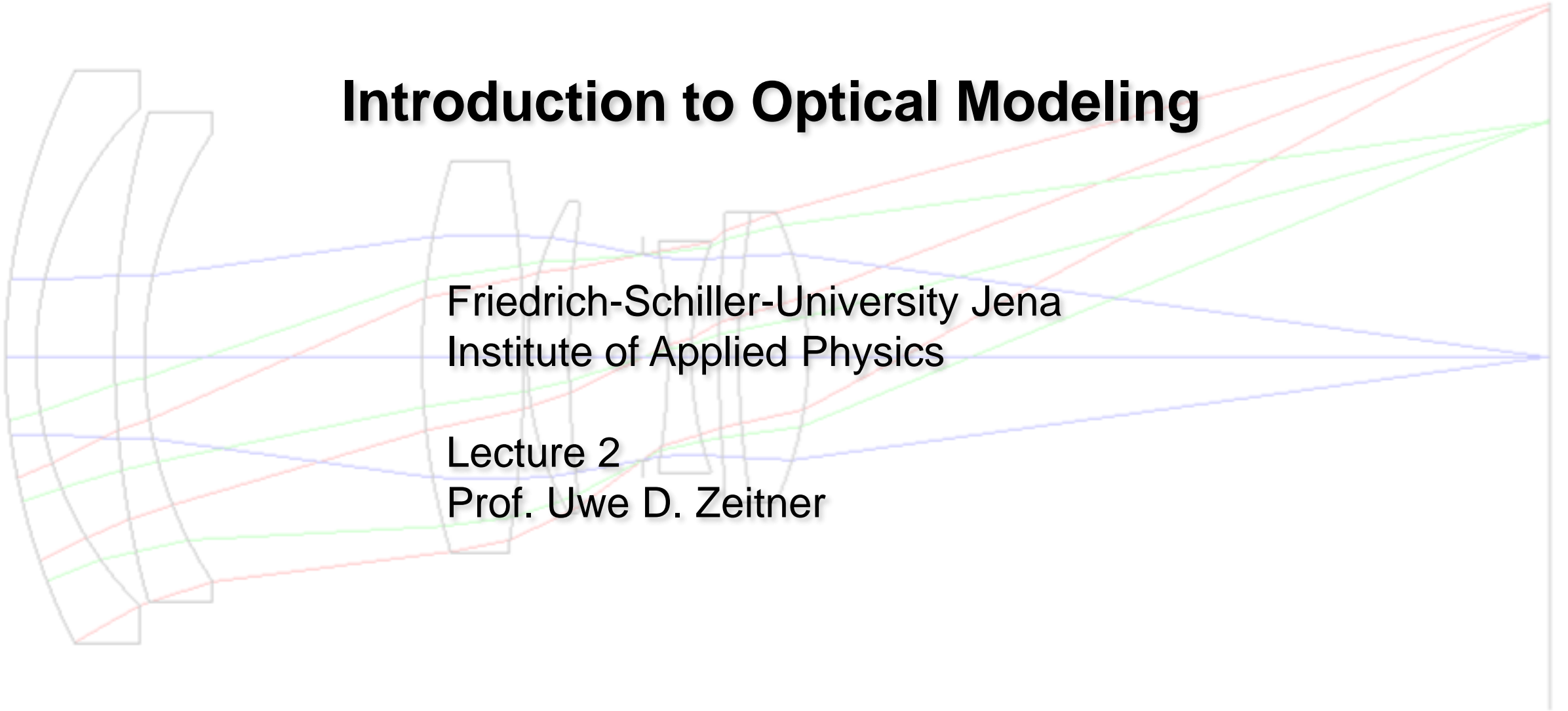


Introduction to Optical Modeling

Friedrich-Schiller-University Jena
Institute of Applied Physics

Lecture 2
Prof. Uwe D. Zeitner



Course Overview

Part 1: Geometrical optics based modeling and design (U.D. Zeitner)

1. Introduction
- 2. Paraxial approximation / Gaussian optics**
- 3. ABCD-matrix formalism**
4. Real lenses
5. Optical materials
 - glass types, dispersion
 - chromatic aberrations
6. Imaging systems
 - apertures/stops, entrance-/exit-pupil
 - wavefront aberrations

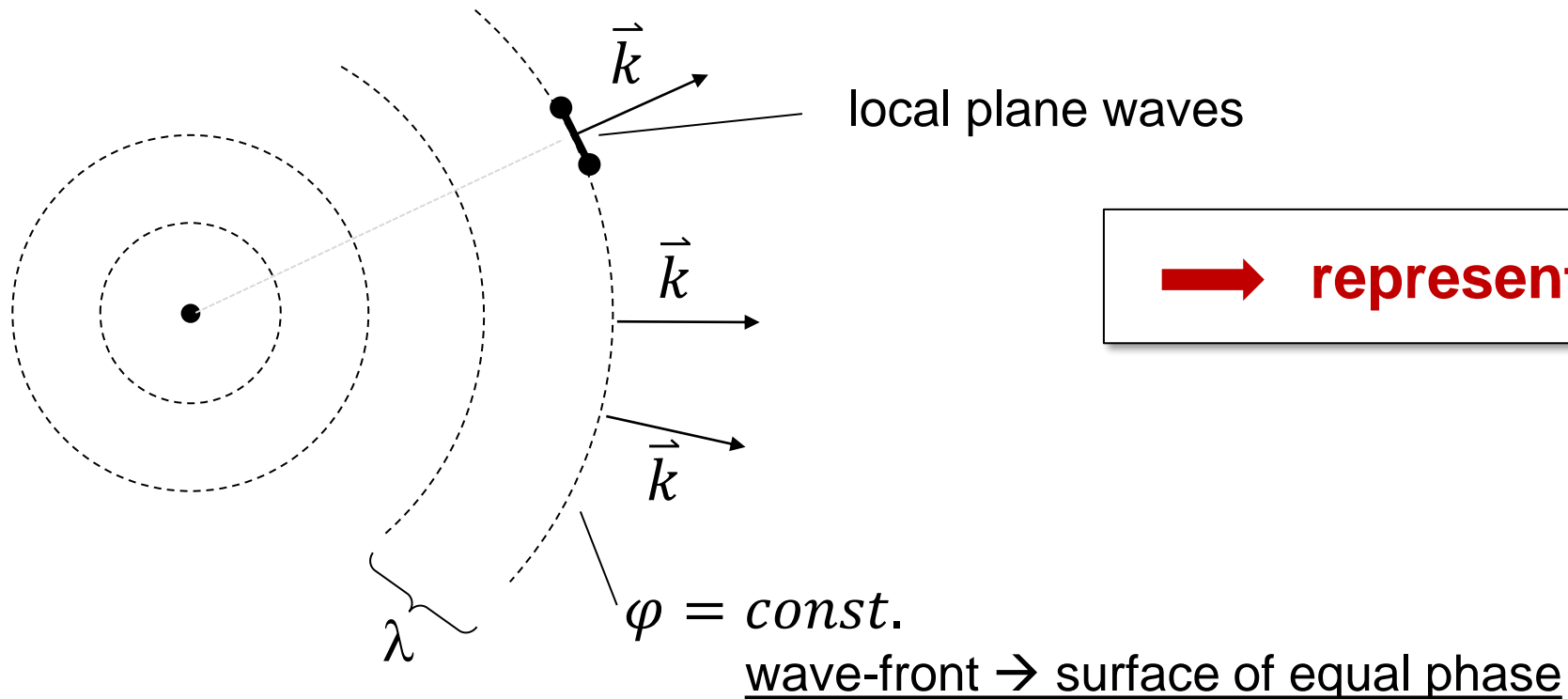
Part 2: Wave-optics based modeling (F. Wyrowski)

Recap: Propagation of Spherical Waves

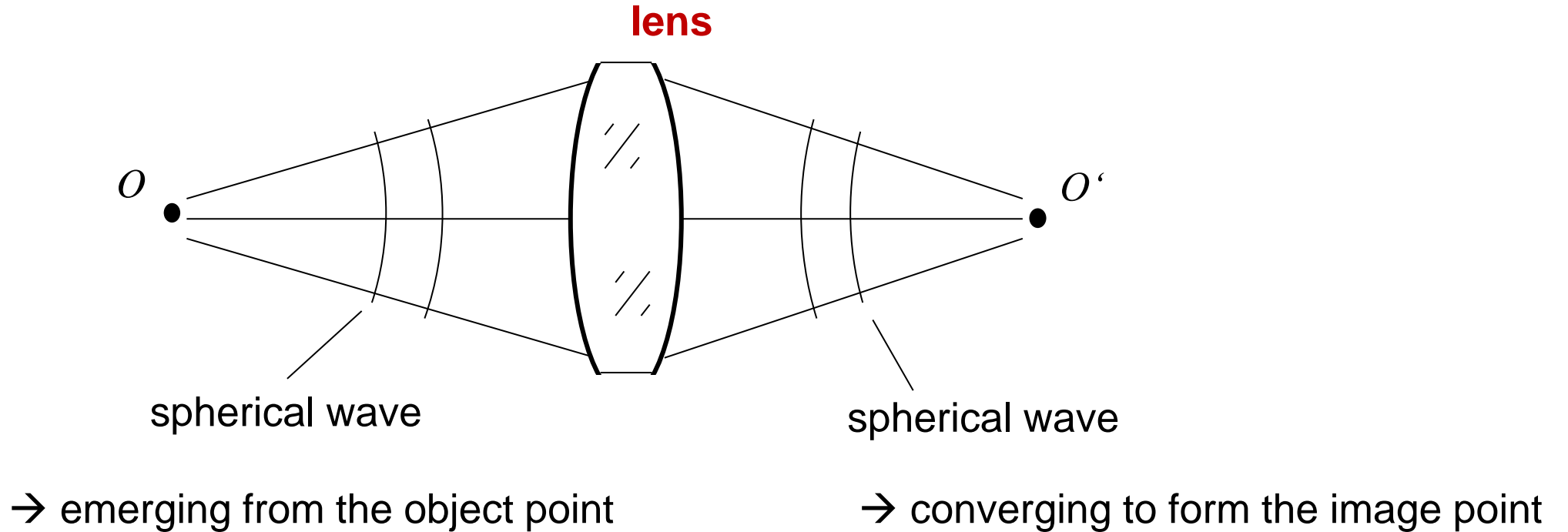
Basic solution of Helmholtz Equation:

$$E(\vec{r}) = \frac{E_0}{r} \cdot e^{ik \cdot r} \quad (1.5)$$

→ spherical wave (from a point source)

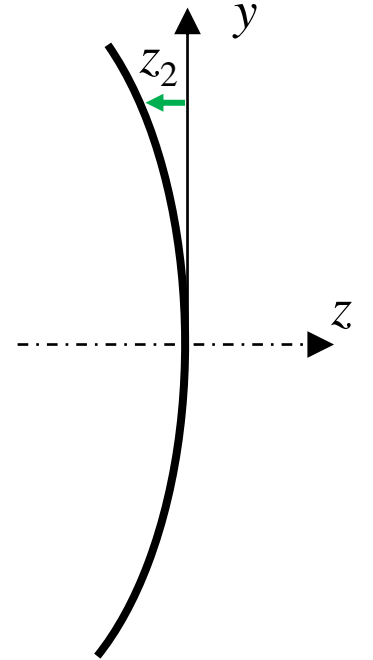
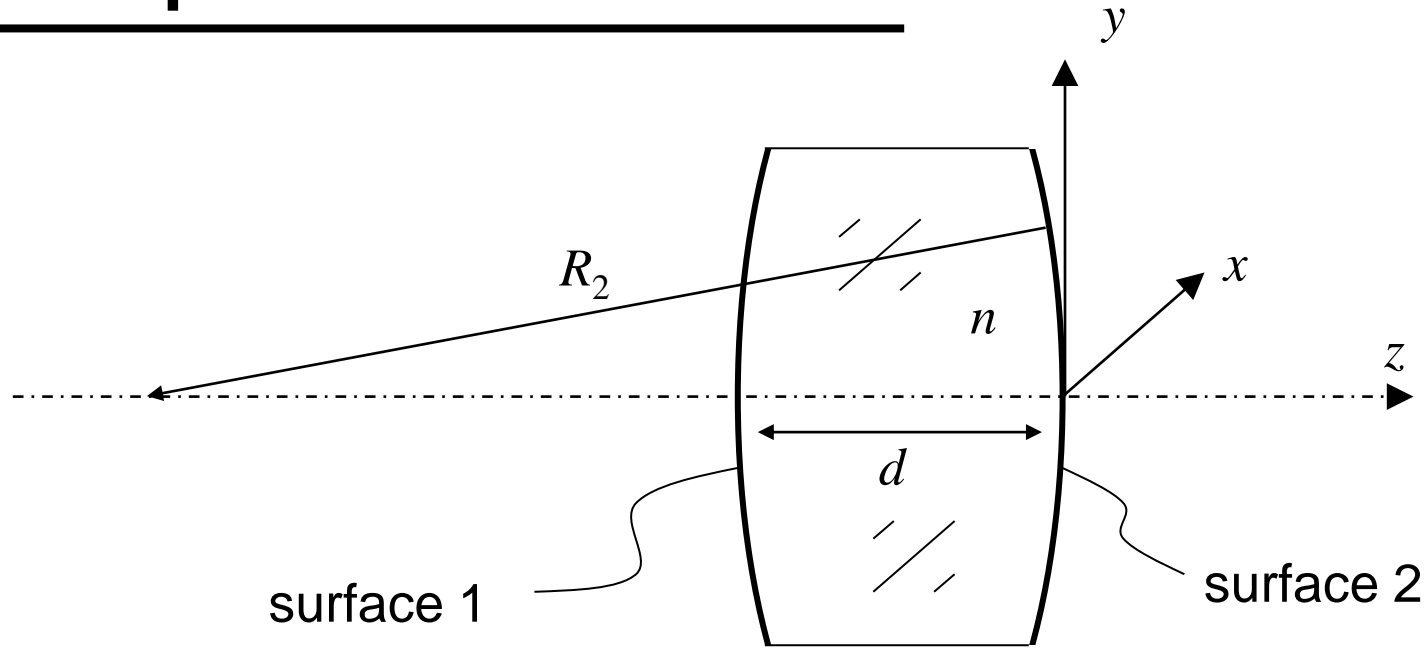


2 Paraxial Imaging / Gaussian Optics



Imaging: **transformation of a spherical wave into another spherical wave**

Description of a Lens



Lens sag (description in the local coordinate system):

$$z_{1/2} = \frac{c_{1/2} \cdot r^2}{1 + \sqrt{1 - (1 + k_{1/2})c_{1/2}^2 r^2}} + a_2 r^2 + a_4 r^4 + \dots \quad (2.1)$$

$r^2 = x^2 + y^2$... radius

$c_{1/2} = \frac{1}{R_{1/2}}$... curvature

R ... radius of curvature

k ... conic constant

d ... lens thickness

(distance of vertex points)

Description of a lens, conic constant

Sag formula:

$$z_{1/2} = \frac{c_{1/2} \cdot r^2}{1 + \sqrt{1 - (1 + k_{1/2})c_{1/2}^2 r^2}} + a_2 r^2 + a_4 r^4 + \dots$$

k ... conic constant \rightarrow describes surface as conic section

$k < -1$ \rightarrow hyperbola

$k = -1$ \rightarrow parabola

$-1 < k < 0$ \rightarrow ellipse

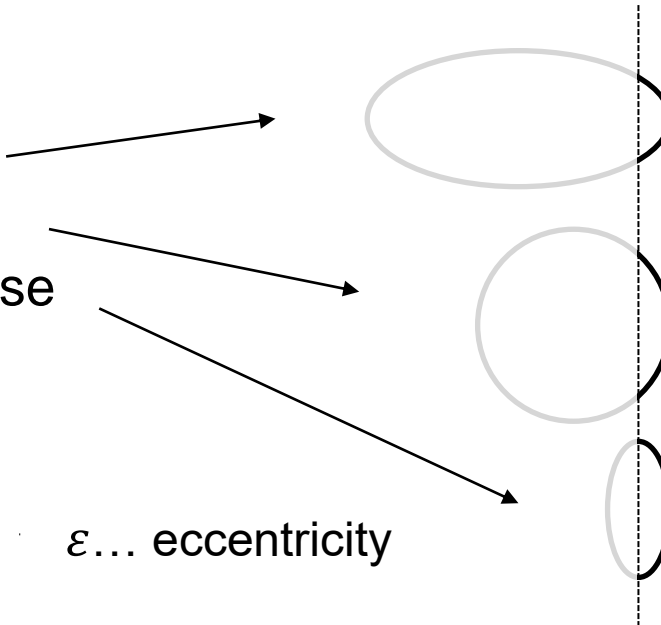
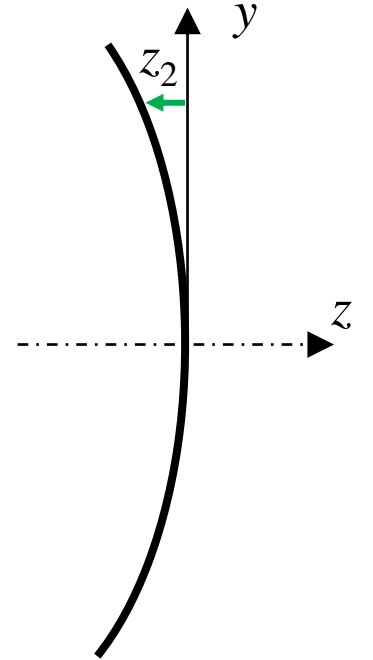
$k = 0$ \rightarrow sphere

$k > 0$ \rightarrow oblate ellipse

$$k = -\varepsilon^2 = -(1 - b^2/a^2)$$

$$1/c = R = \pm b^2/a$$

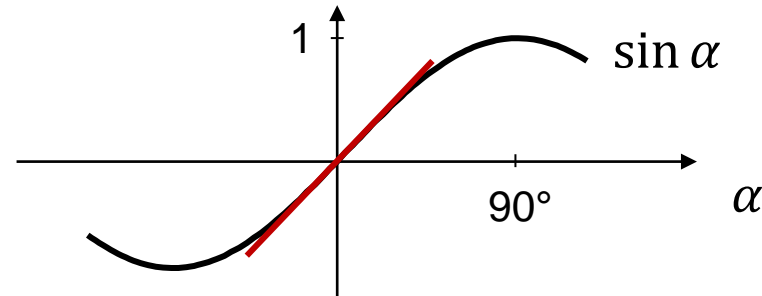
ε ... eccentricity



2.1 Paraxial Approximation

Law of refraction:

$$n \cdot \sin \alpha = n' \cdot \sin \alpha' \quad (1.6)$$



Paraxial approximation:

α, α' small \rightarrow linear approx. of sin-Fct.



$$\begin{aligned} n \cdot \alpha &= n' \cdot \alpha' \\ \cos \alpha &\approx 1 \end{aligned}$$

paraxial law of refraction (2.2)

Simplification used in:

- law of refraction
- corresponding ray angles
- equations describing optical surfaces



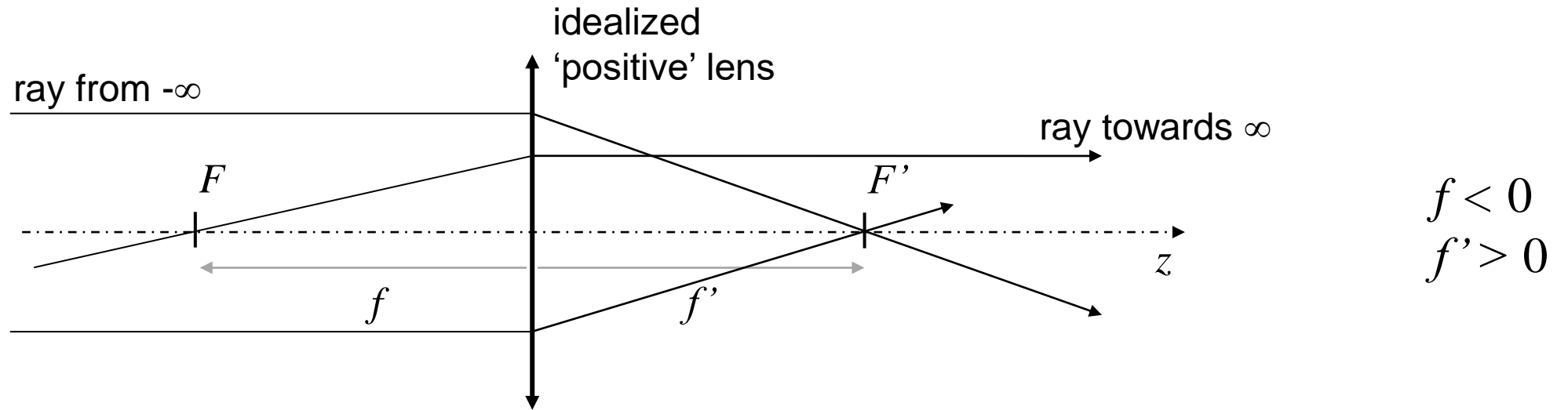
- most **equations** become **linear**
- **no aberrations** during imaging

2.2 Ideal Lens

Optical function: transforming a spherical wave into another spherical wave

Ideal Lens: assumption that optical effect takes place in the “plane of the lens”

Considering two special cases: object / image point at infinite distance from the lens



Sign convention in optics:

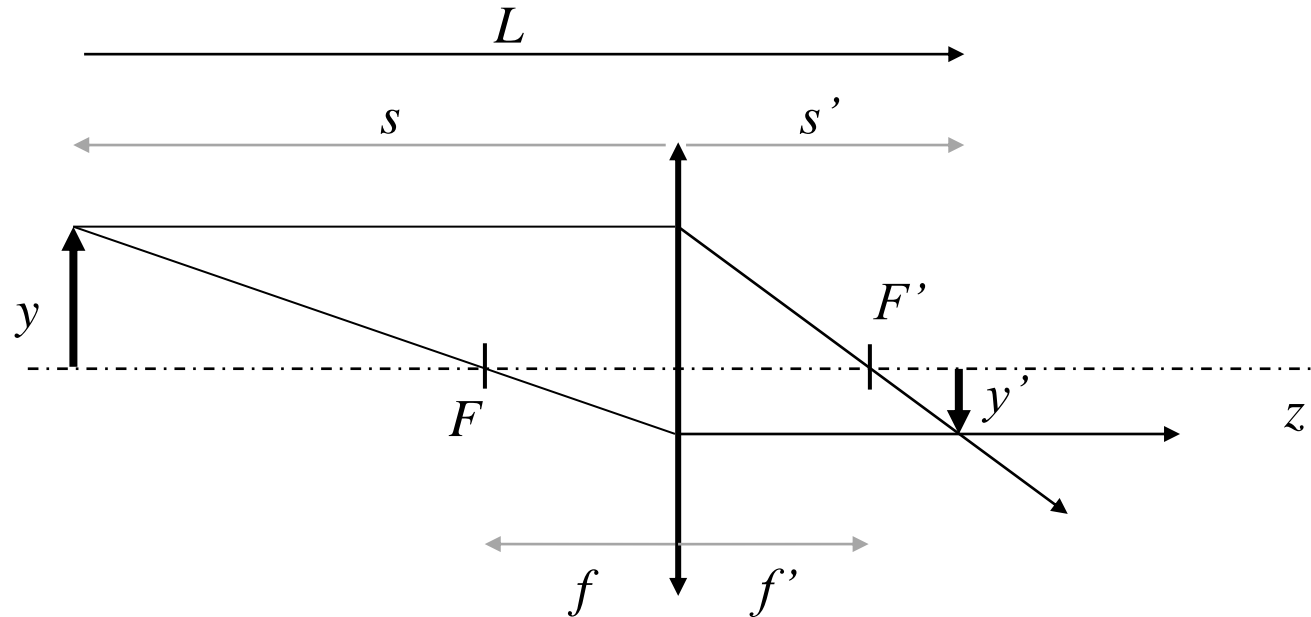
distances are counted positive if oriented in sense of a vector in positive axial direction

Radii: direction is oriented from the surface towards the center of curvature

Simple Image Construction

for a finite object distance

Positive lens: $f' > 0$



y ... object height (> 0)

y' ... image height (< 0)

s ... object distance

s' ... image distance

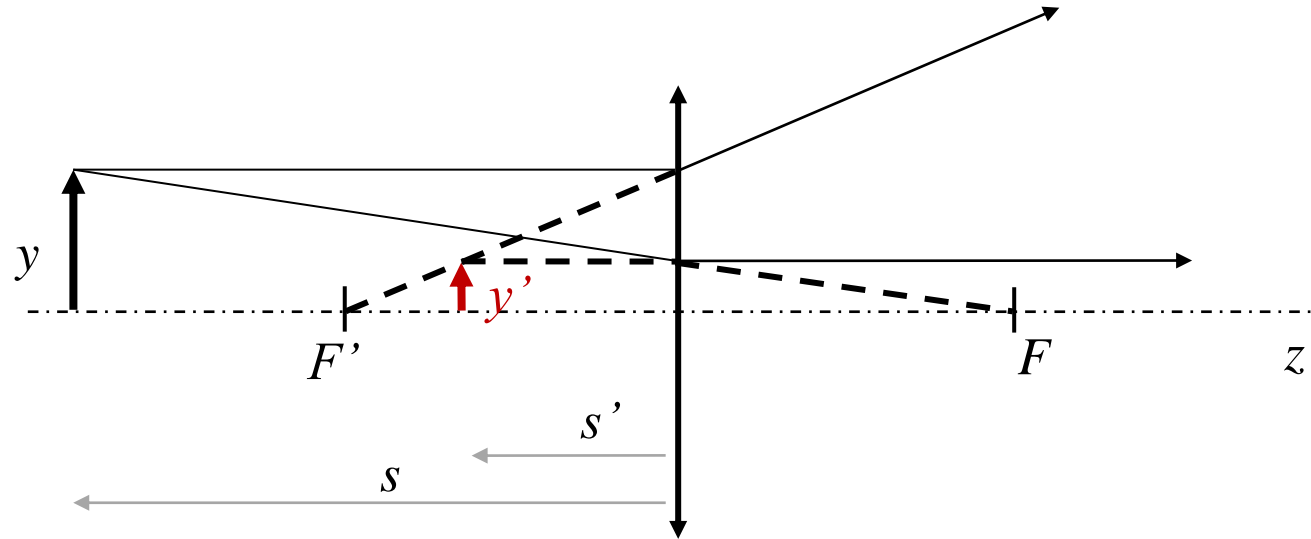
$L = s' - s$... object – image distance

here: equal n before and behind the lens

Simple Image Construction

for a finite object distance

Negative lens: $f' < 0$ → reversed position of F and F'



→ virtual image !!

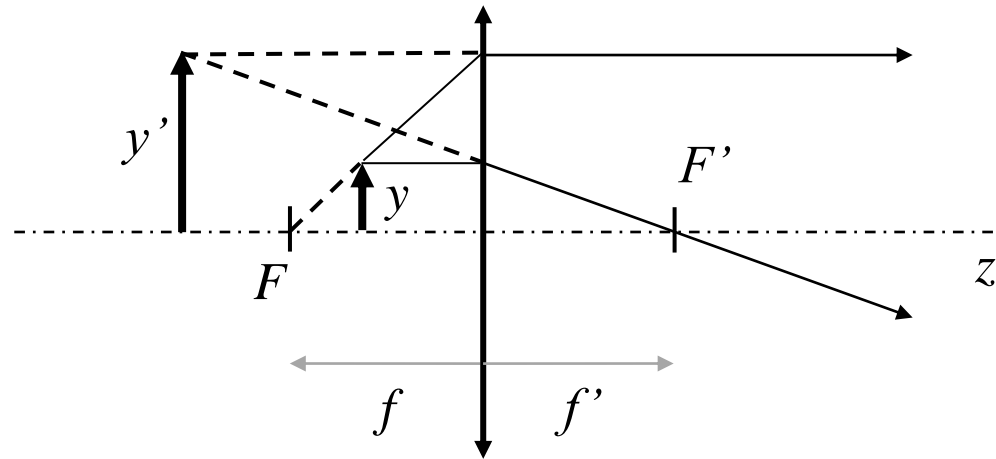
y ... object height (> 0)
 y' ... image height (> 0)
 s ... object distance
 s' ... image distance


here: equal n before and behind the lens

Two important optical “systems”

a) Magnifying Lens

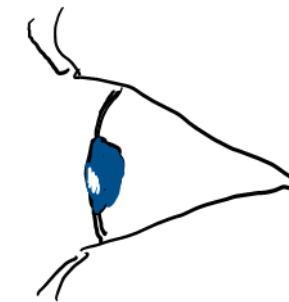
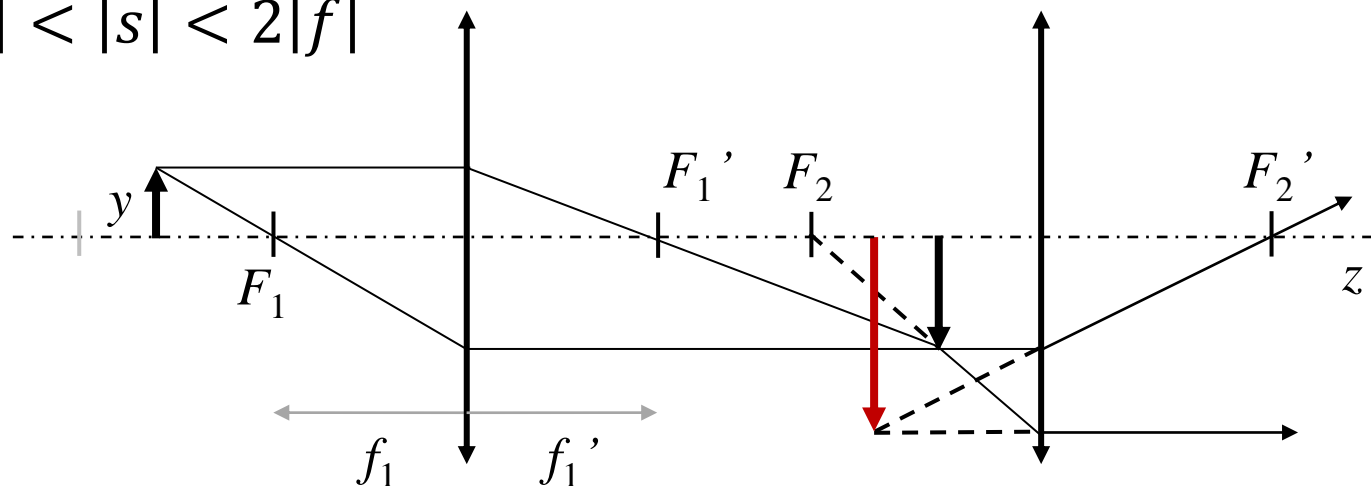
$$|s| < |f|$$



 virtual magnified image
→ observable by eye

b) Microscope → combination of a magnified real image and a magnifying lens

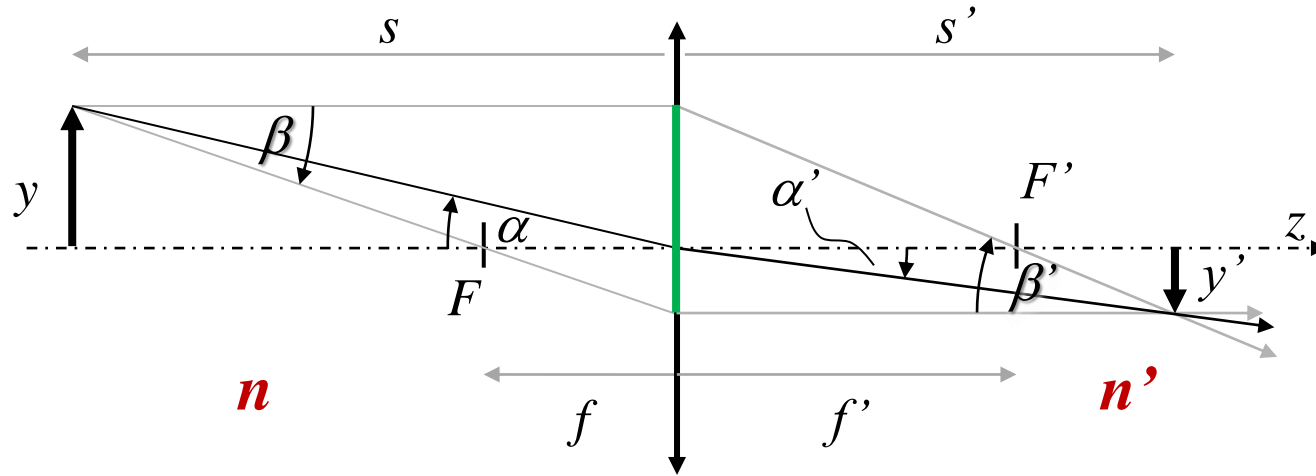
$$|f| < |s| < 2|f|$$



observation
by eye

More general relation between f and f'

more general situation: **different n** in object and image space



$$\Rightarrow f \neq f'$$

$$n \cdot \alpha = n' \cdot \alpha' \quad (2.2)$$

similar consideration:

$$\frac{f'}{s'} + \frac{f}{s} = 1 \quad (2.4)$$

imaging equation

or with $n = n'$

$$\frac{1}{s'} - \frac{1}{s} = \frac{1}{f'} \quad (2.5)$$

“lens makers equation”

$$\left. \begin{array}{l} \beta = \frac{y - y'}{s} \quad \text{and} \quad \beta' = \frac{y - y'}{s'} \\ \Rightarrow s \cdot \beta = s' \cdot \beta' \\ \text{with} \quad \beta = \frac{y'}{f} \quad ; \quad \beta' = -\frac{y}{f'} \\ \text{and} \quad \alpha = \frac{y}{s} \quad ; \quad \alpha' = \frac{y'}{s'} \end{array} \right\} \xRightarrow{(2.2)} \boxed{\frac{f'}{n'} = -\frac{f}{n}} \quad (2.3)$$

Magnification

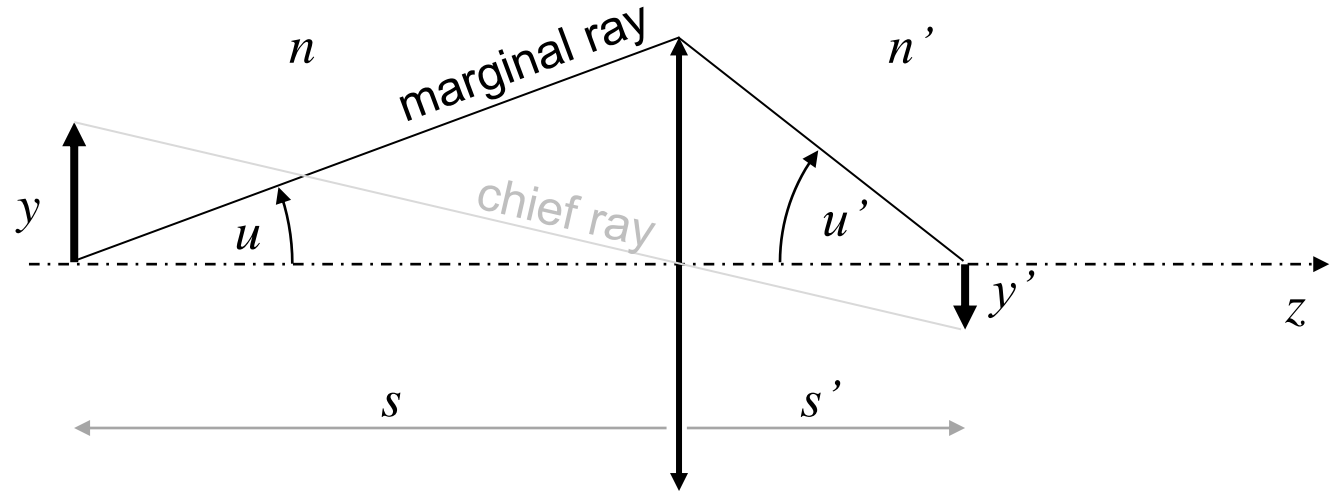
Definition:

$$m = \frac{y'}{y}$$

(2.6)

from sketch: $m = \frac{n \cdot s'}{n' \cdot s}$

with aperture angles u, u' : $m = \frac{n \cdot \tan u}{n' \cdot \tan u'} \approx \frac{n \cdot u}{n' \cdot u'}$

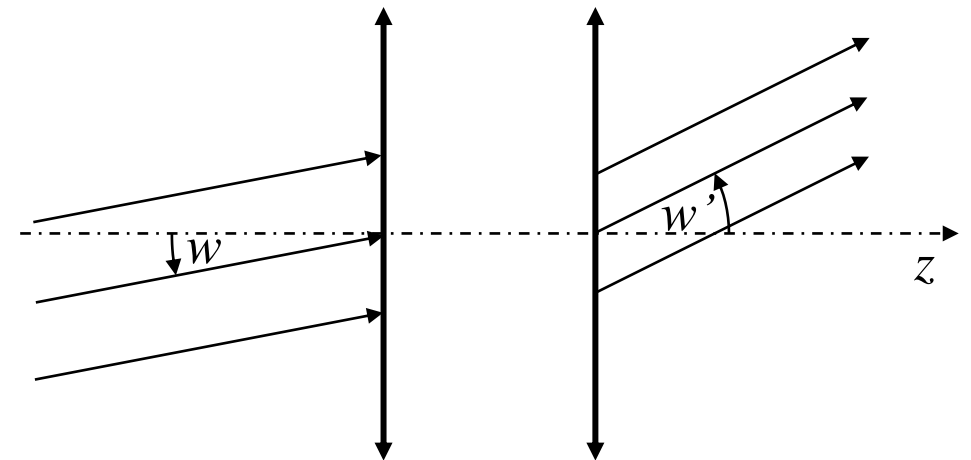


Special case: object and image at infinity (afocal system) ➔ Telescope

➔ y, y' are not defined

➔ definition of the angular magnification by the chief ray angles w, w'

$$\Gamma = \frac{w'}{w}$$



Paraxial Imaging Equations

Quantity to be calculated	Calculation equations		
s	$s = \frac{s'f'}{f' - s'}$	$s = s' - L$	$s = \frac{s'}{m}$
	$s = -\frac{L}{s} \pm \sqrt{\frac{L^2}{4} - f' \cdot L}$	$s = \frac{(1-m) \cdot f'}{m}$	$s = \frac{L}{m-1}$
s'	$s' = \frac{s \cdot f'}{f' + s}$	$s' = s + L$	$s' = \frac{L}{m-1}$
	$s' = \frac{L}{2} \pm \sqrt{\frac{L^2}{4} - f' \cdot L}$	$s' = f' \cdot (1 - m)$	$s' = \frac{L \cdot m}{m-1}$
f'	$f' = \frac{s \cdot s'}{s - s'}$	$f' = -\frac{s \cdot (L + s)}{L}$	$f' = \frac{s \cdot m}{1 - m}$
	$f' = \frac{s' \cdot (L - s')}{L}$	$f' = \frac{s'}{1 - m}$	$f' = -\frac{L \cdot m}{(1 - m)^2}$
L	$L = s \cdot (m - 1)$	$L = s - s'$	$L = -\frac{s^2}{s + f'}$
	$L = \frac{s'^2}{s' - f'}$	$L = \frac{s' \cdot (m - 1)}{m}$	$L = f' \cdot \left(2 - m - \frac{1}{m}\right)$
m	$m = \frac{s'}{s}$	$m = \frac{f'}{s + f'}$	$m = \frac{f' - s'}{f'}$
	$m = \frac{L + s}{s}$	$m = \frac{s'}{s' - L}$	$m = 1 - \frac{L}{2f'} \pm \sqrt{\frac{L}{f'} \cdot \left(\frac{L}{4f'} - 1\right)}$

set of 30 equations
for the calculation of
imaging parameters

2.3 ABCD-Matrix Formalism

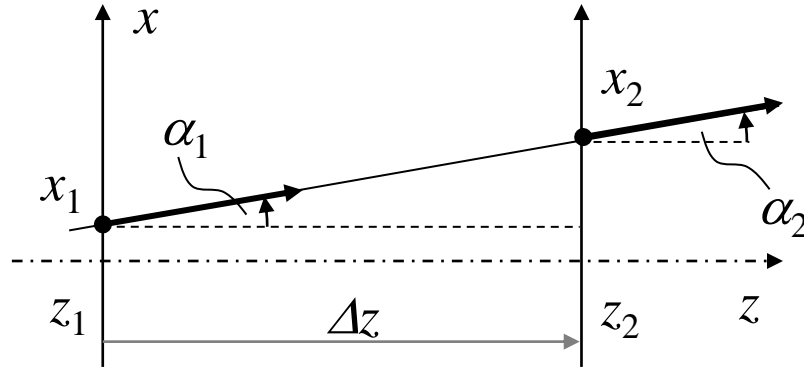
- based on a geometrical optics consideration of field propagation
- simple but powerful method for the (paraxial) treatment of complex optical systems
- ABCD-matrices can also be used for studying paraxial diffraction phenomena!

2.3.1 Derivation of the formalism

→ see blackboard

ABCD-matrices for common optical elements I

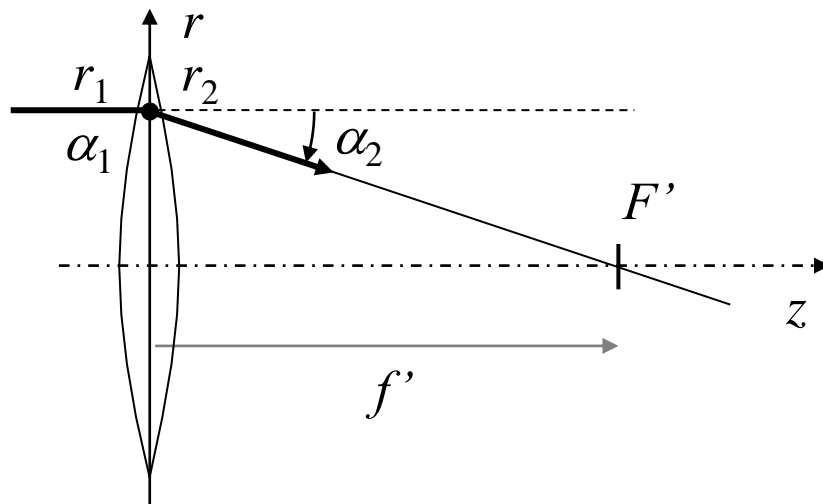
a) Free-space propagation:



$$\begin{aligned}x_2 &= x_1 + \Delta z \cdot \tan \alpha_1 \\ \alpha_2 &= \alpha_1\end{aligned}$$

$$\Rightarrow M_{\Delta z} = \begin{pmatrix} 1 & \Delta z \\ 0 & 1 \end{pmatrix} \quad (2.8)$$

b) Thin lens:

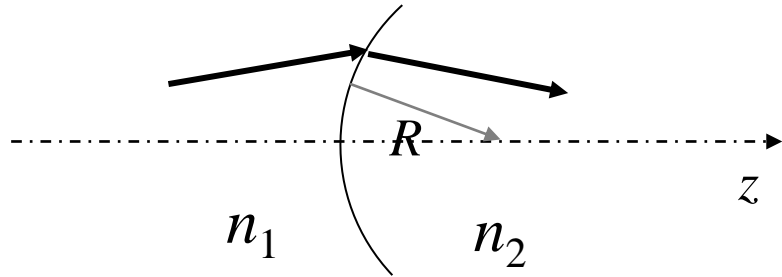


$$\begin{aligned}r_2 &= r_1 \\ \alpha_2 &= -\frac{r_1}{f'} + \alpha_1\end{aligned}$$

$$\Rightarrow M_f = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f'} & 1 \end{pmatrix} \quad (2.9)$$

ABCD-matrices for common optical elements II

c) Transition at a spherical interface $n_1 \rightarrow n_2$; ROC R :

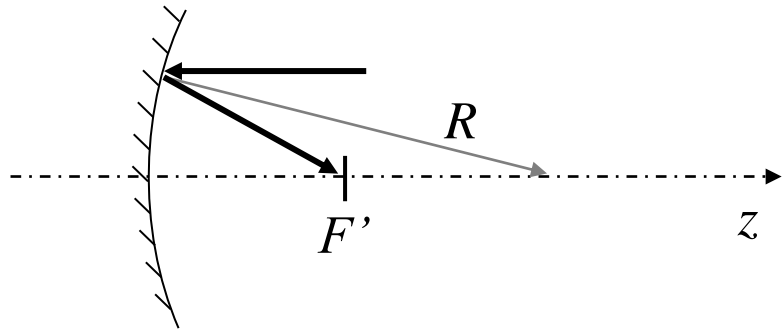


$$\rightarrow M_{SI} = \begin{pmatrix} 1 & 0 \\ \frac{n_1 - n_2}{n_2 \cdot R} & \frac{n_1}{n_2} \end{pmatrix} \quad (2.10)$$

special case: $R \rightarrow \infty$ (plane interface)

$$M_{ref} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{n_1}{n_2} \end{pmatrix} \quad (2.11)$$

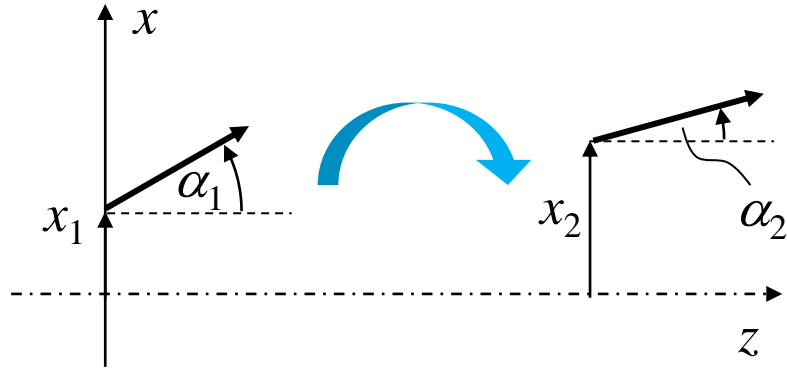
d) Curved mirror, radius $R = 2f$:



$$\rightarrow M_M = \begin{pmatrix} 1 & 0 \\ \frac{2}{R} & 1 \end{pmatrix} \quad (2.12)$$

ABCD-matrices for common optical elements III

e) Magnification m :



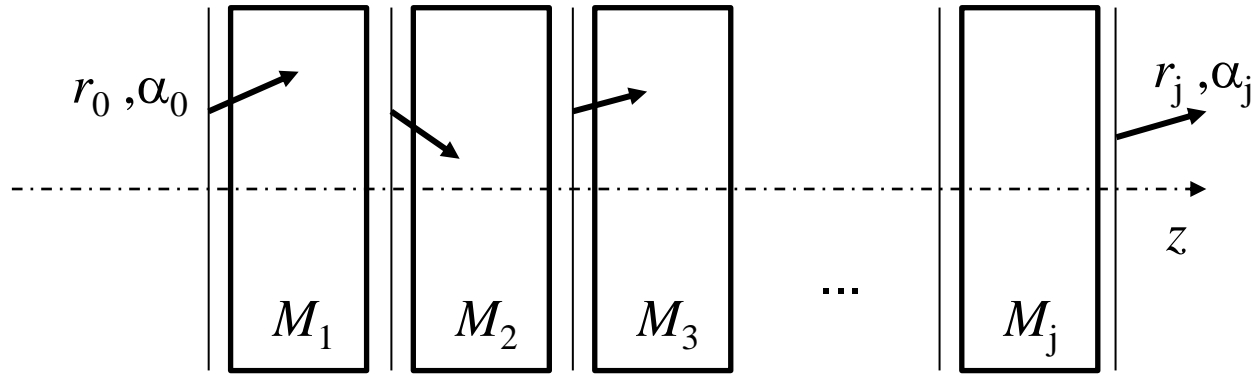
$$m = \frac{x_2}{x_1}$$

$$\rightarrow M_m = \begin{pmatrix} m & 0 \\ 0 & \frac{1}{m} \end{pmatrix} \quad (2.13)$$

→ a number of other matrices for other elements can be found in literature

ABCD-matrices for systems

optical system \rightarrow sequence of single optical elements



$$\rightarrow M_{sys} = M_j \cdot \dots \cdot M_3 \cdot M_2 \cdot M_1 \quad (2.14)$$

The system matrix can be easily obtained by multiplication of the matrices of the single elements the system is composed of.

2.3.2 General properties of ABCD-matrices I

1) Determinant of M

n_1 ... refractive index in front of the system

n_2 ... refractive index behind the system

$$|M| = AD - BC = \frac{n_1}{n_2} \quad (2.15)$$

equal indices: $AD - BC = 1$

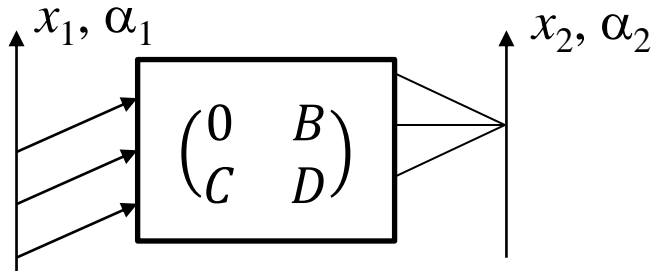
→ only 3 independent variables possible!

2) Inversion of the light direction / backwards propagation through the same system

$$M^{-1} = \begin{pmatrix} D & -B \\ -C & A \end{pmatrix} \quad (2.16)$$

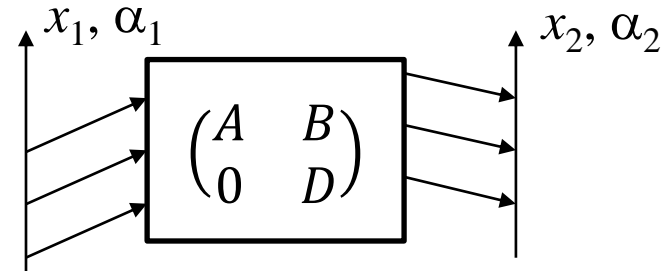
General properties if one matrix element is **zero**

a) $A = 0 \Rightarrow x_2 = B\alpha_1$



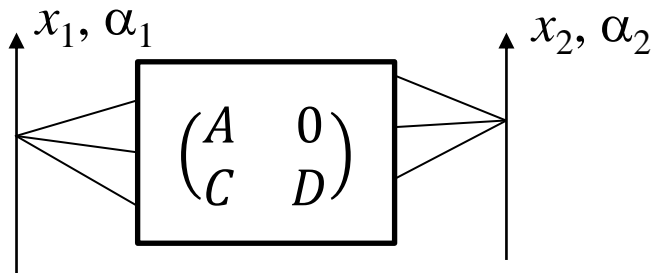
→ focusing

c) $C = 0 \Rightarrow \alpha_2 = D\alpha_1$



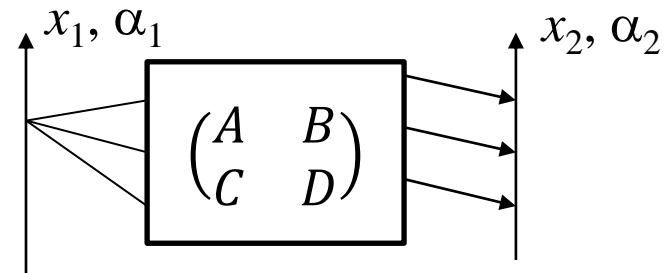
→ deflection

b) $B = 0 \Rightarrow x_2 = Ax_1$



→ imaging

d) $D = 0 \Rightarrow \alpha_2 = Cx_1$



→ defocusing

General properties of ABCD-matrices II

- 3) **Equivalent optical system:** \rightarrow systems having the same ABCD-matrix
 \rightarrow showing the same optical behavior

Decomposition of a given M into a “fixed” series of four basic operations:

- magnification change m
- change of refractive index $n_1 \rightarrow n_2$
- thin lens of optical power $\Phi = 1/f'$
- propagation along distance Δz

equivalent matrix: $M_{eq} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} 1 & \Delta z \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{f' \cdot n_2} & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & \frac{n_1}{n_2} \end{pmatrix} \begin{pmatrix} m & 0 \\ 0 & \frac{1}{m} \end{pmatrix}$

 4 free variables in terms of the given ABCD-matrix:

$$\frac{n_1}{n_2} = AD - BC$$

$$m = \frac{AD - BC}{D}$$

$$\frac{1}{f' \cdot n_2} = \frac{CD}{AD - BC} = -\frac{C}{m}$$

$$\Delta z = \frac{B}{D}$$

(2.17)

Attention: if $D = 0 \rightarrow$ re-arrange the sequence of the 4 operations

General properties of ABCD-matrices III

4) ABCD-matrices and fields, **Collins-Integral**:

- the 4 operations of the equivalent matrix (magnification, material transition, thin lens phase, propagation) can easily be applied on arbitrary fields $\vec{E}(x, y, z)$
- not only a ray-optical consideration
- propagation of wave-optical fields through complex systems possible

Other option: integrate ABCD-matrix directly into a diffraction integral

$$E_2(x_2, y_2) = \frac{i}{\lambda B} e^{-ikL} \iint E_1(x_1, y_1) \cdot e^{-i\frac{\pi}{\lambda B}(A(x_1^2+y_1^2)+D(x_2^2+y_2^2)-2(x_1x_2+y_1y_2))} dx_1 dy_1 \quad (2.18)$$

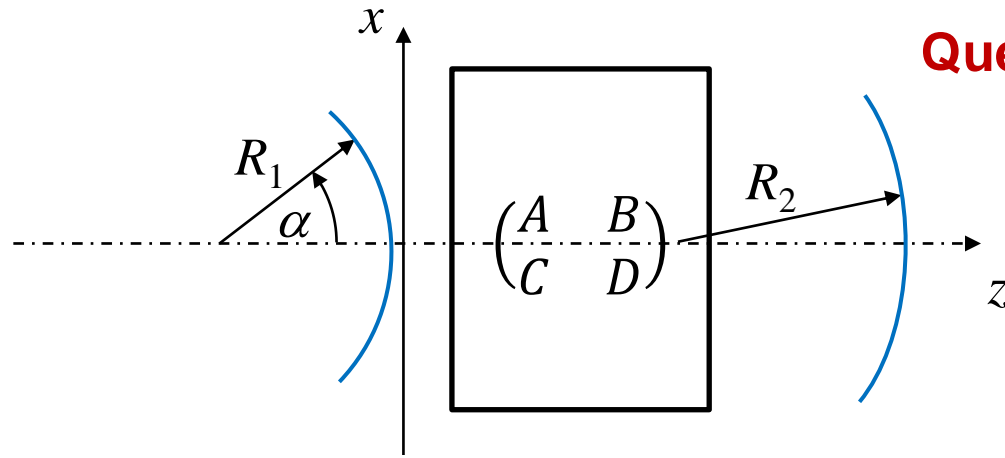
→ **Collins Integral**

E_1 ... field at the input plane of the system
 E_2 ... field at the output plane of the system
 L ... optical path along the optical axis

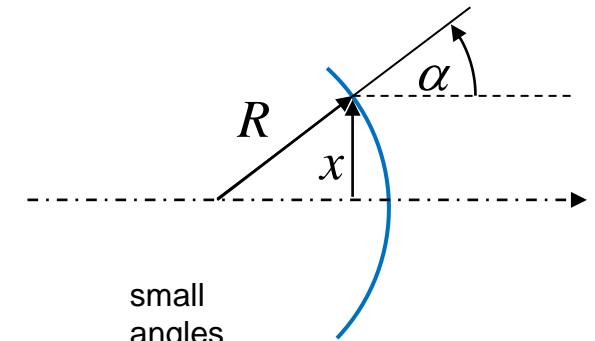
General properties of ABCD-matrices IV

5) Transformation of a spherical wave:

illumination of the system with spherical wave, ROC = R_1



Question: $R_2 = ?$



relation: $\tan \alpha = \frac{x}{R} \stackrel{\text{small angles}}{\approx} \alpha$

relation between x and α

$$\left. \begin{array}{l} x_1 = R_1 \cdot \alpha_1 \\ x_2 = R_2 \cdot \alpha_2 \end{array} \right\} R_2 = \frac{AR_1 + B}{CR_1 + D} \quad (2.19)$$

Example: lens with focal length f'

$$\rightarrow \frac{1}{R_2} = \frac{1}{R_1} - \frac{1}{f'}$$