

FoMO in a nutshell

Important things to memorize for the course »Fundamentals of Modern Optics«*

Thomas Kaiser, Abbe Center of Photonics

WS 16/17

Here, we collect important relations, facts, formulas and other things that one should memorize throughout the course to give you a summary of the most important results. Of course, we cannot list every required formula and we assume that you already have a solid math background. Hence, we cannot guarantee completeness of the list.

FOMO topics

Maxwell's equations

- Macroscopic Maxwell's equations (time domain):

$$\begin{aligned}\nabla \cdot \mathbf{D}(\mathbf{r}, t) &= \rho(\mathbf{r}, t) & \nabla \times \mathbf{E}(\mathbf{r}, t) &= -\frac{\partial \mathbf{B}(\mathbf{r}, t)}{\partial t} \\ \nabla \cdot \mathbf{B}(\mathbf{r}, t) &= 0 & \nabla \times \mathbf{H}(\mathbf{r}, t) &= \mathbf{j}(\mathbf{r}, t) + \frac{\partial \mathbf{D}(\mathbf{r}, t)}{\partial t}\end{aligned}$$

- Macroscopic Maxwell's equations (frequency domain):

$$\begin{aligned}\nabla \cdot \mathbf{D}(\mathbf{r}, \omega) &= \rho(\mathbf{r}, \omega) & \nabla \times \mathbf{E}(\mathbf{r}, \omega) &= i\omega \mathbf{B}(\mathbf{r}, \omega) \\ \nabla \cdot \mathbf{B}(\mathbf{r}, \omega) &= 0 & \nabla \times \mathbf{H}(\mathbf{r}, \omega) &= \mathbf{j}(\mathbf{r}, \omega) - i\omega \mathbf{D}(\mathbf{r}, \omega)\end{aligned}$$

- Constitutive relations (linear material response; in optics usually $\mu(\mathbf{r}, \omega) \equiv 1$ (non-magnetizable):

$$\begin{aligned}\mathbf{D}(\mathbf{r}, \omega) &= \epsilon_0 \mathbf{E}(\mathbf{r}, \omega) + \mathbf{P}(\mathbf{r}, \omega) = \epsilon_0 \epsilon(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega) & \mathbf{B}(\mathbf{r}, \omega) &= \mu_0 \mu(\mathbf{r}, \omega) \mathbf{H}(\mathbf{r}, \omega) \\ \mathbf{P}(\mathbf{r}, \omega) &= \epsilon_0 \chi(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega)\end{aligned}$$

- Time domain material response (response function):

$$\mathbf{P}(\mathbf{r}, t) = \epsilon_0 \int_{-\infty}^t R(\mathbf{r}, t-t') \mathbf{E}(\mathbf{r}, t') dt', \text{ where } R(\mathbf{r}, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \chi(\mathbf{r}, \omega) e^{-i\omega t} d\omega$$

- Complex permittivity:

$$\epsilon(\mathbf{r}, \omega) = 1 + \chi(\mathbf{r}, \omega) + i \frac{\sigma(\mathbf{r}, \omega)}{\epsilon_0 \omega}$$

*Please report any typos or other errors to: thomas.kaiser.1@uni-jena.de

- Continuity equation (conservation of charge):

$$\nabla \cdot \mathbf{j}(\mathbf{r}, t) + \frac{\partial \rho(\mathbf{r}, t)}{\partial t} = 0$$

$$\underbrace{\oint_{\partial V} \mathbf{j} \cdot d\mathbf{S}}_I = - \frac{\partial}{\partial t} \underbrace{\iiint_V \rho \, dV}_Q$$

- Time averaged Poynting vector, loss:

$$\langle \mathbf{S}(\mathbf{r}) \rangle = \frac{1}{2} \operatorname{Re} [\mathbf{E}(\mathbf{r}, \omega) \times \mathbf{H}^*(\mathbf{r}, \omega)] \quad \nabla \cdot \langle \mathbf{S}(\mathbf{r}) \rangle < 0 \Leftrightarrow \text{system is lossy}$$

- Wave equations (vacuum):

$$\nabla \times \nabla \times \mathbf{E}(\mathbf{r}, t) - \frac{1}{c^2} \frac{\partial^2 \mathbf{E}(\mathbf{r}, t)}{\partial t^2} = 0$$

$$\nabla \times \nabla \times \mathbf{H}(\mathbf{r}, t) - \frac{1}{c^2} \frac{\partial^2 \mathbf{H}(\mathbf{r}, t)}{\partial t^2} = 0$$

Normal modes in homogeneous, isotropic matter, non-magnetizable

- Helmholtz equation (wave equation in temporal Fourier domain; homogeneous, isotropic matter):

$$\Delta \mathbf{E}(\mathbf{r}, \omega) + \frac{\omega^2}{c^2} \epsilon(\omega) \mathbf{E}(\mathbf{r}, \omega) = 0$$

- Plane waves are the eigenmodes of free space. They take the form

$$\mathbf{E}(\mathbf{r}, \omega) = \mathbf{E}_0 e^{i\mathbf{k}\mathbf{r}}$$

and their dispersion relation reads as:

$$k_x^2 + k_y^2 + k_z^2 = \frac{\omega^2}{c^2} \epsilon(\omega)$$

- Refractive index:

$$n(\omega) = \frac{k(\omega)}{k_0} = \frac{k(\omega)}{\omega} c$$

- Maxwell relation:

$$n(\omega) = \sqrt{\epsilon(\omega)} \text{ which implies for real and imaginary parts: } (n' + in'')^2 = n'^2 - n''^2 + 2in'n'' = \epsilon' + i\epsilon''$$

- Finding electric or magnetic field from each other in frequency domain in regions without source:

$$\mathbf{H}(\mathbf{r}, \omega) = -\frac{i}{\omega \mu_0} \nabla \times \mathbf{E}(\mathbf{r}, \omega)$$

$$\mathbf{E}(\mathbf{r}, \omega) = \frac{i}{\omega \epsilon_0 \epsilon(\mathbf{r}, \omega)} \nabla \times \mathbf{H}(\mathbf{r}, \omega)$$

- Propagating, lossy and evanescent waves:

$$\langle \mathbf{S} \rangle \neq 0, \nabla \cdot \langle \mathbf{S} \rangle = 0 \Leftrightarrow \text{lossless propagation}$$

$$\langle \mathbf{S} \rangle \neq 0, \nabla \cdot \langle \mathbf{S} \rangle < 0 \Leftrightarrow \text{lossy propagation}$$

$$\langle \mathbf{S} \rangle = 0, \nabla \cdot \langle \mathbf{S} \rangle = 0 \Leftrightarrow \text{evanescence}$$

Beam propagation

- ... see beam propagation scheme handout

Gaussian beams

- Rayleigh length:

$$z_0 = \frac{k}{2} w_0^2$$

- q -parameter definition:

$$\frac{1}{q(z)} = \frac{1}{R(z)} + i \frac{2}{k w^2(z)}$$

- q -parameter at waist position:

$$q_0 = -i z_0$$

- q -parameter after propagation:

$$q(d) = q(0) + d$$

- Condition for finding waist position:

$$\text{Re}[q(z)] = 0$$

- ABCD matrix formalism:

$$\hat{M} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \quad \hat{M}_{\text{tot}} = \hat{M}_n \cdot \dots \cdot \hat{M}_3 \cdot \hat{M}_2 \cdot \hat{M}_1$$

$$q_{n+1} = \frac{A q_n + B}{C q_n + D}$$

Mathematical tools

Miscellaneous math formulas

- Complex exponentials, trigonometric and hyperbolic functions:

$$\begin{aligned}e^{ix} &= \cos x + i \sin x \\ \cos x &= \frac{1}{2} (e^{ix} + e^{-ix}) \\ \sin x &= \frac{1}{2i} (e^{ix} - e^{-ix}) \\ \cosh x &= \frac{1}{2} (e^x + e^{-x}) & \cosh(ix) &= \cos x \\ \sinh x &= \frac{1}{2} (e^x - e^{-x}) & \sinh(ix) &= i \sin x\end{aligned}$$

- Integration:

$$\begin{aligned}\int u(x) \frac{dv(x)}{dx} dx &= u(x)v(x) - \int \frac{du(x)}{dx} v(x) dx \\ \int f(x) dx &= \int f(\xi) \frac{dx}{d\xi} d\xi\end{aligned}$$

- Gaussian functions:

$$\begin{aligned}\int_{-\infty}^{\infty} e^{-x^2} dx &= \sqrt{\pi} \\ \text{FT} \left\{ A \exp \left[-\frac{1}{2} \frac{t^2}{t_0^2} \right] \right\} &= \frac{A t_0}{\sqrt{2\pi}} \exp \left[-\frac{1}{2} \frac{\omega^2}{(1/t_0)^2} \right]\end{aligned}$$

The Fourier transform of a Gaussian function is a Gaussian function.

- Area element in radial coordinates:

$$dA = r dr d\phi$$

Field theory

- Vector identities:

$$\begin{aligned}\nabla \times \nabla \times \mathbf{a} &= \nabla \cdot (\nabla \mathbf{a}) - \Delta \mathbf{a} \\ \Delta \mathbf{a} &= \nabla(\nabla \cdot \mathbf{a}), \text{ but } \Delta \mathbf{a} = (\nabla \cdot \nabla) \mathbf{a} = \nabla^2 \mathbf{a} \neq \nabla(\nabla \cdot \mathbf{a}) \\ \mathbf{a} \times \mathbf{b} \times \mathbf{c} &= \mathbf{b}(\mathbf{a} \cdot \mathbf{c}) - \mathbf{c}(\mathbf{a} \cdot \mathbf{b})\end{aligned}$$

- Integral theorems:

$$\begin{aligned}\iiint_V \nabla \cdot \mathbf{a} dV &= \oint_{\partial V} \mathbf{a} \cdot d\mathbf{S} & (\text{Gauss}) \\ \iint_A \nabla \times \mathbf{a} \cdot d\mathbf{S} &= \oint_{\partial A} \mathbf{a} \cdot d\mathbf{r} & (\text{Stokes})\end{aligned}$$

Fourier transform, δ -function

- In the course, we define the one-dimensional Fourier transform as (these definitions regarding sign and prefactor conventions influence nearly every expression in this document that contains Fourier transforms):

Forward (going to Fourier domain):

$$\tilde{f}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) e^{i\omega t} dt \quad (\text{temporal Fourier domain})$$

$$\tilde{f}(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx \quad (\text{spatial Fourier domain})$$

A Fourier transform gives the strengths of the different plane wave (space) or time harmonic (time) frequency components.

Backward (coming from Fourier domain):

$$f(t) = \int_{-\infty}^{\infty} \tilde{f}(\omega) e^{-i\omega t} d\omega \quad (\text{temporal Fourier domain})$$

$$f(x) = \int_{-\infty}^{\infty} \tilde{f}(k) e^{ikx} dk \quad (\text{spatial Fourier domain})$$

An inverse Fourier transform represents a decomposition of a function into plane wave components (space) or harmonic oscillations (time).

- Fourier shifting theorem:

$$\text{FT} \{f(t - t_0)\} = e^{i\omega t_0} \tilde{f}(\omega)$$

$$\text{FT}^{-1} \{\tilde{f}(\omega - \omega_0)\} = e^{-i\omega_0 t} f(t)$$

$$\text{FT} \{f(x - x_0)\} = e^{-ikx_0} \tilde{f}(k)$$

$$\text{FT}^{-1} \{\tilde{f}(k - k_0)\} = e^{ik_0 x} f(x)$$

The shifting of a function corresponds to a harmonic modulation in Fourier domain.

- Fourier transform of a derivative (I cannot recall an occasion where the backward relations would be needed, so I leave them out):

$$\text{FT} \left\{ \frac{df(t)}{dt} \right\} = -i\omega \tilde{f}(\omega)$$

$$\text{FT} \left\{ \frac{df(x)}{dx} \right\} = ik \tilde{f}(k)$$

A derivative in the real domain corresponds to a simple multiplication in Fourier domain.

- δ -function:

$$\delta(x - x_0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ik(x-x_0)} dk$$

The δ -function is the Fourier transform of a plane wave (space) or harmonic oscillation (time).

$$f(x_0) = \int_{-\infty}^{\infty} f(x) \delta(x - x_0) dx$$

The δ -function »picks out« the value of a function at a particular point. It is just meaningful under an integral.

- Fourier transform of the δ -function:

$$\text{FT}\{\delta(t-t_0)\} = \frac{1}{2\pi} e^{i\omega t_0}$$

$$\text{FT}^{-1}\{\delta(\omega-\omega_0)\} = e^{-i\omega_0 t}$$

$$\text{FT}\{\delta(x-x_0)\} = \frac{1}{2\pi} e^{-ikx_0}$$

$$\text{FT}^{-1}\{\delta(k-k_0)\} = e^{ik_0 x}$$

The Fourier transform of a δ -function is a plane wave (space) of harmonic oscillation (time)

- Convolution:

$$[f \otimes g](t) = \int f(\tau)g(t-\tau) d\tau$$

$$[f \otimes g](x) = \int f(x')g(x-x') dx$$

$$\text{FT}\{[f \otimes g](t)\} = 2\pi \tilde{f}(\omega)\tilde{g}(\omega)$$

$$\text{FT}\{[f \otimes g](x)\} = 2\pi \tilde{f}(k)\tilde{g}(k)$$

Convolution in the real domain corresponds to a simple multiplication in the Fourier domain.