

## Maxwell's Equations

## Poynting Vector

## Normal Modes and Poynting Vector

Tutorial 1, Normal Modes and Poynting Vector.

a).  $\vec{H}(\vec{r}) = -\frac{i}{\omega \mu_0} \text{rot} \vec{E}(\vec{r})$

$$= -\frac{i}{\omega \mu_0} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \partial_x & \partial_y & \partial_z \\ 0 & \vec{E}_x & 0 \end{vmatrix} = -\frac{i}{\omega \mu_0} \begin{pmatrix} -\partial_z \vec{E}_y \\ 0 \\ \partial_x \vec{E}_y \end{pmatrix}$$

$$= -\frac{i}{\omega \mu_0} \begin{pmatrix} i(\beta - \alpha) E_0 \cos[(\beta + i\alpha)x] e^{i(\beta - \alpha)z} \\ 0 \\ -(\beta + i\alpha) E_0 \sin[(\beta + i\alpha)x] e^{i(\beta - \alpha)z} \end{pmatrix} = \frac{i}{\omega \mu_0} \begin{pmatrix} i(\beta - \alpha) E_0 \cos[(\beta + i\alpha)x] e^{i(\beta - \alpha)z} \\ 0 \\ (\beta + i\alpha) E_0 \sin[(\beta + i\alpha)x] e^{i(\beta - \alpha)z} \end{pmatrix}$$

b). Wave equation:  $\Delta \vec{E}(\vec{r}) + \frac{\omega^2}{c^2} \epsilon(\omega) \vec{E}(\vec{r}) = 0$

$$\Delta \vec{E}(\vec{r}) = \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \vec{E}(\vec{r}) = \left[ -(\beta + i\alpha)^2 + (\beta - \alpha)^2 \right] \vec{E}(\vec{r})$$

$$\therefore -(\beta + i\alpha)^2 + (\beta - \alpha)^2 + \frac{\omega^2}{c^2} \epsilon(\omega) = 0$$

$$2\alpha^2 - 2\beta^2 - 4\beta\alpha = -\frac{\omega^2}{c^2} \epsilon(\omega) = -\frac{\omega^2}{c^2} (\epsilon' + i\epsilon'')$$

$$\begin{cases} \beta^2 - \alpha^2 = \frac{\omega^2}{c^2} \epsilon' \\ \beta\alpha = \frac{\omega^2}{c^2} \epsilon'' \end{cases} \Rightarrow \begin{cases} \beta \approx \frac{\omega}{\sqrt{c^2 + \epsilon'(\omega)}} \\ \alpha \approx \frac{1}{2\sqrt{c^2 + \epsilon'(\omega)}} \frac{\omega^2}{c^2} \epsilon''(\omega) \end{cases}$$

c)  $\epsilon'' = 0 \rightarrow \alpha = 0 \quad \beta = \frac{\omega}{\sqrt{c^2 + \epsilon'}} \sqrt{\epsilon'}$

$$\vec{E}(\vec{r}) = E_0 \cos\left(\frac{\omega}{\sqrt{c^2 + \epsilon'}} x\right) \exp\left(i \frac{\omega}{\sqrt{c^2 + \epsilon'}} z\right) \cdot \hat{y}$$

$$= E_0 \cos\beta x \cdot e^{i\beta z} \hat{y}$$

$$\vec{E}(\vec{r}) = \bar{E}_0 \cos \beta x \cdot e^{i\beta z} \hat{y}$$

$$H(\vec{r}) = \frac{1}{\mu_0} \begin{pmatrix} i\beta \bar{E}_0 \cos \beta x \cdot e^{i\beta z} \\ 0 \\ \beta \bar{E}_0 \sin \beta x \cdot e^{i\beta z} \end{pmatrix}$$

$$\langle S(\vec{r}, t) \rangle = \frac{1}{2} \operatorname{Re} [\vec{E} \times H^*]$$

$$= \frac{1}{2} \operatorname{Re} \left[ \begin{pmatrix} 0 \\ \bar{E}_0 \cos \beta x e^{i\beta z} \\ 0 \end{pmatrix} \times \frac{-i}{\mu_0} \begin{pmatrix} -i \bar{E}_0 \cos \beta x \cdot e^{-i\beta z} \\ 0 \\ \beta \bar{E}_0 \sin \beta x \cdot e^{-i\beta z} \end{pmatrix} \right]$$

$$= \frac{1}{2\mu_0} \begin{pmatrix} -i \bar{E}_0^2 \sin \beta x \cos \beta x \\ 0 \\ \beta \bar{E}_0^2 \cos^2 \beta x \end{pmatrix}$$

$$\langle S(\vec{r}, t) \rangle = \frac{1}{2\mu_0} \cdot \beta \bar{E}_0^2 \cos^2 \beta x \cdot \frac{1}{2}$$

## Beam Propagation

### Tutorial Problems.

Beam propagation.

a)  $H(\alpha, \beta, z) = i \sqrt{k^2 - \alpha^2 - \beta^2} \cdot e^z$

b)  $\alpha^2 + \beta^2 \ll k^2$

paraxial approximation:  $\sqrt{k^2 - \alpha^2 - \beta^2} = \sqrt{k^2 - r^2} = f(r)$ .

$$f(0) = k$$

$$f'(r) = \frac{-2r}{2\sqrt{k^2 - r^2}} = -\frac{r}{\sqrt{k^2 - r^2}}, \quad f'(0) = 0$$

$$f''(r) = -\frac{\sqrt{k^2 - r^2} + r \cdot \frac{r}{\sqrt{k^2 - r^2}}}{k^2 - r^2}, \quad f''(0) = -\frac{1}{k}$$

$$\therefore f(r) = f(0) + \frac{1}{2!} f''(0) r^2 = k - \frac{\alpha^2 + \beta^2}{2k} r^2$$

$$\therefore H_p(\alpha, \beta, z) = \exp(i k z) \exp(-i \frac{\alpha^2 + \beta^2}{2k} z^2)$$

c)  $U_0(\alpha, z=0) = \underbrace{A \delta(\alpha)}_{\text{constant}} + \underbrace{B \left[ \delta(\alpha - \frac{\pi}{L}) + \delta(\alpha + \frac{\pi}{L}) \right]}_{\text{pair of poles}}$

$$U(\alpha, z) = \cancel{H(\alpha, z)} \cdot H(\alpha, z) \cdot U_0(\alpha, 0)$$

$$= \cancel{\frac{A \delta(\alpha)}{\alpha}} \exp(i \sqrt{k^2 - \alpha^2} z) + \cancel{\frac{B}{\alpha}} \left[ \delta(\alpha - \frac{\pi}{L}) + \delta(\alpha + \frac{\pi}{L}) \right] \exp(i \sqrt{k^2 - \alpha^2} z)$$

$U(x, z) = \mathcal{F}^{-1}[U(\alpha, z)]$

$$= \cancel{\frac{A}{2\pi} e^{ikz}} + \cancel{\frac{B}{2\pi} \left[ e^{i\sqrt{k^2 - \frac{4\pi^2}{L^2}} z} \cdot e^{i\frac{\pi}{L} x} + e^{i\sqrt{k^2 - \frac{4\pi^2}{L^2}} z} \cdot e^{-i\frac{\pi}{L} x} \right]}$$

$$= \cancel{\frac{A}{2\pi} e^{ikz}} + \cancel{\frac{B}{2\pi} e^{i\sqrt{k^2 - \frac{4\pi^2}{L^2}} z} \cdot \cos \frac{\pi}{L} x}$$

$$= A e^{ikz} + B e^{i\sqrt{k^2 - \frac{4\pi^2}{L^2}} z} \cdot \cos \frac{\pi}{L} x$$

$$d) u_0(x, z) = A e^{i k z} + B e^{i \sqrt{k^2 - \frac{4\pi^2}{L^2}} z} \cdot \cos \frac{2\pi}{L} x$$

$$= e^{i k z} \left[ A + B \exp[i(\sqrt{k^2 - \frac{4\pi^2}{L^2}} - k)z] \cdot \cos \frac{2\pi}{L} x \right]$$

First reappearance:

$$\left( \sqrt{k^2 - \frac{4\pi^2}{L^2}} - k \right) \cdot z = 2\pi$$

$$z = \frac{2\pi}{\sqrt{\frac{4\pi^2}{L^2} - \frac{4\pi^2}{L^2}} - \frac{2\pi}{\lambda_0}}$$

$$z = \frac{1}{\sqrt{\frac{1}{\lambda_0^2} - \frac{1}{L^2}} - \frac{1}{\lambda_0}}$$

$$e) \alpha^2 \ll k^2 \text{ i.e. } \frac{2\pi}{L} \ll \frac{\pi}{\lambda_0} \rightarrow L \gg \lambda_0$$

Under Fresnel approximation.

$$U(x, z) = A [j(x) \exp(i k z) \exp(-i \frac{\alpha^2}{2k} z) + B [j(x - \frac{z}{L}) + j(x + \frac{z}{L})] \exp(i k z) \exp(-i \frac{\alpha^2}{2k} z)]$$

$$U(x, z) = A \exp(i k z) + B \exp(i k z) \left[ e^{-i \frac{4\pi^2}{2k L^2} z} e^{i \frac{2\pi}{L} x} + e^{-i \frac{4\pi^2}{2k L^2} z} e^{-i \frac{2\pi}{L} x} \right]$$

$$= A \exp(i k z) + B \exp(i k z) \exp(-i \frac{4\pi^2}{2k L^2} z) \cos \frac{2\pi}{L} x$$

$$= \exp(i k z) \left[ A + B \exp(-i \frac{4\pi^2}{2k L^2} z) \cos \frac{2\pi}{L} x \right]$$

First reappearance:

$$\frac{4\pi^2}{2k L^2} \cdot z = 2\pi$$

$$z = \frac{2\pi}{\lambda_0}$$

## Gaussian Beams

Tutorial problems

Gaussian Beams.

a).  $q(z) = z - i \frac{w_0^2}{z}$

$$\frac{1}{q(z)} = \frac{1}{z + \frac{w_0^2}{z}} + \frac{\frac{1}{w_0^2}}{1 + (\frac{z}{w_0})^2}$$

$$R(z) = z + \frac{w_0^2}{z}$$

$$w(z) = w_0 \sqrt{1 + (\frac{z}{w_0})^2}$$

b).  $\frac{1}{q(z)} = \frac{1}{R(z)} + \frac{2w}{\pi w^2(z)}$

b)  $\tilde{M} = M_{p_2} M_d M_{p_1}$

$$= \begin{pmatrix} 1 & 0 \\ \frac{n_1 - n_2}{n_1 r} & \frac{n_2}{n_1} \end{pmatrix} \begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{n_2 - n_1}{n_2 r} & \frac{n_1}{n_2} \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ \frac{n_1 - n_2}{n_1 r} & \frac{n_2}{n_1} \end{pmatrix} \begin{pmatrix} 1 - \frac{n_2 - n_1 d}{n_2 r} & \frac{n_1}{n_2} d \\ -\frac{n_2 - n_1}{n_2 r} & \frac{n_1}{n_2} \end{pmatrix}$$

$$= \begin{pmatrix} 1 - \frac{n_2 - n_1}{n_2 r} d & \frac{n_1}{n_2} d \\ \frac{(n_1 - n_2)(2 + \frac{n_1 - n_2}{n_2 r} d)}{n_1 r} & \frac{n_1 - n_2}{n_2 r} d + 1 \end{pmatrix}$$

$$q = \frac{A q_1 + B}{C q_1 + D} = \frac{\left(1 - \frac{n_2 - n_1}{n_2 r}\right) q_1 + \frac{n_1}{n_2} d}{\frac{n_1 - n_2}{n_1 r} (2 + \frac{n_1 - n_2}{n_2 r} d) q_1 + \frac{n_1 - n_2}{n_2 r} d + 1}$$

θ

c) when  $d=0$ ,  $\hat{M}_{\text{lens}} = \begin{pmatrix} 1 & 0 \\ \frac{2(n_1-n_2)}{n_1 r} & 1 \end{pmatrix}$

Set:

$$q(0) = -iz_0$$

$$q(1) = L_1 - iz_0$$

$$q_{R_2}^* = \hat{M}_{\text{lens}} q(1) = \frac{q_1 + 0}{\frac{2(n_1-n_2)}{n_1 r} q_1 + 1}$$

$$= \frac{L_1 - iz_0}{\frac{2(n_1-n_2)}{n_1 r} (L_1 - iz_0) + 1} = \frac{L_1 - iz_0}{C(L_1 - iz_0) + 1} = \frac{L_1 - iz_0}{(L_1 + 1) - i(z_0)}$$

$$= \frac{(L_1 - iz_0) [(CL_1 + 1) + iCz_0]}{(CL_1 + 1)^2 + C^2 z_0^2}$$

$$\operatorname{Re} q_{R_2} = \frac{C \cdot (L_1^2 + z_0^2) + L_1}{(CL_1 + 1)^2 + C^2 z_0^2}, \quad q_2 = q_{R_2} + L_2. \quad \operatorname{Re} q_2 = 0. \text{ if waist.}$$

$$\operatorname{Re} q_{R_2} + L_2 = \frac{C \cdot (L_1^2 + z_0^2) + L_1 + L_2 (CL_1 + 1)^2 + L_2 C^2 z_0^2}{(CL_1 + 1)^2 + C^2 z_0^2} = 0$$

$$L_2 = - \frac{C \cdot (L_1^2 + z_0^2) + L_1}{(CL_1 + 1)^2 + C^2 z_0^2} \quad C = \frac{2(n_1-n_2)}{n_1 r}$$

$$L_2 = \frac{\frac{2(n_2-n_1)}{n_1 r} (L_1^2 + z_0^2) - L_1}{\left(\frac{2(n_1-n_2)}{n_1 r} + 1\right)^2 + \frac{4(n_1-n_2)^2}{n_1^2 r^2} z_0^2}.$$

Pulses

Tutorial Problems.

Pulses.

a)  $k(w) = k(w_0) + \frac{\partial k}{\partial w} \Big|_{w_0} (w - w_0) + \frac{1}{2} \frac{\partial^2 k}{\partial w^2} \Big|_{w_0} (w - w_0)^2$

$$\frac{1}{v_{ph}} = \frac{k(w_0)}{c}, \quad \frac{1}{v_g} = \frac{\partial k}{\partial w}, \quad \frac{\partial^2 k}{\partial w^2} = \frac{2}{\partial w} \left( \frac{1}{v_g} \right).$$

phase velocity      group velocity      dispersion

b)  $\frac{\partial k}{\partial w} = \frac{\partial c}{\partial w} n(w) = \frac{1}{c} \left[ n(w_0) + w_0 \frac{\partial n}{\partial w} \Big|_{w_0} \right] = \frac{1}{v_g}$

$v_g = \frac{c}{(A + 3Bw_0^2) + w_0(2Bw_0)} = \frac{c}{A + 3Bw_0^2}$

$$v_{g1} = \frac{c}{A + 3Bw_1^2}, \quad v_{g2} = \frac{c}{A + 3Bw_2^2}$$

If the 2<sup>nd</sup> pulse can catch the 1<sup>st</sup>,  $v_{g2} > v_{g1} \rightsquigarrow w_2 < w_1$

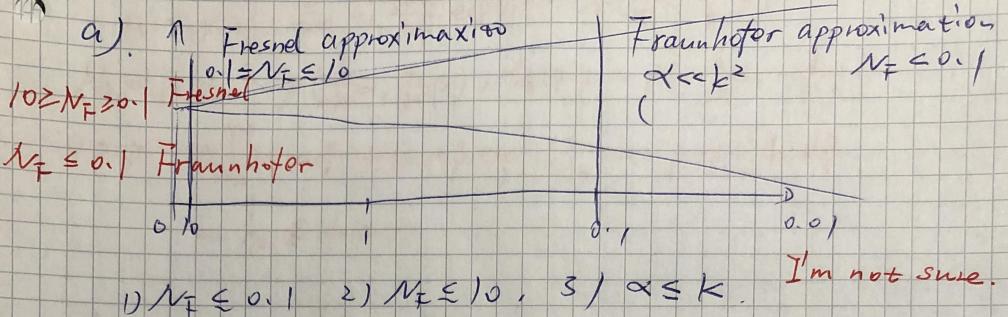
$$t = \frac{v_{g1} \Delta t}{(v_{g2} - v_{g1})} = \frac{\frac{c}{A + 3Bw_1^2} \cdot \Delta t}{\left( \frac{c}{A + 3Bw_2^2} - \frac{c}{A + 3Bw_1^2} \right)}$$

$$= \frac{\Delta t}{\frac{A + 3Bw_1^2}{A + 3Bw_2^2} - 1}$$

## Fraunhofer diffraction

Tutorial problems.

Fraunhofer diffraction.



b)  $\text{rect}(a) \sim \text{sinc}(a)$ .

$$\therefore \text{rect}\left(\frac{b}{z}\right) \sim \text{sinc}\left(\alpha \frac{b}{z}\right) \sim \text{sinc}\left(\frac{kb}{2z}\right).$$

shift theorem:  $f(x-a) \rightarrow \exp(-i\alpha a)$

$$f(x - \frac{a}{z}) \rightarrow \exp\left(-i\frac{k}{z} \cdot \frac{a}{z}\right).$$

$$f(x + \frac{a}{z}) \rightarrow \exp\left(i\frac{k}{z} \cdot \frac{a}{z}\right).$$

$\therefore U(x)$ :

$$= u(x, z) = \text{sinc}\left(\frac{kb}{2z}\right) \cdot \left[ \exp\left(-i\frac{k}{z} \cdot \frac{a}{z}\right) + \exp\left(i\frac{k}{z} \cdot \frac{a}{z}\right) \right]$$

$$= 2 \text{sinc}\left(\frac{kb}{2z}\right) \cdot \cos\left(\frac{ka}{2z}\right).$$

$$I = |u(x, z)|^2 = 4 \text{sinc}^2\left(\frac{kb}{2z}\right) \cos^2\left(\frac{ka}{2z}\right) \quad \text{not precise}$$

c) ~~a~~  $a$  defines Fraunhofer and  $b$  defines Fresnel.

Fraunhofer:  $\frac{a}{z_B} \cdot \frac{1}{z_B} \leq 0.1 \rightsquigarrow a^2 \leq 0.1 \lambda z_B$

Fresnel:  ~~$\frac{b}{z_B} \cdot \frac{1}{z_B} \leq 10 \rightsquigarrow b^2 \leq 10 \lambda z_B$~~

$b > 10\lambda$

Paraxial Approximation:  $|dx|, |dy| > |o\lambda_n|$

Fourier Optics

Tutorial Problems.

Fourier Optics

$$U_1(x, y, z_f) = -i \frac{(2\pi)^2}{\lambda f} \exp(i k z_f) U_0 \left( \frac{k}{f_1} x, \frac{k}{f_1} y \right)$$

$$H_1(\alpha, \beta, z_f) \sim p \left( \frac{f_1}{k} \alpha, \frac{f_1}{k} \beta \right); H_0(x, y) \sim P \left[ -\frac{k}{f_2} x, -\frac{k}{f_2} y \right]$$

$$U_1(x, y, z_f) \sim \int_{-\infty}^{\infty} H_1(\alpha, \beta, z_f) H_0(x, \beta) \exp \left[ i \left( \frac{f_1}{f_2} \alpha x + \frac{f_1}{f_2} y \beta \right) \right] d\alpha d\beta$$

$$\sim h_0 \otimes U_0(x, y).$$

a)  $U_1(x, y, z_f) \sim U_0 \left( \frac{k}{f_1} x, \frac{k}{f_1} y \right)$

$$\bar{T}(\alpha) = \bar{F} \bar{F}^{-1} [t(x)] = \frac{1}{2} \delta(\alpha) + \frac{1}{4} \left[ S(\alpha - \frac{2\pi}{d}) + S(\alpha + \frac{2\pi}{d}) \right]$$

$$\alpha = \frac{k}{f_1} x = \frac{2\pi}{\lambda f_1} x$$

$$U_0 \left( \frac{k}{f_1} x, f_1 \right) = \frac{1}{2} \delta(\alpha) + \frac{1}{2} S \left( \frac{2\pi}{d} x \right) + \frac{1}{4}$$

$$= \frac{1}{2} S \left( \frac{2\pi}{\lambda f_1} x \right) + \frac{1}{4} \left[ S \left( \frac{2\pi x}{\lambda f_1} - \frac{2\pi}{d} \right) + S \left( \frac{2\pi x}{\lambda f_1} + \frac{2\pi}{d} \right) \right]$$

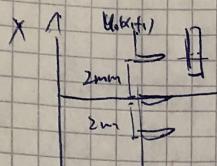
peaks:  $x=0$ ,  $\left| x_1 = \frac{\lambda f_1}{d} \right| \quad \left| x_2 = -\frac{\lambda f_1}{d} \right|$   
 $\underline{x_1 = 0.2 \text{ cm}} \quad \underline{x_2 = -0.2 \text{ cm}}$

b) shift.

$$\Delta s = (n_{\text{glass}} - 1) w = 100.5 \mu\text{m} \quad \rightarrow p(x, y) = \begin{cases} e^{ix} = -1 & 1 \text{ mm} < x < 3 \text{ mm} \\ 1 & \text{elsewhere.} \end{cases}$$

$$H_1(\alpha, z_f) \sim p \left( \frac{f_1}{k} \alpha \right) = \begin{cases} -1 & 1 < \frac{f_1}{k} \alpha < 3 \\ 1 & \text{elsewhere.} \end{cases}$$

$$\bar{T}_0(x) = \frac{1}{2} \delta(\alpha) + \frac{1}{4} \left[ S \left( \alpha - \frac{2\pi}{d} \right) + S \left( \alpha + \frac{2\pi}{d} \right) \right]$$



$\nabla \times \vec{E}_0$

$$\therefore U_0^+ \left( \frac{k}{\lambda} x, -z \right) = \frac{1}{2} \delta \left( \frac{2\pi}{\lambda f_1} x \right) + \frac{1}{4} \left[ -\delta \left( \frac{2\pi x}{\lambda f_1} - \frac{\pi}{2} \right) + \delta \left( \frac{2\pi x}{\lambda f_1} + \frac{\pi}{2} \right) \right] = U_0(x, z)$$

$$U_1(x, z) = \frac{1}{2} \delta \left( \frac{2\pi}{\lambda f_1} x \right) + \frac{1}{4} \left[ \delta \left( \frac{2\pi x}{\lambda f_1} + \frac{\pi}{2} \right) - \delta \left( \frac{2\pi x}{\lambda f_1} - \frac{\pi}{2} \right) \right] = U_1$$

$$U_1(x, z) \sim U_1 \left( \frac{k}{\lambda} x \right)$$

$$\begin{aligned} U_1 \left( \frac{k}{\lambda} x \right) &= \int_{-\infty}^{\infty} \frac{1}{2} \delta \left( \frac{2\pi}{\lambda f_1} x' \right) \cdot \exp \left( -i \frac{k}{\lambda} x' x \right) dx' + \frac{1}{4} \left[ \int_{-\infty}^{\infty} \left[ \delta \left( \frac{2\pi x'}{\lambda f_1} + \frac{\pi}{2} \right) - \delta \left( \frac{2\pi x'}{\lambda f_1} - \frac{\pi}{2} \right) \right] \exp \left( -i \frac{k}{\lambda} x' x \right) dx' \right] \\ &= \frac{1}{2} + \frac{1}{4} \left[ \exp \left( i \frac{1}{f_1} \frac{2\pi}{\lambda} x \right) - \exp \left( -i \frac{1}{f_1} \frac{2\pi}{\lambda} x \right) \right] \\ &= \frac{1}{2} + \frac{1}{2} \sin \left( \frac{1}{f_1} \frac{2\pi}{\lambda} x \right). \end{aligned}$$

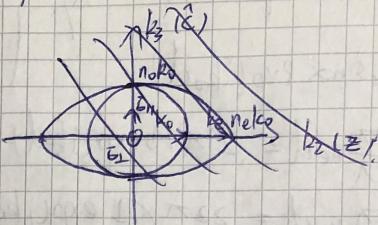
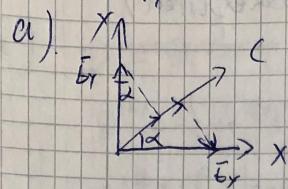
$$\therefore U(x, z) = \frac{1}{2} + \frac{1}{2} \sin \left( \frac{4\pi}{\lambda} x \right)$$

$$C) I = U \cdot U^* = \frac{1}{4} + \frac{1}{4} \sin^2 \left( \frac{4\pi}{\lambda} x \right)$$

## Anisotropy 1

Tutorial problem

Anisotropy 1.

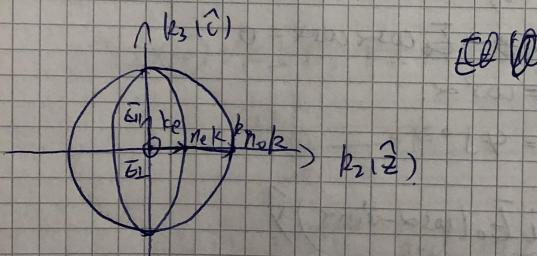


$$x: \begin{cases} \bar{E}_{11} = \bar{E}_x \cos \alpha \exp(i k_0 z) \\ \bar{E}_{\perp} = \bar{E}_x \sin \alpha \exp(i k_0 z) \end{cases}$$

$$y: \begin{cases} \bar{E}_{11} = \bar{E}_y \sin \alpha \exp(i k_0 z) \\ \bar{E}_{\perp} = \bar{E}_y \cos \alpha \exp(i k_0 z), \end{cases}$$

$$k_0 = k \cdot n_0 = \frac{2\pi}{\lambda} \cdot n_0 = 13.82 \text{ } (\mu\text{m}^{-1})$$

$$k_e = k - n_e = \frac{2\pi}{\lambda} \cdot n_e = 13.19 \text{ } (\mu\text{m}^{-1})$$



$$\vec{E} = \bar{E}_x (\cos \alpha \exp(i k_0 z) + \sin \alpha \exp(i k_0 z)) \hat{x} + \bar{E}_y (\sin \alpha \exp(i k_0 z) + \cos \alpha \exp(i k_0 z)) \hat{y}$$

~~$$\Delta \varphi = k \cdot \Delta s = (k_0 - k_e) \cdot \Delta s = \pi$$~~

It is a half wave plate.

$$b) \quad X: \begin{cases} E_{11} = \bar{E}_x \cos \alpha \exp(i k_0 d) \\ E_{\perp} = \bar{E}_x \sin \alpha \exp(i k_0 d) \end{cases} \quad Y: \begin{cases} \bar{E}_{11} = \bar{E}_y \sin \alpha \exp(i k_0 d) \\ \bar{E}_{\perp} = \bar{E}_y \cos \alpha \exp(i k_0 d) \end{cases}$$

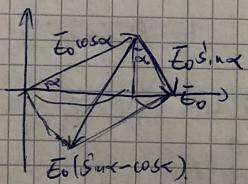
$$k_0 d = \frac{2\pi}{\lambda} n_0 d = 21\pi / \lambda \exp(i k_0 d) = -1$$

$$k_0 d = \frac{2\pi}{\lambda} n_0 d = 22\pi / \lambda \exp(i k_0 d) = 1$$

$$\therefore X: \begin{cases} \bar{E}_{11} = -\bar{E}_x \cos \alpha \\ \bar{E}_{\perp} = \bar{E}_x \sin \alpha \end{cases} \quad Y: \begin{cases} \bar{E}_{11} = -\bar{E}_y \sin \alpha \\ \bar{E}_{\perp} = \bar{E}_y \cos \alpha \end{cases}$$

$$\vec{E}_{\text{out}} = \bar{E}_x (\sin \alpha - \cos \alpha) \hat{x} + \bar{E}_y (\cos \alpha - \sin \alpha) \hat{y}$$

$$c) \quad \vec{E}_{\text{out}} = \bar{E}_0 (\sin \alpha - \cos \alpha) \hat{x}$$



$$\text{if } \vec{E}_{\text{out}} = \bar{E}_0 \hat{y}$$

$$\bar{E}_0 \sin \alpha \sin \alpha - \bar{E}_0 \cos \alpha \cos \alpha = 0$$

$$\sin^2 \alpha = \cos^2 \alpha$$

$$\alpha = 45^\circ$$

$$d) \quad \vec{E}_{\text{out}} = \bar{E}_0 (\sin \alpha - \cos \alpha) \hat{x} + i (\bar{E}_0 (\cos \alpha - \sin \alpha)) \hat{y}$$

$$= \bar{E}_0 \hat{x} - i \bar{E}_0 \hat{y}$$

$$\begin{cases} \sin \alpha - \cos \alpha = 1 \\ \cos \alpha - \sin \alpha = 1 \end{cases} \rightarrow \begin{cases} \sin \alpha = 1 \\ \cos \alpha = 0 \end{cases} \rightarrow \alpha = 90^\circ$$

## Anisotropy 2

Tutorial Problems

Anisotropy 2.

a)  $D_1 = D_1 \exp i\omega t$

$$D^{(a)} = \{D_1 \exp[i(\vec{k}_a \cdot \vec{r} - \omega t)]\} \hat{e}_1$$

$$k_a^2 = \frac{\omega^2}{c^2} n_a^2$$

$$D^{(b)} = \{D_2 \exp[i(\vec{k}_b \cdot \vec{r} - \omega t)]\} \hat{e}_2$$

$$k_b^2 = \frac{\omega^2}{c^2} n_b^2$$

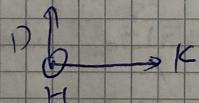
b) Start from Maxwell's equations for plane wave.

$$\textcircled{1} \vec{k} \cdot \vec{D} = 0 \quad \textcircled{3} \vec{c} \times \vec{E} = \omega \mu_0 \vec{H} \quad \text{and} \textcircled{5} \langle S \rangle = \vec{E} \times \vec{H}$$

$$\textcircled{2} \vec{k} \cdot \vec{H} = 0 \quad \textcircled{4} \vec{k} \times \vec{H} = -\omega \vec{D}$$

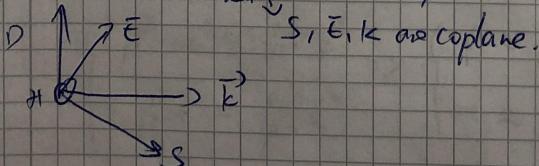
~~$\textcircled{6} \vec{k} \perp \vec{D} + \vec{H}$~~

$$\textcircled{1} \textcircled{2} \textcircled{4}: \vec{k} \perp \vec{D}, \vec{k} \perp \vec{H}, \vec{D} + \vec{H} \rightsquigarrow \vec{H} \perp \vec{D} \vec{k}$$



\textcircled{3}:  $\vec{H} \perp \vec{E} \vec{k}$  and  $\vec{H} \perp \vec{D} \vec{k} \rightsquigarrow \vec{E}, \vec{D}, \vec{k}$  are coplanar.

\textcircled{5}:  $S \perp \vec{E} \vec{H} \vec{k}$ . Generally,  $\vec{E} \perp \vec{D}, \therefore S \perp \vec{K}$ .



c)

