

Optics

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Chapter 1

History of Optics

Chapter 2

Electromagnetic Theory and photons

2.1 Longitudinal and Transverse

- Longitudinal: medium is in the direction of motion of wave.
- Transverse: medium is perpendicular to the motion of wave.

2.2 Wave Equation

$$\psi(x, t) = f(x + vt)$$

$$\begin{aligned} \begin{cases} \frac{\partial}{\partial x} = \frac{\partial}{\partial(x+vt)} \cdot \frac{\partial(x+vt)}{\partial x} = \frac{\partial}{\partial(x+vt)} \\ \frac{\partial}{\partial t} = \frac{\partial}{\partial(x+vt)} \cdot \frac{\partial(x+vt)}{t} = v \cdot \frac{\partial}{\partial(x+vt)} \end{cases} &\Rightarrow \begin{cases} \frac{\partial^2}{\partial x^2} = \frac{\partial^2}{\partial(x+vt)^2} \\ \frac{\partial^2}{\partial t^2} = v^2 \cdot \frac{\partial^2}{\partial(x+vt)^2} \end{cases} \Rightarrow \begin{cases} \frac{\partial^2 \psi}{\partial x^2} = \frac{\partial^2 \psi}{\partial(x+vt)^2} \\ \frac{\partial^2 \psi}{\partial t^2} = v^2 \cdot \frac{\partial^2 \psi}{\partial(x+vt)^2} \end{cases} \\ \Rightarrow \frac{\partial^2 \psi}{\partial x^2} = \frac{1}{v^2} \cdot \frac{\partial^2 \psi}{\partial t^2} &\Rightarrow \nabla^2 \psi(x, y, z) = \frac{1}{v^2} \cdot \frac{\partial^2 \psi(x, y, z)}{\partial t^2} \end{aligned}$$

2.3 Maxwell's Equation

Faraday's Induction Law

$$\oint_c \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \iint \vec{B} \cdot d\vec{S} \Rightarrow \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Gauss's Law

$$\oiint_A \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} \iiint_v \rho dV \Rightarrow \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\oiint_A \vec{B} \cdot d\vec{S} = 0 \Rightarrow \nabla \cdot \vec{B} = 0$$

Ampere's Circuital Law

$$\oint_C \vec{B} \cdot d\vec{l} = \mu_0 \iint_A \left(\vec{J} + \varepsilon_0 \frac{\partial \vec{E}}{\partial t} \right) \cdot d\vec{S} \Rightarrow \nabla \times \vec{B} = \mu_0 \varepsilon_0 \cdot \frac{\partial \vec{E}}{\partial t}$$

We can now take the derivatives of the 4 equations

$$\begin{cases} \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \\ \nabla \times \vec{B} = \mu_0 \varepsilon_0 \cdot \frac{\partial \vec{E}}{\partial t} \\ \nabla \cdot \vec{E} = \frac{\rho}{\varepsilon_0} \\ \nabla \cdot \vec{B} = 0 \end{cases} \Rightarrow \begin{cases} \nabla^2 \vec{E} = \mu_0 \varepsilon_0 \cdot \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{1}{c^2} \cdot \frac{\partial^2 \vec{E}}{\partial t^2} \\ \nabla^2 \vec{B} = \mu_0 \varepsilon_0 \cdot \frac{\partial^2 \vec{B}}{\partial t^2} = \frac{1}{c^2} \cdot \frac{\partial^2 \vec{B}}{\partial t^2} \end{cases}$$

Which indicates the speed of electromagnetic wave is exactly the speed of light.

Furthermore, it can be seen that the electric field and magnetic field are transverse. They are perpendicular to each other. We assume the electric field is parallel to the y-axis.

$$E_y(x, t) = E_0 \cos[\omega(t - x/c) + \varepsilon]$$

According to Faraday's Law

$$\frac{\partial E_y}{\partial x} = -\frac{\partial B_z}{\partial t}$$

We can calculate B_z

$$B_z = \frac{1}{c} \cdot E_0 \cos[\omega(t - x/c) + \varepsilon]$$

So that

$$E_y = v B_z = \begin{cases} \frac{1}{\sqrt{\mu_0 \varepsilon_0}} \cdot B_z & \text{in vacuum} \\ \frac{1}{\sqrt{\mu \varepsilon}} \cdot B_z & \text{not in vacuum} \end{cases}$$

2.4 Energy

$$u_E = \frac{1}{2} \cdot \frac{\varepsilon_0}{1} \cdot E^2 \quad u_B = \frac{1}{2} \cdot \frac{1}{\mu_0} \cdot B^2 \quad u_E = u_B \quad u = u_E + u_B = \varepsilon_0 E^2 = \frac{1}{\mu_0} B^2$$

$$S = uc = \varepsilon_0 c E^2 \quad (\text{Power: Transport of Energy per unit time across a unit area})$$

$$\vec{S} = \frac{1}{\mu} \cdot \vec{E} \times \vec{B} = c^2 \varepsilon \cdot \vec{E} \times \vec{B} \quad (\text{Poynting Vector}) \quad I = \frac{S}{2} = \frac{\varepsilon_0 c}{2} E_0^2 \quad (\text{Irradiance})$$

2.5 Radiation Pressure

$$P(t) = u = u_E + u_B = \frac{S}{c} \quad \text{Radiation Pressure equals energy density of the EM wave}$$

$$\langle P(t) \rangle_T = \frac{1}{2} \cdot \frac{S}{c} = \frac{I}{c} \quad \text{Average Radiation Pressure}$$

$$AP = \frac{\Delta p}{\Delta t} \Rightarrow Ac\Delta t P = c\Delta p \Rightarrow p_V = \frac{P}{c} = \frac{S}{c^2} \quad \text{Momentum per Volume}$$

2.6 Light in Bulk Matter

2.6.1 Speed of light and Dielectric Constant

$$v = \frac{1}{\sqrt{\varepsilon\mu}} \quad n = \frac{c}{v} = \sqrt{\frac{\varepsilon\mu}{\varepsilon_0\mu_0}} = \sqrt{\frac{\varepsilon}{\varepsilon_0}}$$

2.6.2 Dispersion

For gas and solid

$$m_e \frac{d^2x}{dt^2} + \gamma m_e \frac{dx}{dt} + m_e \omega_0^2 x = -eE(t)$$

$$E(t) = E_0 \exp(-i\omega t)$$

Assume

$$x = x_0 \exp(-i\omega t)$$

We got a solution

$$x_0 (\omega_0^2 - \omega^2 - i\gamma\omega) = -\frac{eE_0}{m_e}$$

$$x_0 = -\frac{eE_0}{m_e (\omega_0^2 - \omega^2 - i\gamma\omega)}$$

$$x(t) = -\frac{eE(t)}{m_e (\omega_0^2 - \omega^2 - i\gamma\omega)}$$

$$P(t) = -Nex(t) = \frac{Ne^2 E(t)}{m_e (\omega_0^2 - \omega^2 - i\gamma\omega)}$$

$$\varepsilon_r = \frac{\varepsilon}{\varepsilon_0} = n^2 = 1 + \frac{P}{\varepsilon_0 E} = 1 + \frac{Ne^2}{\varepsilon_0 m_e (\omega_0^2 - \omega^2 - i\gamma\omega)}$$

$$\begin{cases} \text{Re}(\varepsilon_r) = 1 + \frac{Ne^2 (\omega_0^2 - \omega^2)}{\varepsilon_0 m_e [(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2]} \\ \text{Im}(\varepsilon_r) = \frac{Ne^2 \gamma \omega}{\varepsilon_0 m_e [(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2]} \end{cases}$$

When $\gamma = 0$

$$\varepsilon_r = n^2 = 1 + \frac{Ne^2}{\varepsilon_0 m (\omega_0^2 - \omega^2)}$$

For metal

$$m_e \frac{d^2x}{dt^2} + \gamma m_e \frac{dx}{dt} = -eE(t)$$

$$\varepsilon_r = 1 - \frac{Ne^2}{\varepsilon_0 m_e (\omega^2 + i\gamma\omega)} = 1 - \frac{\omega_p^2}{\omega (\omega + i\gamma)}$$

Chapter 3

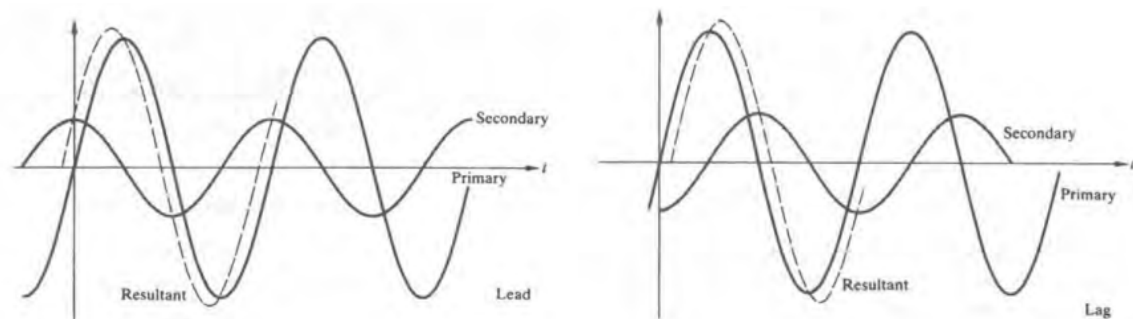
The Propagation of Light

3.1 Scattering and Interference

Destructive interference of the scattering light

- The denser the substance through which light advances, the less the lateral scattering.
- The longer the wavelength, the less the lateral scattering.
- On an overcast day, sky looks white because of large water droplets scatters all lights. On sunny day, sky only scatters blue light. And if there were no atmosphere, sky would be black as it is on moon.
- All molecules have electronic resonances in UV, the closer driving frequency is to a resonance, the more vigorously the oscillator responds. Blue and violet response more than red, sky is blue.

3.2 Speed of Light in Medium



If the phase of light in dielectric lags behind vacuum one, the resultant lags, and vice versa.

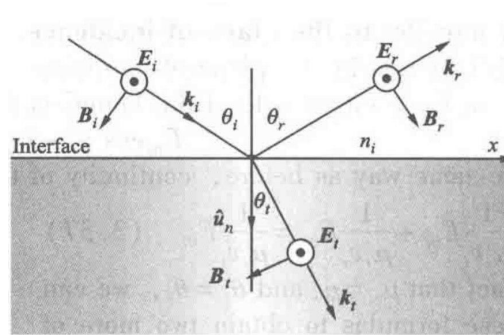
3.3 Internal and External Reflection



Beam I (internal reflection) and Beam II (external reflection) has 180° phase shift, when the gap between right part and left part of the glass in picture b becomes zero, two beams diminishes. This case is the same as picture a where the glass has not been cut.

3.4 The Fresnel Equations

3.4.1 Electric Field Perpendicular to Plane of Incidence

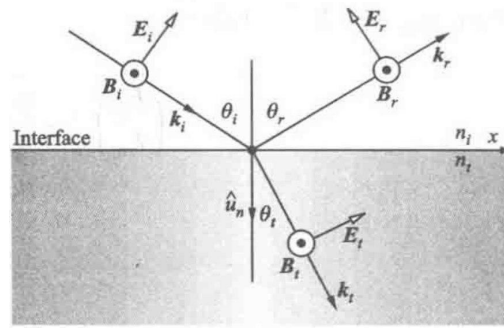


$$\begin{cases} E_i + E_r = E_t \\ B_i \cos \theta_i = B_r \cos \theta_r + B_t \cos \theta_t \end{cases} \Rightarrow \begin{cases} E = vB \\ v = \frac{c}{n} \end{cases}$$

We define the amplitude reflection coefficient r , the amplitude transmission coefficient t

$$\begin{cases} r = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t} \\ t = \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t} \end{cases} + n_i \sin \theta_i = n_t \sin \theta_t \Rightarrow \begin{cases} r_{\perp} = \frac{\sin(\theta_i - \theta_t)}{\sin(\theta_i + \theta_t)} \\ t_{\perp} = \frac{2 \sin \theta_t \cos \theta_i}{\sin(\theta_i + \theta_t)} \end{cases}$$

3.4.2 Electric Field Parallel to Plane of Incidence



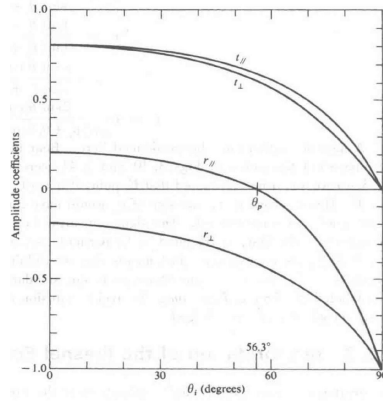
$$\begin{cases} B_i + B_r = B_t \\ E_i \cos \theta_i = E_r \cos \theta_r + E_t \cos \theta_t \end{cases} \Rightarrow \begin{cases} E = vB \\ v = \frac{c}{n} \end{cases}$$

We define the amplitude reflection coefficient r , the amplitude transmission coefficient t

$$\begin{cases} r = \frac{n_t \cos \theta_i - n_i \cos \theta_t}{n_t \cos \theta_i + n_i \cos \theta_t} \\ t = \frac{2n_i \cos \theta_i}{n_t \cos \theta_i + n_i \cos \theta_t} \end{cases} + n_i \sin \theta_i = n_t \sin \theta_t \Rightarrow \begin{cases} r_{\parallel} = \frac{\sin(2\theta_i) - \sin(2\theta_t)}{\sin(2\theta_i) + \sin(2\theta_t)} = \frac{\tan(\theta_i - \theta_t)}{\tan(\theta_i + \theta_t)} \\ t_{\parallel} = \frac{2 \sin \theta_t \theta_i}{\sin(\theta_i + \theta_t) \cos(\theta_i - \theta_t)} \end{cases}$$

3.5 Polarization Angle

$$\theta_p = \arctan\left(\frac{\theta_p}{\theta_i}\right)$$



3.6 Critical Angle

$$\theta_c = \arcsin\left(\frac{n_t}{n_i}\right)$$

3.7 Phase Shift

When $\theta_i = 0$

$$r_{\perp} = -r_{\parallel} = \frac{n_i - n_t}{n_i + n_t}$$

$$t_{\parallel} = t_{\perp} = \frac{2n_i}{n_i + n_t}$$

While $n_i > n_t$ (Inner reflection)

$$r_{\parallel} < 0$$

$$r_{\perp} > 0$$

No phase shift.

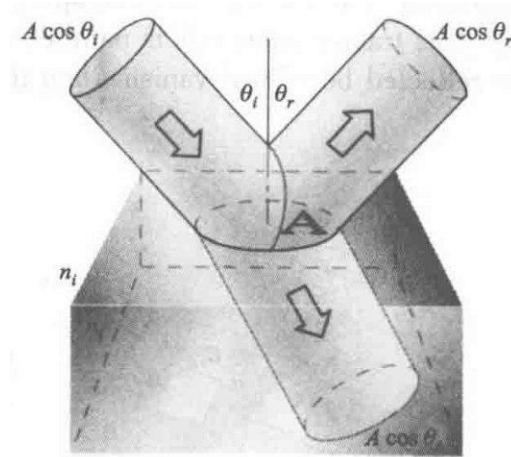
While $n_i < n_t$ (Outer reflection)

$$r_{\parallel} > 0$$

$$r_{\perp} < 0$$

Phase shifted by π .

3.8 Reflectance and Transmittance



$$\begin{cases} R = \frac{I_r A \cos \theta_r}{I_i A \cos \theta_i} = \frac{I_r}{I_i} \\ T = \frac{I_t A \cos \theta_t}{I_i A \cos \theta_i} = \frac{I_t \cos \theta_t}{I_i \cos \theta_i} \end{cases}$$

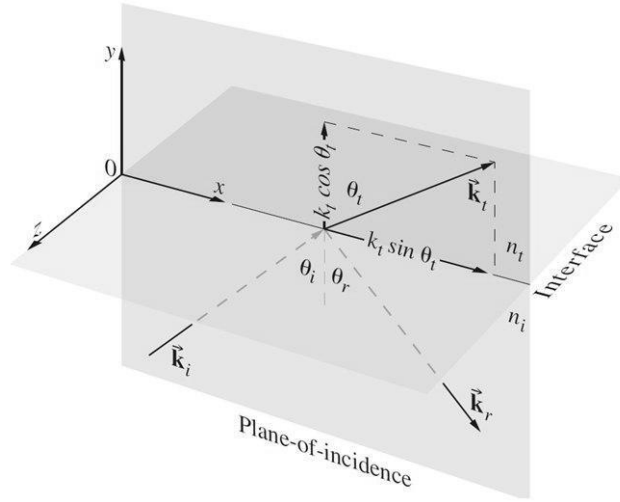
$$\vec{S} = c^2 \epsilon_0 \vec{E} \times \vec{B} \Rightarrow I = \frac{1}{2} \epsilon v E_0^2 = \frac{1}{2} \epsilon_0 \epsilon_r v E_0^2 = \frac{1}{2} \epsilon_0 n^2 v E_0^2 = \frac{1}{2} \epsilon_0 n c E_0^2$$

$$\begin{cases} R = \frac{I_r}{I_i} = \left(\frac{E_{0r}}{E_{0i}} \right)^2 = r^2 \\ T = \frac{I_t \cos \theta_t}{I_i \cos \theta_i} = \left(\frac{n_t \cos \theta_t}{n_i \cos \theta_i} \right) \left(\frac{E_{0t}}{E_{0i}} \right)^2 = \left(\frac{n_t \cos \theta_t}{n_i \cos \theta_i} \right) t^2 \end{cases} \Rightarrow \begin{cases} R_{\perp} = r_{\perp}^2 \\ R_{\parallel} = r_{\parallel}^2 \\ T_{\perp} = \left(\frac{n_t \cos \theta_t}{n_i \cos \theta_i} \right) t_{\perp}^2 \\ T_{\parallel} = \left(\frac{n_t \cos \theta_t}{n_i \cos \theta_i} \right) t_{\parallel}^2 \end{cases} \Rightarrow \begin{cases} R_{\perp} + T_{\perp} = 1 \\ R_{\parallel} + T_{\parallel} = 1 \\ R + T = 1 \end{cases}$$

When $\theta_i = 0$, any distinction between the parallel and perpendicular components of R and T vanishes. Thus

$$\begin{cases} R = R_{\parallel} = R_{\perp} = \left(\frac{n_t - n_i}{n_t + n_i} \right)^2 \\ T = T_{\parallel} = T_{\perp} = \frac{4n_t n_i}{(n_i + n_t)^2} \end{cases}$$

3.9 The Evanescent Wave



$$\vec{E}_t = \vec{E}_{0t} \exp \left[i \left(\vec{k}_t \cdot \vec{r} - \omega t \right) \right]$$

$$\vec{k}_t \cdot \vec{r} = k_{tx}x + k_{ty}y$$

$$k_{tx} = k_t \sin \theta_t = \left(\frac{n_i}{n_t} \right) k_t \sin \theta_i = n_i k_0 \sin \theta_i$$

$$k_{ty} = k_t \cos \theta_t = i k_t \sqrt{\frac{n_i^2 \sin^2 \theta_i}{n_t^2} - 1} = i\beta$$

$$\vec{E}_t = \vec{E}_{0t} \exp(-\beta y) \exp[i(n_i k_0 x \sin \theta_i - \omega t)]$$

3.10 Optical Properties of Metals

The index of refraction of metal is complex

$$\tilde{n} = n_R - in_I$$

$$\nabla \times \vec{H} = \epsilon_0 \epsilon_r \frac{\partial \vec{E}}{\partial t} + \sigma \vec{E} = -i\omega \epsilon_0 \epsilon_r \vec{E} + \sigma \vec{E} = -i\omega \epsilon_0 \tilde{\epsilon}_r \vec{E}$$

Whereas

$$\tilde{\epsilon}_r = \epsilon_r + i \frac{\sigma}{\omega \epsilon_0}$$

$$\tilde{n}^2 = \tilde{\epsilon}_r = \epsilon_r + i \frac{\sigma}{\omega \epsilon_0} = (n_R + i n_I)^2$$

Since $\frac{\sigma}{\omega \epsilon_0 \epsilon_r} \gg 1$

$$n_I \approx n_R = \sqrt{\frac{\sigma}{2\omega \epsilon_0}}$$

Skin depth

$$\delta = \sqrt{\frac{1}{2\omega \mu_0 \sigma}}$$

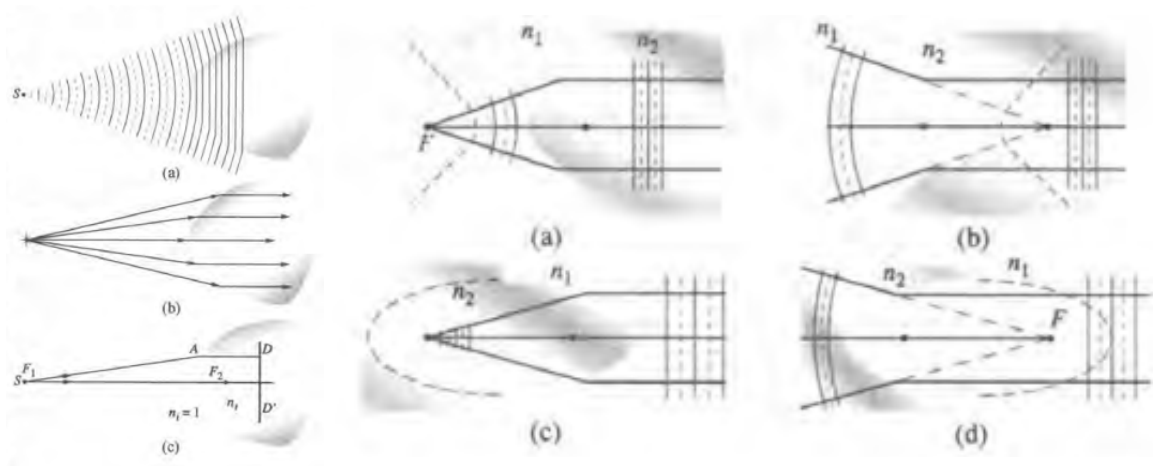
Reflectance

$$R = \left| \frac{n_i - n_t}{n_i + n_t} \right|^2 = \left(\frac{\tilde{n} - 1}{\tilde{n} + 1} \right) \left(\frac{\tilde{n} - 1}{\tilde{n} + 1} \right)^* = \frac{(n_R - 1)^2 + n_I^2}{(n_R + 1)^2 + n_I^2}$$

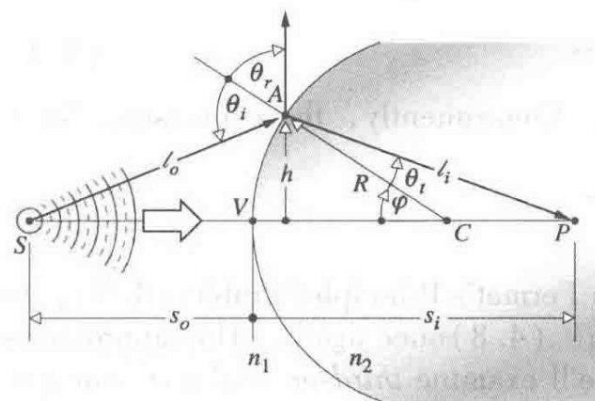
Chapter 4

Geometrical Optics

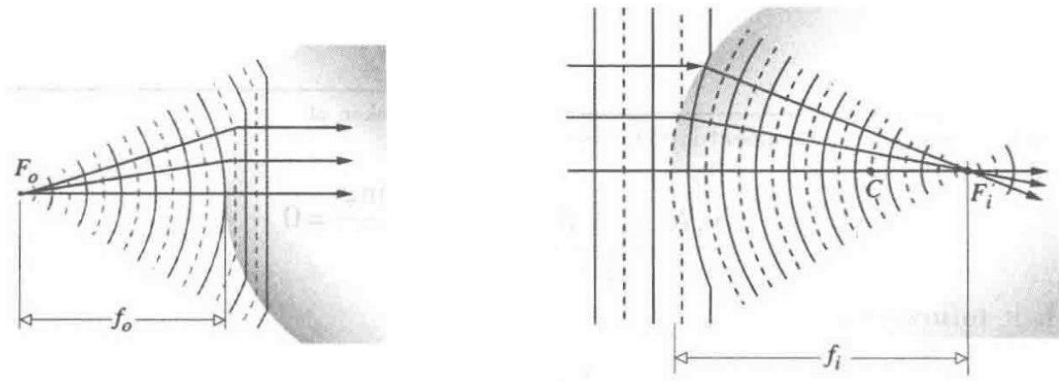
4.1 Aspherical Surface



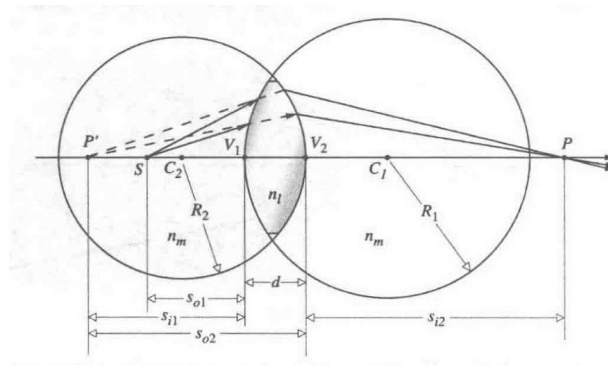
4.2 Refraction at a Spherical Interface



$$\frac{n_1}{s_o} + \frac{n_2}{s_i} = \frac{n_2 - n_1}{R} = \Phi \Rightarrow \begin{cases} f_o = \frac{n_1}{n_2 - n_1} R & (s_i = \infty) \\ f_i = \frac{n_2}{n_2 - n_1} R & (s_o = \infty) \end{cases}$$



4.3 Lenses



$$\frac{n_m}{s_{o1}} + \frac{n_m}{s_{i2}} = (n_l - n_m) \left(\frac{1}{R_1} - \frac{1}{R_2} \right) + \frac{n_l d}{(s_{i1} - d) s_{i1}}$$

For lenses in the air, where $n_m = 1$

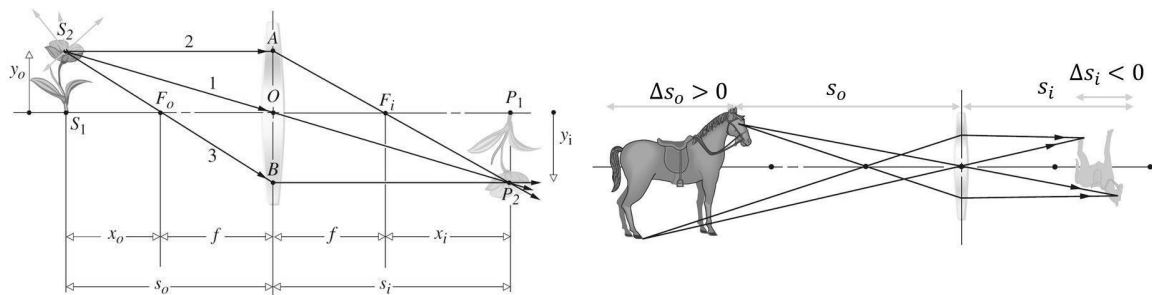
$$\frac{1}{s_{o1}} + \frac{1}{s_{i2}} = (n_l - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right) + \frac{n_l d}{(s_{i1} - d) s_{i1}}$$

For thin lenses, $d \approx 0$

$$\frac{1}{s_{o1}} + \frac{1}{s_{i2}} = (n_l - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right) = \frac{1}{f}$$

Which is the Lensmaker's Formula.

4.4 Magnification



$$\begin{cases} \frac{y_o}{|y_i|} = \frac{f}{x_i} \\ \frac{|y_i|}{y_o} = \frac{f}{x_o} \end{cases} \Rightarrow x_o x_i = f^2 \quad (\text{Newton's formula})$$

Transverse Magnification

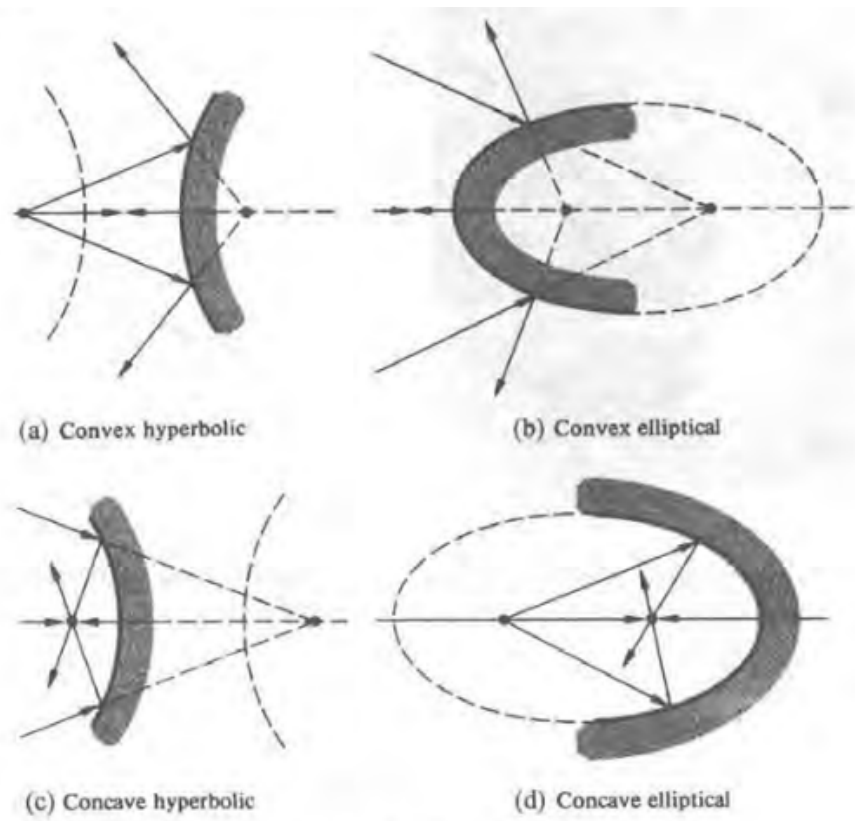
$$M_T = \frac{y_i}{|y_o|} = -\frac{s_o}{s_i} = -\frac{f}{x_o} = -\frac{x_i}{f}$$

Longitudinal Magnification

$$M_L = \frac{dx_i}{dx_o} = \frac{d}{dx_o} \left(\frac{f^2}{x_o} \right) = -\frac{f^2}{x_o^2} = -M_T^2$$

4.5 Mirrors

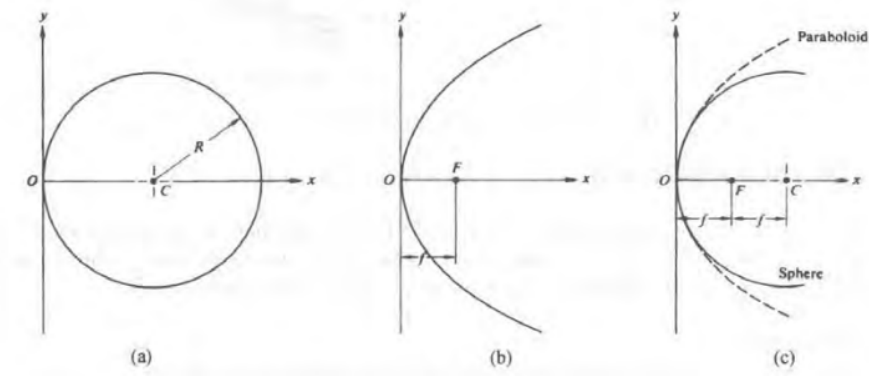
4.5.1 Aspherical Mirrors



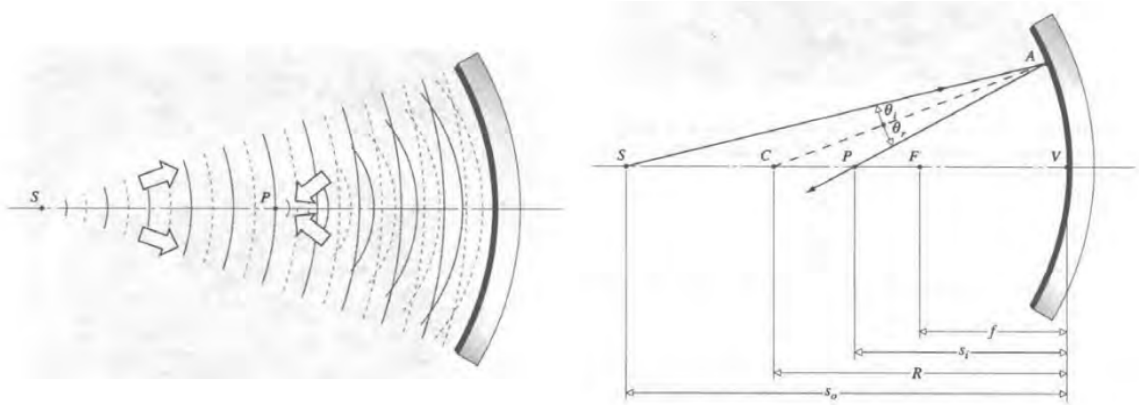
Precise aspheric surface are difficult and expensive to fabricate.

4.5.2 Spherical Mirrors

The difference between spherical and paraboloidal mirror will be appreciable only if y is large.



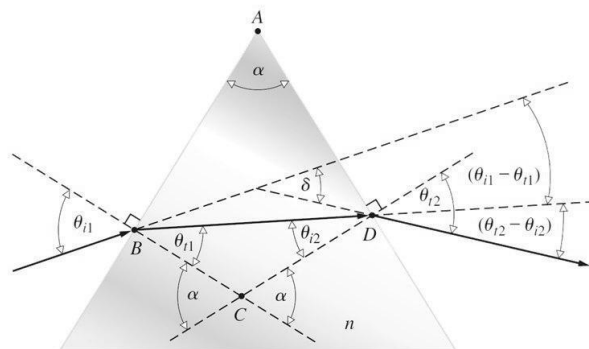
Mirror Formula



$$\frac{\overline{SC}}{\overline{SA}} = \frac{\overline{CP}}{\overline{PA}} + \begin{cases} \overline{SC} = s_o - |R| = s_o + R \\ \overline{CP} = -s_i + |R| = -(s_i + R) \\ \overline{SA} = s_o \\ \overline{PA} = s_i \end{cases} \Rightarrow \frac{s_o + R}{s_o} = -\frac{s_i + R}{s_i} \Rightarrow \frac{1}{s_o} + \frac{1}{s_i} = -\frac{2}{R}$$

$$\begin{cases} s_o = \infty \Rightarrow s_i = f = -R/2 \\ s_i = \infty \Rightarrow s_o = f = -R/2 \end{cases} \Rightarrow \begin{cases} s_o = 2f \Rightarrow s_i = 2f \\ s_o = f \Rightarrow s_i = \infty \\ s_i = f \Rightarrow s_o = \infty \end{cases}$$

4.6 Prism



$$\begin{cases} \delta = (\theta_{i1} - \theta_{t1}) + (\theta_{t2} - \theta_{i2}) \\ \alpha = \theta_{t1} + \theta_{i2} \end{cases} \Rightarrow \delta = \theta_{i1} + \theta_{t2} - \alpha$$

$$\begin{aligned} \theta_{t2} &= \sin^{-1}(n \sin \theta_{i2}) = \sin^{-1}[n \sin(\alpha - \theta_{t1})] = \sin^{-1}[n(\sin \alpha \cos \theta_{t1} - \cos \alpha \sin \theta_{t1})] \\ &= \sin^{-1}\left[n\left(\sin \alpha \sqrt{1 - \sin^2 \theta_{t1}} - \cos \alpha \sin \theta_{t1}\right)\right] = \sin^{-1}\left[\sin \alpha \sqrt{n^2 - \sin^2 \theta_{i1}} - \cos \alpha \sin \theta_{i1}\right] \end{aligned}$$

$$\delta = \theta_{i1} + \sin^{-1}\left[\sin \alpha \sqrt{n^2 - \sin^2 \theta_{i1}} - \cos \alpha \sin \theta_{i1}\right] - \alpha$$

Minimum deviation:

$$\begin{cases} \frac{d\delta}{d\theta_{i1}} = 0 & \Rightarrow & 1 + \frac{d\theta_{t2}}{d\theta_{i1}} = 0 & \Rightarrow & d\theta_{i1} = -d\theta_{t2} \\ d\alpha = 0 & \Rightarrow & d(\theta_{t1} + \theta_{i2}) = 0 & \Rightarrow & d\theta_{t1} = -d\theta_{i2} & \Rightarrow & \frac{\cos \theta_{i1}}{\cos \theta_{t2}} = \frac{\cos \theta_{t1}}{\cos \theta_{i2}} \\ \sin \theta_{i1} = n \sin \theta_{t1} & \Rightarrow & \cos \theta_{i1} d\theta_{i1} = n \cos \theta_{t1} d\theta_{t1} \\ \sin \theta_{t2} = n \sin \theta_{i2} & \Rightarrow & \cos \theta_{t2} d\theta_{t2} = n \cos \theta_{i2} d\theta_{i2} \end{cases}$$

\Rightarrow

$$\frac{1 - \sin^2 \theta_{i1}}{1 - \sin^2 \theta_{t2}} = \frac{1 - n^2 \sin^2 \theta_{i1}}{1 - n^2 \sin^2 \theta_{t2}} \Rightarrow \theta_{i1} = \theta_{t2}$$

This means that the ray for which the deviation is a minimum traverses the prism symmetrically.

$$\theta_{t1} = \theta_{i2} = \alpha/2 \quad \theta_{i1} = \theta_{t2} = \frac{\delta_m + \alpha}{2} \quad n = \frac{\sin[(\delta_m + \alpha)/2]}{\sin \alpha/2}$$

Chapter 5

The Superstition of Waves

5.1 The Addition of Waves

5.1.1 The Algebraic Method

$$E(x, t) = E_0 \sin[\omega t - (kx + \varepsilon)]$$

let

$$\alpha(x, \varepsilon) = -(kx + \varepsilon)$$

Then

$$E(x, t) = E_0 \sin[\omega t + \alpha(x, \varepsilon)]$$

Two waves of the same frequency

$$\begin{cases} E_1 = E_{01} \sin(\omega t + \alpha_1) \\ E_2 = E_{02} \sin(\omega t + \alpha_2) \end{cases}$$

$$\begin{aligned} E = E_1 + E_2 &= E_{01} (\sin \omega t \cos \alpha_1 + \cos \omega t \sin \alpha_1) + E_{02} (\sin \omega t \cos \alpha_2 + \cos \omega t \sin \alpha_2) \\ &= (E_{01} \cos \alpha_1 + E_{02} \cos \alpha_2) \sin \omega t + (E_{01} \sin \alpha_1 + E_{02} \sin \alpha_2) \cos \omega t \\ &= E_0 \cos \alpha \sin \omega t + E_0 \sin \alpha \cos \omega t \\ &= E_0 \sin(\omega t + \alpha) \end{aligned}$$

$$\begin{cases} E_0 \cos \alpha = E_{01} \cos \alpha_1 + E_{02} \cos \alpha_2 \\ E_0 \sin \alpha = E_{01} \sin \alpha_1 + E_{02} \sin \alpha_2 \end{cases} \Rightarrow \begin{cases} E_0^2 = E_{01}^2 + E_{02}^2 + 2E_{01}E_{02} \cos(\alpha_2 - \alpha_1) \\ \tan \alpha = \frac{E_{01} \sin \alpha_1 + E_{02} \sin \alpha_2}{E_{01} \cos \alpha_1 + E_{02} \cos \alpha_2} \end{cases}$$

The phase difference

$$\delta = (kx_1 + \varepsilon_1) - (kx_2 + \varepsilon_2) = \frac{2\pi}{\lambda} (x_1 - x_2) + (\varepsilon_1 - \varepsilon_2)$$

When $E_{01} = E_{02}$ and $\alpha_2 - \alpha_1 = \Delta x$

$$E_0^2 = 2E_{01}^2 + 2E_{01}^2 \cos(k\Delta x) = 2E_{01}^2 [1 + \cos(k\Delta x)]$$

$$\cos 2x = 2 \cos^2 x - 1 \Rightarrow \cos(k\Delta x) = 2 \cos^2\left(\frac{k\Delta x}{2}\right) - 1$$

\Rightarrow

$$E_0^2 = 2E_{01}^2 \cos^2\left(\frac{k\Delta x}{2}\right)$$

Period of the amplitude of addition

$$\frac{k\Delta x}{2} = \frac{\pi}{2} \Rightarrow k(\alpha_2 - \alpha_1) = \pi \Rightarrow \Delta x = \alpha_2 - \alpha_1 = \frac{\lambda}{2}$$

5.1.2 The Complex Method

$$E_1 = E_{01} \cos(kx \pm \omega t) \Rightarrow \tilde{E}_1 = E_{01} \exp[i(kx \pm \omega t)]$$

$$\begin{cases} E_1 = E_{01} \exp[i\alpha_1] \\ E_2 = E_{02} \exp[i\alpha_2] \\ E_0 = E_1 + E_2 \end{cases}$$

$$\begin{aligned} E_0^2 &= (E_{01} \exp[i\alpha_1] + E_{02} \exp[i\alpha_2]) \cdot (E_{01} \exp[-i\alpha_1] + E_{02} \exp[-i\alpha_2]) \\ &= E_{01}^2 + E_{02}^2 + 2E_{01}E_{02} \cos(\alpha_1 - \alpha_2) \end{aligned}$$

5.1.3 Phasor Addition Method

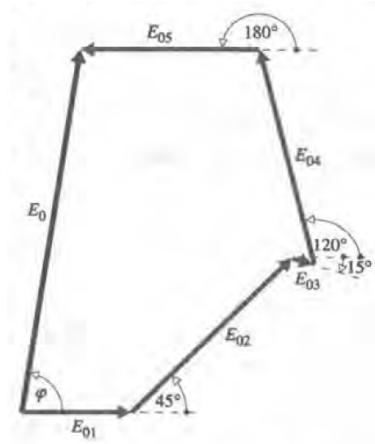


Figure 5.1: Phasor Addition Method

5.2 Standing Waves

$$E_L = E_{0t} \sin(kx - \omega t)$$

$$E_R = E_{0t} \sin(kx - \omega t) \quad E = E_L + E_R$$

$$\sin \alpha + \sin \beta = 2 \sin \frac{\alpha + \beta}{2} \cos \frac{\alpha - \beta}{2}$$

\Rightarrow

$$E = 2E_{0t} \sin kx \cos \omega t$$

5.3 Addition of Waves of Different Frequency

$$\begin{cases} E_1 = E_{01} \cos(k_1 x - \omega_1 t) \\ E_2 = E_{02} \cos(k_2 x - \omega_2 t) \\ E = E_1 + E_2 \end{cases}$$

$$\cos \alpha + \cos \beta = 2 \cos \frac{\alpha + \beta}{2} \cos \frac{\alpha - \beta}{2}$$

\Rightarrow

$$\begin{aligned} E &= E_{01} [\cos(k_1 x - \omega_1 t) + \cos(k_2 x - \omega_2 t)] \\ &= 2E_{01} \cos \frac{1}{2} [(k_1 + k_2)x - (\omega_1 + \omega_2)t] \times \cos \frac{1}{2} [(k_1 - k_2)x - (\omega_1 - \omega_2)t] \end{aligned}$$

Define

$$\begin{aligned} \bar{\omega} &= \frac{1}{2} (\omega_1 + \omega_2) & \omega_m &= \frac{1}{2} (\omega_1 - \omega_2) \\ \bar{k} &= \frac{1}{2} (k_1 + k_2) & k_m &= \frac{1}{2} (k_1 - k_2) \end{aligned}$$

Then

$$E = 2E_{01} \cos(k_m x - \omega_m t) \cos(\bar{k}x - \bar{\omega}t) = E_0(x, t) \cos(\bar{k}x - \bar{\omega}t)$$

Noted that

$$\begin{aligned} \bar{\omega} &= \frac{1}{2} (\omega_1 + \omega_2) & \omega_m &= \frac{1}{2} (\omega_1 - \omega_2) \\ \bar{k} &= \frac{1}{2} (k_1 + k_2) & k_m &= \frac{1}{2} (k_1 - k_2) \end{aligned} \gg$$

$E_0 = 2E_{01} \cos(k_m x - \omega_m t)$ varies far less frequently than $\cos(\bar{k}x - \bar{\omega}t)$

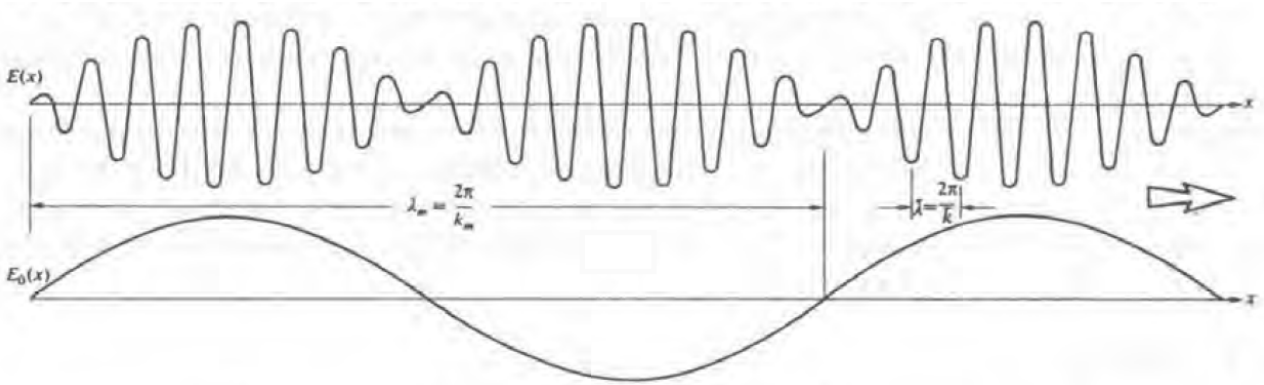


Figure 5.2: Standing Wave

| | |
|------------------------|-------------|
| Beat Frequency (Time) | $2\omega_m$ |
| Beat Frequency (Space) | $2k_m$ |

| | |
|-----------------|------------------------------|
| Group Frequency | $v_g = \omega_m/k_m$ |
| Phase Velocity | $v_p = \bar{\omega}/\bar{k}$ |

5.4 Light in Dispersible Media

| | |
|------------------------|--|
| Average Phase Velocity | $\bar{v}_p = \frac{c}{\bar{n}}$ |
| Group Velocity | $v_g = \frac{c}{\bar{n}} \left(1 + \frac{\bar{\lambda}}{\bar{n}} \frac{\Delta n}{\Delta \lambda} \right)$ |

| | |
|----------------------------|-------------------|
| Normal Dispersion Media | $\bar{v}_p > v_g$ |
| Anomalous Dispersion Media | $\bar{v}_p < v_g$ |

Chapter 6

Polarization

6.1 Circular Polarization

$$\begin{cases} \vec{E}_x(z, t) = \vec{i}E_0 \cos(kx - \omega t) \\ \vec{E}_y(z, t) = \vec{j}E_0 \sin(kx - \omega t) \end{cases} \Rightarrow \vec{E} = E_0 [\vec{i} \cos(kx - \omega t) + \vec{j} \sin(kx - \omega t)] \quad \text{Right-circularly polarized}$$

$$\begin{cases} \vec{E}_x(z, t) = \vec{i}E_0 \cos(kz - \omega t) \\ \vec{E}_y(z, t) = -\vec{j}E_0 \sin(kz - \omega t) \end{cases} \Rightarrow \vec{E} = E_0 [\vec{i} \cos(kx - \omega t) - \vec{j} \sin(kx - \omega t)] \quad \text{Left-circularly polarized}$$

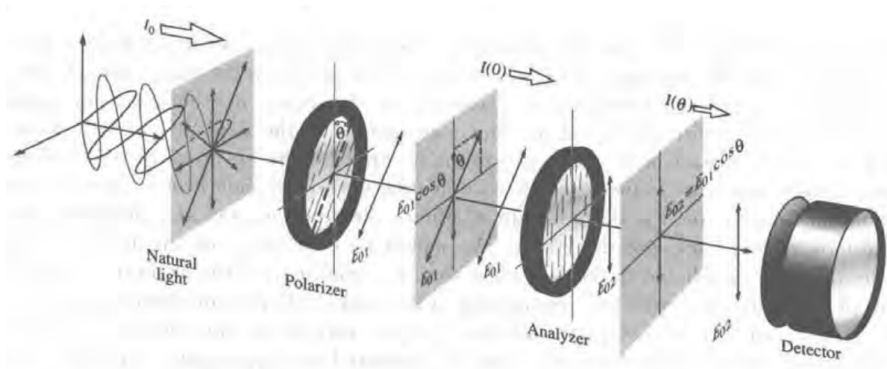
6.2 Elliptical Polarization

$$\begin{cases} \vec{E}_x = E_{0x} \cos(kx - \omega t) \\ \vec{E}_y = E_{0y} \cos(kz - \omega t + \epsilon) \end{cases} \quad \text{Elliptical Polarization}$$

6.3 Angular Momentum

$$\frac{dE}{dt} = \omega M = \omega \frac{dL}{dt} \Rightarrow L = \frac{E}{\omega} = \frac{h\nu}{\omega} = \pm \hbar = \begin{cases} -\hbar & \text{Right-circularly polarized} \\ +\hbar & \text{Left-circularly polarized} \end{cases}$$

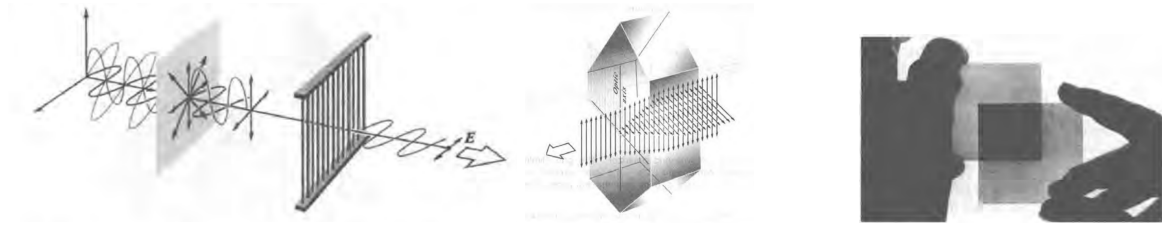
6.4 Malus's Law



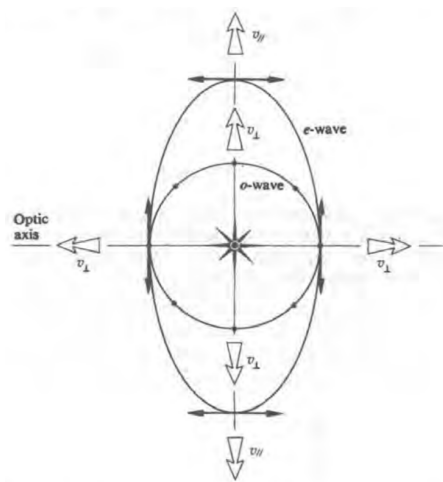
$$\begin{cases} E_{02} = E_{01} \cos \theta \\ I(\theta) = \frac{c\epsilon_0}{2} E_{01}^2 \cos^2 \theta = I(0) \cos^2 \theta \end{cases}$$

6.5 Dichroism

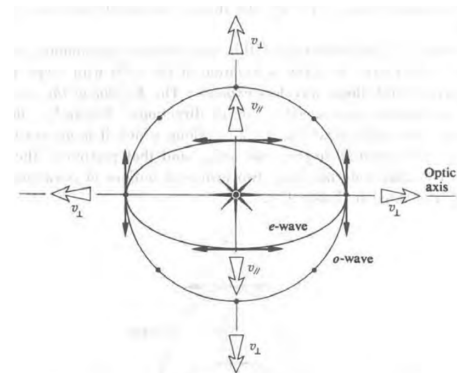
6.5.1 The Wire-Grid Polarizer and Dichroic Crystals, Polaroid



6.6 Birefringent Crystals



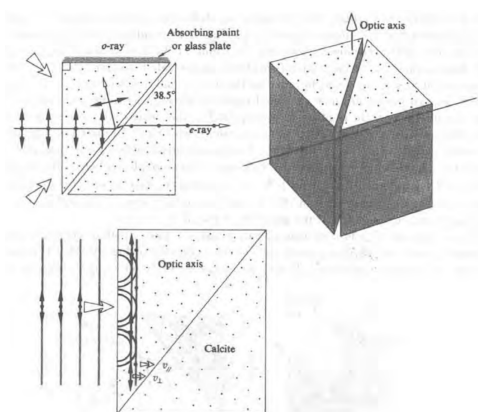
(a) $v_{\perp} < v_{\parallel} \Rightarrow n_o > n_e \Rightarrow$ negative uniaxial



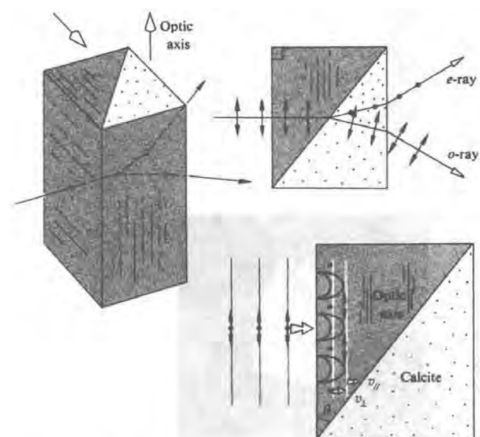
(b) $v_{\perp} > v_{\parallel} \Rightarrow n_o > n_e \Rightarrow$ positive uniaxial

Figure 6.2: negative and positive uniaxial

6.7 Polarizers



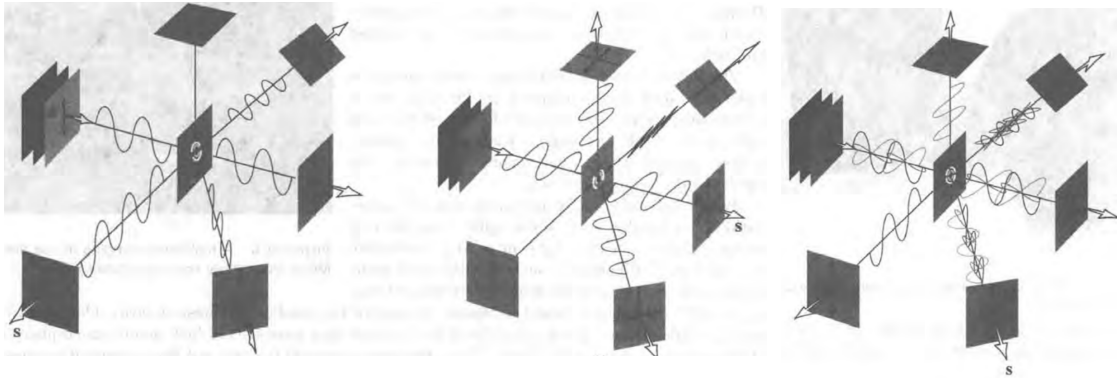
(a) The Glan-Foucault Prism



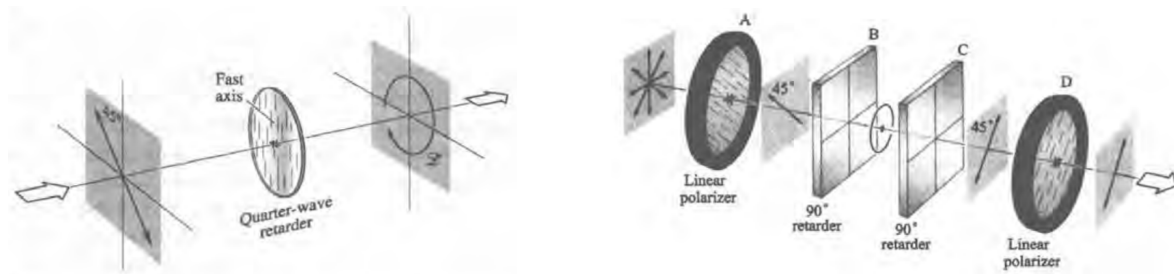
(b) The Wollaston Prism

Figure 6.3: Two Birefringent Polarizers

6.8 Scattering and Polarization



6.9 Retarders



(a) Quarter-wave Retarder

(b) Two Linear Polarizers and Two Quarter-wave Retarders

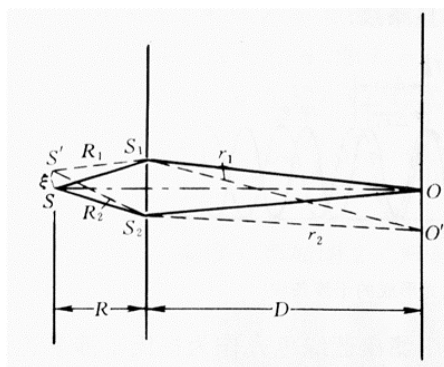
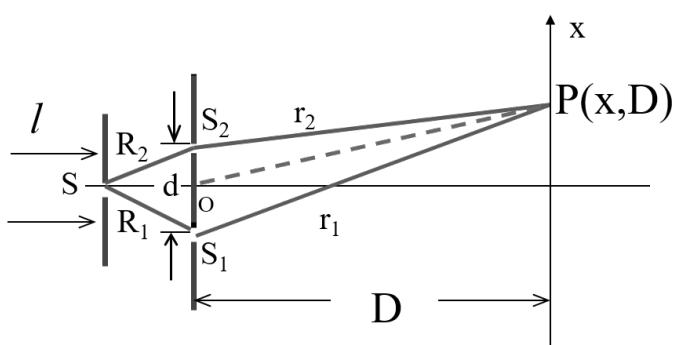
Figure 6.5: Quarter-wave Retarder and its Application

$$d(n_o - n_e) = \frac{4m + 1}{4} \lambda_0$$

Chapter 7

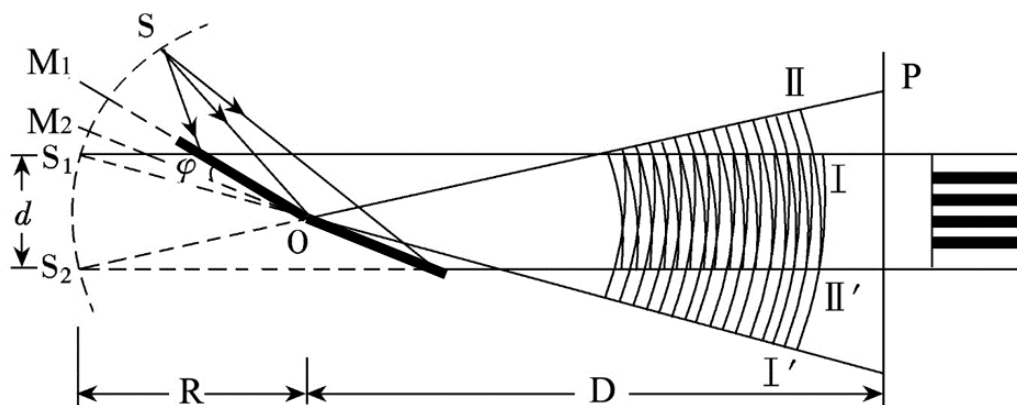
Interference

7.1 Young's Experiment



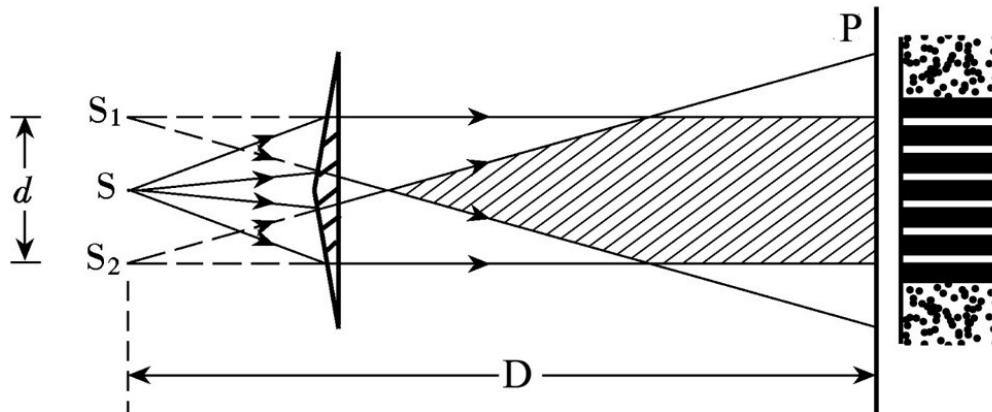
$$\Delta x = \frac{D}{d} \lambda \quad b \leq \lambda R \frac{1}{d} \quad I = I_0 \cos^2 \left(\frac{d\pi}{D\lambda} x \right) \quad x_0 = -\frac{D}{R} \xi$$

7.2 Fresnel's Double Mirror



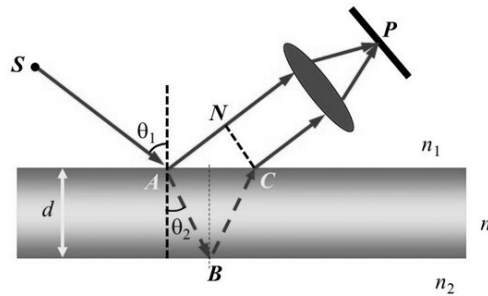
$$x_{white} = k\lambda \frac{D}{d} \quad x_{black} = \frac{2k+1}{2} \lambda \frac{D}{d}$$

7.3 Fresnel's Double Prism



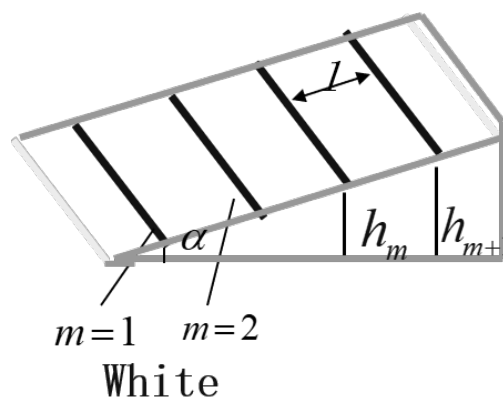
$$x_{white} = k\lambda \frac{D}{d} \quad x_{black} = \frac{2k+1}{2}\lambda \frac{D}{d}$$

7.4 Equal Inclination Interference



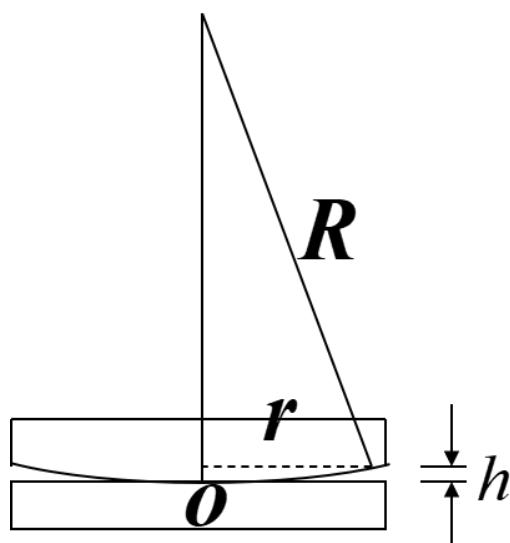
$$\Lambda = \begin{cases} 2nk_0d \cos \theta_2 \pm \pi & n_1 > n_2 < n_3 \text{ OR } n_1 < n_2 > n_3 \\ 2nk_0d \cos \theta_2 & n_1 < n_2 < n_3 \text{ OR } n_1 > n_2 > n_3 \end{cases}$$

7.5 Equal Thickness Interference



$$e = \Delta h = \frac{\lambda}{2n} \quad l = \frac{e}{\sin \alpha} = \frac{\lambda}{2n\alpha} \approx \frac{\lambda}{2n\alpha}$$

7.6 Newton's Rings



$$\Delta = 2nh + \frac{\lambda}{2} = \begin{cases} k\lambda & \text{White} \\ \left(k + \frac{1}{2}\right)\lambda & \text{Black} \end{cases}$$

$$h = R - \sqrt{R^2 - r^2} = R \left[1 - \sqrt{1 - \left(\frac{r}{R}\right)^2} \right] \approx \frac{r^2}{2R}$$

$$r^2 = \begin{cases} \left(k - \frac{1}{2}\right) \frac{R\lambda}{n} & \text{White} \\ \frac{kR\lambda}{n} & \text{Black} \end{cases}$$

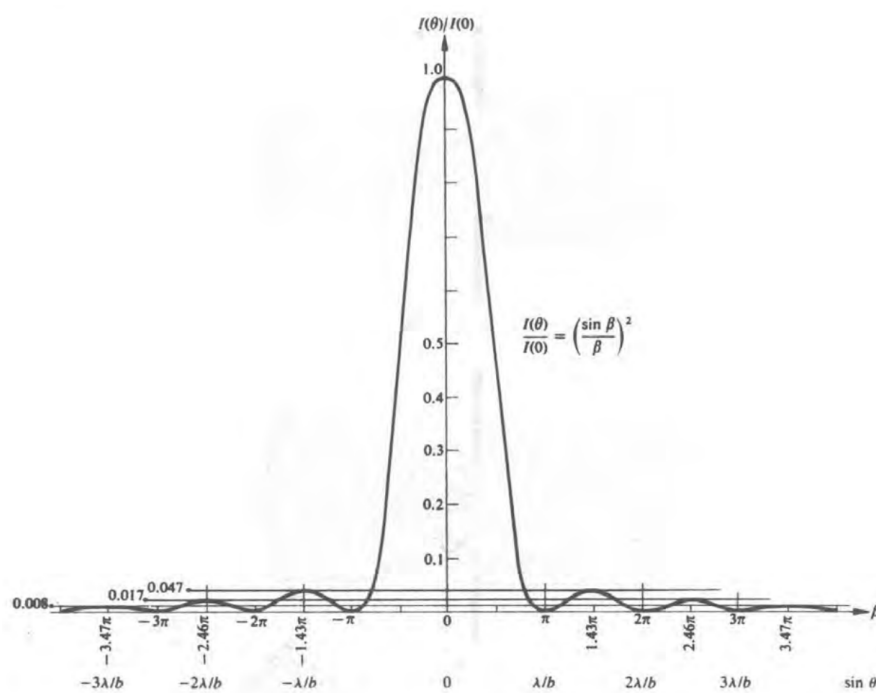
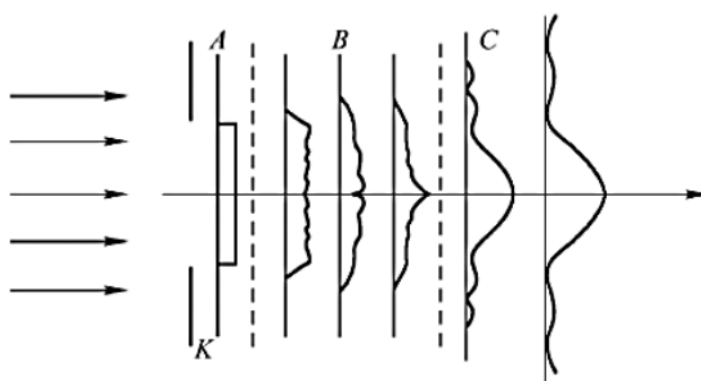
Chapter 8

Diffraction

Chapter 9

Fraunhofer and Fresnel Diffraction

9.1 Fraunhofer Diffraction



White fringes:

$$\left\{ \begin{array}{l} b \sin \theta = 0 \\ \sin \theta = \pm (2m + 1) \cdot \frac{\lambda}{2b} \end{array} \right. \quad \begin{array}{l} \text{Central Fringe} \\ m = 1, 2, 3, \dots \end{array} \quad \left\{ \begin{array}{l} \Delta \theta_0 = 2 \cdot \frac{\lambda}{b} \\ \Delta \theta = \frac{\lambda}{b} \end{array} \right.$$

Dark fringes:

$$\sin \theta = \pm m \cdot \frac{\lambda}{b} \quad m = 1, 2, 3, \dots$$