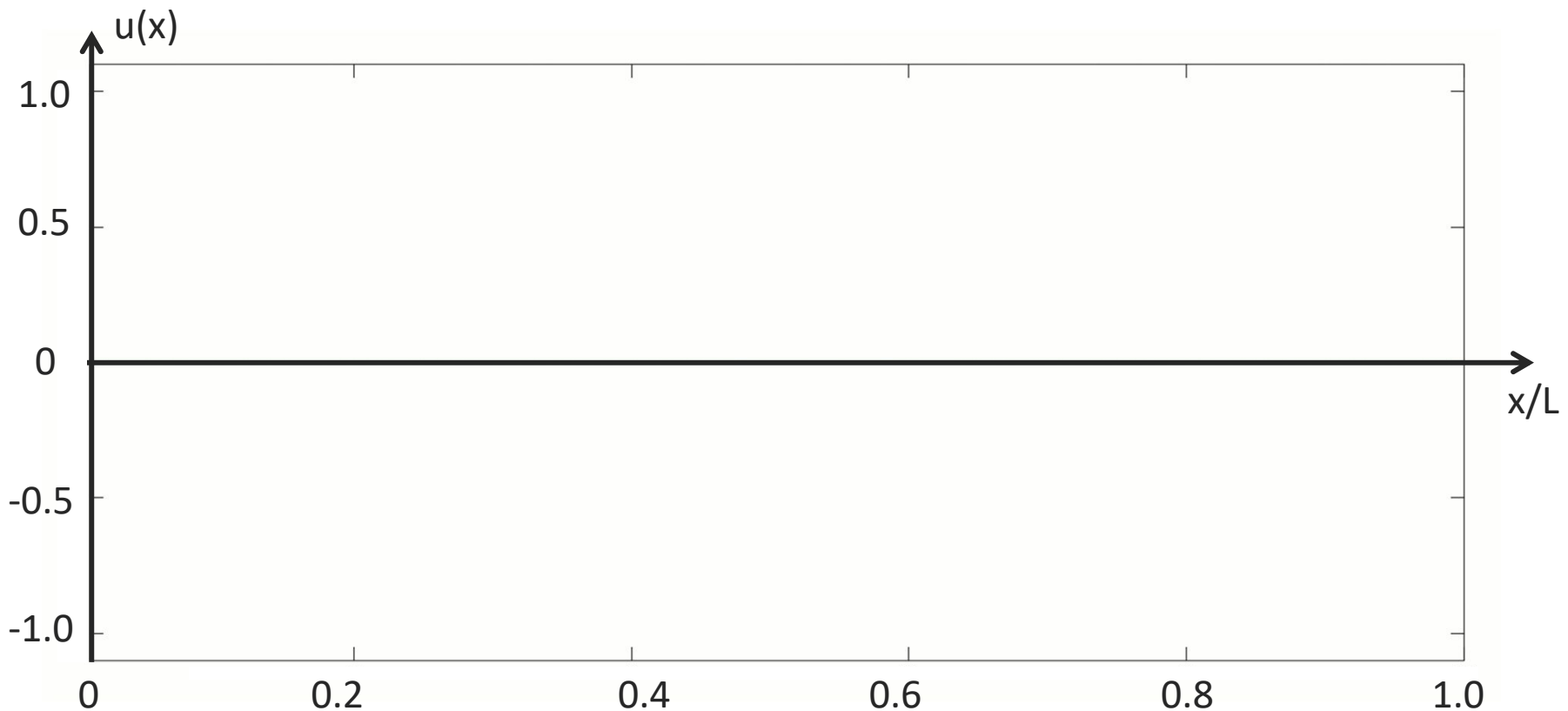
2  $f(T_\epsilon) = \alpha e^{-\beta (T_\epsilon/2)^2} = \epsilon \alpha$   
 $\Rightarrow \beta = -\left(\frac{2}{T_\epsilon}\right)^2 \ln \epsilon$

$$f(t) = \alpha e^{-\beta (t-T_\epsilon/2)^2}$$
$$\beta = -\left(\frac{2}{T_\epsilon}\right)^2 \ln \epsilon$$

Typical value for **sufficient smoothness**:  $\epsilon = 0.001$

# • EXCITATION FUNCTIONS

Run the transient scheme with **Gaussian excitation** at the left end of the string:



Video: Guitar\_string\_with\_Gaussian\_Excitation

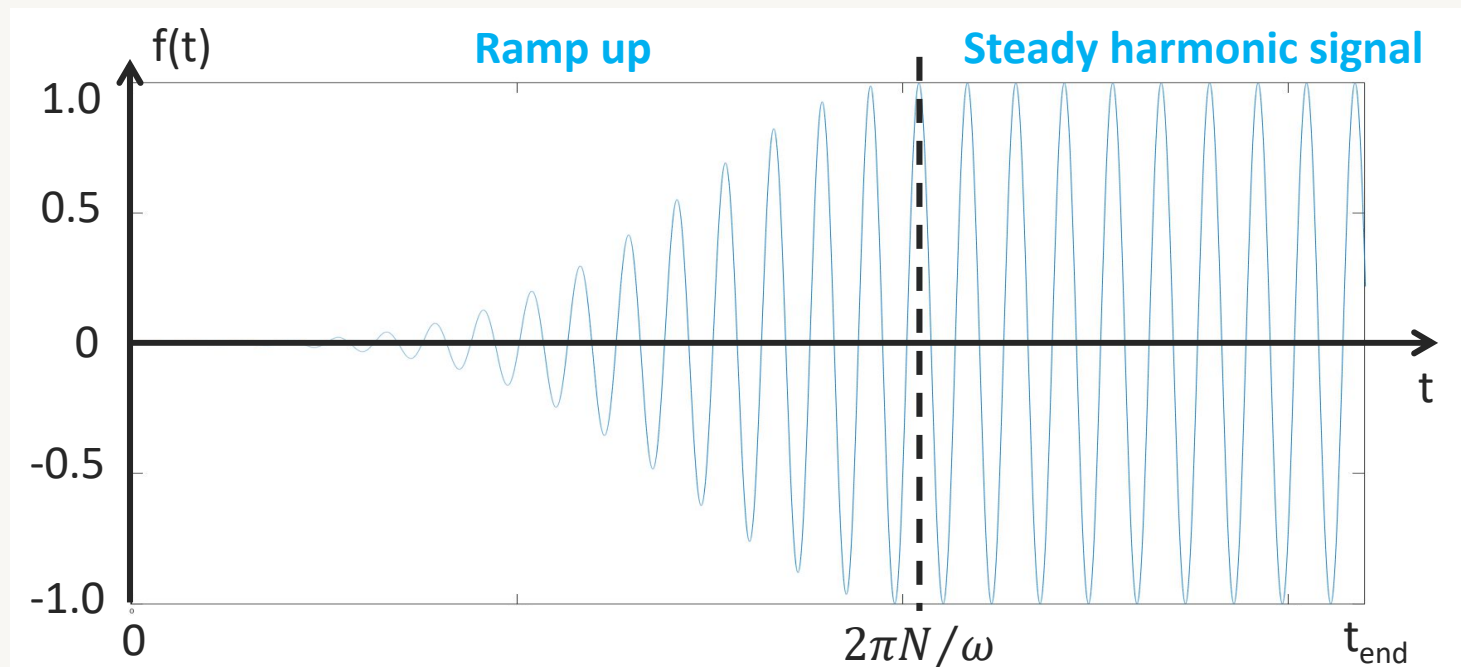


# EXCITATION FUNCTIONS

As an alternative to a pulse shaped excitation, a ramped up harmonic excitation can be used, too. A **ramped up sinusoidal** excitation function can be defined as follows:

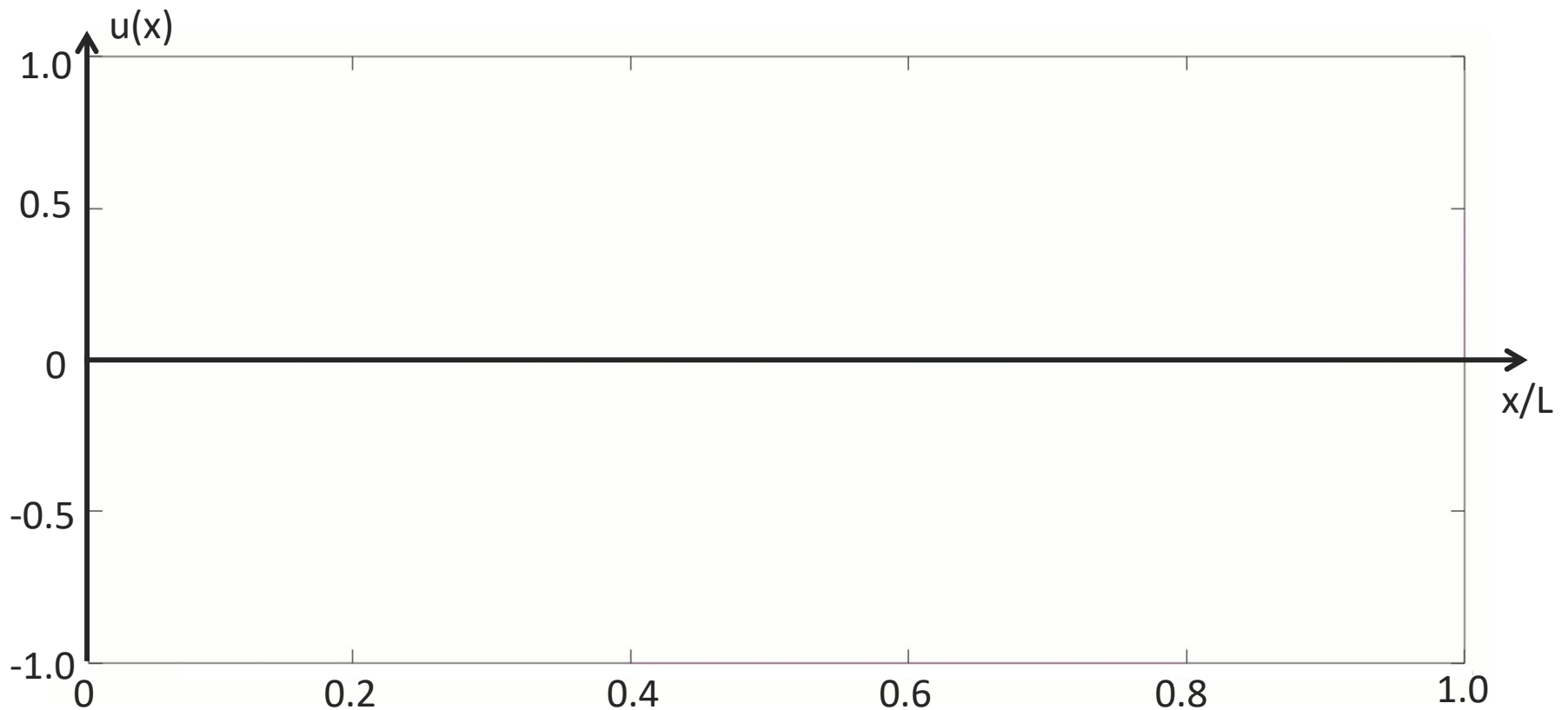
$$f(t) = \begin{cases} \frac{\ln 0.001}{\left(N\frac{2\pi}{\omega}\right)^2} (t - N\frac{2\pi}{\omega})^2 \sin(\omega t), & t < 2\pi N/\omega \\ \sin(\omega t) & , \quad t \geq 2\pi N/\omega \end{cases}$$

**Ramp up for N periods**



# • EXCITATION FUNCTIONS

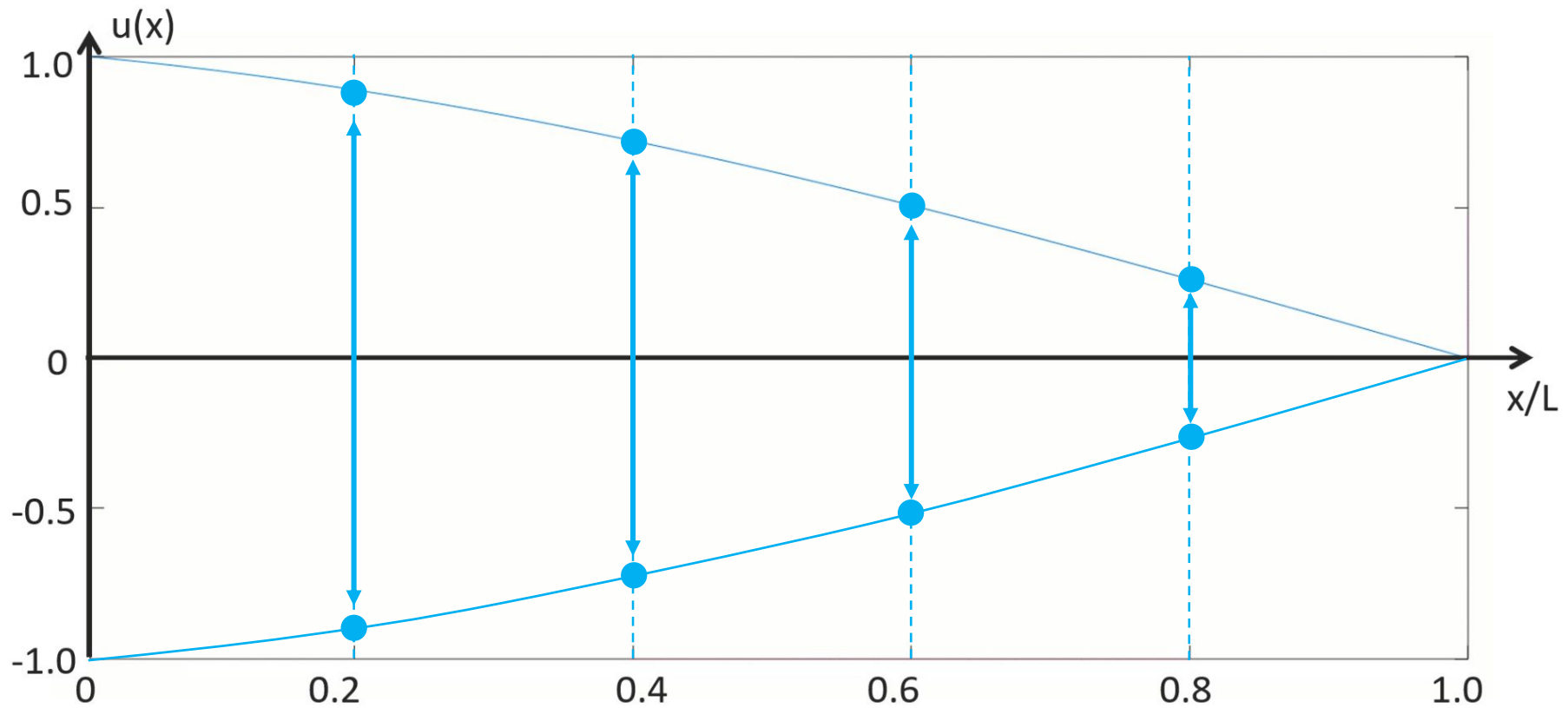
Now run the transient scheme with a **slowly ramped up sinusoidal** excitation at the left end of the string:



Video: Guitar\_string\_with\_harmonic\_excitation

# EXCITATION FUNCTIONS

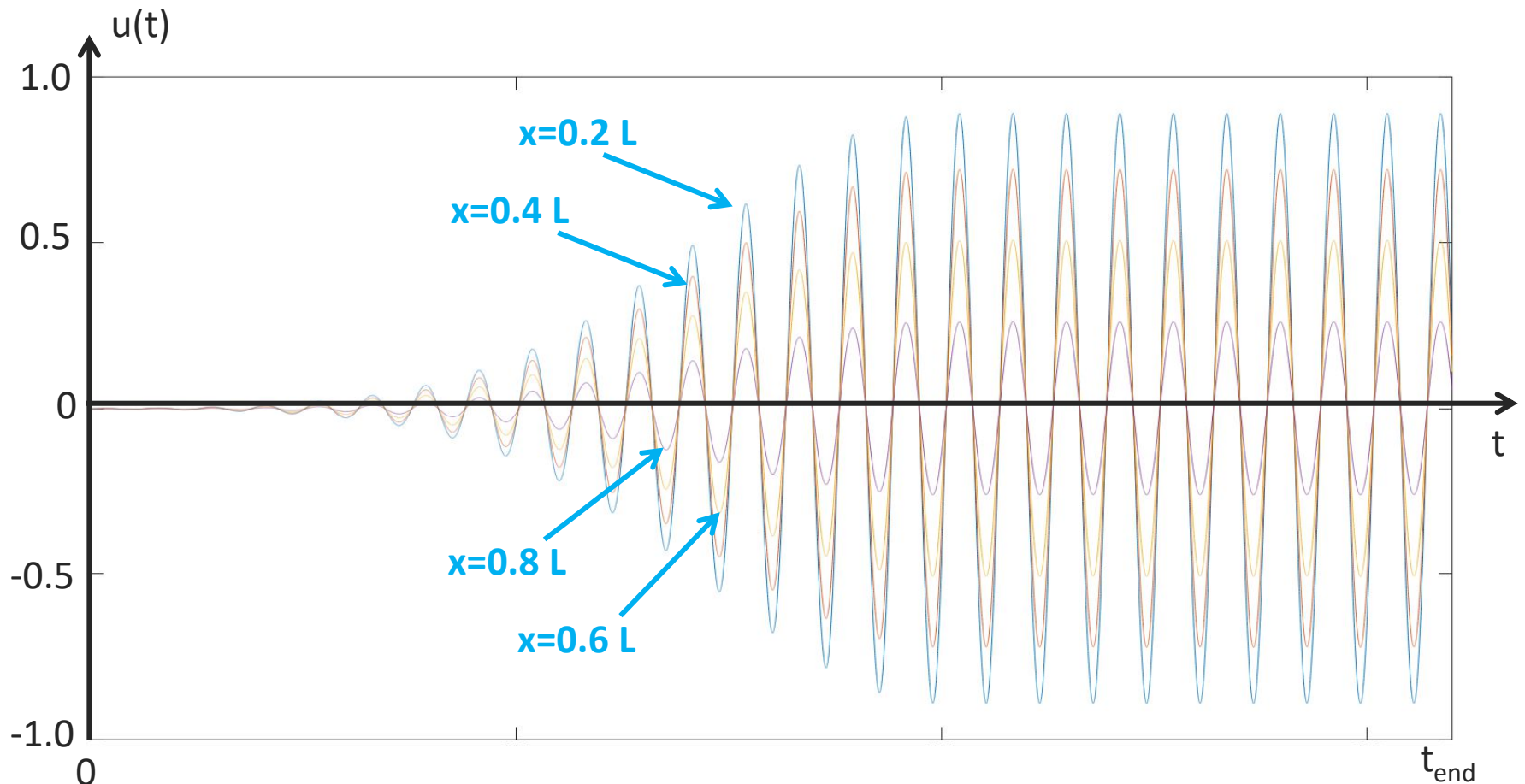
At each **position along the string**, the string performs a **harmonic oscillation** with a **position dependent amplitude**:



Thus, after a while the string performs a **continuous harmonic steady-state oscillation**

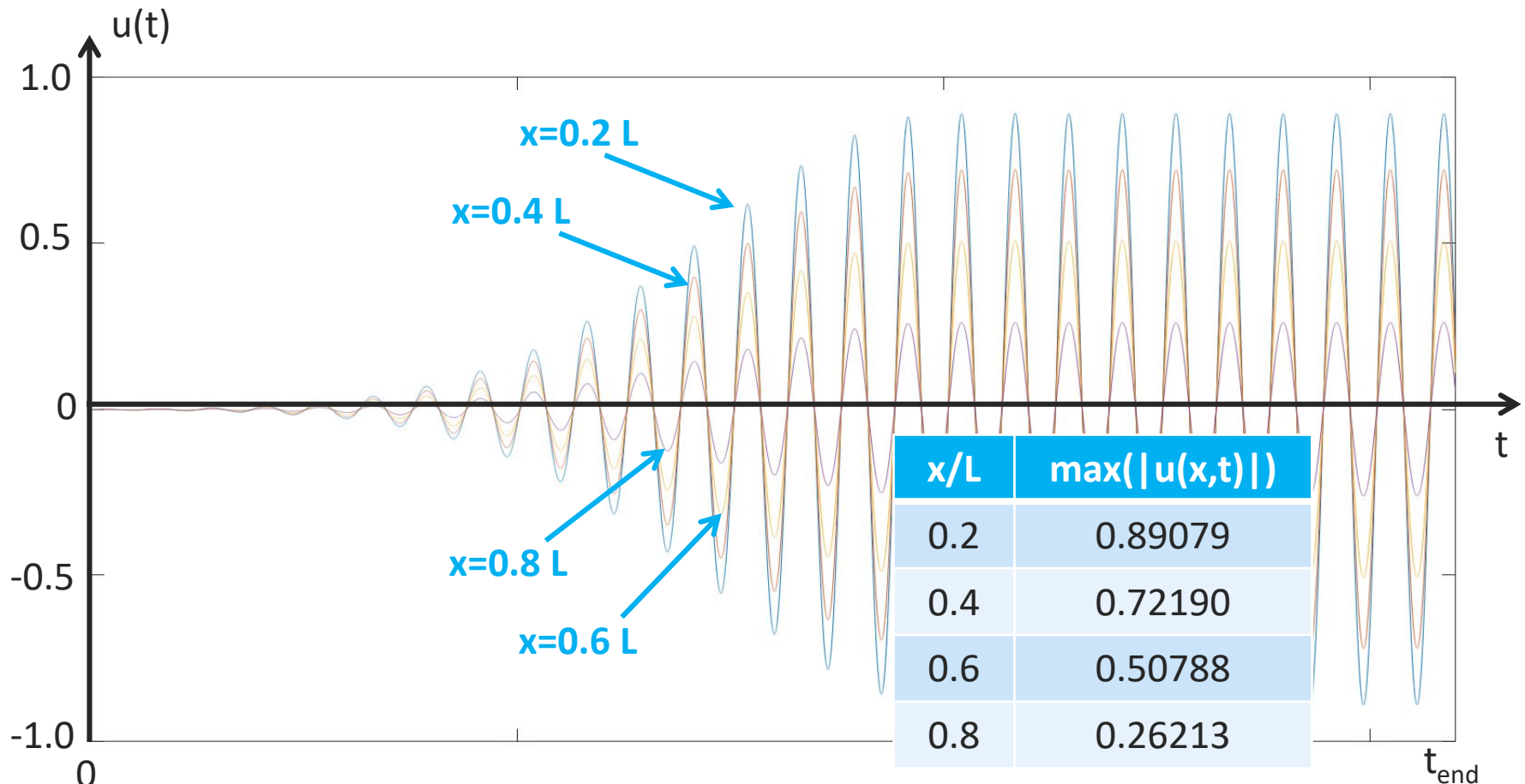
# EXCITATION FUNCTIONS

At each **position along the string**, the string performs a **harmonic oscillation** with a **position dependent amplitude**:



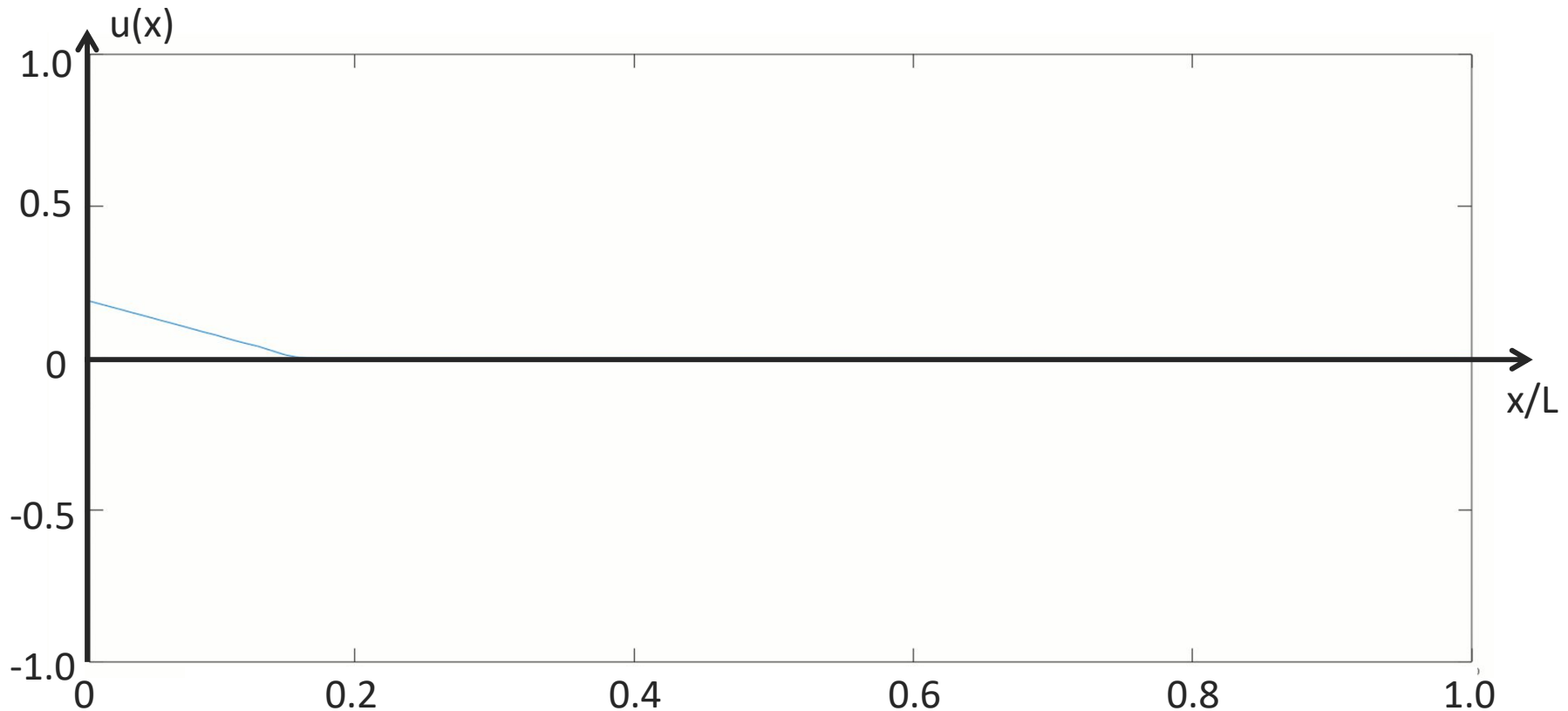
# EXCITATION FUNCTIONS

At each **position along the string**, the string performs a **harmonic oscillation** with a **position dependent amplitude**:



# EXCITATION FUNCTIONS

What happens if we switch on the sinus **without ramping up the amplitude**?

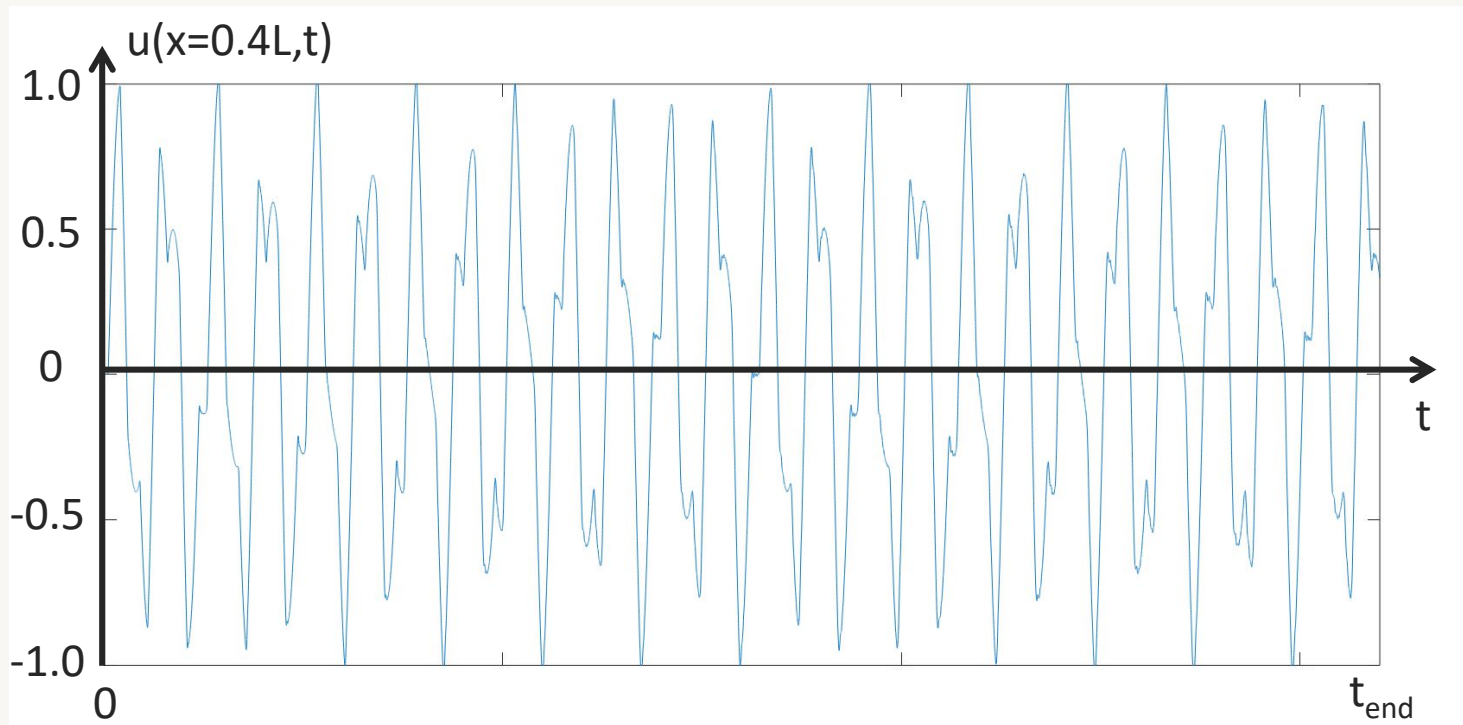


Video: Guitar\_string\_with\_step\_harmonic\_excitation



# EXCITATION FUNCTIONS

The **step function type excitation** creates a lot of **high frequency** content stimulating many **higher order modes** which disturb the harmonic solution:



**Note: Smooth excitation functions** need to be used if **harmonic steady-state** solutions shall be obtained from transient simulations directly

# • FINITE DIFFERENCE METHOD AGENDA

- ✓ 1. Wave equation with excitation
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# DRIVEN PROBLEMS IN FREQUENCY DOMAIN

## FREQUENCY DOMAIN SOLUTION OF THE DRIVEN PROBLEM

Time harmonic **steady-state** solutions can be obtained from the **Frequency Domain** formulation:

$$\frac{\partial^2 \underline{u}(x)}{\partial x^2} = -k \omega^2 \underline{u}(x)$$

Applying a spatial discretization (**central difference**) again, we obtain:

$$\frac{\underline{u}_{i+1} - 2\underline{u}_i + \underline{u}_{i-1}}{\Delta x^2} = -k \omega^2 \underline{u}_i$$

For the **special case  $i = 1$**  we get:

$$\frac{\underline{u}_2 - 2\underline{u}_1 + \underline{u}_0}{\Delta x^2} = -k \omega^2 \underline{u}_1$$

Our **time harmonic excitation function** defines the value of  $\underline{u}_0 = \underline{f}$

$$\frac{\underline{u}_2 - 2\underline{u}_1}{\Delta x^2} = -k \omega^2 \underline{u}_1 - \frac{\underline{u}_0}{\Delta x^2}$$



# DRIVEN PROBLEMS IN FREQUENCY DOMAIN

## FREQUENCY DOMAIN SOLUTION OF THE DRIVEN PROBLEM

In **matrix form**, we can then write (the first row contains the **excitation**):

$$\begin{pmatrix} \ddots & & & \\ & \ddots & & \\ & 1 & -2 & 1 \\ & & \ddots & \\ & & & \ddots \end{pmatrix} \begin{pmatrix} \underline{u}_1 \\ \underline{u}_2 \\ \vdots \\ \underline{u}_{N-1} \end{pmatrix} = -k \omega^2 \Delta x^2 \begin{pmatrix} \underline{u}_1 \\ \underline{u}_2 \\ \vdots \\ \underline{u}_{N-1} \end{pmatrix} + \begin{pmatrix} -\underline{f} \\ 0 \\ \vdots \\ 0 \end{pmatrix}$$

Or by **adding the term**  $k \omega^2 \Delta x^2 \underline{u}$ :

$$\underbrace{\begin{pmatrix} \ddots & & & \\ & \ddots & & \\ & 1 & k \omega^2 \Delta x^2 - 2 & 1 \\ & & \ddots & \\ & & & \ddots \end{pmatrix}}_M \underbrace{\begin{pmatrix} \underline{u}_1 \\ \underline{u}_2 \\ \vdots \\ \underline{u}_{N-1} \end{pmatrix}}_{\underline{u}} = \underbrace{\begin{pmatrix} -\underline{f} \\ 0 \\ \vdots \\ 0 \end{pmatrix}}_b$$



# DRIVEN PROBLEMS IN FREQUENCY DOMAIN

## FREQUENCY DOMAIN SOLUTION OF THE DRIVEN PROBLEM

The resulting **equation system** has the form of:

$$M \underline{u} = b$$

where  $b$  is called the **excitation vector** or often simply **right-hand side**.

The system can be solved by calculating the **matrix inverse**

$$\underline{u} = M^{-1} b$$

or by applying any algorithm to solve this **linear system of equations**:

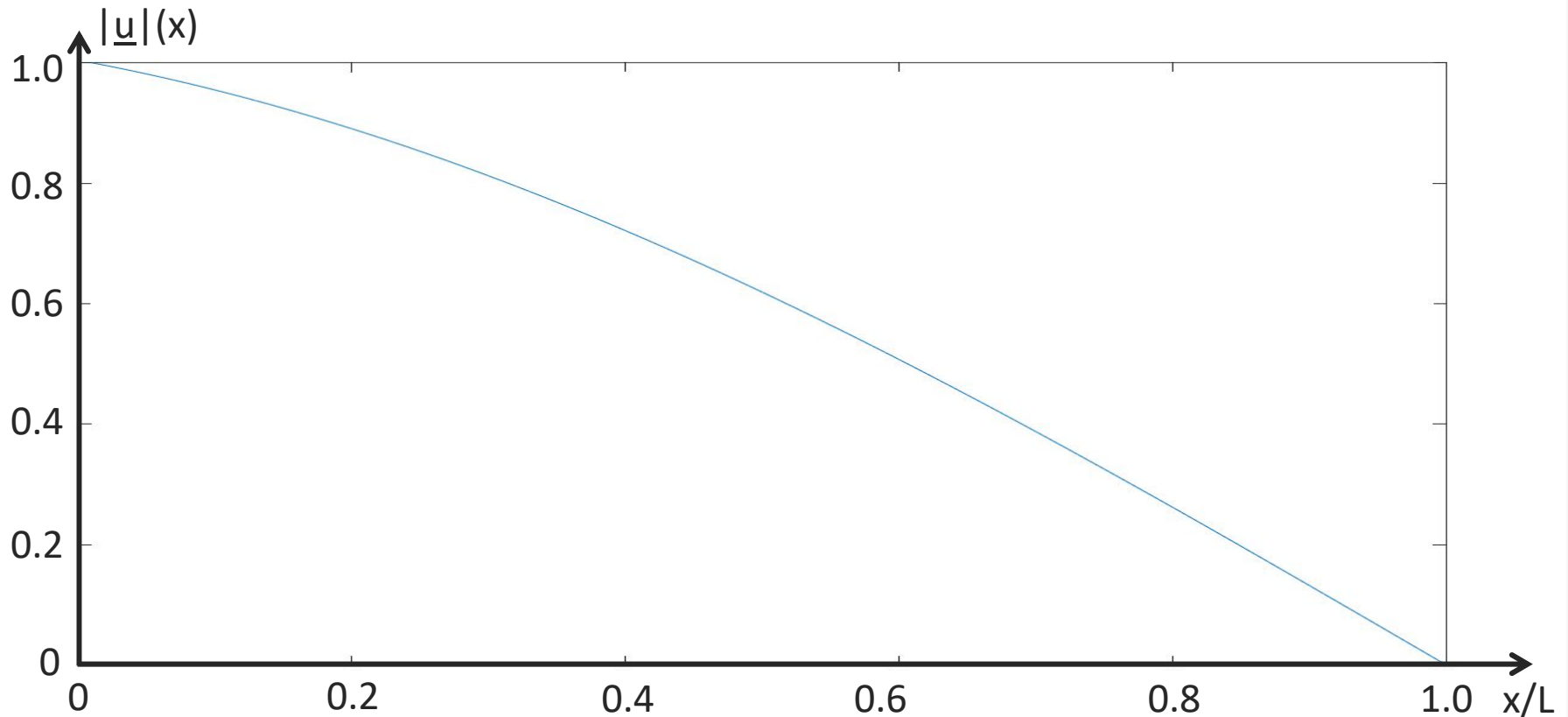
1. **Direct methods**, e.g.
  1. Gaussian elimination
  2. Matrix factorization
2. **Iterative methods**, e.g.
  1. Preconditioned Krylow subspace methods (e.g. CG)
  2. Multigrid methods



# DRIVEN PROBLEMS IN FREQUENCY DOMAIN

## FREQUENCY DOMAIN SOLUTION OF THE DRIVEN PROBLEM

Solving the equation system in MATLAB using `linsolve` yields:





# MATLAB: USEFUL TIPS AND TRICKS

## SOLVING MATRIX EQUATIONS $M u = b$

```
u = inv(M)*b           % or better: u = M\b

u = linsolve(M, b) % direct solution using factorization

% iterative solvers for sparse matrices
S = sparse(M)         % Convert a full matrix into sparse form

u = pcg(S, b)          % preconditioned conjugate gradients method
u = bicg(S,b)          % BiConjugate gradients method
u = bicgstab(S,b)      % BiConjugate gradients stabilized method
u = qmr(S,b)           % Quasi-Minimal Residual (QMR) method
u = tfqmr(S,b)         % Transpose-Free QMR method
```

There is a **large collection** of other solvers and **additional libraries** available for solving **linear equation systems**. **Optimal solver selection** depends on matrix properties and influences **runtime** and **memory** requirements

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# RESULTS IN TIME AND FREQUENCY DOMAIN

Comparing the Time Domain and Frequency Domain results for the harmonic steady-state behavior, we get:

x/L	Transient $\max( u(x,t) )$	Frequency Domain $ \underline{u}(x) $	Deviation   TD from FD
0.2	0.89079	0.89376	0.33%
0.4	0.72190	0.72787	0.82%
0.6	0.50788	0.51672	1.74%
0.8	0.26213	0.27345	4.14%

## Reasons for the differences:

1. Steady-state solution not yet reached in transient simulation
2. Excitation of higher order eigenmodes in transient simulation (smoothness of excitation)
3. Discretization error of the time derivatives (numerical dispersion)
4. Numerical error when solving the linear equation system

**Conclusion:** It is not straight-forward to compute harmonic steady-state solutions from transient simulations, but possible.

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# • SOLUTIONS OF THE WAVE EQUATION

$$\frac{\partial^2 u(x, t)}{\partial x^2} - k \frac{\partial^2 u(x, t)}{\partial t^2} = 0$$

Wave Equation

Time Domain

Frequency Domain

Initial Value Problem

Driven Problem

Eigenmode Problem

Driven Problem

Weather forecast

Automotive  
suspension dynamics

Bridge resonance

Radiation of Wifi  
router



# SOLUTIONS OF THE WAVE EQUATION

## SUMMARY OF FORMULATIONS

### 1. Transient

- Discretization of **spatial** and **temporal derivatives**
- Ideally suited for **transient phenomena**
- **Explicit** scheme, matrix times vector multiplications only
- Can be **less efficient** for calculating **harmonic steady-state** solutions

### 2. Frequency Domain

- Transformation into **frequency domain** and discretization of **spatial derivatives**
- Ideally suited for computing **harmonic steady-state** solutions
- Requires solutions of **equation system**, either direct or iterative
- Can be **less efficient** for calculating **transient solutions** due to large number of required frequency samples

### 3. Eigenmode

- Special variant of the **frequency domain** formulation **without excitation**
- Ideally suited for computing the **eigenmodes** (resonances) of a system



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# • VARYING MATERIAL PROPERTIES

## MATERIAL PROPERTIES

So far, we have solved the wave equation for **constant**  $k$ :

$$\frac{\partial^2 u(x, t)}{\partial x^2} - k \frac{\partial^2 u(x, t)}{\partial t^2} = 0$$

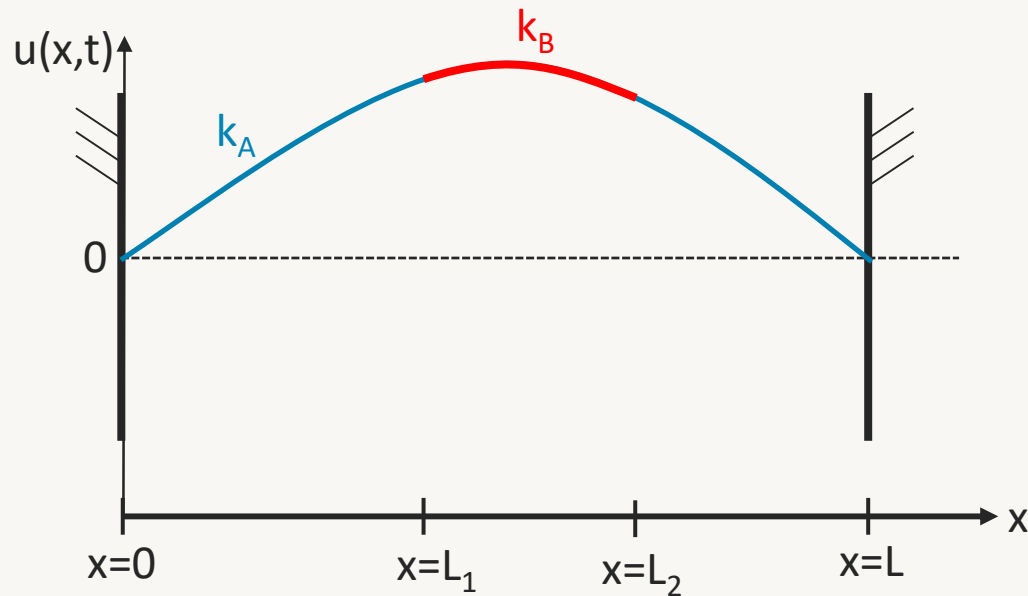
The parameter  $k$  describes the **speed** of the spatial **wave propagation** and is determined by the **properties of the materials** in which the wave propagates

In our example, it is described by the **stiffness of the string**

What happens if the string contains a segment with a **different stiffness**?

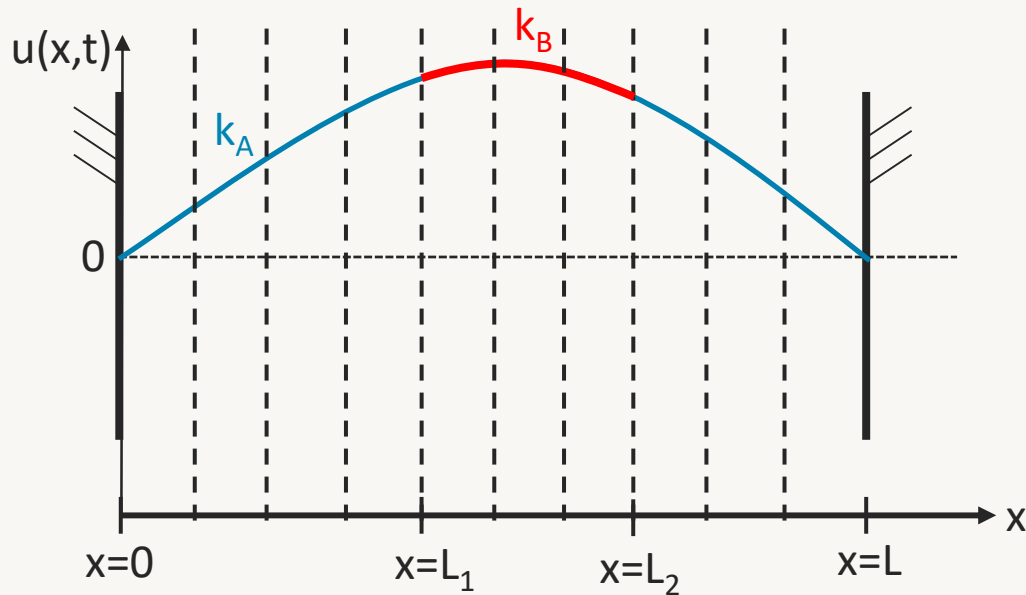
# • VARYING MATERIAL PROPERTIES

## MATERIAL PROPERTIES



# • VARYING MATERIAL PROPERTIES

## MATERIAL PROPERTIES



1. Perform a **spatial discretization** of the string as before
2. Apply a **specific value** of  $k$  for each segment (instead of using a constant  $k$ ):

$$\frac{\underline{u}_{i+1} - 2\underline{u}_i + \underline{u}_{i-1}}{\Delta x^2} = -\underline{k}_i \omega^2 \underline{u}_i$$

# • VARYING MATERIAL PROPERTIES

## EXAMPLE: TRANSIENT SOLUTION OF THE DRIVEN PROBLEM

Based on this approach, we obtain the following **iteration scheme**:  $c_i = \frac{\Delta t^2}{k_i \Delta x^2}$

$$i = 1, \dots, N - 1 \quad : \quad u_{i,0} = 0 \quad \text{Initialization}$$

$$u_{1,1} = f_1 \quad \text{First step}$$

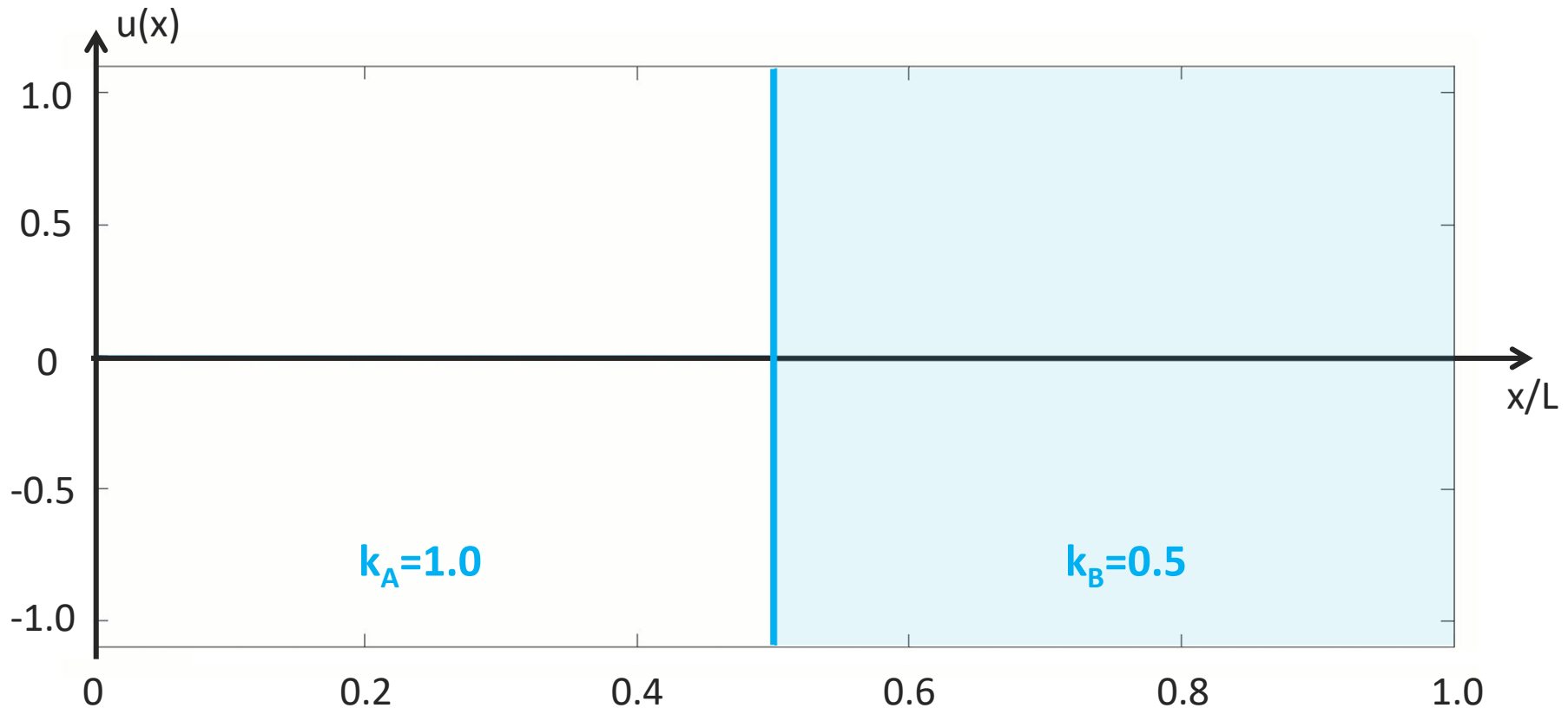
$$i = 2, \dots, N - 1 \quad : \quad u_{i,1} = c_i (u_{i+1,0} + u_{i-1,0} - 2u_{i,0}) + u_{i,0}$$

$$j = 2, \dots : u_{1,j} = f_j \quad \text{Iteration}$$

$$i = 2, \dots, N - 1, j = 2, \dots : u_{i,j} = c_i (u_{i+1,j-1} + u_{i-1,j-1} - 2u_{i,j-1}) + 2u_{i,j-1} - u_{i,j-2}$$

# • VARYING MATERIAL PROPERTIES

Run the transient scheme with **Gaussian** excitation at the left end of the string and **two different values for the string stiffness**:



Video: Guitar\_string\_with\_two\_materials

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# • COMMERCIAL SOFTWARE CST STUDIO SUITE



**CST Studio Suite®** is a high-performance **3D Electromagnetic and Thermal analysis** software package for designing, analyzing and optimizing components and systems. Next week, we will start using this commercial software package to run some simulations

A **academic version** is available in the computer rooms at the university

However, due to the current restrictions it is **highly recommended** (but optional) that you install a student edition of the software on your own computer

Since the confirmation of your registration will take a **few days**, please go ahead and download and install the software **as soon as possible**

The Student Edition can be downloaded free of charge from

<https://www.3ds.com/products-services/simulia/products/cst-studio-suite>

The following slides will explain how to **install the student edition** on your own computer

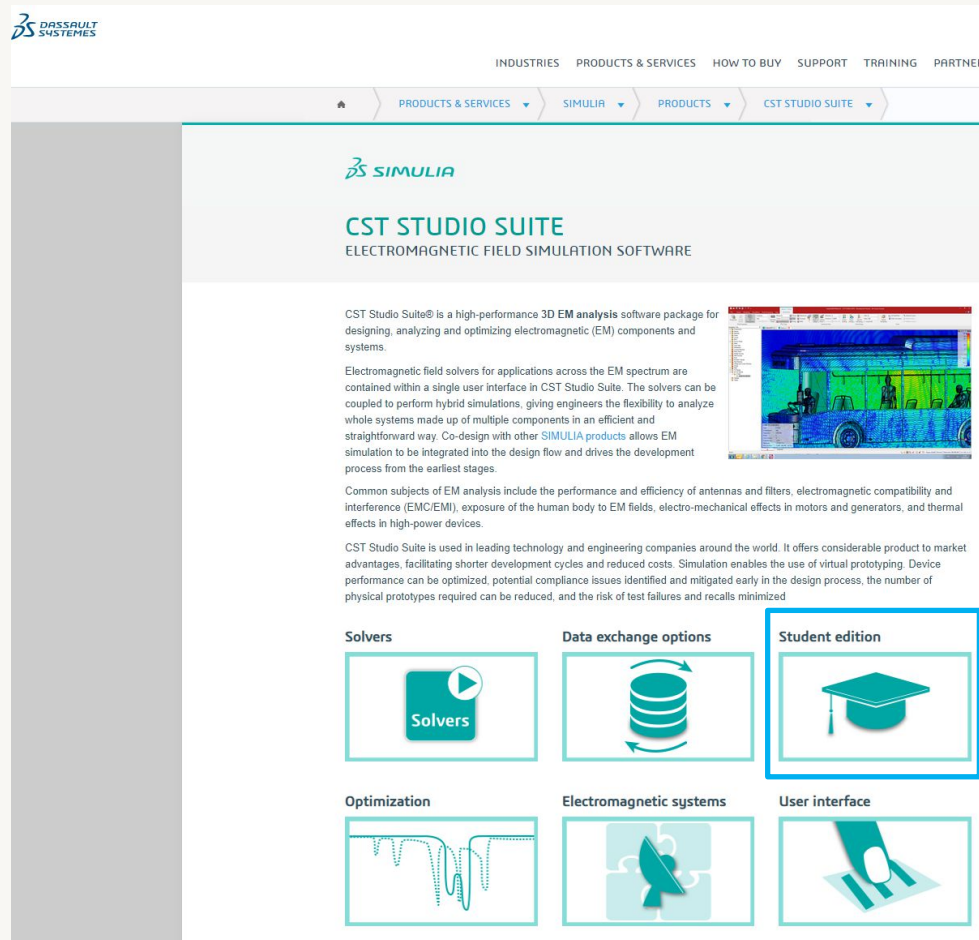




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# COMMERCIAL SOFTWARE CST STUDIO SUITE



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To accompany the CST Studio Suite Student Edition, we have prepared some examples, which are typical of the type of textbook problems you may encounter during your studies of electromagnetic theory or other related courses. Each tutorial includes a descriptive text, a CST Studio Suite file and also a short video showing how to construct each of the models.

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SIMULIA

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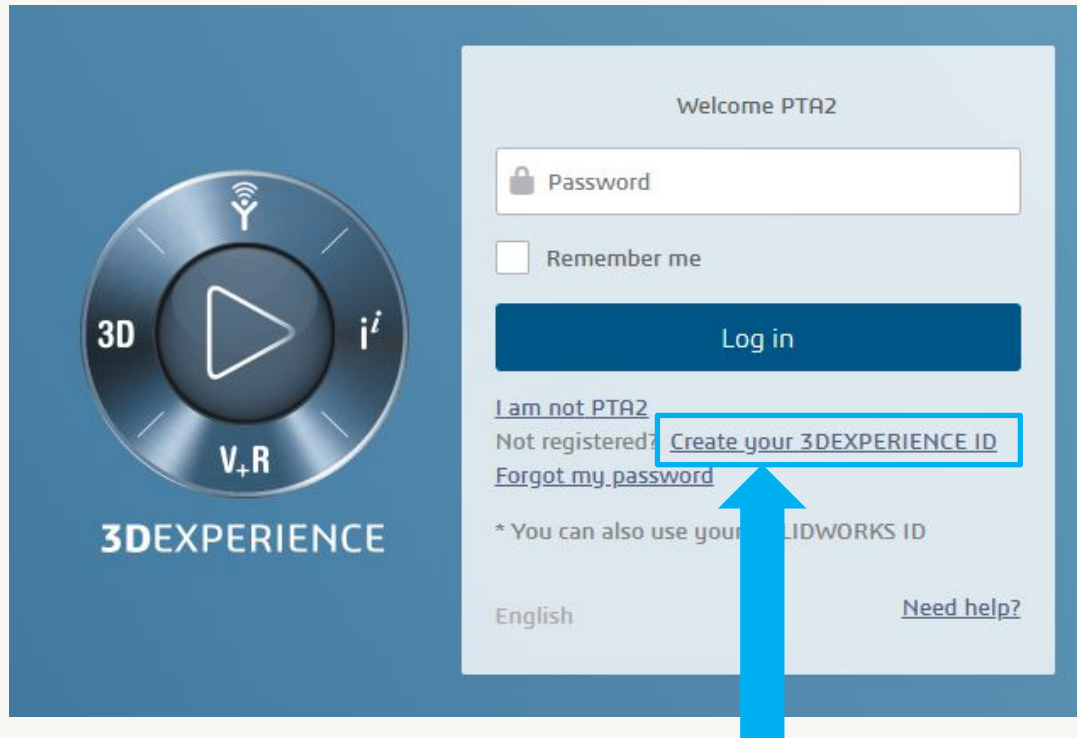
The CST Studio Suite Student Edition introduces you to the world of electromagnetic simulation, making Maxwell's equations easier to understand than ever. With this free edition you have – bar some restrictions – access to our powerful visualization engine and some of the most advanced solvers of CST Studio Suite.

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CST Studio Suite Student Edition is ideal for use as part of coursework as well as for anyone wishing to become more proficient with CST Studio Suite.

Even if you are not affiliated with an academic institution, you can immediately download the CST Studio Suite Student Edition for free and get access to the tutorials at the SIMULIA Learning Community (SLC)

# COMMERCIAL SOFTWARE CST STUDIO SUITE



You need to **create an account** to be able to download the software

The video contains introductions about **how to create an account** (starting at 0:15):

<https://youtu.be/Z8cMPTVeWsl>



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Once you have created an account and logged in, you click on “Free Download” again and follow the installation instructions.

The CST Studio Suite Student Edition introduces you to the world of electromagnetic simulation, making Maxwell's equations easier to understand than ever. With this free edition you have – bar some restrictions – access to our powerful visualization engine and some of the most advanced solvers of CST Studio Suite.

To accompany the CST Studio Suite Student Edition, we have prepared some examples, which are typical of the type of textbook problems you may encounter during your studies of electromagnetic theory or other related courses. Each tutorial includes a descriptive text, a CST Studio Suite file and also a short video showing how to construct each of the models.

CST Studio Suite Student Edition is ideal for use as part of coursework as well as for anyone wishing to become more proficient with CST Studio Suite.

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# • FINITE DIFFERENCE METHOD AGENDA

- ✓ 1. Wave equation with excitation
- ✓ 2. Excitation functions
- ✓ 3. Driven problems in Frequency Domain
- ✓ 4. Results in Time and Frequency Domain
- ✓ 5. Solutions of the wave equation
- ✓ 6. Varying material properties
- ✓ 7. Commercial software CST Studio Suite



# THANK YOU!

See you in the exercise!