

- a) Calculate the intensity of the diffracted field pattern  $I(x, z_B) = |u(x, z_B)|^2$  at  $z = z_B$  in paraxial Fraunhofer approximation for two slits illuminated with a normally incident plane wave (just the diffraction pattern, no prefactors). The width of each slit is  $a$  ( $a > \lambda$ ) and separated by a distance  $d$  ( $d > a$ ):

$$u_0(x, z=0) = \begin{cases} 1, & \text{for } |x| \leq d/2 \\ 0, & \text{elsewhere.} \end{cases}$$

- b) Try to roughly sketch a figure of the intensity and identify the factor due to interference from the one due to slit diffraction.

*Hint: The Fourier transform of a single slit of width  $a$  is  $\propto \text{sinc}(ka)$ .*

$$\text{FT}[f(x, y, t)] = \frac{1}{(2\pi)^2} \iint f(x, y, t) e^{i(kx+ky-wt)} dxdy$$

$$u(x, t) = e^{i(kx-wt)}$$

$$\text{a) } h_F = -\frac{i k e^{ikz}}{2z_B} \exp\left[\frac{ik}{2z_B}(x^2 + y^2)\right] \quad (z=z_B \Rightarrow h_F = -\frac{i e^{ikz_B}}{\lambda z_B} \exp\left[\frac{ik}{2z_B}(x^2 + y^2)\right])$$

$$h_{FR} = -\frac{i \lambda}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B} x^2\right) U_+(\frac{kx}{2z_B}) \quad -\frac{a}{2} \leq x \leq \frac{d}{2} \quad -\frac{a}{2} \leq x \leq \frac{a}{2} - \frac{d}{2} \quad \frac{a-d}{2} \leq x \leq \frac{a+d}{2}$$

$$U_+\left(\frac{kx}{2z_B}\right) = \frac{1}{2\pi} \int_{-\infty}^{\infty} u_+(x', z=0) \exp\left(-\frac{ikx'}{2z_B}\right) dx' = \frac{1}{2\pi} \int_{-\frac{a-d}{2}}^{\frac{a-d}{2}} \exp\left(-\frac{ikx'}{2z_B}\right) dx' + \frac{1}{2\pi} \int_{\frac{a+d}{2}}^{\frac{a+d}{2}} \exp\left(-\frac{ikx'}{2z_B}\right) dx'$$

$$= \frac{1}{2\pi} \left[ -\frac{2z_B}{ikx} \exp\left(-\frac{ikx}{2z_B}\right) \Big|_{-\frac{a-d}{2}}^{\frac{a-d}{2}} - \frac{2z_B}{ikx} \exp\left(-\frac{ikx}{2z_B}\right) \Big|_{\frac{a+d}{2}}^{\frac{a+d}{2}} \right]$$

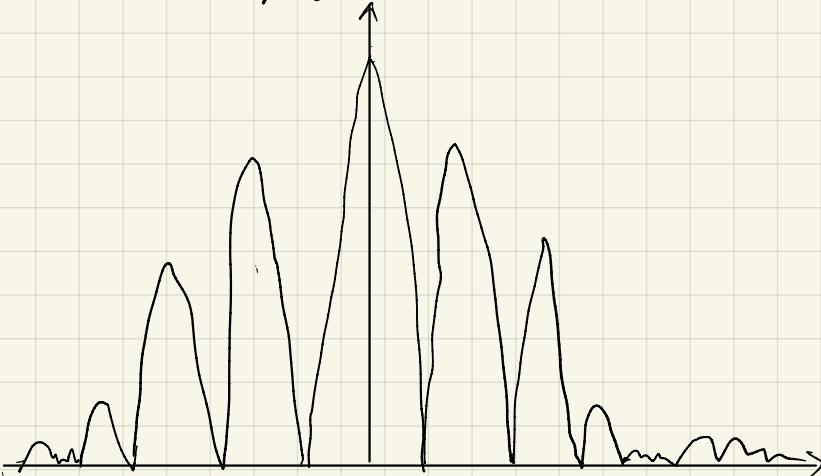
$$= -\frac{2z_B}{i2\pi kx} \left\{ \exp\left(\frac{ikx}{2z_B} \cdot \frac{d}{2}\right) \left[ \exp\left(-\frac{ikx}{2z_B} \frac{a}{2}\right) - \exp\left(\frac{ikx}{2z_B} \frac{a}{2}\right) \right] + \exp\left(\frac{ikx}{2z_B} \cdot \frac{d}{2}\right) \left[ \exp\left(-\frac{ikx}{2z_B} \frac{a}{2}\right) - \exp\left(\frac{ikx}{2z_B} \frac{a}{2}\right) \right] \right\}$$

$$= \frac{2z_B}{\pi kx} \left\{ \left[ \exp\left(\frac{ikx}{2z_B} \frac{d}{2}\right) + \exp\left(\frac{ikx}{2z_B} \frac{d}{2}\right) \right] \sin\left(\frac{ka}{2z_B} x\right) \right\} = \frac{2z_B}{\pi kx} \sin\left(\frac{ka}{2z_B} x\right) \cos\left(\frac{kd}{2z_B} x\right)$$

$$= \frac{2z_B}{\pi kx} \cdot \frac{|ka|}{2z_B} \sin\left(\frac{ka}{2z_B} x\right) \cos\left(\frac{kd}{2z_B} x\right) = \frac{a}{\pi} \sin\left(\frac{ka}{2z_B} x\right) \cos\left(\frac{kd}{2z_B} x\right)$$

$$h_{FR} = -\frac{i 2\pi}{\lambda z_B} \cdot \frac{a}{\pi} e^{ikz_B} \exp\left(\frac{ik}{2z_B} x^2\right) \sin\left(\frac{ka}{2z_B} x\right) \cos\left(\frac{kd}{2z_B} x\right) = -\frac{i 2a}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B} x^2\right) \sin\left(\frac{ka}{2z_B} x\right) \cos\left(\frac{kd}{2z_B} x\right)$$

$$\Rightarrow I = |h_{FR}|^2 = \frac{4a^2}{\lambda^2 z_B^2} \sin^2\left(\frac{ka}{2z_B} x\right) \cos^2\left(\frac{kd}{2z_B} x\right)$$



- a) Name two different approximations that can be made in diffraction theory and shortly explain their meaning in two sentences.

- b) Calculate the intensity of the diffracted field pattern  $I(x, z_B) = |u(x, z_B)|^2$  at  $z = z_B$  in paraxial Fraunhofer approximation for one slit illuminated with a normally incident plane wave (just the diffraction pattern, no prefactors). The width of the slit is  $a$  ( $a > \lambda$ ):

$$u_0(x, z=0) = \begin{cases} 1, & \text{for } |x| \leq a/2 \\ 0, & \text{elsewhere.} \end{cases}$$

(a) Paraxial Approximation: Approximation when the angles between light and axis are very small

Far field Approximation: The amplitude in the far field is given by the Fourier Transform of field behind aperture.

$$(b) h_{FR}(x, y, z_B) = -i \frac{(2\pi)^2}{\lambda z_B} e^{ikz_B} \exp\left[\frac{ik}{2z_B}(x^2 + y^2)\right] U_+\left(\frac{kx}{2z_B}, \frac{ky}{2z_B}\right)$$

$$U_f\left(\frac{kx}{2z_B}, \frac{ky}{2z_B}\right) = \left(\frac{1}{2\pi}\right)^2 \iint_{-\infty}^{\infty} u_0(x', z=0) \exp\left[-i\left(\frac{kx}{2z_B}x' + \frac{ky}{2z_B}y'\right)\right] dx' dy'$$

$$U_{FR}(x, z_B) = -i \frac{2\lambda}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B}x^2\right) U_f\left(\frac{kx}{2z_B}\right)$$

$$\begin{aligned} U_f\left(\frac{kx}{2z_B}, z_B\right) &= \frac{1}{2\pi} \int_{-\frac{a}{2}}^{\frac{a}{2}} u_0 \exp\left[-i\left(\frac{kx}{2z_B}x'\right)\right] dx' = \frac{1}{2\pi} \int_{-\frac{a}{2}}^{\frac{a}{2}} \exp\left(-i\frac{kx}{2z_B}x'\right) dx' = \frac{1}{2\pi} \cdot \frac{z_B}{-ikx} \exp\left(-i\frac{kx}{2z_B}x'\right) \Big|_{-\frac{a}{2}}^{\frac{a}{2}} \\ &= \frac{1}{2\pi} \cdot \frac{z_B}{-ikx} \left[ \exp\left(-i\frac{kx}{2z_B}\frac{a}{2}\right) - \exp\left(i\frac{kx}{2z_B}\frac{a}{2}\right) \right] = \frac{z_B}{2\pi kx} \sin\left(\frac{kxa}{2z_B}\right) = \frac{z_B}{\pi} \frac{\sin(kxa)}{\frac{kxa}{2z_B}} \cdot \frac{a}{2z_B} \\ &= \frac{a}{2\pi} \sin\left(\frac{kxa}{2z_B}\right) \end{aligned}$$

$$\Rightarrow U_{FR} = -i \frac{a}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B}x^2\right) \sin\left(\frac{kxa}{2z_B}\right) \quad | = |U_{FR}|^2 = \frac{a^2}{(\lambda z_B)^2} \sin^2\left(\frac{kxa}{2z_B}\right)$$

A one-dimensional optical field directly behind a specific optical element is given as

$$u_0(x, 0) = \begin{cases} A \exp[i\Phi_0(\frac{x}{a} + 1)] & |x| \leq a \\ 0 & \text{otherwise} \end{cases}$$

- a) Find the optical intensity  $I(x, z)$  in the case that the paraxial approximation holds and the distance  $z \gg a$ . You may omit possible prefactors.

- b) How large does  $\Phi_0$  have to be so that the intensity on the optical axis  $x = 0$  vanishes?

$$(a) U_{FR} = -i \frac{2\lambda}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B}x^2\right) U_f\left(\frac{kx}{2z_B}\right)$$

$$\begin{aligned} U_f\left(\frac{kx}{2z_B}\right) &= \frac{1}{2\pi} \int_{-a}^a u_0(x') \exp\left(-i\frac{kx}{2z_B}x'\right) dx' = \frac{A}{2\pi} \int_{-a}^a \exp[i\Phi_0(\frac{x'}{a} + 1)] \exp\left(-i\frac{kx}{2z_B}x'\right) dx' \\ &= \frac{A}{2\pi} e^{i\Phi_0} \int_{-a}^a \exp[i(\frac{\Phi_0}{a} - \frac{kx}{2z_B})x'] dx' = \frac{A}{2\pi} e^{i\Phi_0} \frac{\exp[i(\frac{\Phi_0}{a} - \frac{kx}{2z_B})x'] \Big|_{-a}^a}{i(\frac{\Phi_0}{a} - \frac{kx}{2z_B})} \\ &= \frac{A}{2\pi} e^{i\Phi_0} \frac{\exp[i(\frac{\Phi_0}{a} - \frac{kx}{2z_B})a] - \exp[-i(\frac{\Phi_0}{a} - \frac{kx}{2z_B})a]}{i(\frac{\Phi_0}{a} - \frac{kx}{2z_B})} = \frac{A}{\pi} e^{i\Phi_0} \frac{\sin[(\frac{\Phi_0}{a} - \frac{kx}{2z_B})a]}{(\frac{\Phi_0}{a} - \frac{kx}{2z_B})} \\ &= \frac{Au}{\pi} e^{i\Phi_0} \sin\left[\left(\frac{\Phi_0}{a} - \frac{kx}{2z_B}\right)a\right] \end{aligned}$$

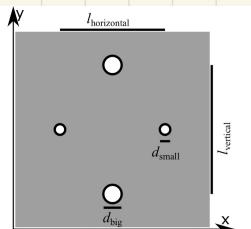
$$\Rightarrow U_{FR} = -i \frac{2\Phi_0 a}{\lambda z_B} e^{i\Phi_0} e^{ikz_B} \exp\left(\frac{ik}{2z_B}x^2\right) \sin\left[\left(\frac{\Phi_0}{a} - \frac{kx}{2z_B}\right)a\right]$$

$$I = |U_{FR}|^2 = \frac{4A^2 a^2}{\lambda^2 z_B^2} \sin^2\left[\left(\frac{\Phi_0}{a} - \frac{k}{2z_B}x\right)a\right]$$

$$\text{if } x \neq 0 \Rightarrow I = \frac{4A^2 a^2}{\lambda^2 z_B^2} \sin^2(\Phi_0) = \frac{4A^2 a^2}{\lambda^2 z_B^2} \cdot \frac{\sin \Phi_0}{\Phi_0} \quad \Phi_0 = n\pi \quad n = \pm 1, \pm 2, \dots \Rightarrow I = 0$$

In Figure 1, a transmission mask is shown with the dimensions, indicated by  $d_{big}$  and  $d_{small}$  for the diameters of the holes and  $l_{vertical}$  and  $l_{horizontal}$  for the distances between the holes. They obey the following relation:  $l_{vertical} > l_{horizontal} \gg d_{big} > d_{small}$ .

- b) Which of these dimensions defines the applicability of the Fresnel approximation and the Fraunhofer approximation, respectively? Explain the condition(s) for both approximations.



$$N_F = \frac{a^2}{\lambda z_B} \quad z_0 = \frac{\pi a^2}{\lambda}$$

$0.1 < N_F < 10 \quad z_0 < z_B < 3z_0 \rightarrow \text{Fresnel diffraction}$

$N_F < 0.1 \quad z_0 > 3z_0 \rightarrow \text{Fraunhofer diffraction}$

$l_{horizontal}$  and  $l_{vertical}$  defines the applicability of Fresnel Approximation.

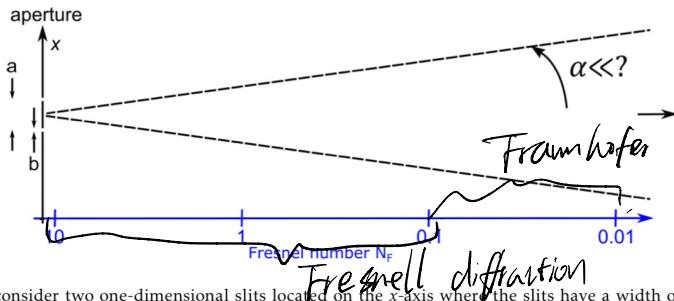
When  $z_B \ll l_{horizontal}$  and  $z_B \ll l_{vertical}$ , Fresnel Approximation can apply

$d_{big}$  and  $d_{small}$  defines the applicability of Fraunhofer Approximation. When  $d_{big} \ll z_B$

and  $d \ll z_B$ . Fraunhofer Approximation can apply.

And if  $\alpha \ll \text{horizontal}$  and  $z_B \ll \text{vertical}$  are also met, it's Paraxial Fraunhofer diffraction. If they are not met, it's non-paraxial Fraunhofer diffraction.

- a) Write down the conditions where 1) the Fresnel approximation, 2) the paraxial Fraunhofer approximation, and 3) the non-paraxial Fraunhofer approximation are valid. Mark the regions where the different diffraction approximations are valid in the following figure.



- b) We consider two one-dimensional slits located on the x-axis where the slits have a width of  $b$  and are separated by a distance of  $a$ . Calculate the resulting far-field intensity when the field directly after the aperture is

$$u(x, z=0) = \begin{cases} 1 & \text{for } |x| < b/2 \\ 1 & \text{for } a - b/2 < x < a + b/2 \\ 0 & \text{otherwise} \end{cases}$$

Hint: You may leave out the prefactors.

- c) When  $a \gg b$ , which of these dimensions defines the applicability of the Fresnel approximation and the Fraunhofer approximation, respectively? For both approximations, state the condition(s) in relation with the wavelength  $\lambda$  and the observation distance  $z_B$ .

$$(b) U_{FR}(x, z) = -\frac{i\lambda}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B} x^2\right) U_0\left(\frac{kx}{z_B}\right)$$

$$\begin{aligned} U_0\left(\frac{kx}{z_B}\right) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} U_0(x, z=0) e^{-i\frac{kx}{z_B} x'} dx' = \frac{1}{2\pi} \int_{-\frac{b}{2}}^{\frac{b}{2}} \exp(-i\frac{kx}{z_B} x') dx' + \frac{1}{2\pi} \int_{a-\frac{b}{2}}^{a+\frac{b}{2}} \exp(-i\frac{kx}{z_B} x') dx' \\ &= -\frac{z_B}{i2\pi kx} \left[ \exp\left(-i\frac{kx}{z_B} \frac{b}{2}\right) - \exp\left(i\frac{kx}{z_B} \frac{b}{2}\right) \right] - \frac{z_B}{i2\pi kx} \left\{ \exp\left[-i\frac{kx}{z_B}(a+\frac{b}{2})\right] - \exp\left[i\frac{kx}{z_B}(a-\frac{b}{2})\right] \right\} \\ &= \frac{z_B}{\pi kx} \sin\left(\frac{kb}{2z_B} x\right) + \frac{z_B}{i2\pi kx} \exp\left(-i\frac{ka}{z_B}\right) \left[ \exp\left(\frac{ikb}{2z_B} x\right) - \exp\left(-\frac{ikb}{2z_B} x\right) \right] \\ &= \frac{z_B}{\pi kx} \sin\left(\frac{kb}{2z_B} x\right) \left[ 1 + \exp\left(-i\frac{ka}{z_B} x\right) \right] = \frac{b}{2\pi} \operatorname{sinc}\left(\frac{kb}{2z_B} x\right) \left[ 1 + \exp\left(-i\frac{ka}{z_B} x\right) \right] \end{aligned}$$

$$U_{FR} = -\frac{ib}{\lambda z_B} e^{ikz_B} \exp\left(\frac{ik}{2z_B} x^2\right) \operatorname{sinc}\left(\frac{kb}{2z_B} x\right) \left[ 1 + \exp\left(-i\frac{ka}{z_B} x\right) \right]$$

$$\begin{aligned} I = |U_{FR}|^2 &= \left(\frac{b}{\lambda z_B}\right)^2 \operatorname{sinc}^2\left(\frac{kb}{2z_B} x\right) \left[ 1 + \exp\left(-i\frac{ka}{z_B} x\right) \right] \left[ 1 + \exp\left(i\frac{ka}{z_B} x\right) \right] \\ &= \left(\frac{b}{\lambda z_B}\right)^2 \operatorname{sinc}^2\left(\frac{kb}{2z_B} x\right) \left[ 2 + \exp\left(-i\frac{ka}{z_B} x\right) + \exp\left(i\frac{ka}{z_B} x\right) \right] = 2 \left(\frac{b}{\lambda z_B}\right)^2 \operatorname{sinc}^2\left(\frac{kb}{2z_B} x\right) \left[ 1 + \cos\left(\frac{ka}{z_B} x\right) \right] \\ &= \left(\frac{2b}{\lambda z_B}\right)^2 \operatorname{sinc}^2\left(\frac{kb}{2z_B} x\right) \cos^2\left(\frac{ka}{2z_B} x\right) \end{aligned}$$

(c) a defines the applicability of Fresnel Approximation.

$$a \ll z_B \quad \text{and} \quad 0.1 < N_F < 10 \quad N_F = \frac{b^2}{\lambda z_B} \quad \frac{b^2}{\lambda z_B} < 10 \quad z_B > \frac{b^2}{10\lambda} \Rightarrow \frac{b}{10\lambda} < z_B < \frac{10b^2}{\lambda}$$

b defines the applicability of Fraunhofer Approximation

$$N_F < 0.1 \Rightarrow \frac{b^2}{\lambda z_B} < 0.1 \Rightarrow z_B > \frac{10b^2}{\lambda}$$