

# Problem 1

a) in time domain

$$\nabla \cdot \mathbf{H} = 0 \quad \nabla \times \mathbf{H} = \mathbf{j} + \frac{\partial \mathbf{D}}{\partial t}$$

$$\nabla \cdot \mathbf{D} = 0 \quad \nabla \times \mathbf{\Sigma} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t}$$

in freq domain

$$\nabla \cdot \bar{\mathbf{H}} = 0 \quad \nabla \times \bar{\mathbf{H}} = \bar{\mathbf{j}} - i\omega \bar{\mathbf{D}}$$

$$\nabla \cdot \bar{\mathbf{D}} = 0 \quad \nabla \times \bar{\mathbf{\Sigma}} = \mu_0 i\omega \bar{\mathbf{H}}$$

b)

in time domain

$$\mathbf{D} = \epsilon_0 \mathbf{\Sigma} + \epsilon_0 \int_{-\infty}^t \mathbf{R}(\mathbf{r}, t-t') \mathbf{\Sigma}(\mathbf{r}, t') dt'$$

in Freq domain

$$\bar{\mathbf{D}} = \epsilon_0 \bar{\mathbf{\Sigma}} + \epsilon_0 \chi(\omega) \bar{\mathbf{\Sigma}} = \epsilon_0 \epsilon(\omega) \bar{\mathbf{\Sigma}}$$

c)

$$\nabla \times \nabla \times \bar{\mathbf{H}} = \nabla \times \bar{\mathbf{j}} - i\omega \nabla \times \bar{\mathbf{D}}$$

$$\nabla(\nabla \cdot \bar{\mathbf{H}}) - \Delta \bar{\mathbf{H}} = \sigma(\omega) \nabla \times \bar{\mathbf{\Sigma}} - i\omega \epsilon_0 \epsilon(\omega) \nabla \times \bar{\mathbf{\Sigma}}$$

$$-\Delta \bar{\mathbf{H}} = \sigma(\omega) i\omega \mu_0 \bar{\mathbf{H}} + \frac{\omega^2}{c^2} \epsilon_0 \epsilon(\omega) \bar{\mathbf{H}}$$

$$\Delta \bar{\mathbf{H}} + \frac{\omega^2}{c^2} \epsilon(\omega) \bar{\mathbf{H}} = -\sigma i\omega \mu_0 \bar{\mathbf{H}}$$

Problem 2

(a) since  $k^2 = \frac{\omega}{c} \epsilon(\omega) \Rightarrow k = \frac{\omega}{c} \sqrt{\epsilon(\omega)}$

$k = \frac{\omega}{c} \sqrt{\epsilon' + i\epsilon''}$  since  $\epsilon' \gg \epsilon''$

$k = \frac{\omega}{c} \sqrt{\epsilon'} + \frac{i\epsilon''}{2\sqrt{\epsilon'}} \Rightarrow k' = \frac{\omega}{c} \sqrt{\epsilon'}, k'' = \frac{\epsilon''}{2\sqrt{\epsilon'}}$

(b)  $\Sigma_r(r,t) = \Sigma_0 e^{-k''z} \cos(\underbrace{k'z - \omega t + \phi}_A) \hat{x}$

$\nabla \times \Sigma_r = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \Sigma_x & 0 & 0 \end{vmatrix} = (\Sigma_0 k'' e^{-k''z} \cos A + \Sigma_0 e^{-k''z} k' \sin A) \hat{y}$

$H_y = \frac{1}{\mu_0 \omega} [\Sigma_0 k'' e^{-k''z} \sin A - \Sigma_0 e^{-k''z} k' \cos A] \hat{y}$

(c)  $S_r(r,t) = \Sigma_r \times H_r$

(d)  $\langle S_r(r,t) \rangle \Rightarrow \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \Sigma_x & 0 & 0 \\ 0 & H_y & 0 \end{vmatrix} = \Sigma_x H_y \hat{z}$

$\Rightarrow \frac{\Sigma_0^2 k''^2}{\mu_0 \omega} e^{-2k''z} \cos A \sin A - \frac{\Sigma_0^2 k' k'}{\mu_0 \omega} e^{-2k''z} (\cos^2 A) \hat{z}$

by taking the average of both

$\langle S_r(r,t) \rangle = \frac{-\Sigma_0^2 k' k'}{2\mu_0 \omega} e^{-2k''z}$

Problem 3

(a)  $U_0(x) = A \left[ 1 + \cos \frac{2\pi x}{G} \right]$  call  $\frac{2\pi}{G} = k$

$$U_0(\alpha) = \frac{A}{2\pi} \int_{-\infty}^{\infty} (1 + \cos kx) e^{i\alpha x} dx$$

$$= A \left[ \delta(\alpha) + \frac{1}{2} \delta(\alpha - k) + \delta(\alpha + k) \right]$$

(b)

$$U(\alpha; z) = A \left[ \delta(\alpha) + \frac{1}{2} [\delta(\alpha - k) + \delta(\alpha + k)] \right] e^{i\sqrt{k_0^2 - \alpha^2} z}$$

$$U(x; z) = A \int_{-\infty}^{\infty} U(\alpha; z) e^{-i\alpha x} d\alpha$$

$$= A \left[ e^{ik_0 z} + \frac{1}{2} \left[ e^{i\sqrt{k_0^2 - k^2} z + ikx} + e^{i\sqrt{k_0^2 - k^2} z - ikx} \right] \right]$$

$$\Rightarrow A \left[ e^{ik_0 z} + e^{i\sqrt{k_0^2 - k^2} z} \cos kx \right]$$

(c) since in Talbot effect  $U(x, z=L_T) = U(x, z=0) e^{ik_0 z + 2\pi i n}$

$\Rightarrow$  if we ignored the phase term that destroy the periodicity

$$e^{i\sqrt{k_0^2 - k^2} L_T - ik_0 L_T - 2\pi n i} = 1$$

$$\Rightarrow \sqrt{k_0^2 - k^2} L_T - k_0 L_T - 2\pi n = 0$$

$$L_T = \frac{2\pi n}{\sqrt{k_0^2 - k^2} - k_0}, \text{ The shortest repetition } n=1$$

$$L_T = \frac{2\pi}{\sqrt{\frac{(2\pi)^2}{\lambda^2} - \frac{(2\pi)^2}{G^2}} - \frac{2\pi}{\lambda}} = \frac{1}{\sqrt{\frac{1}{\lambda^2} - \frac{1}{G^2}} - \frac{1}{\lambda}}$$

① under farfield condition that  $k_0^2 \gg \alpha^2$

which cause  $H_f = e^{i k_0 z} \frac{e^{-i \alpha^2 z}}{e^{2 k_0 z}}$  where  $\alpha^2 = \left(\frac{2\pi}{G}\right)^2$

Then

$$L_T = \frac{2\pi n}{k_0 - \frac{k^2}{2k_0} - k_0} = \frac{2\pi n (2k_0)}{k^2}$$

$$= \frac{4\pi n 2\pi}{\lambda \frac{4\pi^2}{G^2}} = \frac{2n G^2}{\lambda}$$

$$(a) \quad \frac{1}{V_{ph}} = \frac{K_0}{\omega_0} = \frac{n(\omega_0)}{c_0} \Rightarrow V_{ph} = \frac{c_0}{n(\omega_0)}$$

$$V_{ph} = \frac{c_0}{2 + (10^{-32} \times 4 \times 10^{30})} = \frac{c_0}{2 + (0.04)} = \frac{c_0}{2.04}$$

$$(*) \quad \frac{1}{V_g} = \frac{\partial K}{\partial \omega} = \frac{1}{c_0} \left[ n(\omega_0) + \omega \frac{\partial n}{\partial \omega} \Big|_{\omega_0} \right]$$

$$\Rightarrow V_g = \frac{c_0}{2 + 3(10^{-32} \times 4 \times 10^{30})} = \frac{c_0}{2.12}$$

where  $\frac{\partial n}{\partial \omega} = 2\omega c$

$$(b) \quad * \quad q(0) = \frac{T_0^2 (C_0 + i)}{2D(1 + C_0^2)} \quad \text{since } C_0 = 0 \text{ (Flat phase)}$$

$$q(0) = \frac{T_0^2 i}{2D}$$

\* Propagate distance  $L$

$$q(L) = \frac{T_0^2 i}{2D} + L = \frac{T_0^2 i + 2LD}{2D}$$

$$\frac{1}{q(L)} = \frac{2D}{T_0^2 i + 2LD} \Rightarrow \frac{-2DT_0^2 i + 4D^2 L}{4L^2 D^2 + T_0^4}$$

$$\Rightarrow \frac{4D^2 L}{4L^2 D^2 + T_0^4} - \frac{i 2DT_0^2}{4L^2 D^2 + T_0^4}$$

$$\frac{1}{q(\lambda)} = \frac{4\lambda^2 D^2 (T_0^2)}{[4\lambda^2 D^2 + T_0^4] (T_0^2)} - \frac{i 2 D T_0^2 (T_0^2)}{[4\lambda^2 D^2 + T_0^4] (T_0^2)}$$

$$= \frac{2 D}{T_0^2} \left[ \underbrace{\frac{2 D \lambda T_0^2}{4\lambda^2 D^2 + T_0^4}}_{C(2)} - \underbrace{\frac{i T_0^4}{4\lambda^2 D^2 + T_0^4}}_{\frac{T_0^2}{T^2(2)}} \right]$$

For the evolved chirp  $C(2) = 0$

$$\Rightarrow \frac{T_0^4}{4\lambda^2 D^2 + T_0^4} = \frac{T_0^2}{T^2(2)} \Rightarrow T^2(\lambda) = \frac{4\lambda^2 D^2 + T_0^4}{T_0^2}$$

$$T(\lambda) = \sqrt{\frac{4\lambda^2 D^2 + T_0^4}{T_0^2}}$$

$$\text{since } D = \frac{\partial^2 \chi^2}{\partial \omega^2} = \frac{1}{-v_g^2} \frac{\partial v_g}{\partial \omega} = \frac{\partial}{\partial \omega} \left( \frac{1}{v_g} \right)$$

$$\text{and } \frac{1}{v_g} = \frac{1}{c_0} [n + 2\omega^2 C] = \frac{1}{c_0} [B + 3\omega^2 C]$$

$$\frac{\partial}{\partial \omega} \left( \frac{1}{v_g} \right) = \frac{1}{c_0} 6\omega C = D$$

$$\text{Then } D = 6 \times 2 \times 10^{15} \times 10^{-32} = \frac{1.2 \times 10^{-16}}{c_0}$$

$$\text{Then } T_1(\lambda) = \sqrt{\frac{4 \times 400 \times \left( \frac{1.2 \times 10^{-16}}{c_0} \right)^2 + (4 \times 10^{-12})^4}{(4 \times 10^{-12})^2}} =$$

©

$$C=0 \Rightarrow D=0$$

$$T_2(\lambda) = \sqrt{\frac{(4 \times 10^{-12})^4}{(4 \times 10^{-12})^2}} = \sqrt{4 \times 10^{-12}} = 2 \times 10^{-6} \text{ s}$$



Singularity of a black hole

$$\text{The difference} = |T_1 - T_2| = ?$$



# blem 5

(7)

$$a) \quad z_0 = \pi w_0^2 / \lambda$$

$$\frac{1}{q(z)} = \frac{1}{R(z)} + \frac{i}{z_0} \quad \text{since we start at the waist } R=\infty$$

$$\frac{1}{q(0)} = +\frac{i}{z_0} \Rightarrow q(0) = -iz_0$$

after distance  $z$

$$q(z) = -iz_0 + z$$

$$\Rightarrow \frac{1}{q(z)} = \frac{1}{z - iz_0} \Rightarrow \frac{z^*}{z^2 + z_0^2} + \frac{-iz_0}{z^2 + z_0^2}$$

$$\Rightarrow \underbrace{\frac{1}{z(1 + \frac{z_0^2}{z^2})}}_{R(z)} + i \underbrace{\frac{1}{z_0(1 + \frac{z^2}{z_0^2})}}_{R(z)}$$

$$\Rightarrow \frac{\lambda}{\pi w^2(z)} = \frac{\lambda}{\pi w_0^2(1 + \frac{z^2}{z_0^2})} \Rightarrow w^2(z) = w_0^2(1 + \frac{z^2}{z_0^2})$$

$$b) \quad M = \begin{pmatrix} 1 & z' \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -1/f & 1 \end{pmatrix} = \begin{pmatrix} 1 - z'/f & z' \\ 1/f & 1 \end{pmatrix}$$

$$q = \frac{-iz_0(1 - z'/f) + z'}{-iz_0/f + 1} \Rightarrow \frac{1}{q} = \frac{1 - iz_0/f}{z' - iz_0(1 - z'/f)}$$

$$\frac{1}{q} = \frac{z' - iz_0(1 - z'/f) - \frac{iz_0 z'}{f} + \frac{z_0^2}{f}(1 - z'/f)}{z'^2 + z_0^2(1 - z'/f)^2}$$

$$\frac{1}{q} = \frac{Z' - iZ_0 + \frac{i2Z_0 Z'}{f} - \frac{i2Z_0 Z'}{f} + \frac{Z_0^2}{f} + \frac{Z_0^2 Z'}{f}}{Z'^2 + Z_0^2 - \frac{2Z_0^2 Z'}{f} + \frac{Z_0^2 Z'^2}{f^2}}$$

$$= \frac{Z' + \frac{Z_0^2}{f} + \frac{Z_0^2 Z'}{f}}{Z'^2 + Z_0^2 + \frac{Z_0^2 Z'^2}{f^2} - \frac{2Z_0^2 Z'}{f}} - \frac{iZ_0}{Z'^2 + Z_0^2 + \frac{Z_0^2 Z'^2}{f^2} - \frac{2Z_0^2 Z'}{f}}$$

at  $\omega_0'$

The Real Part  $\Rightarrow Z' + \frac{Z_0^2}{f} + \frac{Z_0^2 Z'}{f} = 0$

$$Z' \left( 1 + \frac{Z_0^2}{f} \right) = - \frac{Z_0^2}{f}$$

$$Z' = \frac{\frac{Z_0^2}{f}}{1 + \frac{Z_0^2}{f}}, \text{ ~~not~~}$$

Then the Imaginary part

$$\frac{Z_0}{\left( \frac{Z_0^2/f}{1 + Z_0^2/f} \right)^2 + Z_0^2 + \frac{Z_0^2 \left( \frac{Z_0^2/f}{1 + Z_0^2/f} \right)^2 - 2Z_0^2 \left( \frac{Z_0^2/f}{1 + Z_0^2/f} \right)}{f^2} = \frac{\lambda}{\pi \omega_0'^2}$$

when  $Z_0 \gg f$

$$\Rightarrow Z' = \frac{1}{\frac{f}{Z_0^2} \left( 1 + \frac{Z_0^2}{f} \right)} = \frac{1}{\frac{f}{Z_0^2} + \frac{1}{f}} = f$$

Then  $\omega_0'^2 = \frac{\lambda}{\pi} \frac{Z_0 f^2 + Z_0^2 + Z_0^2 - 2Z_0^2}{Z_0} = \frac{\lambda f^2}{\pi Z_0} = \frac{\lambda^2 f^2}{\pi^2 \omega_0'^2}$

$$\omega_0'^2 = \frac{\lambda f}{\pi \omega_0'}$$



