

Task 1 Jinsong Lin

a)

Solution:

From the 2f-setup equation:

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$$U(x, y, z) = -i \frac{(2\pi)^3}{\lambda f} \exp(2ikf) U_0\left(\frac{k}{f}x, \frac{k}{f}y\right)$$

$$\therefore U(x) = -i \frac{2\pi}{\lambda f_1} \exp(2ikf_1) U_0\left(\frac{k}{f_1}x\right)$$

$$U^+(x) = U(x) \cdot H(x) = -i \frac{2\pi}{\lambda f_1} \exp(2ikf_1) U_0\left(\frac{k}{f_1}x\right) \cdot H(x)$$

$$U_1(x) = -i \frac{2\pi}{\lambda f_2} \exp(2ikf_2) U^+\left(\frac{k}{f_2}x\right)$$

$$= -i \frac{2\pi}{\lambda f_2} \exp(2ikf_2) \frac{1}{2\pi} \int_{-\infty}^{\infty} -i \frac{2\pi}{\lambda f_1} \exp(2ikf_1) U_0\left(\frac{k}{f_1}x'\right) H(x') \exp(-i \frac{k}{f_2}xx') dx'$$

$$= -i \frac{1}{\lambda f_2} \exp(2ikf_2) \left[-i \frac{2\pi}{\lambda f_1} \exp(2ikf_1) \right] \int_{-\infty}^{\infty} U_0\left(\frac{k}{f_1}x'\right) H(x') \exp(-i \frac{k}{f_2}xx') dx'$$

$$= -\frac{2\pi}{\lambda^2 f_1 f_2} \exp[2ik(f_1 + f_2)] \int_{-\infty}^{\infty} U_0\left(\frac{k}{f_1}x'\right) H(x') \exp(-i \frac{k}{f_2}xx') dx'$$

b)

Solution:

$$U_0(x) = e^{iA \cos(\omega_0 x)} = 1 + iA \cos(\omega_0 x)$$

$$U_0\left(\frac{k}{f_1}x\right) = \frac{1}{2\pi} \int_{-\infty}^{\infty} [1 + iA \cos(\omega_0 x')] \exp(-i \frac{k}{f_1}xx') dx'$$

$$= \delta(\omega) + i \frac{A}{2} [\delta(\omega + \omega_0) + \delta(\omega - \omega_0)]$$

$$\text{Let } C = -\frac{2\pi}{\lambda^2 f_1 f_2} \exp[2ik(f_1 + f_2)] \quad \varphi(\omega) = \frac{A}{2} [\delta(\omega + \omega_0) + \delta(\omega - \omega_0)]$$

$$U_1(x) = C \int_{-\infty}^{\infty} U_0(\omega) H\left(\frac{f_1}{k}\omega\right) \exp\left[i \frac{k}{f_2}xx'\right] dx'$$

$$= C \int_{-\infty}^{\infty} U_0(\omega) H\left(\frac{f_1}{k}\omega\right) \exp\left(i \frac{f_1}{f_2}x\omega\right) \frac{f_1}{k} d\omega$$

$$= \frac{C f_1}{k} \int_{-\infty}^{\infty} [\delta(\omega) + i\varphi(\omega)] H\left(\frac{f_1}{k}\omega\right) \exp\left(i \frac{f_1}{f_2}x\omega\right) d\omega$$

$$\therefore H(x) = \begin{cases} \exp(i\varphi_0) & -a \leq x \leq a \\ 1 & \text{else} \end{cases}$$

$$\therefore H\left(\frac{f_1}{k}\omega\right) = \begin{cases} \exp(i\varphi_0) & -\frac{k}{f_1}a \leq \omega \leq \frac{k}{f_1}a \\ 1 & \text{else} \end{cases}$$

In order to convert the phase modulation into an amplitude modulation, we need to multiply $\exp(i\varphi_0)$ with $\delta(\omega)$ which is zero-frequency component. $\therefore \frac{k}{f_1}a < \omega_0 = 2\pi/\Lambda$, $a < \frac{2\pi f_1}{k\Lambda}$. And $\varphi_0 = \pm \frac{\pi}{2}$, $\exp(i\varphi_0) = \pm i$, when we calculate the $|U_1(x)|^2$, the phase modulation can convert into an amplitude modulation.

Task 2 Jinsong Liu

a)

Solution:

$$FT[g(x)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{2} [1 + \cos(\frac{2\pi}{d}x')] \exp(-i\frac{\alpha}{2}xx') dx'$$

$$= \frac{1}{2} \delta(\alpha) + \frac{1}{4} [\delta(\alpha + \frac{2\pi}{d}) + \delta(\alpha - \frac{2\pi}{d})]$$

$$FT[n(x', y')] = \frac{1}{(2\pi)^2} \iint_{-\infty}^{\infty} \exp[-\frac{x'^2}{W^2} - \frac{(y' - \frac{R}{2})^2}{l^2}] \exp[-i(\alpha x' + \beta y')] dx' dy'$$

$$= \frac{W \cdot l}{4\pi} \exp[-\frac{W^2 \alpha^2 + l^2 \beta^2}{4}] \cdot \exp(-i\beta \frac{R}{2})$$

$$FT[m(x', y')] = \frac{lW}{4\pi} \exp[-\frac{l^2 \alpha^2 + W^2 \beta^2}{4}] \cdot \exp(i\beta \frac{R}{2})$$

$$FT[e_1(x', y')] = \frac{W^2}{4\pi} \exp[-\frac{W^2 \alpha^2 + W^2 \beta^2}{4}] \cdot \exp(-i\alpha \frac{R}{2}) \cdot \exp(-i\beta \frac{R}{2})$$

$$FT[e_2(x', y')] = \frac{W^2}{4\pi} \exp[-\frac{W^2 \alpha^2 + W^2 \beta^2}{4}] \exp(i\alpha \frac{R}{2}) \exp(-i\beta \frac{R}{2})$$

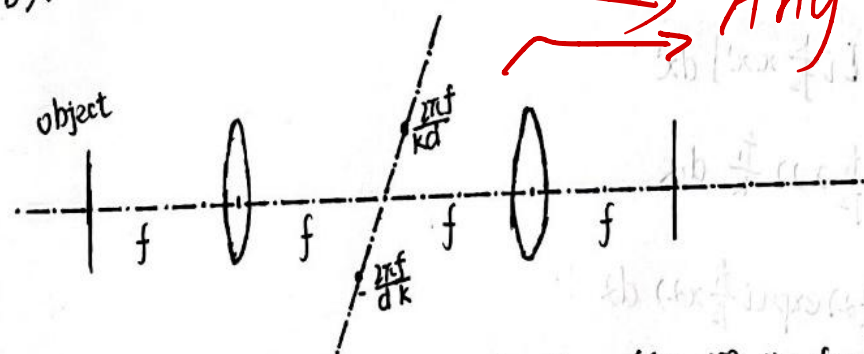
$$\therefore R > l > w \Rightarrow d$$

$$\therefore \frac{2\pi}{d} \gg \frac{1}{W} > \frac{1}{l}$$

So the $\delta(\alpha + \frac{2\pi}{d})$ and $\delta(\alpha - \frac{2\pi}{d})$ will not make too much influence over the Gaussian spectrum.

To remove the prism, we can add filter at $\alpha = 0, \pm \frac{2\pi}{d}$
 $(x = 0, \pm \frac{2\pi f}{kd})$

b):



Any real filter should have a width

The initial optical field is the object, after the focal length of the first lens, it does a Fourier transform and the field in between. The final optical field is inverse Fourier transformation of the field between 2 lenses

of the field between 2 lenses

c)

Solution:

∴ The prism has zero-frequency component which is involved in the Gaussian spectrum.

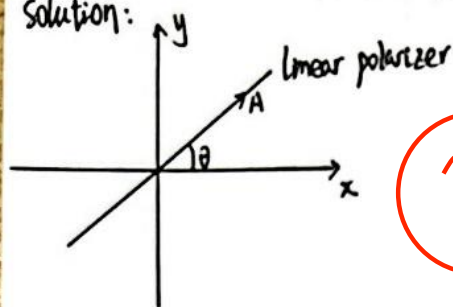
To minimize the effect, we can multiply a very small number with $\delta(x)$.

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Task 3

a)

Solution:



incident: $J_{in} = \begin{bmatrix} A_1 \\ B_1 \end{bmatrix}$

$$A = A_1 \cos \theta + B_1 \sin \theta$$

After the polarizer:

$$J_{out} = \begin{bmatrix} A_2 \\ B_2 \end{bmatrix}$$

$$A_2 = A \cos \theta = A_1 \cos^2 \theta + B_1 \cos \theta \sin \theta$$

$$B_2 = A \sin \theta = A_1 \cos \theta \sin \theta + B_1 \sin^2 \theta$$

$$\therefore T_\theta = \begin{bmatrix} \cos^2 \theta & \cos \theta \sin \theta \\ \cos \theta \sin \theta & \sin^2 \theta \end{bmatrix}$$

b)

Solution:

incident: $J_{in} = \begin{bmatrix} A_1 \\ B_1 \end{bmatrix}$

From the problem, we can get

$$T J_{in} = \lambda J_{in}$$

$$\therefore (T - \lambda) J_{in} = 0$$

$$\det(T - \lambda) = \begin{vmatrix} 1 - 2\lambda & \frac{i}{2} \\ -\frac{i}{2} & 1 - 2\lambda \end{vmatrix} = 0$$

$$\therefore (1 - 2\lambda)^2 - 1 = 0$$

$$\lambda = 0, 1 \quad (\lambda = 0 \text{ is wrong})$$

when $\lambda = 1$

$$\begin{cases} A_1 + i B_1 = 2 A_1 \\ -i A_1 + B_1 = 2 B_1 \end{cases}$$

$$\therefore A_1 = 1 \quad B_1 = -i$$

The normalized Jones vector is $J = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix}$

This state is right circular polarized light.