

Task 1

a)

Solution:

$$I = |\langle \vec{S}(\vec{r}, t) \rangle|$$

$$\langle \vec{S}(\vec{r}, t) \rangle = \frac{1}{2} \text{Re} [\vec{E}(\vec{r}) \times \vec{H}(\vec{r})]$$

$$= \frac{1}{2} \text{Re} \left[\frac{\partial A_0}{\omega_0 \mu_0} \vec{e}_2 + \frac{\beta A_1}{\omega_0 \mu_0} \vec{e}_2 \right]$$

$$= \frac{\partial A_0 + \beta A_1}{2 \omega_0 \mu_0} \vec{e}_2$$

b)

$$\vec{E}(z, t) = A_0 \exp[i(\alpha z - \omega_0 t)] \vec{e}_x + A_1 \exp[i(\beta z - \omega_0 t)] \vec{e}_y$$

$$I = \vec{E}(\vec{r}) \cdot \vec{E}(\vec{r})$$

$$= A_0^2 + A_1^2$$

$$\therefore \vec{E}(\vec{r}, \omega) = A_0 e^{i\alpha z} \delta(\omega - \omega_0) \vec{e}_x + A_1 e^{i\beta z} \delta(\omega - \omega_0) \vec{e}_y$$

$$\text{From Maxwell's equations } \vec{H}(\vec{r}, \omega) = -\frac{i}{\omega \mu_0} \text{rot } \vec{E}(\vec{r}, \omega)$$

$$\therefore \vec{H}(\vec{r}, \omega) = -\frac{i}{\omega \mu_0} [i\alpha A_0 e^{i\alpha z} \delta(\omega - \omega_0) \vec{e}_y - i\beta A_1 e^{i\beta z} \delta(\omega - \omega_0) \vec{e}_x]$$

$$= \frac{\alpha}{\omega \mu_0} A_0 e^{i\alpha z} \delta(\omega - \omega_0) \vec{e}_y - \frac{\beta}{\omega \mu_0} A_1 e^{i\beta z} \delta(\omega - \omega_0) \vec{e}_x$$

$$\vec{H}(\vec{r}, t) = \frac{\alpha A_0}{\omega_0 \mu_0} e^{[i(\alpha z - \omega_0 t)]} \vec{e}_y - \frac{\beta A_1}{\omega_0 \mu_0} e^{[i(\beta z - \omega_0 t)]} \vec{e}_x$$

Task 2

a)

Solution:

$$U_0(\alpha, \beta) = \frac{1}{4\pi^2} \iint_{-\infty}^{\infty} u_0(x, y) \exp[-i(\alpha x + \beta y)] dx dy$$

$$= \frac{1}{4\pi^2} \iint_{-\infty}^{\infty} A e^{-\frac{x^2+y^2}{W^2}} e^{-i(\alpha x + \beta y)} dx dy$$

$$= \frac{A}{4\pi^2} \int_{-\infty}^{\infty} e^{-\frac{x^2}{W^2} - i\alpha x} dx \int_{-\infty}^{\infty} e^{-\frac{y^2}{W^2} - i\beta y} dy$$

$$= \frac{A}{4\pi^2} \int_{-\infty}^{\infty} e^{-\frac{1}{W^2}(x + \frac{i\alpha W^2}{2})^2} e^{-\frac{\alpha^2 W^2}{4}} dx \int_{-\infty}^{\infty} e^{-\frac{1}{W^2}(y + \frac{i\beta W^2}{2})^2} e^{-\frac{\beta^2 W^2}{4}} dy$$

$$= \frac{A}{4\pi^2} \left[e^{-\frac{\alpha^2 W^2}{4}} \int_{-\infty}^{\infty} e^{-\frac{1}{W^2}(x + \frac{i\alpha W^2}{2})^2} dx \right] \left[e^{-\frac{\beta^2 W^2}{4}} \int_{-\infty}^{\infty} e^{-\frac{1}{W^2}(y + \frac{i\beta W^2}{2})^2} dy \right]$$

$$= \frac{A}{4\pi^2} \left[e^{-\frac{\alpha^2 W^2}{4}} \int_{-\infty}^{\infty} e^{-\frac{1}{W^2}(x + \frac{i\alpha W^2}{2})^2} d[\frac{1}{W}(x + \frac{i\alpha W^2}{2})] \right] \left[e^{-\frac{\beta^2 W^2}{4}} \int_{-\infty}^{\infty} e^{-\frac{1}{W^2}(y + \frac{i\beta W^2}{2})^2} d[\frac{1}{W}(y + \frac{i\beta W^2}{2})] \right]$$

From Gaussian integral

$$\int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi}$$

$$\therefore U_0(\alpha, \beta) = \frac{A}{4\pi^2} (W\sqrt{\pi} e^{-\frac{\alpha^2 W^2}{4}} \cdot W\sqrt{\pi} e^{-\frac{\beta^2 W^2}{4}}) = \frac{AW^2}{4\pi} e^{-\frac{(\alpha^2 + \beta^2)W^2}{4}}$$

b)

Solution: $\because k_r = \sqrt{k_x^2 + k_y^2}$

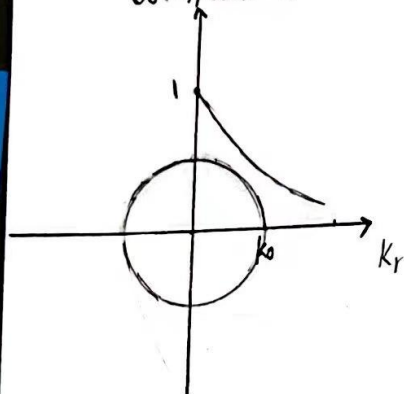
$$U_0(k_r = k_0) / U_0(k_r = 0) = \frac{\frac{AW^2}{4\pi} e^{-\frac{k_0^2 W^2}{4}}}{\frac{AW^2}{4\pi}} = e^{-\frac{k_0^2 W^2}{4}} = e^{-\frac{4\pi^2 W^2}{4\lambda^2}} = e^{-\frac{\pi^2 W^2}{\lambda^2}}$$

$$W = \lambda / \pi$$

$$U_0(k_r) / U_0(k_r = 0) = e^{-\frac{k_r^2 \lambda^2}{4\pi^2}}$$

$$U_0(k_r = k_0) / U_0(k_r = 0) = \frac{1}{e}$$

$U_0(k_r) / U_0(k_r = 0)$



(c)

Solution:

Draw a circle with radius = k_0

The area in the circle ($k_x^2 + k_y^2 \leq k_0^2$) contributes to propagating wave (real part of k_z)

The area beyond the circle ($k_x^2 + k_y^2 > k_0^2$) contributes to evanescent wave (imaginary part of k_z)

Task 3

a)

Solution:

In Talbot Effect, we can get

$$f(x, z=L_T) = f(x, z=0) \exp(ikL_T + i2\pi m_L)$$

$$f(x+a) = f(x)$$

Express $f(x)$ as a Fourier series

$$f(x) = \sum_{n=-\infty}^{+\infty} C_n e^{in\omega_0 x} \quad (n \in \mathbb{Z})$$

$$\because T=a \therefore \omega_0 = \frac{2\pi}{T} = \frac{2\pi}{a}$$

$$f(x) = \sum_{n=-\infty}^{+\infty} C_n e^{in\frac{2\pi}{a}x}$$

$$\begin{aligned} F\{f(x, z=0)\} &= F(\alpha, z=0) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \sum_{n=-\infty}^{+\infty} C_n e^{in\frac{2\pi}{a}x} e^{-i\alpha x} dx \\ &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} \sum_{n=-\infty}^{+\infty} C_n e^{-i(\alpha - n\frac{2\pi}{a})x} dx \\ &= \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} C_n \delta(\alpha - n\frac{2\pi}{a}) \end{aligned}$$

$$F(\alpha, z) = F(\alpha, z=0) \cdot H(\alpha, z)$$

$$H(\alpha, z) = \exp[i\sqrt{k^2 - \alpha^2} z]$$

$$\therefore F(\alpha, z) = \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} C_n \delta(\alpha - n\frac{2\pi}{a}) e^{i\sqrt{k^2 - \alpha^2} z}$$

$$f(x, z) = \int_{-\infty}^{+\infty} F(\alpha, z) e^{i\alpha x} d\alpha$$

$$= \int_{-\infty}^{+\infty} \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} C_n \delta(\alpha - n\frac{2\pi}{a}) e^{i\sqrt{k^2 - \alpha^2} z} e^{i\alpha x} d\alpha$$

$$= \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} C_n e^{i\sqrt{k^2 - (\frac{2n\pi}{a})^2} z} e^{i\frac{2n\pi}{a}x}$$

$$\therefore f(x, z=L_T) = f(x, z=0) e^{i(kL_T + 2\pi m_L)}$$

$$\therefore \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} e^{i\sqrt{k^2 - (\frac{2n\pi}{a})^2} L_T} e^{i\frac{2n\pi}{a}x} = \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} C_n e^{i\frac{2n\pi}{a}x} e^{i(kL_T + 2\pi m_L)}$$

$$\therefore e^{i\sqrt{k^2 - (\frac{2n\pi}{a})^2} L_T} = e^{i(kL_T + 2\pi m_L)}$$

$$\sqrt{k^2 - (\frac{2n\pi}{a})^2} L_T = kL_T + 2\pi m_L$$

$$L_T = \frac{2\pi m_L}{\sqrt{k^2 - (\frac{2n\pi}{a})^2} - k}$$

\therefore The paraxial is valid

$$\therefore \alpha^2 \ll k^2$$

$$\sqrt{k^2 - (\frac{2n\pi}{a})^2} \approx k - \frac{(\frac{2n\pi}{a})^2}{2k}$$

$$\therefore L_T \approx \frac{-k a^2 m_L}{n^2 \pi}$$

b) solution:

From (a) we can get

$$L_T = -\frac{\kappa a^2 m_L}{n^2 \pi}$$

$$\because m_L, n \in \mathbb{Z}$$

$$\therefore \text{let } m_L = -n^2$$

$$L_T = \frac{\kappa a^2}{\pi} = \frac{2a^2}{\lambda}$$

$$\therefore L_T = 25.7 \mu\text{m} \quad a = 3 \times 10^{-3} \text{ m}$$

$$\therefore \lambda = 0.7 \times 10^{-6} \text{ m}$$

c)

Solution:

$$f(x, z=0) = A \cos(x 2\pi/a_1) \cos(x 2\pi/a_2)$$

$$= \frac{A}{2} \left[\cos\left(\frac{2\pi}{a_1}x + \frac{2\pi}{a_2}x\right) + \cos\left(\frac{2\pi}{a_1}x - \frac{2\pi}{a_2}x\right) \right]$$

$$F(\omega, z=0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{A}{2} \left[\cos\left(\frac{2\pi}{a_1}x + \frac{2\pi}{a_2}x\right) + \cos\left(\frac{2\pi}{a_1}x - \frac{2\pi}{a_2}x\right) \right] e^{-i\omega x} dx$$

$$= \frac{A}{4\pi} \int_{-\infty}^{\infty} \left[\frac{e^{i(\frac{2\pi}{a_1} + \frac{2\pi}{a_2})x} + e^{-i(\frac{2\pi}{a_1} + \frac{2\pi}{a_2})x}}{2} + \frac{e^{i(\frac{2\pi}{a_1} - \frac{2\pi}{a_2})x} + e^{-i(\frac{2\pi}{a_1} - \frac{2\pi}{a_2})x}}{2} \right] e^{-i\omega x} dx$$

Suppose $\frac{2\pi}{a_1} + \frac{2\pi}{a_2} = M$ $\frac{2\pi}{a_1} - \frac{2\pi}{a_2} = N$

$$F(\omega, z=0) = \frac{A}{4\pi} \int_{-\infty}^{\infty} \frac{1}{2} \left[e^{i(M-\omega)x} + e^{-i(M-\omega)x} + e^{i(N-\omega)x} + e^{-i(N-\omega)x} \right] dx$$

$$= \frac{A}{8\pi} \left[\delta(M-\omega) + \delta(M+\omega) + \delta(N-\omega) + \delta(N+\omega) \right]$$

$$F(\omega, z) = F(\omega, z=0) e^{i(k^2 - \omega^2)z}$$

$$f(x, z) = \int_{-\infty}^{\infty} F(\omega, z=0) e^{i(k^2 - \omega^2)z} e^{i\omega x} d\omega$$

$$= \int_{-\infty}^{\infty} \frac{A}{8\pi} \left[\delta(M-\omega) + \delta(M+\omega) + \delta(N-\omega) + \delta(N+\omega) \right] e^{i(k^2 - \omega^2)z} e^{i\omega x} d\omega$$

$$= \frac{A}{8\pi} \left[e^{i(k^2 - M^2)z} e^{iMx} + e^{i(k^2 - M^2)z} e^{-iMx} + e^{i(k^2 - N^2)z} e^{iNx} + e^{i(k^2 - N^2)z} e^{-iNx} \right]$$

$$= \frac{A}{4\pi} \left[e^{i(k^2 - M^2)z} \cos(Mx) + e^{i(k^2 - N^2)z} \cos(Nx) \right]$$

$$\boxed{f(x, z=L_T) = f(x, z=0) \exp(ikL_T)}$$

$\therefore f(x, z=L_T) = f(x, z=0) \exp(ikL_T + i2\pi m_e)$ the Talbot effect still takes place outside the paraxial regime

$$\therefore \frac{A}{4\pi} \left[e^{i(k^2 - M^2)L_T} \cos(Mx) + e^{i(k^2 - N^2)L_T} \cos(Nx) \right] = \frac{A}{4\pi} \left[\cos(Mx) + \cos(Nx) \right] e^{i(kL_T + 2\pi m_e)}$$

$$\therefore \boxed{e^{i(k^2 - M^2)L_T} = e^{i(kL_T + 2\pi m_e)}} \quad \boxed{e^{i(k^2 - N^2)L_T} = e^{i(kL_T + 2\pi m_e)}}$$

$$L_T =$$