

Final Test: 2021.

Multiple-test:

1. The wavelength 488 nm belongs to Visible.
2. The momentum operator in quantum mechanics has real eigenvalues, is self-adjoint.
3. Solid materials have a bulk modulus different from zero.
4. Metals are characterized by the presence of free electrons, A good electrical conductivity.
5. The effective mass of an electron in a crystal lattice may be negative or positive.
6. The rotational-vibrational absorption spectrum of gas of diatomic molecules may be observed in the infrared spectral region, (show a P-branch).
7. The atomic state 1S_1 is impossible, $(l=0, s=0, j=1)$.
8. All crystals are _____.

True or Wrong.

- a) From $[\hat{Q}, \hat{Q}] = 0$ and $[\hat{Q}, \hat{Q}] = 0$ it follows that $[\hat{Q}, \hat{Q}] = 0$ Wrong.
- b) The operator of discrete translations \hat{T}_n defined by $\hat{T}_n \psi(r) = \psi(r + na)$ (a - lattice vector) with the eigenvalues $\hat{T}_n = e^{i2\pi n}$ is self-adjoint _____.
- c) All self-adjoint operators commute with each other Wrong.
- d) Amorphous solids lack any short-range order in their atomic configuration Wrong.
- e) If a crystal lattice has inversion symmetry, the corresponding solid must be a metal Wrong.

5. Imagine that you have measured the IR absorption spectrum of a gas of $(CF_3)_2CH$ molecules. You observe the fundamental transition wavenumber of the strength stretching vibration of the CH-group as $\nu_p = 2992 \text{ cm}^{-1}$. You also register the transition wavenumber corresponding to the first overtone as $\nu_{2p} = 5882 \text{ cm}^{-1}$. From these data, assuming a Morse potential and neglecting any rotations, estimate the transition wavenumber corresponding to the second overtone ν_{3p} !

$$\begin{aligned} \nu_{nm} &= G(V=n) - G(V=m) = \nu_0(1-x_e) - \nu_0(1-x_e) \\ &= \nu_0(1-x_e)(n-m) - \nu_0x_e[n^2 - m^2] \end{aligned}$$

$$m=0 \rightarrow \nu_{n0} = \nu_0(1-x_e)n - \nu_0x_en^2$$

$$\text{Assume } a = \frac{\nu_{n0}}{\nu_{n0}} = \frac{[1-x_e)n - x_en^2]}{[1-x_e)n - x_en^2]} \Rightarrow x_e = \frac{a(n-1)}{a(n-1)^2 - (n-1)^2}$$

$$\text{When } n=1, n=2 \quad a = \frac{\nu_{20}}{\nu_{10}} = \frac{5882}{2992} \approx 1.966$$

$$\Rightarrow x_e \approx 0.0164$$

$$\text{from } \nu_{n0} = \nu_0(1-x_e)n - \nu_0x_en^2 \rightarrow \nu_0 \approx 3092.5 \text{ cm}^{-1}$$

$$x_e = \frac{\nu_e}{4D_e} \rightarrow D_e \approx 47156.5 \text{ cm}^{-1}$$

$$\therefore \nu_{30} = \nu_0(1-x_e)3 - \nu_0x_e3^2$$

$$= 9670 \text{ cm}^{-1}$$

$$\nu_{10} = \nu_0(1-x_e) - \nu_0x_e = 2992$$

$$\nu_{20} = 2\nu_0(1-x_e) - 4\nu_0x_e = 5882$$

$$\therefore \text{We know } \begin{cases} \nu_0(1-x_e) = 3093 \\ \nu_0x_e = 51 \end{cases}$$

$$\therefore \nu_{30} = 3\nu_0(1-x_e) - 9\nu_0x_e = 3 \times 3093 - 9 \times 51 = 8670 \text{ cm}^{-1}$$

3. In a 1-D harmonic oscillator, obtain the oscillator strength for the quantum transitions.

from the eigenstate with quantum number $n=99$ to $n=100$?

$$f_{100,99} = \frac{2m}{\hbar} \frac{W_{100,99}}{W_{100,99}} = \frac{n\hbar}{2m\omega_0} = \frac{100\hbar}{2m\omega_0}$$

$$W_{100,99} = W_0 \quad (W_{100,99} = W_0 \text{ for } m=10-1)$$

$$\therefore f_{100,99} = \frac{2m}{\hbar} \frac{W_{100,99}}{W_{100,99}} = \frac{2m}{\hbar} \frac{W_{100,99}}{W_{100,99}} = 100$$

4. From Planck's formula, find an expression for the energy density of radiation in equilibrium conditions per wavelength interval i.e. $u(\lambda, T) = \frac{1}{V} \frac{dE}{d\lambda}$.

$$\begin{aligned} u(\lambda, T) &= \frac{dE}{V d\lambda} = \frac{\hbar \omega^3}{c^2 \lambda^3} \cdot \frac{1}{e^{\frac{\hbar \omega}{kT}} - 1} \quad \omega = \frac{2\pi c}{\lambda} \\ u(\lambda, T) &= \frac{1}{V} \frac{dE}{d\lambda} = \frac{1}{V} \frac{dE}{d\omega} \frac{d\omega}{d\lambda} = -\frac{8\pi}{\lambda^3} \cdot \frac{1}{e^{\frac{\hbar c}{\lambda kT}} - 1} \cdot \left(-\frac{2\pi c}{\lambda^2}\right) \\ &= \frac{16\pi^2 c}{\lambda^5} \cdot \frac{1}{e^{\frac{\hbar c}{\lambda kT}} - 1} \end{aligned}$$

$$\omega = \frac{2\pi c}{\lambda} \quad d\omega = -\frac{2\pi c}{\lambda^2} d\lambda \Rightarrow u(\omega, T) d\omega = -u(\lambda, T) d\lambda \quad u(\lambda, T) = \frac{8\pi h c}{\lambda^5} \cdot \frac{1}{e^{\frac{\hbar c}{\lambda kT}} - 1}$$

$$u(\lambda, T) = \frac{1}{V} \frac{dE}{d\lambda} \quad \therefore \frac{E}{V} = \int_0^\infty u(\lambda, T) d\lambda = 8\pi h c \int_0^\infty \frac{1}{\lambda^5} \frac{1}{e^{\frac{\hbar c}{\lambda kT}} - 1} d\lambda$$

$$x = \frac{\hbar c}{kT \lambda} \Rightarrow d\lambda = -\frac{\hbar c}{kT} \frac{1}{x^2} dx$$

$$\therefore \frac{dE}{d\lambda} = \frac{8\pi h c V}{\lambda^5} \frac{1}{e^{\frac{\hbar c}{\lambda kT}} - 1}$$

6. Consider the hydrogen $2p_z$ state ($=|210\rangle$) with $n=2, l=1, m=0$. Calculate the variance of r defined as $\text{Var}(r) \equiv \langle r^2 \rangle - \langle r \rangle^2$ in this state!

$$\int_0^\infty r^n e^{-r} dr = \frac{n!}{(n+1)!}$$

$$\text{make use of } |210\rangle = \frac{1}{4\sqrt{20} a_0^{3/2}} \frac{r}{a_0} e^{-\frac{r}{2a_0}} \cos\theta$$

$$\begin{aligned} \langle r^2 \rangle &= \int_0^\infty \int_0^\pi \int_0^{2\pi} \psi^* r^2 \psi dV = \frac{1}{32\pi a_0^5} \int_0^\infty r^4 e^{-\frac{r}{a_0}} \cos^2\theta dV \\ &= \frac{1}{32\pi a_0^5} \int_0^\infty r^5 e^{-\frac{r}{a_0}} dr \int_0^\pi \cos^2\theta \sin\theta d\theta \int_0^{2\pi} d\phi \\ &= \frac{1}{32\pi a_0^5} \cdot \frac{1}{120 a_0^6} \cdot \frac{2}{3} \cdot 2\pi = 5a_0 \end{aligned}$$

$$\langle r \rangle = \int_0^\infty \int_0^\pi \int_0^{2\pi} \psi^* r \psi dV = \frac{1}{32\pi a_0^5} \int_0^\infty r^3 e^{-\frac{r}{a_0}} dr \int_0^\pi \cos^2\theta \sin\theta d\theta \int_0^{2\pi} d\phi$$

$$= 1$$

$$\therefore \text{Var}(r) \equiv \langle r^2 \rangle - \langle r \rangle^2 = 5a_0 - 1$$