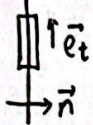


Task 1

a)

Solution:

medium 1 medium 2



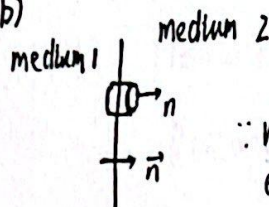
∴ We do not consider surface current and surface charges

$$\therefore \int \nabla \times \vec{H} \cdot d\vec{S} = \oint \vec{H} \cdot d\vec{l} = \hat{n} \times (\vec{H}_2 - \vec{H}_1) = 0$$

Conclusion is right, but almost no derivation, so I have to write some point.

-0.5

b)



∴ We do not consider surface current and surface charges

$$\therefore \int \text{div } \vec{D} \, dv = \oint \vec{D} \cdot d\vec{S} = \hat{n} \cdot (\vec{D}_2 - \vec{D}_1) = 0$$

E and B ? -2

Task 2

a) (3)

The definition of potential $\varphi(\vec{r}, t)$ and $\vec{A}(\vec{r}, t)$:

$$\vec{H}(\vec{r}, t) = \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t)$$

$$\vec{E}(\vec{r}, t) = -\nabla \varphi(\vec{r}, t) - \frac{\partial \vec{A}(\vec{r}, t)}{\partial t}$$

Maxwell's equations:

$$① \nabla \cdot \vec{D}(\vec{r}, t) = \rho(\vec{r}, t)$$

$$② \nabla \cdot \vec{B}(\vec{r}, t) = 0$$

$$③ \nabla \times \vec{E}(\vec{r}, t) = -\frac{\partial \vec{B}(\vec{r}, t)}{\partial t}$$

$$④ \nabla \times \vec{H}(\vec{r}, t) = \vec{j}(\vec{r}, t) + \frac{\partial \vec{D}(\vec{r}, t)}{\partial t}$$

put $\vec{H}(\vec{r}, t) = \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t)$ into equation ④

$$\nabla \cdot (\nabla \times \vec{A}(\vec{r}, t)) = 0$$

put $\vec{E}(\vec{r}, t) = -\nabla \varphi(\vec{r}, t) - \frac{\partial \vec{A}(\vec{r}, t)}{\partial t}$ into equation ③

$$\nabla \times (-\nabla \varphi - \frac{\partial \vec{A}(\vec{r}, t)}{\partial t}) = -\frac{\partial \vec{B}(\vec{r}, t)}{\partial t}$$

$$\therefore \vec{B}(\vec{r}, t) = \nabla \times \vec{A}(\vec{r}, t)$$

$$\therefore \nabla \times (-\nabla \varphi) = 0$$

Wave equations:

put $\vec{H}(\vec{r}, t) = \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t)$ into equation ④

$$\nabla \times \nabla \times \vec{A}(\vec{r}, t) = \mu_0 \vec{j}(\vec{r}, t) + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

put $\vec{E}(\vec{r}, t) = -\nabla \varphi(\vec{r}, t) - \frac{\partial \vec{A}(\vec{r}, t)}{\partial t}$ into this equation

$$\nabla (\nabla \cdot \vec{A}(\vec{r}, t)) - \nabla^2 \vec{A}(\vec{r}, t) = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} (-\nabla \varphi - \frac{\partial \vec{A}}{\partial t})$$

$$\nabla (\nabla \cdot \vec{A}(\vec{r}, t)) - \nabla^2 \vec{A}(\vec{r}, t) = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} (-\nabla \varphi - \frac{\partial \vec{A}}{\partial t})$$

$$\text{when } \nabla \cdot \vec{A}(\vec{r}, t) = -\frac{1}{c^2} \frac{\partial \varphi(\vec{r}, t)}{\partial t}$$

The wave equation is

$$\nabla^2 \vec{A}(\vec{r}, t) - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}(\vec{r}, t)}{\partial t^2} = -\mu_0 \vec{j}(\vec{r}, t)$$

Put $\vec{E}(\vec{r}, t) = -\nabla\varphi(\vec{r}, t) - \frac{\partial\vec{A}(\vec{r}, t)}{\partial t}$ into equation ①

$$\nabla \cdot (-\nabla\varphi(\vec{r}, t) - \frac{\partial\vec{A}}{\partial t}) = \frac{\rho(\vec{r}, t)}{\epsilon_0} \quad \checkmark$$

$$\text{When } \nabla \cdot \vec{A}(\vec{r}, t) = -\mu_0\epsilon_0 \frac{\partial\varphi(\vec{r}, t)}{\partial t}$$

The wave equation is

$$\nabla^2\varphi(\vec{r}, t) - \mu_0\epsilon_0 \frac{\partial^2\varphi(\vec{r}, t)}{\partial t^2} = -\frac{\rho(\vec{r}, t)}{\epsilon_0}$$

b)

Solution:

$$\vec{E}(\vec{r}, t) = -\nabla\varphi(\vec{r}, t) - \partial_t\vec{A}(\vec{r}, t)$$

$$= -\nabla\varphi - \partial_t\lambda(\vec{r}, t) - \partial_t\vec{A}(\vec{r}, t) + \partial_t\nabla\lambda(\vec{r}, t)$$

$$= -\nabla\varphi - \partial_t\vec{A}(\vec{r}, t) \quad \checkmark$$

$$\vec{H}(\vec{r}, t) = \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t)$$

$$= \frac{1}{\mu_0} \nabla \times (\vec{A}(\vec{r}, t) - \nabla\lambda(\vec{r}, t))$$

$$= \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t) - \frac{1}{\mu_0} \nabla \times \nabla\lambda(\vec{r}, t)$$

$$= \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t) - \boxed{0} \quad (\because \lambda(\vec{r}, t) \text{ is an arbitrary scalar, } \therefore \nabla \times \nabla\lambda(\vec{r}, t) = 0)$$

$$= \frac{1}{\mu_0} \nabla \times \vec{A}(\vec{r}, t) \quad \checkmark$$

$\therefore \vec{E}(\vec{r}, t)$ and $\vec{H}(\vec{r}, t)$ are invariant under the transformation

\therefore Maxwell's equations are invariant under the transformation too.

c)

d) Solution:

The wave equation for $\varphi(\vec{r}, t)$:

$$\nabla^2\varphi(\vec{r}, t) - \mu_0\epsilon_0 \frac{\partial^2\varphi(\vec{r}, t)}{\partial t^2} = -\frac{\rho(\vec{r}, t)}{\epsilon_0}$$

$$\nabla^2\varphi(\vec{r}, t) - \frac{1}{c^2} \frac{\partial^2\varphi(\vec{r}, t)}{\partial t^2} = -\frac{\rho(\vec{r}, t)}{\epsilon_0} \quad \checkmark$$

The wave equation for $\vec{A}(\vec{r}, t)$

$$\nabla^2\vec{A}(\vec{r}, t) - \mu_0\epsilon_0 \frac{\partial^2\vec{A}(\vec{r}, t)}{\partial t^2} = -\mu_0\vec{j}(\vec{r}, t)$$

$$\nabla^2\vec{A}(\vec{r}, t) - \frac{1}{c^2} \frac{\partial^2\vec{A}(\vec{r}, t)}{\partial t^2} = -\mu_0\vec{j}(\vec{r}, t) \quad \checkmark$$

e)

Solution:

$$\nabla \cdot \vec{A}(\vec{r}, t) = \nabla \cdot \vec{A}(\vec{r}, t) - \nabla \cdot \nabla\lambda(\vec{r}, t)$$

$$= \frac{1}{c^2} \partial_t\varphi(\vec{r}, t) = -\frac{1}{c^2} \frac{\partial\varphi(\vec{r}, t)}{\partial t} - \frac{1}{c^2} \frac{\partial^2\lambda(\vec{r}, t)}{\partial t^2}$$

\therefore the condition to preserve the Lorenz gauge is

$$\nabla \cdot \nabla\lambda(\vec{r}, t) = \frac{\partial^2\lambda(\vec{r}, t)}{\partial t^2} \quad \checkmark$$

Task 3:

1+2+2+0

a) solution:

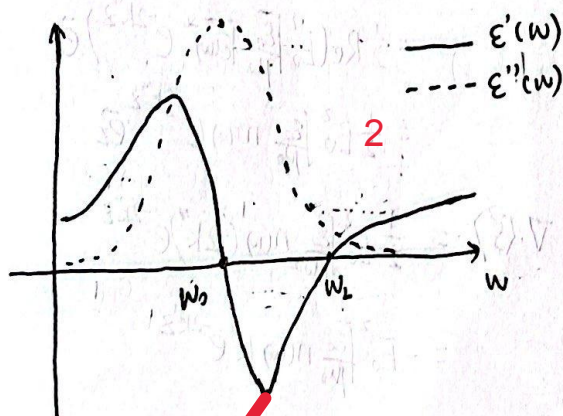
$$\begin{aligned}\chi(\omega) &= \int_{-\infty}^{\infty} R(t) e^{i\omega t} dt \\ &= \int_0^{\infty} \frac{f}{\Omega} e^{-\delta t} \frac{e^{i\Omega t} - e^{i\omega t}}{2i} e^{i\omega t} dt \\ &= \frac{f}{2i\Omega} \left(\int_0^{\infty} e^{[-\delta + i(\Omega + \omega)]t} dt - \int_0^{\infty} e^{[-\delta + i(\Omega - \omega)]t} dt \right) \\ &= \frac{f}{2i\Omega} \left(-\frac{1}{-\delta + i(\Omega + \omega)} + \frac{1}{-\delta + i(\Omega - \omega)} \right) \\ &= \frac{f}{2i\Omega} \frac{-[-\delta + i(\Omega - \omega)] + [-\delta + i(\Omega + \omega)]}{\delta^2 - 2i\delta\omega + \Omega^2 - \omega^2} \\ &= \frac{f}{2i\Omega} \frac{2i\delta}{\delta^2 - 2i\delta\omega + \Omega^2 - \omega^2} \\ &= \frac{f}{\delta^2 - 2i\delta\omega + \omega_0^2 - \delta^2 - \omega^2} \\ &= \frac{f}{\omega_0^2 - \omega^2 - 2i\delta\omega}\end{aligned}$$

b) solution: $\epsilon(\omega) = \epsilon'(\omega) + i\epsilon''(\omega)$

$$\begin{aligned}\epsilon(\omega) &= 1 + \chi(\omega) = 1 + \frac{f}{\omega_0^2 - \omega^2 - 2i\delta\omega} \\ &= 1 + \frac{(\omega_0^2 - \omega^2)f + 2i\delta f\omega}{(\omega_0^2 - \omega^2)^2 + 4\delta^2\omega^2}\end{aligned}$$

$$\epsilon'(\omega) = 1 + \frac{(\omega_0^2 - \omega^2)f}{(\omega_0^2 - \omega^2)^2 + 4\delta^2\omega^2}$$

$$\epsilon''(\omega) = \frac{2\delta f\omega}{(\omega_0^2 - \omega^2)^2 + 4\delta^2\omega^2}$$



c) solution:

$$\bar{p}(\vec{r}, \omega) = \epsilon_0 \chi(\omega) \bar{E}(\vec{r}, \omega)$$

$$\bar{p}(\vec{r}, \omega) = \epsilon_0 \chi(\omega) \bar{E}(\vec{r}) \bar{F}\{\cos(\omega_0 t + \omega t)\}$$

$$\begin{aligned}\bar{F}\{\cos(\omega_0 t + \omega t)\} &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \cos(\omega_0 t + \omega t) e^{i\omega t} dt \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{i(\omega_0 + \omega)t} + e^{i(\omega - \omega_0)t}}{2} e^{i\omega t} dt \\ &= \frac{1}{4\pi} \int_{-\infty}^{\infty} e^{i(\omega_0 + \omega + \omega)t} + e^{i(\omega - \omega_0 - \omega)t} dt \\ &= \frac{1}{4\pi} [2\pi\delta(\omega_0 + \omega) + 2\pi\delta(\omega - \omega_0)] \\ &= \frac{1}{2} [\delta(\omega_0 + \omega) + \delta(\omega - \omega_0)]\end{aligned}$$

$$\bar{p}(\vec{r}, \omega) = \frac{\bar{E}(\vec{r})\epsilon_0 f}{2} [\delta(\omega_0 + \omega) + \delta(\omega - \omega_0)] \frac{f}{\omega_0^2 - \omega^2 - 2i\delta\omega}$$

$$\begin{aligned}\bar{p}(\vec{r}, t) &= \int_{-\infty}^{\infty} \bar{p}(\vec{r}, \omega) e^{-i\omega t} d\omega \\ &= \frac{\bar{E}(\vec{r})\epsilon_0 f}{2} \left[\int_{-\infty}^{\infty} \delta(\omega_0 + \omega) \frac{e^{-i\omega t}}{\omega_0^2 - \omega^2 - 2i\delta\omega} d\omega + \int_{-\infty}^{\infty} \delta(\omega - \omega_0) \frac{e^{-i\omega t}}{\omega_0^2 - \omega^2 - 2i\delta\omega} d\omega \right] \\ &= \frac{\bar{E}(\vec{r})\epsilon_0 f}{2} \left(\frac{e^{i\omega_0 t}}{\omega_0^2 - \omega_0^2 - 2i\delta\omega_0} + \frac{e^{-i\omega_0 t}}{\omega_0^2 - \omega_0^2 - 2i\delta\omega_0} \right) \\ &= \frac{\bar{E}(\vec{r})\epsilon_0 f}{2} \left[\frac{(\omega_0^2 - \omega_0^2)(e^{i\omega_0 t} + e^{-i\omega_0 t})}{(\omega_0^2 - \omega_0^2)^2 + 4\delta^2\omega_0^2} - \frac{2i\delta\omega_0(e^{i\omega_0 t} - e^{-i\omega_0 t})}{(\omega_0^2 - \omega_0^2)^2 + 4\delta^2\omega_0^2} \right] \\ &= \bar{E}(\vec{r})\epsilon_0 f \left[\frac{(\omega_0^2 - \omega_0^2) \cos(\omega_0 t)}{(\omega_0^2 - \omega_0^2)^2 + 4\delta^2\omega_0^2} + \frac{4\delta\omega_0 \sin(\omega_0 t)}{(\omega_0^2 - \omega_0^2)^2 + 4\delta^2\omega_0^2} \right]\end{aligned}$$

2

frequency range

frequency range with normal dispersion: $\frac{\partial \epsilon'(\omega)}{\partial \omega} > 0$

frequency range with anomalous dispersion: $\frac{\partial \epsilon'(\omega)}{\partial \omega} < 0$

Strong absorption occur: $\omega = \omega_0$

Task 4

a)

Solution:

9

The wave equation is

$$\text{rot rot } \vec{E}(\vec{r}, \omega) = \frac{\omega^2}{c^2} \epsilon(\omega) \vec{E}(\vec{r}, \omega)$$

The wave is a monochromatic plane wave with the complex representation $\vec{E}(\vec{r}) = \vec{E}_0 \exp[i\vec{k} \cdot \vec{r}]$

$$\text{div } \vec{E}(\vec{r}, \omega) = 0 \quad \vec{E}(\vec{r}, \omega) = \vec{E}(\omega) \exp[i\vec{k} \cdot \vec{r}]$$

$$\Delta \vec{E}(\vec{r}, \omega) + \frac{\omega^2}{c^2} \epsilon(\omega) \vec{E}(\vec{r}, \omega) = 0$$

$$(ik)^2 \vec{E}(\omega) + \frac{\omega^2}{c^2} \epsilon(\omega) \vec{E}(\omega) = 0$$

$$[-k^2 + \frac{\omega^2}{c^2} \epsilon(\omega)] \vec{E}(\omega) = 0$$

$$\vec{E}(\omega) \neq 0 \quad \therefore -k^2 + \frac{\omega^2}{c^2} \epsilon(\omega) = 0 \quad (\text{dispersion relation})$$

$$k(\omega) = \frac{\omega}{c} \sqrt{\epsilon(\omega)} = \frac{\omega}{c} [n(\omega) + iK(\omega)]$$

$$\sqrt{\epsilon(\omega)} = n(\omega) + iK(\omega)$$

$$\epsilon(\omega) = [n(\omega) + iK(\omega)]^2 = n^2(\omega) + 2in(\omega)K(\omega) - K^2(\omega)$$

$$\text{Re } \epsilon(\omega) = n^2(\omega) - K^2(\omega)$$

$$\text{Im } \epsilon(\omega) = 2n(\omega)K(\omega)$$

$$\therefore n(\omega) = \frac{1}{\sqrt{2}} \sqrt{[\text{Re } \epsilon(\omega)]^2 + [\text{Im } \epsilon(\omega)]^2} + \text{Re } \epsilon(\omega)$$

b)

Solution:

$$\vec{E}(\vec{r}) = \vec{E}_0 \exp[i\vec{k} \cdot \vec{r}]$$

The plane wave has a linear polarization along the y-direction and \vec{k} -vector is pointing in the z-direction

$$\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i(kz - \omega t)} \vec{e}_y \quad \vec{E}(\vec{r}) = \vec{E}_0 e^{ikz} \vec{e}_y$$

$$\vec{H}(\vec{r}, \omega) = \frac{1}{\omega \mu_0} [\vec{k} \times \vec{E}(\omega)] = \frac{1}{\omega \mu_0} \vec{k} \times \vec{E}_0 e^{ikz} = \frac{1}{\omega \mu_0} \text{rot } \vec{E}(\vec{r}, \omega)$$

$$\vec{E}(\vec{r}, \omega) = \mathcal{F}\{\vec{E}(\vec{r}, t)\} = \vec{E}_0 e^{ikz} \delta(\omega' - \omega) \vec{e}_y$$

$$\text{rot } \vec{E}(\vec{r}, \omega) = \nabla \times \vec{E}(\vec{r}, \omega)$$

$$\vec{e}_z \times \vec{e}_y = -\vec{e}_x$$

$$= i\vec{k} \times \vec{E}_0 e^{ikz} \delta(\omega' - \omega) \vec{e}_y$$

$$= i \frac{\omega}{c} \sqrt{\epsilon(\omega)} \vec{E}_0 e^{ikz} \delta(\omega' - \omega) \vec{e}_x$$

$$\therefore \vec{H}(\vec{r}, \omega) = -\frac{i}{\omega \mu_0} \left[i \frac{\omega}{c} \sqrt{\epsilon(\omega)} \vec{E}_0 e^{ikz} \delta(\omega' - \omega) \right] \vec{e}_x$$

$$= \sqrt{\frac{\epsilon_0 \epsilon(\omega)}{\mu_0}} \vec{E}_0 e^{ikz} \delta(\omega' - \omega) \vec{e}_x$$

$$- \vec{H}(\vec{r}, t) = \int_{-\infty}^{\infty} \vec{e}_x \sqrt{\frac{\epsilon_0 \epsilon(\omega)}{\mu_0}} \vec{E}_0 e^{ikz} \delta(\omega' - \omega) e^{-i\omega t} d\omega$$

$$= \sqrt{\frac{\epsilon_0 \epsilon(\omega)}{\mu_0}} \vec{E}_0 e^{ikz} e^{-i\omega t} \vec{e}_x$$

$$- \vec{H}(\vec{r}) = \sqrt{\frac{\epsilon_0 \epsilon(\omega)}{\mu_0}} \vec{E}_0 e^{ikz} \vec{e}_x$$

$$c) \langle \vec{S}(\vec{r}, t) \rangle = \frac{1}{2} \text{Re} [\vec{E}(\vec{r}) \times \vec{H}(\vec{r})^*]$$

$$\vec{E}(\vec{r}) = \vec{E}_0 e^{ikz} \vec{e}_y$$

$$\vec{H}(\vec{r}) = \sqrt{\frac{\epsilon_0}{\mu_0}} \sqrt{\epsilon(\omega)} \vec{E}_0 e^{ikz} \vec{e}_x$$

$\therefore \vec{k}$ is a complex wavevector

$$\therefore \vec{k} = \vec{k}' + i\vec{k}''$$

$$\vec{H}(\vec{r})^* = \sqrt{\frac{\epsilon_0}{\mu_0}} \sqrt{\epsilon(\omega)^*} \vec{E}_0 e^{-k''z} e^{ik'z} \vec{e}_x$$

$$\therefore \langle \vec{S}(\vec{r}, t) \rangle = \frac{1}{2} \left(\vec{E}_0 e^{ik'z} e^{ik''z} \sqrt{\frac{\epsilon_0}{\mu_0}} \sqrt{\epsilon(\omega)^*} \vec{E}_0 e^{-k''z} e^{ik'z} \right) \vec{e}_z$$

$$\vec{E}_0 \sqrt{\frac{\epsilon_0}{\mu_0}} \sqrt{\epsilon(\omega)}$$

$$\therefore \langle \vec{S}(\vec{r}, t) \rangle = \frac{1}{2} \text{Re} \left(\vec{E}_0 e^{ik'z} e^{ik''z} \sqrt{\frac{\epsilon_0}{\mu_0}} \sqrt{\epsilon(\omega)^*} \vec{E}_0 e^{-k''z} e^{ik'z} \right) \vec{e}_z$$

$$= \frac{1}{2} \text{Re} \left(\vec{E}_0^2 \sqrt{\frac{\epsilon_0}{\mu_0}} \sqrt{\epsilon(\omega)^*} e^{-2k''z} \right) \vec{e}_z$$

$$= \frac{1}{2} \vec{E}_0^2 \sqrt{\frac{\epsilon_0}{\mu_0}} n(\omega) e^{-2k''z} \vec{e}_z$$

$$\nabla \cdot \langle \vec{S} \rangle = \frac{1}{2} \vec{E}_0^2 \sqrt{\frac{\epsilon_0}{\mu_0}} n(\omega) (-2k'') e^{-2k''z}$$

$$= -\vec{E}_0^2 \sqrt{\frac{\epsilon_0}{\mu_0}} n(\omega) k'' e^{-2k''z}$$

d)

Solution:

i propagating waves ~~loss~~ without loss

$$\text{div} \langle \vec{S} \rangle = 0 \quad \therefore \epsilon'' = 0$$

ii propagating waves with loss

$$dw \langle \vec{S} \rangle > 0 \quad \therefore \epsilon'' > 0$$

iii nonpropagating waves ~~with loss~~ without loss

$$\text{div} \langle \vec{S} \rangle = 0 \quad \therefore \epsilon'' = 0$$

2

missing the ranges for ϵ'