

Andrew
Pepper

1. $\frac{d\psi}{dt} = F(\psi, t) = \psi\psi = (2 + i\omega)\psi$

a) Backward-Euler differencing order of accuracy:

① $\frac{d\psi}{dt}(t_n) \approx \frac{\psi(t_n) - \psi(t_n - \Delta t)}{\Delta t}$

Now we can Taylor expand $\psi(t_n - \Delta t)$:

② $\psi(t_n - \Delta t) = \psi(t_n) - \Delta t \frac{d\psi}{dt} + \frac{(\Delta t)^2}{2} \frac{d^2\psi}{dt^2} + \text{h.o.t.}$

Combining ①/②:

$$\frac{\psi(t_n) - \psi(t_n - \Delta t)}{\Delta t} - \frac{d\psi}{dt}(t_n) = \frac{\Delta t}{2} \frac{d^2\psi}{dt^2}$$

↓ (cancel out Δt)

$$\therefore \boxed{\text{Backward-Euler differencing has order of Accuracy} = 1}$$

($\Delta t \propto \text{power of } 1$)

Trapezoidal Differencing:

Since our goal is to obtain the order of accuracy, we can look at the general trapezoidal method equation:

$$\frac{\psi_{n+1} - \psi_n}{\Delta t} = \frac{(2 + i\omega)\psi_{n+1} + \psi_n}{2}$$

and if we Taylor expand ψ_{n+1} :

$$\psi_{n+1} = \psi(t_n + \Delta t) = \psi(t_n) + \Delta t \frac{d\psi}{dt}(t_n) + \frac{(\Delta t)^2}{2} \frac{d^2\psi}{dt^2} + \text{h.o.t.}$$

Now if we examine the 2 previous equations

we see on the RHS that the term $\frac{(\Delta t)^2}{2} \frac{d^2\psi}{dt^2}$

survives $\therefore \boxed{\text{order of accuracy} = 2}$

1. b) $|A| = \left| \frac{\phi_{n+1}}{\phi_n} \right| \leq 1 + \eta \Delta t$

For backward Euler we will have to simplify:

$$|A| = \left| \frac{\phi_{n+1}}{\phi_n} \right| \quad \begin{matrix} \phi_{n+1} \xrightarrow{n-1} \phi_n \\ \phi_n \xrightarrow{n-1} \phi_{n-1} \end{matrix}$$

$$\therefore \left| \frac{\phi_n}{\phi_{n-1}} \right| = \left| \frac{1}{1 - (\gamma \Delta t)} \right| \quad \begin{matrix} \text{From} \\ \phi_{n-1} \propto \phi_n - \Delta t \frac{d\phi}{dt} + \frac{(\Delta t)^2}{2} \frac{d^2\phi}{dt^2} + \dots \end{matrix}$$

$$\begin{aligned} (1 - 2\gamma \Delta t)^2 + \omega^2 \Delta t^2 &\leq 1 \\ 1 - 2\gamma \Delta t + (\omega^2 + 2\gamma^2) \Delta t^2 &\leq 1 \\ \Delta t &\leq \frac{2\gamma}{(\omega^2 + 2\gamma^2)} \end{aligned}$$

$$\frac{1}{(1 - (2\gamma \Delta t)^2 + (\omega \Delta t)^2)^{1/2}}$$

$$|1 - \gamma \Delta t| =$$

$$\begin{aligned} &\sqrt{(1 - 2\gamma \Delta t)^2} \\ &= \sqrt{1 + (2\gamma \Delta t)^2 - 2\gamma \Delta t} \end{aligned}$$

Comparing to Figure 1, we find

that as long as $2\gamma \Delta t \leq 0$ we have A-stability and can use finite Δt without worry of model blowing up. As soon as $2\gamma \Delta t > 0$ we can only avoid model blowing up if we have a large Δt .

For trapezoidal we have:

$$|A| = \left| \frac{\phi_{n+1}}{\phi_n} \right| \Rightarrow \frac{\phi_{n+1} - \phi_n}{\Delta t} = \frac{\gamma \phi_{n+1} + \phi_n}{2}$$

$$\therefore |A| \leq 1$$

$$\frac{2}{\phi_{n+1} + \phi_n} \cdot \frac{\phi_{n+1} - \phi_n}{\Delta t} = \gamma$$

this means for any $2\gamma \Delta t < 0$ we have absolute stability.

$$\begin{aligned} &\downarrow \\ &\frac{1 + \frac{\gamma \Delta t}{2}}{1 - \frac{\gamma \Delta t}{2}} \leq 1 \end{aligned}$$

Figure 5 reinforces this claim, and tells us $2\gamma \Delta t > 0$ does not

have any absolute stability, which if we consider this makes sense because we would arrive at an impossible statement

$$[1 + (-) \neq 1 - (-)]$$

2. Starting with $A = 1 - \frac{1}{2}(\omega \Delta t)^2 \leq 1$

we can solve for Δt and get

$$\Delta t \leq \frac{2}{\omega^2} \quad \text{which corresponds to a parameter in the code.}$$

It seems after adjusting Δt values that when $2\Delta t$ is positive, the higher time difference ($2\Delta t$) does a better modeling job which is expected because the numerical solution should be approaching 0 as the ~~real~~ true solution increases exponentially, but the large Δt slows this from happening. When $2\Delta t < 0$, no true solution was found, and the smaller Δt (2.5) did the best job. Overall, for positive $2\Delta t$, we want high Δt steps, and for negative $2\Delta t$, we want lower (more accurately) Δt steps.

The true solution for $z = -0.5$ shouldn't be 0, but I couldn't get the correct graph to show.

	<u>200</u>	<u>270</u>
2:	worst	best
1.25:	okay	okay
0.5:	best	worst

3. While I could not get graphical data to support this, I know that choosing a higher Δt step will ^{the} cause divergence from true solution. In general, trapezoidal method does a much better job at avoiding amplitude error for the case of $\lambda \Delta t \ll 1$. A phase error is still present, but being able to properly model essentially any frequencies for a given Δt (assuming $\lambda \ll 1$) is powerful.

4. I think the largest differences between these 3 approaches can be identified by considering what they do the best. If we need as little amplification error as possible, Trapezoidal method is the easy choice. If instead we need to amplify our model due to some little signal, then forward time method is the best choice. Damping of course comes from backward time method, and both FT/RT have areas that are good at minimizing amplitude error. FT when our $\lambda \Delta t$ is $\ll 1$ and our frequencies are relatively low. RT on the other hand does better at higher frequencies ^{or time steps} and larger Δt steps provided $\lambda \Delta t \ll 1$ (like FT).

RT and Trapezoidal are perfectly A-stable. FT is not.