

Continuous System: Waves

Continuous String as a Limiting Case of the Loaded String

wish to allow the number of masses to become infinite $n \rightarrow \infty$

$$\text{let } m \rightarrow 0, d \rightarrow 0 \quad \frac{m}{d} = \text{constant} \equiv \rho$$

$$n \rightarrow \infty, d \rightarrow 0 \quad (n+1)d = L$$

$$q_r(t) = \sum_r n_r(t) \sin\left(\frac{jr\pi}{n+1}\right)$$

$$B_r e^{i\omega_r t} = r\pi \frac{j d}{(n+1)d} = r\pi \frac{x}{L}$$

$$\Rightarrow q(x, t) = \sum_r n_r(t) \sin\left(\frac{r\pi x}{L}\right)$$

$$= \sum_r B_r e^{i\omega_r t} \sin\left(\frac{r\pi x}{L}\right)$$

$$q(x, 0) = \sum_r \mu_r \sin\left(\frac{r\pi x}{L}\right) \quad \frac{dq(x, 0)}{dt} = -\sum_r (\omega_r v_r) \sin\left(\frac{r\pi x}{L}\right)$$

$$\int_0^L \sin\left(\frac{r\pi x}{L}\right) \sin\left(\frac{s\pi x}{L}\right) dx = \frac{L}{2} \delta_{rs}$$

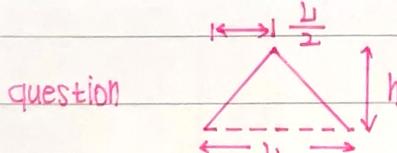
x sin $\frac{s\pi x}{L}$ ↓ integrate

$$\mu_r = \frac{2}{L} \int_0^L q(x, 0) \sin\left(\frac{r\pi x}{L}\right) dx \quad v_r = -\frac{2}{(\omega_r L)} \int_0^L \frac{dq(x, 0)}{dt} \sin\left(\frac{r\pi x}{L}\right) dx$$

$$\omega_r = 2 \sqrt{\frac{\Gamma}{md}} \sin\left[\frac{r\pi}{2(n+1)}\right]$$

$$= \frac{2}{d} \sqrt{\frac{\Gamma}{\rho}} \sin\left(\frac{r\pi}{2L}\right)$$

when $d \rightarrow 0 \quad \omega_r = \frac{r\pi}{L} \sqrt{\frac{\Gamma}{\rho}}$



question

$$q(x, 0) = \begin{cases} \frac{2h}{L} x & 0 \leq x \leq \frac{L}{2} \\ \frac{2h}{L} (L-x) & \frac{L}{2} \leq x \leq L \end{cases} \quad \frac{dq(x, 0)}{dt} = 0$$

$$\mu_r = \frac{4h}{L^2} \int_0^{\frac{L}{2}} x \sin\left(\frac{r\pi x}{L}\right) dx + \frac{4h}{L^2} \int_{\frac{L}{2}}^L (L-x) \sin\left(\frac{r\pi x}{L}\right) dx$$

because released from rest, all the v_r vanish

$$= \frac{8h}{r^2 \pi^2} \sin\left(\frac{r\pi}{2}\right) \quad \begin{cases} 0 & \text{even} \\ \frac{8h}{r^2 \pi^2} (-1)^{\frac{1}{2}(r-1)} & r \text{ odd} \end{cases}$$

$$q(x, t) = \frac{8h}{\pi^2} \sin\left(\frac{r\pi x}{L}\right)$$

Energy of a Vibrating String

$$E_k = \frac{1}{2} \rho \int_0^L \left(\frac{\partial \eta}{\partial t} \right)^2 dx = \frac{1}{2} \rho \int_0^L \left[\sum_r \frac{d\eta_r}{dt} \sin\left(\frac{r\pi x}{L}\right) \right]^2 dx$$

$$= \frac{1}{2} \rho \sum_{rs} \frac{d\eta_r}{dt} \frac{d\eta_s}{dt} \int_0^L \sin\left(\frac{r\pi x}{L}\right) \sin\left(\frac{s\pi x}{L}\right) dx$$

$$= \frac{\rho L}{4} \sum_{rs} \frac{d\eta_r}{dt} \frac{d\eta_s}{dt} \delta_{rs} = \frac{\rho L}{4} \sum_r \left(\frac{d\eta_r}{dt} \right)^2$$

$$\text{Re}\left[\left(\frac{d\eta_r}{dt}\right)^2\right] = \text{Re} \left\{ \frac{d}{dt} [\mu_r + i\nu_r] (\cos\omega_r t + i\sin\omega_r t) \right\}^2$$

$$= (-\omega_r \mu_r \sin\omega_r t - \omega_r \nu_r \cos\omega_r t)^2$$

$$\Rightarrow E_k = \frac{\rho L}{4} \sum_r \omega_r^2 (\mu_r \sin\omega_r t + \nu_r \cos\omega_r t)^2$$

$$U = \frac{1}{2} \frac{I}{d} \sum_j (q_{j+1} - q_j)^2 = \frac{1}{2} I \sum_j \left(\frac{q_{j+1} - q_j}{d} \right)^2 d \quad \text{and } d \rightarrow 0$$

$$= \frac{1}{2} I \int_0^L \left(\frac{\partial \eta}{\partial x} \right)^2 dx = \frac{1}{2} I \int_0^L \left[\sum_r \frac{r\pi}{L} \eta_r \cos\left(\frac{r\pi x}{L}\right) \right]^2 dx$$

$$= \sum_r \frac{r\pi}{L} \eta_r \cos\left(\frac{r\pi x}{L}\right)$$

$$= \frac{I}{2} \sum_{rs} \frac{r\pi}{L} \frac{s\pi}{L} \eta_r \eta_s \int_0^L \cos\left(\frac{r\pi x}{L}\right) \cos\left(\frac{s\pi x}{L}\right) dx$$

$$= \frac{I}{2} \sum_{rs} \frac{r\pi}{L} \frac{s\pi}{L} \eta_r \eta_s \frac{L}{2} \delta_{rs}$$

$$= \frac{I}{2} \sum_r \frac{r^2 \pi^2}{L^2} \frac{L}{2} \eta_r^2 = \frac{\rho L}{4} \sum_r \omega_r^2 \eta_r^2$$

$$= \frac{\rho L}{4} \sum_r \omega_r^2 (\mu_r \cos\omega_r t - \nu_r \sin\omega_r t)^2$$

$$E = E_k + U = \frac{\rho L}{4} \sum_r \omega_r^2 (\mu_r^2 + \nu_r^2) = \frac{\rho L}{4} \sum_r \omega_r^2 |B_r|^2$$

the average over one complete period
of the fundamental vibration $r=1$

$$\langle E_k \rangle = \frac{\rho L}{4} \sum_r \omega_r^2 \langle (\mu_r \sin\omega_r t + \nu_r \cos\omega_r t)^2 \rangle = \frac{\rho L}{8} \sum_r \omega_r^2 (\mu_r^2 + \nu_r^2)$$

$$\langle U \rangle = \frac{\rho L}{4} \sum_r \omega_r^2 \langle \mu_r \cos\omega_r t - \nu_r \sin\omega_r t \rangle^2 = \frac{\rho L}{8} \sum_r \omega_r^2 (\mu_r^2 + \nu_r^2)$$

$$\Rightarrow \langle T \rangle = \langle U \rangle$$

General Solutions of the Wave Equation

$\Psi = \Psi(x, t)$ time-dependent wave equation

$\psi = \psi(x)$ time-independent wave equation

$$\frac{\partial^2 \Psi}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 \Psi}{\partial t^2} = 0, \quad \frac{\partial^2 \Psi}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} = 0$$

$$\begin{aligned} \varsigma &= x + vt \\ \eta &= x - vt \end{aligned}$$

$$\frac{\partial \Psi}{\partial x} = \frac{\partial \Psi}{\partial \varsigma} \frac{\partial \varsigma}{\partial x} + \frac{\partial \Psi}{\partial \eta} \frac{\partial \eta}{\partial x} = \frac{\partial \Psi}{\partial \varsigma} + \frac{\partial \Psi}{\partial \eta}$$

$$\begin{aligned} \frac{\partial^2 \Psi}{\partial x^2} &= \frac{\partial}{\partial x} \frac{\partial \Psi}{\partial x} = \frac{\partial}{\partial x} \left(\frac{\partial \Psi}{\partial \varsigma} + \frac{\partial \Psi}{\partial \eta} \right) = \frac{\partial}{\partial \varsigma} \left(\frac{\partial \Psi}{\partial \varsigma} + \frac{\partial \Psi}{\partial \eta} \right) + \frac{\partial}{\partial \eta} \left(\frac{\partial \Psi}{\partial \varsigma} + \frac{\partial \Psi}{\partial \eta} \right) \frac{\partial \eta}{\partial x} \\ &= \frac{\partial^2 \Psi}{\partial \varsigma^2} + 2 \frac{\partial^2 \Psi}{\partial \varsigma \partial \eta} + \frac{\partial^2 \Psi}{\partial \eta^2} \end{aligned}$$

and

$$\begin{aligned} \frac{1}{v} \frac{\partial \Psi}{\partial t} &= \frac{\partial \Psi}{\partial \varsigma} - \frac{\partial \Psi}{\partial \eta}, \quad \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} = \frac{1}{v} \frac{\partial}{\partial t} \left(\frac{1}{v} \frac{\partial \Psi}{\partial t} \right) = \frac{1}{v} \frac{\partial}{\partial t} \left(\frac{\partial \Psi}{\partial \varsigma} - \frac{\partial \Psi}{\partial \eta} \right) \\ &= \frac{\partial^2 \Psi}{\partial \varsigma^2} - 2 \frac{\partial^2 \Psi}{\partial \varsigma \partial \eta} + \frac{\partial^2 \Psi}{\partial \eta^2} \end{aligned}$$

$$\Rightarrow \frac{\partial^2 \Psi}{\partial \varsigma \partial \eta} = 0$$

$$\Psi = f(\varsigma) + g(\eta) = f(x+vt) + g(x-vt)$$

$$\text{as time } t=0 \quad \Psi(x, 0) = f(x) + g(x)$$