

# Electromagnetism II: Electricity

Chapters 3 and 4 of Purcell cover the material presented here, as does chapter 6 of Wang and Ricardo, volume 2. Image charges are covered in more detail in section 3.2 of Griffiths. For an array of interesting physical examples, see chapters II-6 through II-9 of the Feynman lectures. There is a total of **78** points.

## 1 The Method of Images

### Idea 1

The method of images can be used in some highly symmetric situations to compute the electric field in the vicinity of a conductor. Specifically, consider any configuration of static charges and take any equipotential surface containing some of the charges. Then the resulting field configuration outside that surface is the field configuration we would have if that surface bounded a conductor. This is simply because it has constant potential on the conductor surface, so it must be the right answer by the uniqueness theorem.

[4] **Problem 1.** The simplest application of the method of images is the case of a charge  $q$  a distance  $a$  from an infinite grounded conducting plane. This problem explores some of its subtleties, assuming you've already read the basic treatment in section 3.4 of Purcell.

- (a) Find the force on the charge.
- (b) Find the work needed to move the charge out to infinity. (Answer:  $q^2/16\pi\epsilon_0 a$ .)
- (c) Find the total potential energy of the charges on the conducting plane, i.e. the potential energy associated only with their interaction with each other. (Answer:  $q^2/16\pi\epsilon_0 a$ .)
- (d) Now suppose there is another parallel grounded conducting plane on the other end of the charge, a distance  $b$  away. How many image charges are needed now? Draw some of them.
- (e) A conducting plane forces the electric field to be perpendicular to it. Suppose we modified the plane so that the electric field was instead always *parallel* to it. Find the force on the charge.

**Solution.** (a) The field above the plate from the screening charges on the plate can be mimicked by placing an image charge  $-q$  below the plane. This follows because in both cases, the plane of the plate is an equipotential of potential 0. Thus, the force on the charge is  $q^2/(16\pi\epsilon_0 a^2)$  pointing towards the plane.

- (b) We just integrate the expression from before to get  $q^2/(16\pi\epsilon_0 a)$ .

It's tempting to just write down  $q^2/(8\pi\epsilon_0 a)$  because this is the energy of two charges  $\pm q$  separated by  $2a$ , but we must remember that the image charge isn't a real charge. It instead describes the effects of all the screening charges on the plane. You might think it takes work to move the screening charges too, but they can move "for free", since the electric field is perpendicular to the plane throughout this entire process.

- (c) Suppose we freeze the plane's screening charges in place, then move the point charge out to infinity. In this case the image charge is stationary, so the work needed is  $q^2/(8\pi\epsilon_0 a)$ .

Let's think about what this means. There are two components to the initial potential energy: the energy  $U_1$  of the point charge interacting with the screening charges, and the energy  $U_2$  of the screening charges interacting with each other.

In part (b), we showed that it takes work  $q^2/(16\pi\epsilon_0 a)$  to move the point charge out, after which point all charges in the problem are widely separated. So

$$U_1 + U_2 = -\frac{q^2}{16\pi\epsilon_0 a}.$$

By the other argument we just made, if we freeze the screening charges and move the point charge away, we are left with just the screening charges. So

$$U_1 = -\frac{q^2}{8\pi\epsilon_0 a}.$$

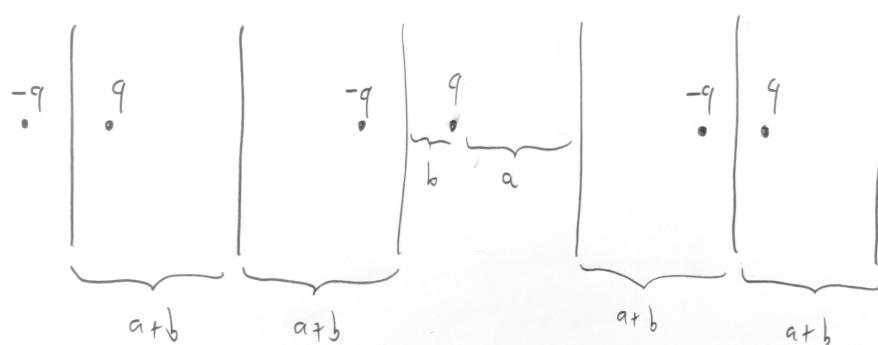
The quantity we're looking for is  $U_2$ , which is thus

$$U_2 = \frac{q^2}{16\pi\epsilon_0 a}.$$

Of course this energy is positive, because the screening charges repel each other.

(At first glance this argument may seem to contradict the statement made in part (b), which states that the screening charges cost zero energy to move, as we move the point charge. However, that statement is only true if the screening charges are always allowed to move, so that they preserve the boundary condition  $E_{\parallel} = 0$ . But above we instead considered an artificial situation where the screening charges were frozen, and this reasoning no longer holds.)

- (d) We actually need infinitely many.



This is the same reason you see infinitely many images of yourself when between two mirrors.

- (e) This boundary condition can be met by placing a charge  $+q$  at  $z = -d$ . Then, the force on the charge would be  $kq^2/(4d^2)$  pointing *away* from the plane.

[4] **Problem 2.** In this problem you'll develop the method of images for spheres.

- (a) A point charge  $-q$  is located at  $x = a$  and a point charge  $Q$  is located at  $x = A$ . Show that the locus of points with  $\phi = 0$  is a circle in the  $xy$  plane, and hence a spherical shell in space.
- (b) What relations must hold among  $q$ ,  $Q$ ,  $a$ , and  $A$  so the center of the sphere is at the origin? In this case, find the radius  $r$  of the sphere, and an expression for  $q$  in terms of  $Q$ ,  $r$ , and  $A$ . These results will be used many times throughout the problem set.
- (c) Now suppose a point charge  $Q$  is a distance  $b$  from the center of a spherical grounded conducting shell of radius  $r$ . Find the force on the charge, considering both the cases  $b < r$  and  $b > r$ .
- (d) The case  $b < r$  is a bit confusing. On one hand, argue that the total charge is  $Q$  plus the image charge, and hence nonzero. On the other hand, argue that the total charge must be zero, by considering an appropriate Gaussian surface. One of these arguments is wrong – which one?

As an aside, the fundamental reason the method of images works for spheres is that electromagnetism has conformal symmetry, a symmetry under any local rescaling of space which preserves angles. (One example of a conformal transformation is inversion in Euclidean geometry.) The setup here is related to the conducting plane by such a transformation.

**Solution.** If you happen to know Euclidean geometry, you can very quickly solve parts (a) and (b) using properties of [Apollonian circles](#). However, for everyone else, we'll present a straightforward solution using coordinates.

- (a) The condition for the potential to vanish is

$$\frac{q}{\sqrt{(x-a)^2 + y^2}} = \frac{Q}{\sqrt{(x-A)^2 + y^2}}.$$

Squaring both sides and clearing denominators, we find

$$(q^2 - Q^2)(x^2 + y^2) + (q^2 A^2 - Q^2 a^2) + 2x(aQ^2 - Aq^2) = 0.$$

This has the form of a conic section, and since the coefficients of  $x^2$  and  $y^2$  are equal, it's a circle. (Strictly speaking, it could also be the empty set, since, for example,  $x^2 + y^2 = -1$  has no solutions. But we know there have to be places where  $\phi = 0$ , because it's positive near the positive charge and negative near the negative charge, so it must cross zero by continuity.)

- (b) The center of the sphere is at the origin if the coefficient of  $x$  vanishes,

$$aQ^2 = Aq^2.$$

Note that this forces  $a$  and  $A$  to have the same sign. Plugging this in and simplifying, we find

$$x^2 + y^2 = Aa$$

from which we conclude the radius is  $r = \sqrt{Aa}$ . If you know some Euclidean geometry, at this point you can also recognize that the two point charges are inversions of each other with respect to the sphere. By combining this with the first equation, conclude

$$q = \frac{Qr}{A}.$$

- (c) We can treat both cases at once. The image charge is a distance  $b' = r^2/b$  from the center of the shell, and its charge is  $q' = -q\sqrt{b'/b} = -qr/b$ . Therefore, the force on the charge is

$$F = \frac{qq'}{4\pi\epsilon_0(b-b')^2} = \frac{q^2rb}{4\pi\epsilon_0(b^2-r^2)^2}.$$

It always attracts the point charge towards the nearest point on the surface of the sphere.

- (d) The second argument is right. Since the conductor shields the details of the charges inside, the field outside the sphere must be spherically symmetric. But we also know the sphere is at zero potential, so the field outside must be exactly zero, so by using a spherical Gaussian surface, the total charge is zero.

On the other hand, the first argument seems to work for all the other examples we've seen: whenever we have an image charge  $Q'$ , it seems to represent the effects of a total charge  $Q'$  on the conducting surface. But in this case, the method of images doesn't determine the charge distribution on the surface, because adding a constant charge density to the sphere's surface creates absolutely *no* field inside. So in reality, this constant density takes whatever value is necessary to make the sphere overall neutral, in accordance with the second argument.

- [2] **Problem 3** (Purcell 3.50). A point charge  $q$  is located a distance  $b > r$  from the center of a *nongrounded* conducting spherical shell of radius  $r$ , which also has charge  $q$ . When  $b$  is close to  $r$ , the charge is attracted to the shell because it induces negative charge; when  $b$  is large the charge is clearly repelled. Find the value of  $b$  so that the point charge is in equilibrium. (Hint: you should have to solve a difficult polynomial equation. You can either use a computer or calculator, or use the fact that it contains a factor of  $1 - x - x^2$ .)

**Solution.** The image charge has value  $-q' = -qr/b$ , and is at position  $r^2/b$ . As of now, the spherical shell has total charge  $-q'$  (by Gauss's law), so to compensate, we add a charge of value  $q + q'$  at the center. Thus, to balance forces we have

$$\frac{q'}{(b - r^2/b)^2} = \frac{q + q'}{b^2}.$$

Defining  $x = r/b$ , this simplifies to

$$x = (1 + x)^3(1 - x)^2$$

which is a rather intimidating quintic equation. At this point you can pull out your calculator, but amazingly, it factors as

$$(1 - x - x^2)(1 + x - x^3) = 0.$$

The only root with  $0 < x < 1$  is from the quadratic,  $x = (\sqrt{5} - 1)/2$ , giving

$$b = \frac{1 + \sqrt{5}}{2} r.$$

- [2] **Problem 4.** Consider a spherical conductor of radius  $R$  placed in a uniform external field  $\mathbf{E}_0$ . By using the method of images, argue that the field created by the sphere is the same as the field of an infinitesimal dipole at the center of the sphere, and find its dipole moment. (Hint: suppose the external field is sourced by two point charges that are very far away.)

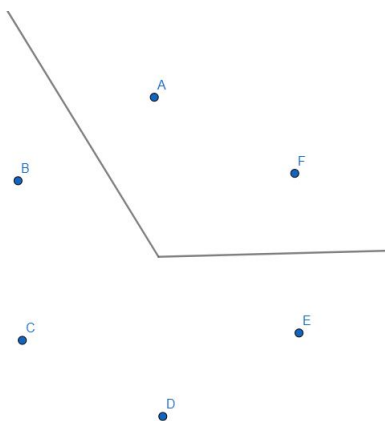
**Solution.** Suppose  $\mathbf{E}_0 = E_0 \hat{\mathbf{z}}$ . Let the conductor be centered at the origin. Now, consider a point charge  $Q$  at  $z = -L$ , and  $-Q$  at  $z = L$  for  $L \gg R$ . This clearly generates a uniform electric field near the origin (for  $r \ll L$ ) with magnitude  $2\frac{Q}{4\pi\epsilon_0 L^2}$ . So we simply choose  $Q = 2\pi\epsilon_0 L^2 E_0$ . Now, the image charges are very close to the origin and form a dipole with moment  $2Q'L' = 2(QR/L)(R^2/L) = \boxed{4\pi\epsilon_0 R^3 E_0}$ .

[3] **Problem 5.** Two grounded conducting half-planes intersect, so that in cylindrical coordinates, the equations describing the planes are  $\theta = 0$  and  $\theta = 2\pi/3$ .

- A charge  $q$  is placed a distance  $r$  from the origin at  $\theta = \pi/3$ . Can the method of images be used to find the force on the charge? If so, do it; if not, explain why not.
- Do the same for  $\theta = \pi/2$ .
- In general, given that the second plane is at  $\theta = \theta_p$  and the charge is at  $\theta = \theta_q$ , when does the method of images work?

**Solution.** (a) It can't be done. When we reflect about the  $\theta = 0$  half plane, we get a point that is collinear with  $\theta = 2\pi/3$ . Thus, the new image charge cannot be further reflected in the plane  $\theta = 2\pi/3$  to get the whole set of image charges.

- It can't be done.



Here,  $A$  is the original charge, and  $B, C, D, E, F$  are image charges with signs alternating (so  $q_B = -q$ ,  $q_C = q$ , and so on). We see that it's impossible to satisfy all the image charge constraints unless we include the image charge  $F$ . But  $F$  is in the same region as  $A$ . That isn't allowed for image charges: the point of image charges is to provide an easy way of calculating the effects of screening charges on conducting surfaces on a given set of *real* charges, so it's not legal to introduce *new* real charges in the process.

- If we reflect the charge in the second plane, we get

$$\theta \rightarrow 2\theta_p - \theta.$$

If we reflect in the first plane again, we get

$$2\theta_p - \theta = \theta - 2\theta_p.$$

That is, in general the composition of two reflections is just a rotation by  $-2\theta_p$ . In order to avoid an infinite number of image charges,  $\theta_p$  must be a rational multiple of  $2\pi$ . If  $\theta_p$  is of the form  $\pi/n$  for integer  $n$ , there is no problem for any  $\theta_q$ . But if  $\theta_p$  is not of this form, then after a sufficient number of these operations we will arrive at an image charge inside the physical region, which is unacceptable. So the procedure works if and only if  $\theta_p = \pi/n$ .

[5] **Problem 6** (Purcell 3.45). [A] Consider a point charge  $q$  located between two parallel infinite grounded conducting planes. The planes are a distance  $\ell$  apart, and the point charge is a distance  $b$  from the left plane. The goal of this problem is to find the total charge induced on each plane.

- (a) Argue that the total charge on each plane would not change if we replaced the point charge  $q$  with two point charges  $q/2$ , both a distance  $b$  from the left plane. By iterating this process, convert the point charge into a uniformly charged plane, and use this to get the answer.
- (b) Alternatively, using image charges, show that the electric field on the inside surface of the left plane, perpendicular to the plane, at a point a distance  $r$  from the axis containing all the image charges, satisfies

$$4\pi\epsilon_0 E_{\perp} = \sum_{n=-\infty}^{\infty} \frac{2q(2n\ell + b)}{((2n\ell + b)^2 + r^2)^{3/2}}.$$

- (c) Since  $\sigma = -\epsilon_0 E_{\perp}$ , we can integrate both sides to find the total charge on the left plane. However, the integral of each term by itself is simply  $q$ , so the series doesn't converge. To get the result, do the following steps in this specific order: group the terms  $\pm n$  together, then integrate them only out to a distance  $R \gg b$ , then sum over the values of  $|n|$ , then take the limit  $R \rightarrow \infty$ . You should get a finite result that matches that of part (a). As you'll probably see in the process, if you do the steps in any other order, you'll get a nonsensical answer.

Those concerned with mathematical rigor might be bothered by the many choices made in part (c). You might ask, couldn't we have gotten a different result by changing how we did the computation? In fact, by the Riemann rearrangement theorem, we could have gotten almost *any* result. But the way we did it is the physically correct way. It roughly sums the terms "in to out", which respects the fact that real plates are finite. Closely related ideas are used to "cancel infinities" in quantum field theory, in a process known as renormalization.

**Solution.** (a) Consider the induced charge distribution for one point charge  $q$ . By the superposition principle, the boundary conditions in this case are also satisfied if we take half that charge distribution, and center it about each of the charges  $q/2$ . By uniqueness, this is the solution.

Extrapolating to infinitely many charges, we get a uniform plane of charge. For the two plates to be at the same voltage, the electric fields above and below the plane have the ratio  $b/(\ell - b)$ . Then the charges have the ratio  $b/(\ell - b)$  and sum to  $-q$ , so the charge on the near plate is  $-q(\ell - b)/\ell$ , while the charge on the far plate is  $-qb/\ell$ .

- (b) We do this by summing over image charges. As we can see from the figure in the solution to problem 1, there are infinitely many image charges. The  $n = 0$  term in the sum corresponds to the image charge and real charge closest to the left plane. The  $n = 1$  term corresponds to the two image charges a distance  $2\ell + b$  from the left plane, while the  $n = -1$  term corresponds to the image charges a distance  $2\ell - b$  from the left plane, and so on.

- (c) Applying Gauss's law,  $\sigma = \epsilon_0 E_\perp$ , and grouping the terms  $\pm n$  together, we have a total charge on the left plane of

$$\begin{aligned} Q &= \int_0^\infty (-\epsilon_0 E_\perp) \cdot 2\pi r \, dr \\ &= q \int_0^\infty \left[ -\frac{br}{(b^2 + r^2)^{3/2}} + \sum_{n=1}^\infty \left( \frac{(2n\ell - b)r}{((2n\ell - b)^2 + r^2)^{3/2}} - \frac{(2n\ell + b)r}{((2n\ell + b)^2 + r^2)^{3/2}} \right) \right] dr. \end{aligned}$$

The first term integrates to 1, so we will deal with the sum. Consider one term for some given  $n$ , and say we integrate out to some finite but large  $R$ . Note that we can only assume  $b \ll R$ , not  $n\ell \ll R$ . We will drop all factors of  $b^2$ . The term integrates out to

$$\begin{aligned} & -\frac{2n\ell - b}{\sqrt{4n^2\ell^2 + R^2 - 4n\ell b}} + \frac{2n\ell + b}{\sqrt{4n^2\ell^2 + R^2 + 4n\ell b}} \\ & \approx -\frac{2n\ell - b}{\sqrt{R^2 + 4n^2\ell^2}} \left( 1 + \frac{1}{2} \frac{4n\ell b}{R^2 + 4n^2\ell^2} \right) + \frac{2n\ell + b}{\sqrt{R^2 + 4n^2\ell^2}} \left( 1 - \frac{1}{2} \frac{4n\ell b}{R^2 + 4n^2\ell^2} \right) \\ & = \frac{1}{\sqrt{R^2 + 4n^2\ell^2}} \left( b - 2n\ell + 2n\ell + b + \frac{-2n\ell b(2n\ell - b) - 2n\ell b(2n\ell + b)}{R^2 + 4n^2\ell^2} \right) \\ & = \frac{2R^2 b}{(R^2 + 4n^2\ell^2)^{3/2}} \\ & = \frac{2b}{R} \frac{1}{(1 + (2n\ell/R)^2)^{3/2}}. \end{aligned}$$

Since  $R \rightarrow \infty$ , this sum in this limit can be written as an integral, giving

$$\frac{2b}{R} \int_0^\infty \frac{1}{(1 + (2n\ell/R)^2)^{3/2}} dn = \frac{2b}{R} \frac{R}{2\ell} \int_0^\infty \frac{dx}{1 + x^2} = \frac{b}{\ell}.$$

Therefore, the total charge is simply  $\boxed{-q(1 - b/\ell)}$ .

If you found the analysis in part (c) quite tricky, know that you're in good company. About 10 papers have been written about this exact system in the American Journal of Physics alone!

## 2 Capacitors

### Idea 2

For a single, isolated conductor with charge  $Q$ , the self-capacitance is defined as

$$Q = C\phi$$

where  $\phi$  is the potential difference between the conductor and infinity. For a set of two isolated conductors with charges  $\pm Q$ , the capacitance is defined as

$$Q = C\phi$$

where  $\phi$  is the potential difference between the two conductors.

**Idea 3**

The definitions of  $C$  above are only useful when you have only one or two conductors in the problem, respectively. In a situation with more than two, it's very tricky to use the above definitions, because all the conductors will affect each other; even a neutral conductor will have an effect since there will be induced charges on its surface.

Instead, it's better to revert to more general principles. The underlying principle behind capacitance is linearity: the charges are linearly related to the fields. For multiple capacitors, the most general possible linear relation is

$$Q_i = \sum_j C_{ij} \phi_j$$

where conductor  $i$  has charge  $Q_i$  and potential  $\phi_i$ , the potential is taken to be zero at infinity, and the  $C_{ij}$  are called general coefficients of capacitance, or in electrical engineering, the Maxwell capacitance matrix. Similarly, inverting this relation,

$$\phi_i = \sum_j p_{ij} Q_j$$

where the  $p_{ij}$  are called coefficients of potential.

In Olympiad physics, you'll almost never want to compute the  $C_{ij}$  or  $p_{ij}$  explicitly. Instead, the point here is that if you're given the charges and want the potentials, or vice versa, you can build up the answer you want using the principle of superposition, computing all the fields you need, e.g. using Gauss's law.

**Remark**

General capacitance coefficients are discussed further in section 3.6 of Purcell. One nontrivial fact is that  $C_{ij} = C_{ji}$ , which is proven by energy conservation in problem 3.64 of Purcell. Capacitance coefficients can be clunky to work with. For example, suppose you want to compute the familiar capacitance of a system of two conductors. By definition, we have

$$Q_1 = C_{11}\phi_1 + C_{12}\phi_2, \quad Q_2 = C_{21}\phi_1 + C_{22}\phi_2.$$

An ordinary two-plate capacitor corresponds to the special case of opposite charges on the plates, so we write  $Q = Q_1 = -Q_2$ . There is a potential difference  $V$  across the plates, so  $\phi_1 = \phi_2 + V$ , and plugging this in gives

$$Q = (C_{11} + C_{12})\phi_2 + C_{11}V, \quad -Q = (C_{22} + C_{21})\phi_2 + C_{21}V.$$

Eliminating  $\phi_2$  from the system of equations above, we find the familiar capacitance

$$C = \frac{Q}{V} = \frac{C_{11}C_{22} - C_{12}^2}{C_{11} + C_{22} + 2C_{12}}$$

where we used  $C_{12} = C_{21}$ . This is quite an inconvenient formula, so as a result we won't consider general capacitance coefficients any further, except briefly for practice in problem 8.



- [2] **Problem 7** (Purcell 3.21). Consider a capacitor made of four parallel plates with large area  $A$ , evenly spaced with small separation  $s$ . The first and third are connected by a wire, as are the second and fourth. What is the capacitance of the system?

**Solution.** Say we charge the conductors to equal and opposite charges. Then by symmetry, the surface charges are

$$\sigma_1, -\sigma_2, \sigma_2, -\sigma_1$$

reading left to right (1 to 4). The field in between the first and the second plates is  $\sigma_1/\epsilon_0$ , and the field between the second and the third plates is  $(\sigma_1 - \sigma_2)/\epsilon_0$ . Since 1 and 3 are connected, and 2 and 4 are, we must have that

$$\sigma_1 s = (\sigma_2 - \sigma_1)s \implies \sigma_2 = 2\sigma_1.$$

Thus, the potential difference is  $\sigma_1 s/\epsilon_0$ , so the capacitance is  $C = Q/\phi = 3\sigma_1 A/(\sigma_1 s/\epsilon_0) = 3\epsilon_0 A/s$ .

- [3] **Problem 8.** Consider two concentric spherical metal shells, with radii  $a < b$ .

- Compute their capacitance using Gauss's law.
- Compute their capacitance by computing the four capacitance coefficients, verifying that  $C_{12} = C_{21}$  along the way, and using the result for  $C$  above.

**Solution.** (a) Let the shells have charge  $\pm Q$ . The field between the shells is  $(Q/4\pi\epsilon_0 r^2)\hat{r}$ , so

$$V = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{b} \right).$$

Thus the capacitance is

$$C = \frac{Q}{V} = 4\pi\epsilon_0 \frac{ab}{b-a}.$$

- Let the first capacitor be the inner shell. If only the outer shell is charged, with charge  $Q$ , then  $\phi_1 = \phi_2 = Q/4\pi\epsilon_0 b$ . The general capacitance equations in this case are

$$0 = C_{11}\phi_1 + C_{12}\phi_2, \quad Q = C_{21}\phi_1 + C_{22}\phi_2$$

from which we see that

$$C_{11} + C_{12} = 0, \quad C_{21} + C_{22} = 4\pi\epsilon_0 b.$$

Now suppose only the inner shell is charged, with charge  $Q$ . In this case we have  $\phi_1 = Q/4\pi\epsilon_0 a$  while  $\phi_2 = Q/4\pi\epsilon_0 b$ , so

$$Q = C_{11}\phi_1 + C_{12}\phi_2, \quad 0 = C_{21}\phi_1 + C_{22}\phi_2$$

from which we see that

$$\frac{C_{11}}{a} + \frac{C_{12}}{b} = 4\pi\epsilon_0, \quad \frac{C_{21}}{a} + \frac{C_{22}}{b} = 0.$$

Solving these four equations for the capacitance coefficients gives

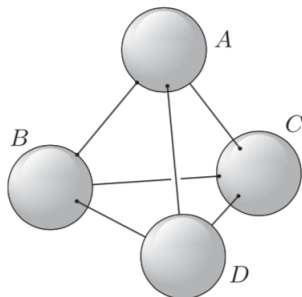
$$C_{11} = 4\pi\epsilon_0 \frac{ab}{b-a}, \quad C_{12} = C_{21} = -4\pi\epsilon_0 \frac{ab}{b-a}, \quad C_{22} = 4\pi\epsilon_0 \frac{b^2}{b-a}.$$

Plugging into the general formula, we have

$$C = \frac{4\pi\epsilon_0}{b-a} \frac{(ab)b^2 - (ab)^2}{ab + b^2 - 2ab} = \frac{4\pi\epsilon_0}{b-a} \frac{b^2 a(b-a)}{b(b-a)} = 4\pi\epsilon_0 \frac{ab}{b-a}.$$

This is certainly a longer route to get to the same conclusion!

- [3] **Problem 9** (MPPP 152). Four identical non-touching metal spheres are positioned at the vertices of a regular tetrahedron, as shown.



A charge  $4q$  given to sphere  $A$  raises it to a potential  $V$ . Sphere  $A$  can also be raised to potential  $V$  if it and one of the other spheres are each charged with  $3q$ . What must be the size of equal charges given to  $A$  and two other spheres for the potential of  $A$  to again be raised to  $V$ ? What if all four spheres are used?

**Solution.** Suppose sphere  $A$  has charge  $4y + 3x + 3x$ , and spheres  $B$  and  $C$  have charges of  $3x$ . We can write this a superposition of three cases,  $(4y, 0, 0, 0) + (3x, 3x, 0, 0) + (3x, 0, 3x, 0)$ . The potential at  $A$  is then just  $(y/q)V + 2(x/q)V$  because of linearity (one can use the capacitance coefficients to formalize this). Therefore, we want to set  $V = (y/q) + 2(x/q)V$ , so  $2x + y = q$ . We also want  $4y + 3x + 3x = 3x$ , so  $y = -3x/4$ . Thus,  $2x - 3x/4 = q$ , so  $x = 4q/5$  and  $y = -3q/5$ . Thus, the charges are  $\boxed{12q/5}$ .

We now consider the case of all four spheres being used. Suppose sphere  $A$  has charge  $4y + 3x + 3x + 3x$ , and all others have charge  $3x$ . The potential at  $A$  is then  $(y/q)V + 3(x/q)V = V$ , so  $3x + y = q$ . Also,  $4y + 9x = 3x$ , so  $y = -3x/2$ . Thus,  $3x - 3x/2 = q$ , so  $x = 2q/3$  and  $y = -q$ . Thus, the charge is  $3x = \boxed{2q}$ .

- [3] **Problem 10.** ⌚ USAPhO 2008, problem A1.

#### Idea 4

The energy stored in a capacitor is

$$U = \frac{1}{2}QV = \frac{1}{2}CV^2$$

where  $V$  is the voltage difference and  $C$  is the (mutual) capacitance. The mutual capacitance  $C$  adds in parallel, while  $1/C$  adds in series.

- [2] **Problem 11** (Purcell 3.24). Some estimates involving capacitance.

- Estimate the capacitance of the Earth.
- Make a rough estimate of the capacitance of the human body.
- By shuffling over a nylon rug on a dry winter day, you can easily charge yourself up to a couple of kilovolts, as shown by the length of the spark when your hand comes too close to a grounded conductor. How much energy would be dissipated in such a spark?

**Solution.** (a) The Earth is a sphere of radius of order  $10^7$  m, so

$$C = 4\pi\epsilon_0 r \sim 10^{-3} \text{ F}.$$

Notably, we can make “bigger” capacitors in the lab! Nonetheless, a huge amount of charge can be delivered to the Earth, such as by lightning strikes. This is because the voltage of the Earth is also huge, which is possible because its huge size means the corresponding electric fields aren’t that big.

(b) A human is approximately a sphere of radius 0.5 m. Then,  $C = 4\pi\epsilon_0 r \sim 5 \times 10^{-11} \text{ F}$ .

(c) We have  $U = \frac{1}{2}CV^2 \sim 10^{-4} \text{ J}$ .

[2] **Problem 12.** The total energy can also be found by integrating the electric field energy,

$$U = \frac{\epsilon_0}{2} \int E^2 dV.$$

(a) Show that this agrees with  $U = CV^2/2$  for a parallel plate capacitor.

(b) Show that this agrees with  $U = CV^2/2$  for a capacitor made of concentric spheres.

The general proof is more advanced, but if you’re interested, one slick method is given in problem 1.33 of Purcell.

**Solution.** (a) Let the plate area be  $A$  and the distance between them be  $d$ . Then

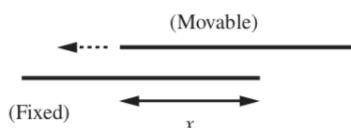
$$U = \frac{\epsilon}{2} E^2 (Ad) = \frac{\sigma^2}{2\epsilon_0} Ad = \frac{C}{2} \frac{\sigma^2 d^2}{\epsilon_0^2} = \frac{CV^2}{2}.$$

(b) Let the radii be  $R_1$  and  $R_2$  and the charges be  $\pm Q$ . The field is  $Q/(4\pi\epsilon_0 r^2)$ , so

$$U = \frac{\epsilon_0}{2} \int \frac{Q^2}{16\pi^2\epsilon_0^2} \frac{dV}{r^4} = \frac{Q^2}{32\pi^2\epsilon_0} \int_{R_1}^{R_2} \frac{4\pi r^2 dr}{r^4} = \frac{Q^2}{8\pi\epsilon_0} \left( \frac{1}{R_1} - \frac{1}{R_2} \right).$$

On the other hand, this should be equal to  $U = QV/2$ , which follows directly from the result of problem 8.

[3] **Problem 13** (Purcell 3.26). A parallel-plate capacitor consists of a fixed plate and a movable plate that is allowed to slide in the direction parallel to the plates. Let  $x$  be the distance of overlap.



The separation between the plates is fixed. Let  $C(x)$  be the capacitance.

(a) Assume the plates are electrically isolated, so that their charges  $\pm Q$  are constant. By differentiating the energy, find the leftward force on the movable plate in terms of  $Q$  and  $C(x)$ .

(b) Now assume the plates are connected to a battery, so that their potential difference  $\phi$  is held constant. Find the leftward force on the movable plate, in terms of  $\phi$  and  $C(x)$ .

- (c) If the movable plate is held in place, the two answers above should be equal because nothing is moving. Verify that this is the case, being careful with signs.
- (d) In terms of electric fields, why is there a force on the movable plate? Does the effect invoked in the answer to this part change the conclusion of parts (a) through (c) at all?

**Solution.** (a) The energy as a function of  $x$  is

$$U(x) = \frac{Q^2}{2C}$$

where we understand that  $C$  is also a function of  $x$ . Thus, the force on the plate is

$$F = -\frac{dU}{dx} = \frac{Q^2}{2} \frac{d}{dx} \left( -\frac{1}{C} \right) = \frac{Q^2}{2C^2} \frac{dC}{dx}.$$

- (b) Here, the energy is  $U(x) = \frac{1}{2}C\phi^2$ , so naively we have

$$F = -\frac{dU}{dx} = -\frac{\phi^2}{2} \frac{dC}{dx}.$$

This is negative, while the answer to part (a) is positive. The reason is that  $U$  should reflect the total energy of the system – and in this case, the system must include the battery that does work to maintain the potential difference  $\phi$ .

Say  $x$  increases by  $dx$ . Let the change in capacitance be  $dC$ . Then,  $dQ = \phi dC$ . Thus, the work the battery does is

$$dW = \phi dQ = \phi^2 dC.$$

Therefore, if  $F$  is the net force the plate feels, we have that

$$dW = F dx + dU \implies F = \frac{1}{2}\phi^2 \frac{dC}{dx}.$$

- (c) Let  $F_Q$  be the first force, and  $F_\phi$  the second. We have that

$$F_Q/F_\phi = \frac{Q^2}{\phi^2 C^2} = 1.$$

If we didn't account for the subtlety in part (b), we would have gotten  $-1$  here.

- (d) The force is due to the fringe fields, as the non-fringe field is perfectly vertical. Thus, we have managed to compute a fringe field effect using energy, even though we were able to completely ignore the fringe fields in the energy calculation! This is yet another example of how conservation laws can hand you information that's very hard to get otherwise.

### Idea 5: Dielectrics

While we're on the subject of capacitors, it's useful to introduce dielectrics, which will be important later. All electrical insulators are dielectrics. In the presence of an electric field, a dielectric will polarize, with positive charges displaced slightly along the field. The resulting electric dipoles distributed throughout the material in turn create a field that tends to

weaken the original applied field.

Each part of a dielectric polarizes based on the local electric field, but that electric field depends on the applied field, and the polarization of *every other* piece of the dielectric. Thus, solving for the electric field for a general dielectric geometry is very difficult, and usually not possible in closed form, just like how it's usually not possible to solve for the field of a charged conductor. In Olympiad physics, you will almost always consider highly symmetric situations, where a dielectric simply reduces the applied electric field everywhere by a factor of  $\kappa$ , called the dielectric constant. (We'll consider some trickier situations in **E8**.)

Consider a parallel plate capacitor with charge  $\pm Q$  on each plate. If a dielectric is inserted with the charge kept the same, then the field inside is reduced by a factor of  $\kappa$ . Thus, the capacitance  $C = Q/V$  increases by a factor of  $\kappa$ . Dielectrics may increase the amount of energy that can be stored in a capacitor, which is typically limited by the voltage  $V_0$  where electrical breakdown occurs. So if  $V_0$  stays the same, the maximal stored energy  $U = CV_0^2/2$  goes up by a factor of  $\kappa$ .

Plugging in the definition of  $C$ , this result implies that the energy density in the capacitor is  $\kappa\epsilon_0 E^2$ . But we showed in **E1** that the energy density of the electric field is only  $\epsilon_0 E^2$ . The extra energy is stored in the dielectric material itself: it takes energy to separate positive and negative charges within the dielectric, as if we were stretching many microscopic springs. This potential energy is released when the capacitor is discharged.

### 3 Tricky Problems

#### Example 1: PPP 151

A closed body with conducting surface  $F$  has self-capacitance  $C$ . The surface is now dented so that the new surface  $F^*$  is entirely inside  $F$ . Prove that the capacitance has decreased.

#### Solution

The energy stored in the capacitor is  $U = Q^2/2C$ . Therefore, if we give the capacitor a fixed charge  $Q$ , proving that  $F^*$  has lower  $C$  is equivalent to showing that it takes positive work to dent the foil from  $F$  to  $F^*$ . It's cleaner to show the other direction, i.e. that starting from  $F^*$ , we can get to  $F$  while only lowering the energy.

Suppose without loss of generality that  $F$  is infinitesimally larger than  $F^*$ . (We can break any finite change into infinitesimal stages and repeat this argument.) We can go from  $F^*$  to  $F$  by just taking each charge on the surface and moving it outward until it hits  $F$ . This lowers the energy because the electric field is always directed outward, as we proved in **E1**.

At this point, the charges lie on  $F$ , but they don't have the right distribution, i.e.  $F$  is not an equipotential. Now we let the charges spontaneously redistribute themselves so that  $F$  is

again an equipotential. This again lowers the energy, proving the desired result.

### Example 2

Are there charge distributions that aren't spherically symmetric, but which produce an *exact*  $\hat{\mathbf{r}}/r^2$  field outside of them?

### Solution

If you know a bit about the multipole expansion, this might seem like a daunting question. To make the field exactly  $\hat{\mathbf{r}}/r^2$ , you need to make sure the charge distribution has no dipole moment, no quadrupole moment, no octupole moment, and so on to infinity, and it seems impossible to satisfy all of these constraints without spherical symmetry. But we have *already* seen an example of such a charge distribution earlier in the problem set!

Recall that when we treated the method of images for spheres, we found that in some situations, the complicated charge densities on conducting spheres were exactly the same as those produced by a fictitious image charge inside the sphere, and generally away from its center. If we place the origin at that image charge, then we have an example of a charge distribution that is perfectly  $\hat{\mathbf{r}}/r^2$  far away, but which isn't spherically symmetric. (The general solution is given [here](#).)

- [2] **Problem 14.** Consider a set of  $n$  conducting, very large parallel plates, placed in zero external electric field. The plates are given charges  $Q_i$ . If the left ends of the plates are at locations  $x_i$ , and the plates have thickness  $d_i$ , what is the total charge on the left end of the leftmost plate, and the right end of the rightmost plate?

**Solution.** If you try to solve this one directly, the calculations can get pretty messy, but there's a simple solution using basic facts about capacitors. First, consider the two leftmost plates. Since they're conductors, the electric field must vanish inside them. Now split the charge in the problem into two parts: (1) the charge on the right side of the leftmost plate, and the left side of the second leftmost plate, and (2) all other charge.

The electric field due to (2) has some uniform value  $E_{\text{ext}}$  within the first and second plates. Therefore, in order for the electric field to vanish in both of those plates, the electric field due to (1) has to be the *same* in both plates. This is only possible if the two charges in (1) are *opposite*, which corresponds to those sides forming a parallel plate capacitor.

We can then repeat this reasoning, which shows that the charges on the right side of plate  $i$  must cancel with the charges on the left side of plate  $i + 1$ , producing no field anywhere except in between plates  $i$  and  $i + 1$ . All that is left is the left end of the leftmost plate and the right end of the rightmost plate, which must have total charge  $\sum_i Q_i$ . Finally, in order to ensure no field in any plate, these two charges must be equal, so the answer is that both have total charge  $(\sum_i Q_i)/2$ .

- [2] **Problem 15** (Purcell 3.9). A conducting spherical shell has charge  $Q$  and radius  $R_1$ . A larger concentric conducting spherical shell has charge  $-Q$  and radius  $R_2$ .

(a) If the outer shell is grounded, explain why nothing happens to the charge on it.

- (b) If instead the inner shell is grounded, e.g. by connecting it to ground by a very thin wire that passes through a very small hole in the outer shell, find its final charge.
- (c) It's not so clear why charge would leave the inner shell in part (b), thinking in terms of forces. A small bit of positive charge will certainly want to hop on the wire and follow the electric field across the gap to the larger shell. But when it gets to the larger shell, it seems like it has no reason to keep going to infinity, because the field is zero outside. And, even worse, the field will point *inward* once some positive charge has moved away from the shells. So it seems like the field will drag back any positive charge that has left. Does charge actually leave the inner shell? If so, what's wrong with the above reasoning?

**Solution.** (a) The potential at the outer shell due to itself is  $-Q/4\pi\epsilon_0 R_2$  and the potential due to the inner shell is  $Q/4\pi\epsilon_0 R_2$ , so it is zero overall. Thus, the outer shell is already effectively grounded.

- (b) The potential at the inner shell due to itself is  $Q'/4\pi\epsilon_0 R_1$  and the potential to the outer shell is  $-Q/4\pi\epsilon_0 R_2$ . Since the total must be zero,  $Q' = R_1 Q/R_2$ .
- (c) The key mistake is that the positive charges are repelled also by the charges behind it in the wire. So yes, eventually the field due to the shells may even become inward, there is a whole line of plus charge behind a given charge that force it forward.

Another way of saying this is that a wire has negligible capacitance; like a thin metal pipe of water, it cannot store extra net charge but can only let charge move rigidly through the entire thing. It is energetically favorable for this to happen, so even if some charges don't want to move forward, their neighbors will push them forward.

- [2] **Problem 16.** The usual expression for the capacitance of a parallel-plate capacitor is  $A\epsilon_0/d$ . However, in reality the field within the capacitor is not perfectly uniform, and there are fringe fields outside. Taking these effects into account, is the true capacitance higher or lower than  $A\epsilon_0/d$ ?

**Solution.** For concreteness, suppose the plates are circular disks, with charge  $\pm Q$ , and consider the electric field along the axis of symmetry. In the naive derivation, we assume the charge density on the plates is uniform. Then we approximate the plates as infinite in order to use Gauss's law to conclude that the field inside is  $\sigma/\epsilon_0 = Q/A\epsilon_0$ . This is inaccurate for two reasons:

- The plates are not actually infinite, so the field on the symmetry axis should actually be smaller.
- The charge distribution is not actually uniform. Instead, since the like charges on each plate repel each other, some charge gets pushed outward. This further decreases the field on the symmetry axis.

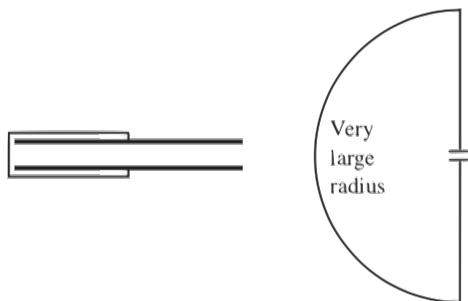
Therefore, the field on the symmetry axis is less than  $\sigma/\epsilon_0$ , so the voltage drop is lower, which implies the capacitance is higher.

- [2] **Problem 17** (Purcell 4.16). In a parallel plate capacitor, the quantity  $\int \mathbf{E} \cdot d\mathbf{s}$  should be equal to  $V$  for any path that connects the two plates.

A charged capacitor can be discharged by attaching a wire to the external surfaces of the plates. No matter how one attaches the wire,  $\int \mathbf{E} \cdot d\mathbf{s}$  along the wire should be equal to  $V$ . And as we've argued in **E2**, this is sufficient to cause charges to move along the wire, even if the electric field

points in the “wrong” direction at some points along the wire, because the wire has negligible capacitance: charges within it move rigidly, each pushing the next one and pulling the previous one.

But it’s puzzling how this works for a capacitor, because the electric field is supposed to be essentially zero just outside it. Consider two possible limiting cases for the wire’s shape.



In each case, explain qualitatively how  $\int \mathbf{E} \cdot d\mathbf{s}$  can be equal to  $V$ . In particular, how large are the contributions from the distinct segments of the wire (the horizontal and vertical parts in the first case, and the straight and curved parts in the second)?

**Solution.** In the first case, the horizontal parts of the wire contribute almost nothing. That’s because the radial part of the electric field vanishes within the conductor plates themselves (since they must be equipotentials), and the horizontal path is right next to the plates. Therefore, the contribution is almost entirely from the vertical segment.

You might be wondering how this is possible, because at the edge of the plates, there seems to be less charge nearby, so the electric field should be smaller than near the middle of the plates. The resolution is that the surface charge density near the edge of plates is a lot higher than the surface charge density near the middle, because like charges repel.

It’s interesting to compare this to the case of two parallel plates with *uniform* charge density. In that case, the vertical segment contributes roughly  $V/2$ . To see why, consider putting a second, identical parallel plate capacitor directly to the left of the first one. Now the vertical segment is in the middle of a big capacitor, and has voltage drop  $V$ . So each of the two halves of that big capacitor contributes  $V/2$  to the vertical segment. However, the total voltage drop is still  $V$ , because for uniform charge density there are substantial horizontal fields, so that the horizontal segments contribute roughly  $V/4$  each.

In the second case, the result is due to the far-field behavior. When you zoom out, the capacitor looks like a dipole, so the field at long distances is a dipole field. Now, the dipole field falls off as  $1/r^3$ , and the circumference of the curved part is proportional to  $r$ , so the contribution of this part goes as  $r/r^3 \rightarrow 0$  as  $r \rightarrow \infty$ . So all the contribution is from the straight part.

To see how this can be the case, note that the vertical field just above the capacitor plates is negligible; the dipole field only kicks in once we’re far enough so that the plates look small, i.e. subtending a small angle from our perspective. If the plates are squares of side length  $a$ , this occurs at a distance of order  $a$ . Then

$$\int \mathbf{E} \cdot d\mathbf{s} \sim 2 \int_a^\infty \frac{p}{2\pi\epsilon_0 r^3} dr = \frac{2}{\pi} \frac{p}{\epsilon_0 a^2}$$

where  $p$  is the electric dipole moment. If the plates are separated by a distance  $d \ll a$ , then  $p = Qd = \sigma a^2 d$ , giving

$$\frac{2}{\pi} \frac{p}{\epsilon_0 a^2} = \frac{2}{\pi} \frac{\sigma d}{\epsilon_0} = \frac{2}{\pi} V$$



which is on the order of the voltage across the capacitor plates. Of course, we didn't get precisely  $V$  because we made a lot of approximations in the calculation, but this illustrates the conceptual point: the full integral of  $\mathbf{E} \cdot d\mathbf{s}$  can indeed be equal to  $V$ , and most of the contribution to this integral comes from the part of the vertical wire which is a distance of order  $a$  from the capacitor.

[3] **Problem 18.** Consider two conducting spheres of radius  $r$  separated by a distance  $a \gg r$ . The spheres can be thought of as the two plates of a parallel plate capacitor.

- (a) By assuming the surface charge density on each sphere is uniform, estimate the capacitance.
- (b) In reality, the surface charges on each sphere will be distorted by the other sphere. The surface charges assumed in part (a) will induce changes in the surface charges, represented by an image charge for each sphere. But these image charges will induce further image charges, and so on, yielding an infinite series for the charge distribution and hence the capacitance. Find the first two terms in the series for the capacitance.
- (c) If  $a/r = 10$ , roughly estimate the error in neglecting the other terms.

**Solution.** (a) The potential of a single sphere, relative to infinity, is  $V = Q/4\pi\epsilon_0 r$ . Hence the two spheres, at this level of approximation, have potentials  $\pm Q/4\pi\epsilon_0 r$ , and the capacitance is

$$C = \frac{Q}{V_1 - V_2} = 2\pi\epsilon_0 r.$$

This is ignoring any interaction between the charges on different spheres.


- (b) The uniform charge on one sphere looks like a point charge to the other sphere, by the shell theorem, and this induces an image charge. The same applies for the spheres in reverse. The image charges are smaller than the original ones by a factor of  $r/a$ , by the result of problem 2. Now, these “first-order” image charges on each sphere in turn induce second-order image charges in the other sphere, which are smaller by another factor of  $r/a$ , and so on. Once we sum up an infinite series of image charges, we get the true charge configuration, from which we can compute the exact capacitance, which is now a power series in  $r/a$ .

In this part we're only interested in the first two terms. We found the first term in part (a), and the next term comes from accounting for the “first-order” image charges. Let the spheres have charges  $\pm Q$ . The negative sphere induces an image charge of  $Qr/a$  within the positive sphere, so that the negative sphere plus the image charge yield a total of zero potential on the positive sphere. So the total potential of the positive sphere is still  $V = Q/4\pi\epsilon_0 r$ , except that the total charge on the sphere is not  $Q$ , but  $Q(1 + r/a)$ . Applying the same reasoning in reverse to the other sphere,

$$C = \frac{Q(1 + r/a)}{V_1 - V_2} = 2\pi\epsilon_0 r(1 + r/a + O((r/a)^2)).$$

It makes sense the true capacitance is higher than the naive guess of part (a), because in part (a) we demanded a uniform charge density, and charges can lower their energy further if we let them spread out.

- (c) As we've seen, the answer is a power series in  $r/a$ . So if we include the zeroth and first order terms, the second order and higher terms are suppressed by at least two powers,  $(r/a)^2 = 1\%$ , so we expect our approximation to be good within a few percent.

[3] **Problem 19.**  USAPhO 2022, problem A2. A computational problem involving surface tension.

**Example 3**

Find the leading interaction force between a dipole of dipole moment  $p$  and a grounded conducting sphere of radius  $r$ , separated by a distance  $R \gg r$ . What if the sphere is electrically neutral instead?

**Solution**

Place the origin at the center of the sphere and orient the  $z$ -axis to pass through the dipole. We can regard the dipole  $p = qd$  as a combination of two charges

$$q \text{ at } z = R, \quad -q \text{ at } z = R + d$$

where  $d$  is very small. In the grounded case, this induces two image charges in the sphere,

$$\frac{qr}{R} \text{ at } z = \frac{r^2}{R}, \quad -\frac{qr}{R+d} \text{ at } z = \frac{r^2}{R+d}$$

approximately separated by  $r^2/R^2$ . We can now use Coulomb's law four times, but that's a bit tedious. Instead, decompose the image charges into a dipole moment and a net charge,

$$p' = \frac{pr^3}{R^3}, \quad Q' = \frac{qr}{R} - \frac{qr}{R+d} \approx \frac{pr}{R^2}.$$

We can place both of these at the origin, because this slight displacement will only affect the answer by subleading terms in  $r/R$ . Then the corresponding fields, far along the  $z$ -axis, are

$$E_{p'}(z) = \frac{2kpr^3}{R^3 z^3}, \quad E_{Q'}(z) = \frac{kpr}{R^2 z^2}.$$

The first term is negligible compared to the first, due to the many powers of  $R$  and  $z$  in the denominator. Thus, keeping only the second term, the force on the original dipole is

$$F = p \frac{d}{dz} E(z) \Big|_{z=R} = -\frac{2kp^2 r}{R^5}$$

which falls off very quickly with distance. This derivation illustrates a common subtlety: it might not always be obvious how far to approximate. We threw away terms subleading in  $r/R$ , because we only wanted the leading contribution. But if we had applied that principle to the image charges at the first step, we would have thrown out the tiny net charge  $Q'$ , which actually provides the dominant contribution to the force, because of how tiny  $p'$  is.

Now, the situation for a neutral sphere is completely different. By the logic of problem 3, there's a third image at the center of the sphere to enforce neutrality,

$$-\frac{pr}{R^2} \text{ at } z = 0.$$

The image charges can now be decomposed into a combination of two dipole moments. We already saw the first one  $p'$  above, while the second is, to leading order

$$p'' \approx \frac{pr}{R^2} \frac{r^2}{R} = \frac{pr^3}{R^3}$$

with the same magnitude and direction as  $p'$ . Thus, this system of image charges has approximate dipole moment  $2p'$ . The corresponding force is

$$F = p \frac{d}{dz} \frac{4kp r^3}{R^3 z^3} \Big|_{z=R} = -\frac{12kp^2 r^3}{R^7}$$

which falls off even more quickly with distance. In this derivation, we didn't have to worry too much about getting  $p''$  exactly right, because there was no net charge ("monopole") term that could've overwhelmed the dipole field, so all other field contributions are safely suppressed by more powers of  $r/R$ . (Of course, if  $p''$  had come out pointing the opposite direction to  $p'$ , so that the two almost cancelled, we would've had to be more careful.)

The lesson of this example is *not* to just use exact expressions and Taylor expand at the end. Here, that brute force approach would have required Taylor expanding six Coulomb's law forces out to order  $1/R^7$ , which is extraordinarily tedious. Instead, to approximate properly, we have to think carefully in every case.

#### Example 4

Estimate the interaction force between a point charge  $q$  and a thin conducting rod of length  $\ell$ , which is a distance  $L \gg \ell$  from the charge and oriented along the separation between them.

#### Solution

The interaction occurs because the point charge induces negative charges on the near end of the rod, and positive charges on the far end. These charges are then acted on by the electric field of the point charge, causing a force.

To get a very crude estimate, let's just suppose that charge  $Q$  appears on the far end and charge  $-Q$  appears on the near end. The resulting field produced in the middle is

$$E \sim \frac{kQ}{\ell^2}.$$

On the other hand, this needs to cancel a field from the point charge of

$$E \sim \frac{kq}{L^2}$$

which tells us that  $Q \sim (\ell/L)^2 q$ . The force on the induced charges is

$$F \sim kqQ \left( \frac{1}{L^2 + \ell^2} - \frac{1}{L^2} \right) \sim -\frac{kqQ\ell^2}{L^4} \sim -\frac{kq^2\ell^4}{L^6}.$$

Again, the force is attractive, and falls off quickly with distance.

- [3] **Problem 20** (Physics Cup 2017). Estimate the interaction force between a point charge  $q$  and an infinitely thin circular neutral conducting disc of radius  $r$  if the charge is at the axis of the disc, and the distance between the disc and the charge is  $L \gg r$ .

**Solution.** The interaction is because charges redistribute on the disc to keep it an equipotential. As an extremely rough approximation, suppose that charge  $Q$  appears at the center of the disc and charge  $-Q$  appears on the rim. Then essentially by dimensional analysis, the electric field in the disc is

$$E \sim \frac{kQ}{r^2}.$$

On the other hand, the electric field due to the point charge along the disc is of order

$$E \sim \frac{kQ}{L^2} \frac{r}{L}$$

where the  $r/L$  factor is from projecting the field along the disc. Then

$$Q \sim -q \frac{r^3}{L^3}.$$

The force is, by Coulomb's law,

$$F \sim kqQ \left( \frac{1}{L^2} - \frac{L}{(L^2 + r^2)^{3/2}} \right) \sim \frac{kq^2 r^5}{L^7}.$$

A more accurate estimate would get the numeric prefactors.

### Example 5

Find the charge distribution on a conducting disc of radius  $R$  and total charge  $Q$ .

### Solution

In general, there are very few situations where the charge distribution on a conductor can be found explicitly. As you've seen, some of the simplest examples can be solved with image charges. Some more complex, two-dimensional examples can be solved with a mathematical technique called conformal mapping. And this special example can be solved with a neat trick.

Consider a uniformly charged spherical shell centered on the origin, and consider a point  $P$  inside the shell, on the  $xy$  plane. The electric field at point  $P$  is zero, by the shell theorem. Recall that in the usual proof of the shell theorem, one draws two cones opening out of  $P$  in opposite directions. The charges contained in each cone produce canceling electric fields.

Now imagine shrinking the spherical shell towards the  $xy$  plane, so it becomes elliptical. The crucial insight is that the shell theorem argument above still works, for points on the  $xy$  plane. When we squash the shell all the way down to the  $xy$  plane, it becomes a disc, with zero electric field on it. This is thus a valid charge distribution for a disc-shaped conductor, and by the uniqueness theorem, it's the only one.

By keeping track of how much charge gets squashed to radius  $[r, r + dr]$ , we find  $\sigma(r) \propto R/\sqrt{R^2 - r^2}$ , and fixing the proportionality constant gives

$$\sigma(r) = \frac{Q}{4\pi R \sqrt{R^2 - r^2}}.$$

## 4 Electrical Conduction

We now leave the world of electrostatics and consider magnetostatics, the study of steady currents.

### Idea 6

In a conductor with conductivity  $\sigma$ , the current density is

$$\mathbf{J} = \sigma \mathbf{E}.$$

Alternatively,  $\mathbf{E} = \rho \mathbf{J}$  where  $\rho$  is the resistivity. The current and charge density satisfy

$$\nabla \cdot \mathbf{J} = -\frac{\partial \rho}{\partial t}.$$

The current passing through a surface  $S$  at a given time is

$$I = \int_S \mathbf{J} \cdot d\mathbf{S}.$$

Since  $\mathbf{J} \propto \mathbf{E}$ , we have Ohm's law  $V = IR$ , where  $V$  is the voltage drop across the resistor. The power dissipated in a resistor is  $P = IV$ . The resistance  $R$  adds in series, while  $1/R$  adds in parallel.

[2] **Problem 21** (HRK). A battery causes a current to run through a loop of wire.

- Suppose the wire makes a sharp corner. How do the charges know to turn around there?
- A copper wire with conductivity  $\sigma$  is joined to an iron wire with conductivity  $\sigma' < \sigma$ . For the current in both sections to be the same, the electric field in the iron wire must be higher. How does that happen?

In general, the surface charge distribution in a DC circuit can be quite complex; the aspects shown in these two short questions are just the beginning. For some more about this, see [this paper](#).

**Solution.** (a) The first charges to make it there don't; they just stop at the surface of the wire, due to the attraction from the protons. Once this charge builds up at the kink, it repels the next electrons so that they automatically turn around. This typically occurs extremely quickly, as the relevant timescale is the  $RC$  of the wire and  $C$  is tiny. The amount of charge required is very small, less than a few hundred electrons.

- It's the same story as part (a). The first charges to reach the iron will start moving slower, because the fields are the same. This then causes a buildup of charge at the interface between them, which increases the field in the iron and decreases the field in the copper. In the steady state, the current densities in both are equal. Again, this occurs very quickly and requires very little charge.

[2] **Problem 22** (PPP 22). Two students, living in neighboring rooms, decided to economize by connecting their ceiling lights in series. They agreed they would each install a 100 W bulb in their own rooms and that they would pay equal shares of the electricity bill. However, both tries to get better lighting at the other's expense. The first student installed a 200 W bulb, while the second student installed a 50 W bulb. Which student subsequently failed the end-of-term examinations?

**Solution.** A standard bulb is designed to be hooked up in parallel with other bulbs, across some fixed voltage  $V$ . Since  $P = V^2/R$ , higher wattage bulbs have lower resistance.

Since the bulbs are in series, then have the same current through them. Since  $P = I^2 R$ , that means the bulb with the *higher* wattage rating draws *less* power, so the first student fails. The second student's bulb burns four times brighter. (The fact that the brightness of a bulb depends on what else is attached is also a reason why hooking lights up in series is impractical.)

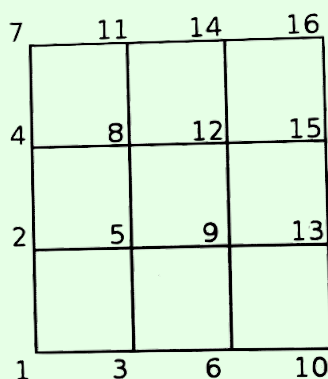
To warm up for DC circuits, we'll consider some resistor network problems.

### Idea 7

If any two points in a resistor network are at the same potential, nothing will change if the two points are connected together and treated as one. More generally, the resistance of any resistor directly connecting the two points may be changed freely.

### Example 6

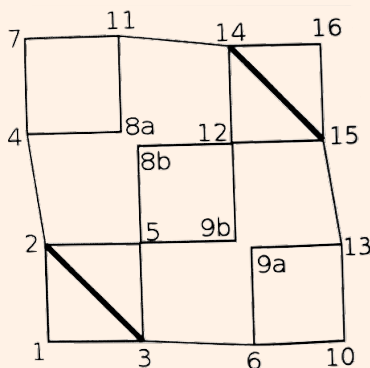
Consider the  $3 \times 3$  grid below, where every edge is a resistor  $R$ .



Find the equivalent resistance between nodes 1 and 16.

### Solution

By the above idea, we can short together nodes 2/3, and 14/15, by the diagonal symmetry of the network. Next, we can break nodes 8 and 9 into two pieces.



This is valid because the separated nodes  $8a/8b$  and  $9a/9b$  still have the same potential in

the new network, by the diagonal symmetry. (This is using the above idea in reverse.) Now, the circuit has been reduced to combinations of series and parallel resistors. The resistance between 1 and 2/3 is  $R/2$ . The resistance between 2/3 and 14/15 is the combination of three networks in parallel, and finally the resistance of 14/15 and 16 is  $R/2$ . Thus,

$$R_{\text{eq}} = \left( \frac{1}{2} + \left( \frac{1}{3} + \frac{1}{3} + \frac{1}{2} \right)^{-1} + \frac{1}{2} \right) R = \frac{13}{7} R.$$

You won't see any resistor problems as complicated as this one for the rest of the training, because they're kind of contrived; the point of this example was just to show multiple uses of symmetry techniques.

### Example 7: PPP 23

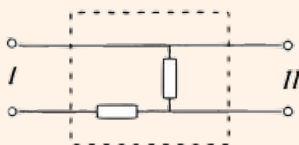
A black box contains a resistor network and has two output terminals.



If a battery of voltage  $V$  is connected across the first terminal, the voltage across the second terminal is  $V/2$ . If a battery of voltage  $V$  is connected across the second terminal, the voltage across the first terminal is  $V$ . Find one possible configuration of the resistors inside the box.

### Solution

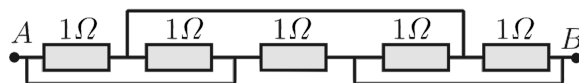
A simple configuration with two equal resistors works.



When a battery is connected across II, the horizontal resistor doesn't do anything. When a battery is connected across I, the two resistors comprise a voltage divider.

[2] **Problem 23.** ⌚ USAPhO 2007, problem A1.

[2] **Problem 24** (IPhO 1996). Consider the following resistor network.



Find the equivalent resistance between A and B.

**Solution.** See the official solutions [here](#).

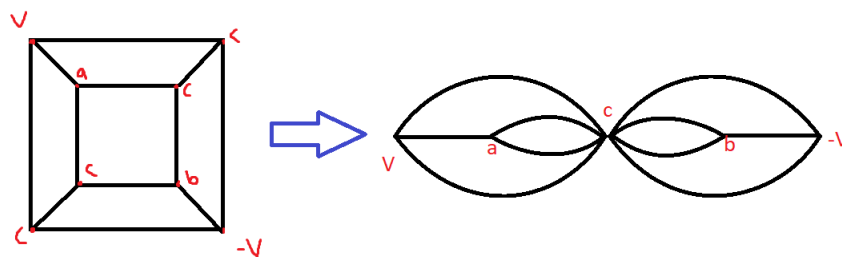
[3] **Problem 25.** Consider a cube of side length  $L$  whose edges are resistors of resistance  $R$ .

- (a) Compute the resistance between two vertices a distance  $\sqrt{3}L$  apart.
- (b) Compute the resistance between two vertices a distance  $\sqrt{2}L$  apart.
- (c) Compute the resistance between two vertices a distance  $L$  apart.
- (d) Generalize to vertices  $\sqrt{n}L$  apart on an  $n$ -dimensional cube. (Give your answer in the form of a summation.)

**Solution.** (a) Let the two vertices be  $A$  and  $B$ . Let the vertices distance 1 from  $A$  be labeled  $a$  and distance two labeled  $b$ . Note that all the vertices labeled  $a$  have the same potential by symmetry, and same for  $b$ . Thus, we can treat all the vertices labeled the same thing as one vertex.

We have 3 connections from  $A$  to  $a$ , 6 from  $a$  to  $b$ , and 3 from  $b$  to  $B$ . Thus, our resistance is  $R/3 + R/6 + R/3 = \boxed{5R/6}$ .

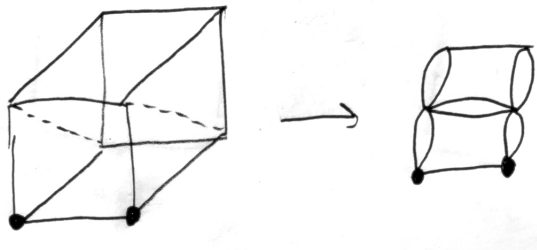
- (b) Apply potential  $V$  and  $-V$  to the vertices, and label the rest of the vertices as shown.



We claim that all the vertices labeled  $c$  have potential 0. The idea is then that negating the potentials  $V$  and  $-V$  must negate the potentials at all vertices. But negating is equivalent to a simple reflection that preserves the locations of the vertices labeled  $c$ . Thus the only option is that their potential is 0. Therefore, all the  $c$ s can be treated as one vertex, and we have the drawn equivalent circuit. This is a combination of parallel and series, and we compute the answer to be

$$2 \cdot \frac{1}{1 + 1 + \frac{1}{1+1/2}} R = \boxed{3R/4}.$$

- (c) In a similar fashion, the points labeled the same below have the same potential, and on the right is the equivalent circuit.





This is again just a series/parallel problem, and we compute the answer to be

$$R \frac{1}{\frac{1}{\frac{1}{\frac{1}{2} + \frac{1}{\frac{1}{\frac{1}{2} + \frac{1}{\frac{1}{2} + 1 + \frac{1}{2}}}} + \frac{1}{2}} + \frac{1}{1}}} = \boxed{7R/12}.$$

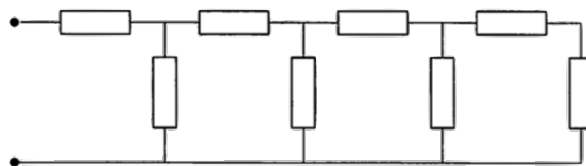
- (d) The coordinates take the form  $(x_1, x_2, \dots, x_n)$  where  $x_i$  is zero or one. We consider the vertices  $(0, 0, \dots, 0)$  and  $(1, 1, \dots, 1)$ . The first vertex is connected to all the vertices with one 1, of which there are  $n$ . By symmetry, these are all at the same voltage. Next, these vertices are connected to all the vertices with two 1's, of which there are  $\binom{n}{2}$ , and so on.

We hence have  $n + 1$  effective vertices of different voltages. Consider the vertex representing points with  $k$  1's. The number of connections to points with  $k + 1$  1's is  $\binom{n}{k}(n - k)$ . Then by adding series and parallel resistances,

$$R_{\text{eq}} = R \sum_{k=0}^n \left( \binom{n}{k}(n - k) \right)^{-1}.$$

For example, this recovers the result of part (a) for  $n = 3$ .

[2] **Problem 26** (PPP 158). Consider the circuit below, where every resistor is  $1 \Omega$ .



- (a) Find the equivalence resistance between the input terminals.  
 (b) Do the same in the case where the chain is infinitely long.

**Solution.** (a) Suppose the rightmost resistor has unit current flowing down, and let  $V$  be the potential difference across  $A$  and  $B$ .

Number the resistors, starting from the back, and going to the left. Let  $I_k$  be the current in  $R_k$ . We claim that

$$I_k = F_k$$

where the  $F_k$  are the Fibonacci numbers  $F_1 = F_2 = 1$  and  $F_n = F_{n-1} + F_{n-2}$ . This is true by induction. Note that

$$I_{2k} = I_{2k-1} + I_{2k-2}$$

by junction law. By the loop law,

$$I_{2k+1} = I_{2k} + I_{2k-1}.$$

Note that all currents are either to the right or down. Let the number of vertical resistors be  $n$ . Then, the potential difference is

$$V = I_{2n-1} + I_{2n} = I_{2n+1} = F_{2n+1}.$$

The equivalent resistance is

$$R = (F_{2n+1}/F_{2n}) \Omega.$$

- (b) By taking the limit  $n \rightarrow \infty$  above, we get the golden ratio,

$$R = \frac{1 + \sqrt{5}}{2} \Omega.$$

There's another, slicker way to do this. In the infinite case, if we let the answer be  $R$ , then we have

$$R = 1 + \frac{1}{1 + \frac{1}{R}}$$

which is equivalent to the quadratic

$$R^2 - R - 1 = 0, \quad R = \frac{1 \pm \sqrt{5}}{2}.$$

Then taking the positive root gives the answer.

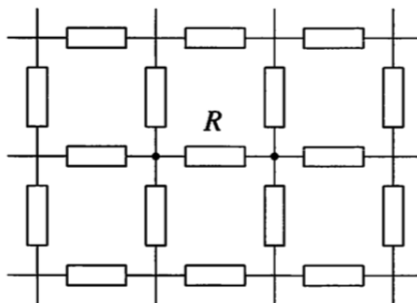
This presents a question: what's the deal with the negative root? Both the positive and negative roots are technically equally good. In fact, there exist circuit elements with negative resistance, as you'll see in **E6**, though they need an active source of power to maintain.

However, the positive root is still better in an important way. Suppose you start out with something on the right end, with (possibly negative) equivalent resistance  $R_0$ . If we attach two  $1\Omega$  resistors on the left, as in the diagram, then we get some new resistance  $R_1$ , and attaching two more gives  $R_2$ , and so on. The two roots for  $R$  found above are the two fixed points for this iteration. However, the positive root is more physical because it's the *stable* fixed point. If you start out with any resistance other than  $(1 - \sqrt{5})/2$ , the recursion will converge to  $(1 + \sqrt{5})/2$ . (For example, in part (a), we showed this happens if you start with  $1\Omega$ .) Only if you start out with the exact negative resistance  $(1 - \sqrt{5})/2$  will you stay there.

Therefore, even though a mathematically infinite network doesn't have anything "on the right end" (since it has no end at all), it's still meaningful to say that "the" resistance is  $(1 + \sqrt{5})/2$ . When we introduce infinite objects in physics, we usually do so just to get a mathematically tractable approximation for a real, finite object. And for almost any long, but finite chain, you'll get an answer near  $(1 + \sqrt{5})/2$ , so that's the useful answer in the infinite case.

- [3] **Problem 27** (PPP 159-161). Superposition can be a useful trick to analyze circuit networks.

- (a) Consider an infinite two-dimensional grid of identical resistors  $R$ .



Find the equivalent resistance between two neighboring points by considering the superposition of 1 A flowing into one point, and 1 A flowing out the other.

- (b) What would the equivalent resistance be if the resistor directly connecting the two neighboring points was removed?

- (c) Now consider an icosahedron of identical resistors  $R$ . By superposing appropriate current distributions, find the equivalent resistance between two neighboring vertices.

**Solution.** (a) We see that if just  $I$  flowed in, then the current in the joining resistor would be  $I/4$ , so in this case its  $I/2$ . Thus, the potential drop is  $IR/2$  for a current of  $I/2$ , so the resistance is  $\boxed{R/2}$ .

- (b) Suppose the answer is  $R'$ . We have that  $R'$  and  $R$  attached in parallel give  $R/2$ . Thus,

$$\frac{1}{R'} + \frac{1}{R} = \frac{2}{R},$$

so  $R' = \boxed{R}$ .

- (c) Consider the current distribution where  $I$  flows into a vertex, and  $I/11$  leaves from each other vertex. We see that the current in each edge coming out from the source vertex is  $I/5$ . So superposing a similar but with signs flipped current distribution for an adjacent vertex, we see that  $(12/11)IR_{\text{eff}} = 2RI/5$ , so  $R_{\text{eff}} = \boxed{11R/30}$ .

### Idea 8

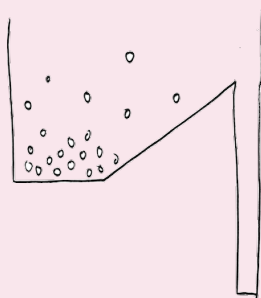
In a circuit of resistors and batteries, Kirchoff's loop rule states that the sum of the voltage drops around a loop is zero. Kirchoff's junction rule states that the net current flowing into a vertex is zero. (This is technically nonzero, because of the effect of problem 21, but negligible because wires have tiny capacitance.)

### Remark

If the sum of the voltage drops around a loop is zero, then why would current ever want to flow? After all, if you had a circular tube of water, the water would never flow, because the net drop in height along the circle is zero. The reason current flows in circuits with batteries is that within the battery, charges are moved from lower to higher electric potential energy, just like how a pump could be used to move water upward to start a liquid circuit, by an "electromotive force".

But this immediately raises the question: what *is* this specific force? It can't be the electric force, because we just established that it's pointing the wrong way. It's not a magnetic effect. For some setups, it is literally a mechanical force like a pump: in the Van der Graaff generator, a motor drives the charges on a statically charged conveyor belt to higher potential. But that's not how batteries work.

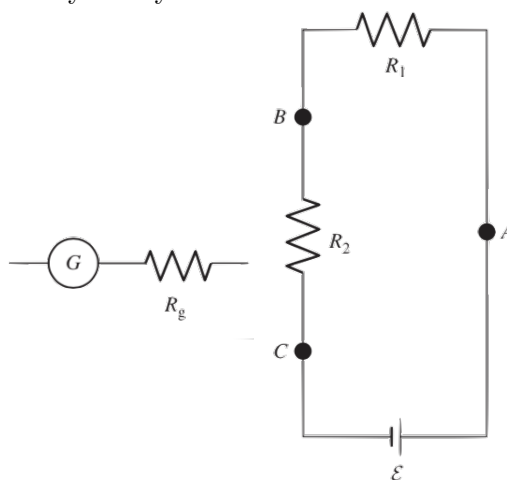
In a battery, there is no specific force pushing charges from low to high electric potential. Instead, the charges just jiggle around randomly, and the result emerges from the effects of their many collisions. To understand this, consider a gravitational analogy.



Consider an ideal gas at temperature  $T$  released in the trough shown above. The gas molecules will randomly collide, sometimes being propelled upward by chance. Sometimes, a gas molecule will climb the hill and fall into the deep hole, at which point it is unlikely to come out again. Thus, if the hole begins empty, it is energetically favorable for gas molecules to fill it. But there is no attractive force pulling molecules up along the slope! Gravity always points down; molecules go up the slope when they are randomly bounced that way.

This is essentially how the potential in an initially neutral battery is set up. The hole corresponds to the lower energy state an electron can reach inside the anode, but there is no long-range force pushing it there, just the average effect of random collisions.

- [2] **Problem 28** (Purcell 4.10). The basic ingredient in older voltmeters and ammeters is the galvanometer, a device to measure very small currents. (It works via magnetic effects, but the exact mechanism isn't important here.) Inherent in any galvanometer is some resistance  $R_g$ , so a physical galvanometer can be represented by the system shown below.



Consider a circuit such as the one shown, with all quantities unknown. We want to measure the current flowing across point A and the voltage difference between points B and C. Given a galvanometer with known  $R_g$ , and also a supply of known resistors (ranging from much smaller to much larger than  $R_g$ ), how can you accomplish these two tasks? Explain how to construct your two devices (called an ammeter and voltmeter), and also how you should insert them in the given circuit. You will need to make sure that you (a) affect the given circuit as little as possible, and (b) don't destroy your galvanometer by passing more current through it than it can handle.

**Solution.** We explain how to create an ammeter and a voltmeter. As usual, attach the ammeter in series, and the voltmeter in parallel.


To create an ammeter, connect a resistor  $R \ll R_g$  in parallel with the galvanometer (including  $R_g$  obviously). Suppose a current  $I$  passes through the device. If  $N = R_g/R$ , then it's not hard to see that the current through galvanometer is about  $I/N$ , and the system acts very much like a wire. Therefore, to get  $I$ , multiply the reading by  $N$ .

To create a voltmeter, attach a resistor  $R \gg R_g$  in series to the galvanometer (again, attach this contraption in parallel to the circuit). If the voltage drop we want to measure is  $V$ , then the current will be about  $V/R$ , so we can use the current to find  $V$ .

### Remark

Occasionally, you might see Olympiad problems where a voltmeter is connected in series. The most common voltmeters are handheld digital multimeters, where the voltmeter setting presents a resistance of about  $10\text{ M}\Omega$ . Thus, for such problems, you should just treat the voltmeter like a high-resistance resistor.

Is this realistic? Well, it certainly happens every day, in almost every introductory physics lab in the world. But no professional would ever do this on purpose, because voltmeters aren't designed to be used this way. There is no guarantee that the resistance of the voltmeter is a constant. Instead, for most digital multimeters, there is a complex circuit inside that adjusts the internal resistance depending on the input and the configuration settings. You probably won't break the voltmeter when you put it in series, but you won't get reliable results either.

- [2] **Problem 29.**  USAPhO Quarterfinal 2009, problems 3 and 4.
- [3] **Problem 30.** [INPhO 2021, problem 1](#). A nice problem on practical circuit measurements. Note that the question statement is a bit vague. You are supposed to keep track of quantities of order  $R_A/R$  and  $R/R_V$ , but you are allowed to neglect quantities as small as  $R_A/R_V$ .

**Solution.** See the official solutions [here](#).