# Semiconductor Fundamentals – (III)

2.5 Boltzmann approximation & E<sub>F</sub>, n, p

2.6 Carrier drift and diffusion

Material developed by Prof. C. Z. Zhao

#### Last lecture:

- Negative charges:
  - Conduction electrons (density = n) mobile
  - Ionized acceptor atoms (density =  $N_A^{-1}$ )
- Positive charges:
  - Holes (density = p) mobile
  - > Ionized donor atoms (density =  $N_D^+$ ) immobile
- The net charge density (C/cm³) in a semiconductor is

$$\rho = q(p - n + N_D^+ - N_A^-)$$

• Law of Mass Action:  $n \cdot p = n_i^2$ 

#### 质量作用定律

How to deduce the relationship between  $E_F$  and n/p?

## **2.5** Boltzmann approximation & $E_F$ , n, p

Fermi function and Fermi level

Density of States

Boltzmann Approximation

Electron and hole Concentrations

# Thermal Equilibrium

- No external forces are applied:
  - electric field = 0, magnetic field = 0
  - mechanical stress = 0
  - > no light
- Dynamic situation in which every process is balanced by its inverse process
  - Electron-hole pair (EHP) generation rate = EHP recombination rate
- Thermal agitation → electrons and holes exchange energy with the crystal lattice and each other
  - → Every energy state in the conduction band and valence band has a certain probability of being occupied by an electron

## Statistical Thermodynamics: Fermi energy

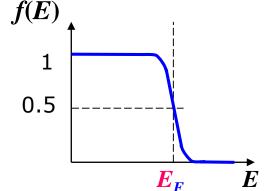
- The Fermi energy,  $E_F$ , is the energy associated with a particle, which is in thermal equilibrium with the system of interest. The energy is strictly associated with the particle and does not consist even in part of heat or work. This same quantity is called the electrochemical potential, m, in most thermodynamics texts.
  - http://hyperphysics.phyastr.gsu.edu/Hbase/solids/fermi.html#c2
  - http://hyperphysics.phyastr.gsu.edu/Hbase/solids/fermi.html#c1

## PL

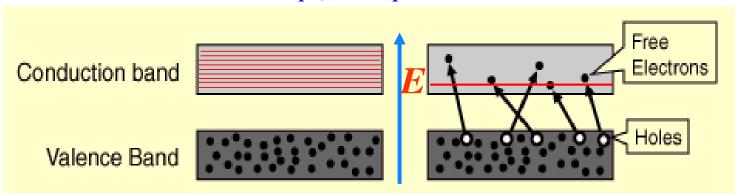
#### Fermi function and Fermi level

 Probability that a state at energy level, E, is occupied by one electron is,

$$f(E) = \frac{1}{1 + \exp(\frac{E - E_F}{kT})}$$



- f(E): Fermi-Dirac function
- An increase in E will reduce f(E)
- $E_F$  --- Fermi-level
  - > When  $E = E_F$ ,  $f(E = E_F) = 0.5$ .

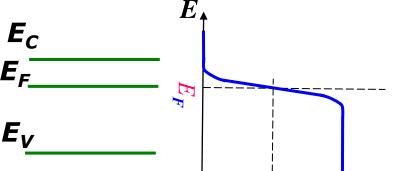


textbook P.66

#### Fermi function and Fermi level

 Probability that a state at energy level, E, is occupied by one electron is,

$$f(E) = \frac{1}{1 + \exp(\frac{E - E_F}{kT})}$$



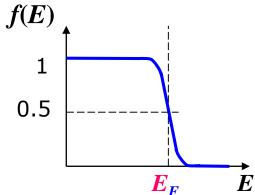
- f(E): Fermi-Dirac function
- An increase in E will reduce f(E)
- $E_F$  --- Fermi-level • When  $E = E_F$ ,  $f(E = E_F) = 0.5$ .
- 1. Need simplify the Fermi-Dirac function
- 2. What is the states' density?

## PL

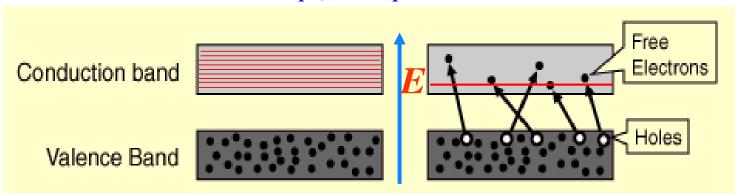
#### Fermi function and Fermi level

 Probability that a state at energy level, E, is occupied by one electron is,

$$f(E) = \frac{1}{1 + \exp(\frac{E - E_F}{kT})}$$



- f(E): Fermi-Dirac function
- An increase in E will reduce f(E)
- $E_F$  --- Fermi-level
  - When  $E = E_F$ ,  $f(E = E_F) = 0.5$ .



textbook P.66

## **2.5** Boltzmann approximation & $E_F$ , n, p

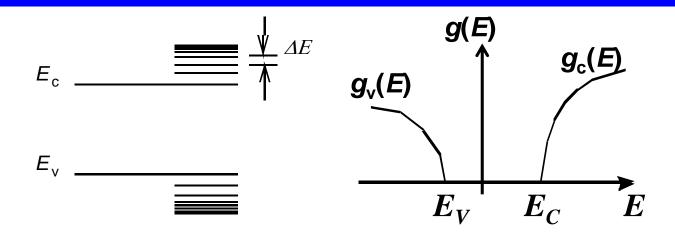
Fermi function and Fermi level

Density of States 态密度

Boltzmann Approximation

Electron and hole Concentrations

#### **Density of States**



 $g(E)\Delta E$  = number of states per cm<sup>3</sup> in the energy range between E and E+ $\Delta E$ 

Near the band edges:

$$g_c(E) = \frac{m_n^* \sqrt{2m_n^* (E - E_c)}}{\pi^2 \hbar^3}$$
  $E \ge E_c$ 

$$E \ge E_c$$

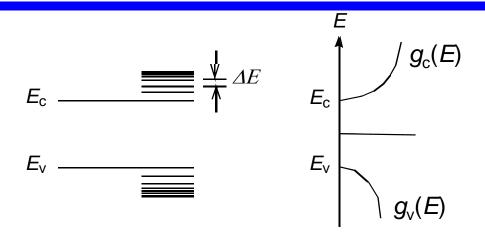
density of states in the conduction band

$$g_{v}(E) = \frac{m_{p}^{*} \sqrt{2m_{p}^{*}(E_{v} - E)}}{\pi^{2}\hbar^{3}} \qquad E \leq E_{v}$$

$$E \leq E_{v}$$

density of states in the valence band

#### **Density of States**



g(E)dE = number of states per cm<sup>3</sup> in the energy range between E and E+dE

Near the band edges:

$$g_c(E) = \frac{m_n^* \sqrt{2m_n^* (E - E_c)}}{\pi^2 \hbar^3} \quad E \ge E_c$$

$$g_{v}(E) = \frac{m_{p}^{*} \sqrt{2m_{p}^{*}(E_{v} - E)}}{\pi^{2}\hbar^{3}}$$
  $E \leq E_{v}$ 

## **2.5** Boltzmann approximation & $E_F$ , n, p

Fermi function and Fermi level

Density of States

Boltzmann Approximation

Electron and hole Concentrations

#### **Boltzmann Approximation**

$$f(E) = \frac{1}{1 + \exp(\frac{E - E_F}{kT})}$$

If 
$$E - E_F > 3kT$$
,  $f(E) \cong e^{-(E - E_F)/kT}$ 

 $E_{C}$ 

because of  $exp[(E-E_F)/(kT)] >> 1$ 

E.,

If 
$$E_F - E > 3kT$$
,  $f(E) \cong 1 - e^{(E - E_F)/kT}$ 

E<sub>C\_\_\_\_\_\_</sub>

E<sub>F\_\_\_\_\_</sub>

*E<sub>V</sub>*\_\_\_\_\_

Probability that a state is **empty**:

$$1-f(E) \cong e^{(E-E_F)/kT} = e^{-(E_F-E)/kT}$$

Probability that a state is occupied by a hole

## **2.5** Boltzmann approximation & $E_F$ , n, p

Fermi function and Fermi level

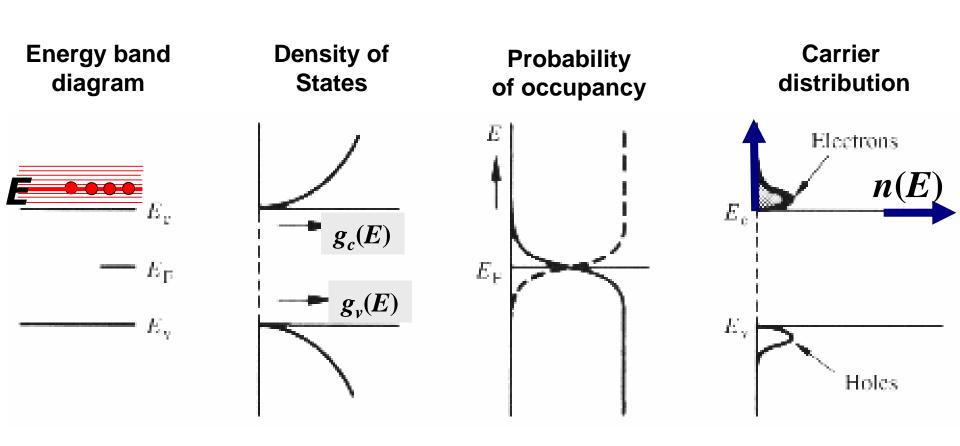
Density of States

Boltzmann Approximation

Electron and hole Concentrations

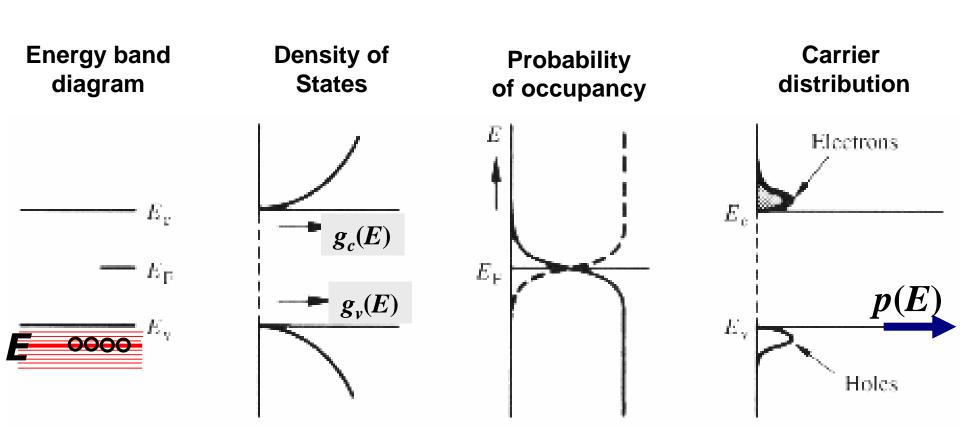
#### **Equilibrium Distribution of Electrons**

• Obtain n(E) by multiplying  $g_c(E)$  and f(E)



#### **Equilibrium Distribution of Holes**

• Obtain p(E) by multiplying  $g_{V}(E)$  and 1-f(E)



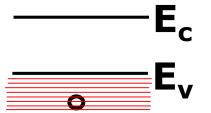
#### **Equilibrium Hole Concentrations**

 Integrate p(E) over all the energies in the valence band to obtain p:

$$p = \int_{-\infty}^{E_v} g_v(E) [1 - f(E)] dE$$

 By using the Boltzmann approximation, and extending the integration limit to -∞, we obtain

$$p = N_{v}e^{-(E_{F}-E_{v})/kT}$$
 where  $N_{v} = 2\left(\frac{2\pi m_{p}^{*}kT}{h^{2}}\right)^{3/2}$ 



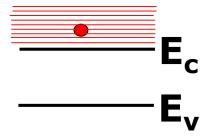
#### **Equilibrium Electron Concentrations**

 Integrate n(E) over all the energies in the conduction band to obtain n:

top of conduction band
$$n = \int_{E_c}^{\text{top of conduction band}} g_c(E) f(E) dE$$

 By using the Boltzmann approximation, and extending the integration limit to ∞, we obtain

$$n = N_c e^{-(E_c - E_F)/kT}$$
 where  $N_c = 2 \left( \frac{2\pi m_n^* kT}{h^2} \right)^{3/2}$ 



#### **Intrinsic Carrier Concentration**

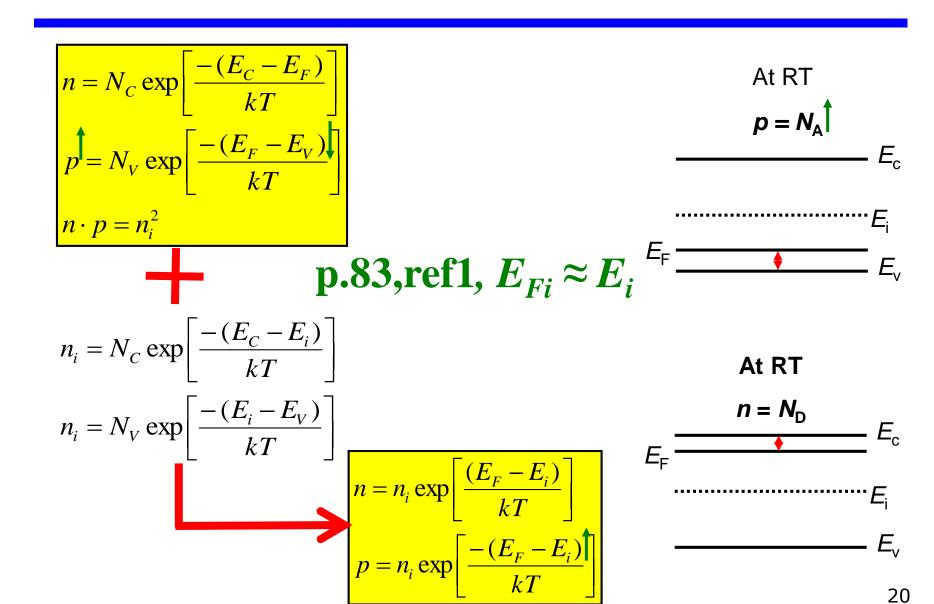
$$np = \left(N_c e^{-(E_c - E_F)/kT}\right) \left(N_v e^{-(E_F - E_v)/kT}\right)$$

$$= N_c N_v e^{-(E_c - E_v)/kT} = N_c N_v e^{-E_g/kT}$$

$$= n_i^2 \quad \text{Law of Mass Action}$$

$$n_i = \sqrt{N_c N_v} e^{-E_g/2kT}$$

#### **Electron and hole concentrations**



# HW3: Energy-band diagram

Question: Where is  $E_F$  for  $n = 10^{17}$  cm<sup>-3</sup>?

#### 2.6 Carrier drift and diffusion

Carrier scattering

载流子散射

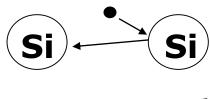
- Carrier drift:
  - Carrier mobility
  - Conductivity &Resistivity
  - Energy band model

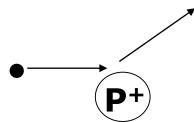
Carrier diffusion

## **Thermal Motion**

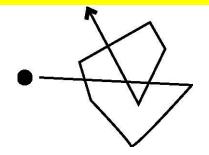
#### 载流子

- In thermal equilibrium, carriers are not sitting still:
  - undergo collisions with vibrating Si atoms (Brownian motion)
  - electrostatically interact with charged dopants and with each other
- Characteristic time constant of thermal motion
  - ho mean free time between collisions:  $\tau_c \equiv$  collision time [s]
  - In between collisions, carriers acquire high velocity: v<sub>th</sub> ≡ thermal velocity [cm/s]
  - ...but get nowhere! (on avearge)
- Characteristic length of thermal motion:
  - $\lambda \equiv$ mean free path [cm],  $\lambda = v_{th} \tau_{c}$





#### 平均自由时间



平均自由程

# Carrier Scattering

- random motion<sup>4</sup>

  electron

  5
- Mobile electrons and atoms in the Si lattice are always in random thermal motion.
  - Average velocity of thermal motion for electrons in Si: ~10<sup>7</sup> cm/s @ 300K
  - Electrons make frequent "collisions" with the vibrating atoms
     晶格散射或声子散射
    - "lattice scattering" or "phonon scattering"
  - Other scattering mechanisms:
    - deflection by ionized impurity atoms
    - deflection due to <u>Coulombic</u> force between carriers
- The average current in any direction is zero, if no electric field is applied.

#### **Effective Mass**

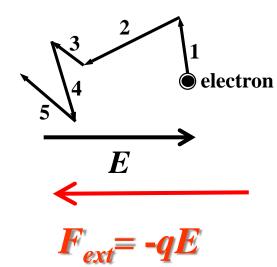
- Under an externally applied force, F<sub>ext</sub>, the movement of electrons (or holes) is influenced by the positively charged protons and by negatively charged electron in the lattice. So, the movement in the crystal is different from that in vacuum.
- The total force F<sub>total</sub>

$$F_{total} = F_{ext} + F_{int} = ma$$

where a is the acceleration,  $F_{\text{int}}$  is the internal force. We can write

$$F_{ext} = m^* a$$

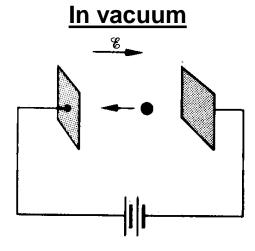
where  $m^*$  is called **effective mass**.



**Notation:**  $m_n^*$  for electrons,  $m_p^*$  for holes,

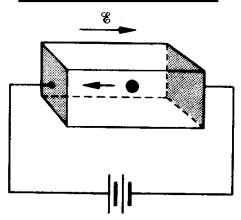
有效质量

# **Electrons as Moving Particles**



$$F = (-q)E = m_0 a$$

#### In semiconductor



$$F = (-q)E = m_0 a$$
  $F_{\text{ext}} = (-q)E = m_n *a$ 

where  $m_n^*$  is the electron effective mass.

If  $\tau_{cn}$  is electron mean free time between collisions,

$$|a| = \frac{dv}{dt} \approx \frac{v_e}{\tau_{cn}}$$

$$\Rightarrow v_e = \frac{q\tau_{cn}E}{m_n^*}, \quad v_h = \frac{q\tau_{cp}E}{m_p^*}$$

漂移速度: average drift velocity 26

## 2.6 Carrier drift and diffusion

Carrier scattering

载流子漂移

- Carrier drift:
  - Carrier mobility

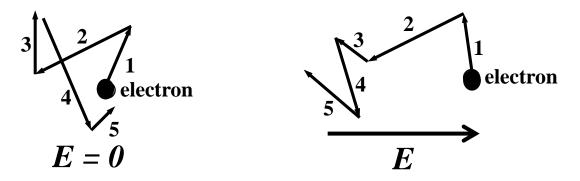
载流子迁移率

- Conductivity & Resistivity
- Energy band model

Carrier diffusion

## **Carrier Drift**

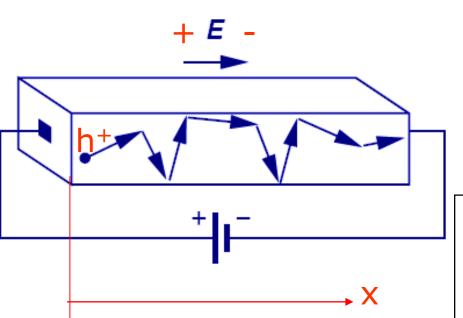
When an electric field (e.g., due to an externally applied voltage) is applied to a semiconductor, mobile charge-carriers will be accelerated by the electrostatic force. This force superimposes on the random motion of electrons:



- Electrons drift in the direction opposite to the E-field
   → Current flows
- ❖ Because of scattering, electrons in a semiconductor do not achieve constant acceleration. However, they can be viewed as classical particles moving at a constant average drift velocity.

## **Carrier Drift**

- The process in which charged particles move because of an electric field is called *drift*.
- Charged particles within a semiconductor move with an average velocity proportional to the electric field.
  - > The proportionality constant is the carrier *mobility*.



迁移率

Hole velocity  $\vec{v_h} = \mu_{_D} \, \vec{E}$ 

Electron velocity  $\stackrel{\rightarrow}{v_e} = -\mu_{\scriptscriptstyle n}\stackrel{\rightarrow}{E}$ 

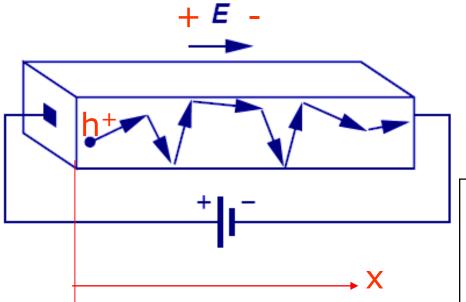
#### **Notation:**

 $\mu_{\rm p} \equiv \text{hole mobility (cm}^2/\text{V}\cdot\text{s})$ 

 $\mu_n \equiv \text{electron mobility (cm}^2/\text{V}\cdot\text{s})$ 

## **Carrier Drift**

$$v_e = \frac{q\tau_{cn}E}{m_n^*}, \quad v_h = \frac{q\tau_{cp}E}{m_p^*} \quad \Rightarrow \quad \mu_n = \frac{q\tau_{cn}}{m_n^*}, \quad \mu_p = \frac{q\tau_{cp}}{m_p^*}$$



Hole velocity 
$$\stackrel{
ightarrow}{v_{_h}}=\stackrel{
ightarrow}{\mu_{_p}}\stackrel{
ightarrow}{E}$$

Electron velocity  $\stackrel{\rightarrow}{v_e} = -\mu_{\scriptscriptstyle n}\stackrel{\rightarrow}{E}$ 

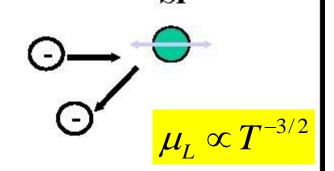
#### **Notation:**

 $\mu_{\rm p}$  = hole mobility (cm<sup>2</sup>/V·s)

 $\mu_n \equiv \text{electron mobility (cm}^2/\text{V}\cdot\text{s})$ 

## Carrier Mobility $1/\mu = 1/\mu_L + 1/\mu_I$

- Mobile carriers are always in random thermal motion. If no electric field is applied, the average current in any direction is zero.
  - Mobility is reduced by
  - 1) collisions with the vibrating atoms "phonon scattering"



2) deflection by ionized impurity atoms "Coulombic



 $\mu_I \propto T^{+3/2} / N_I$ 

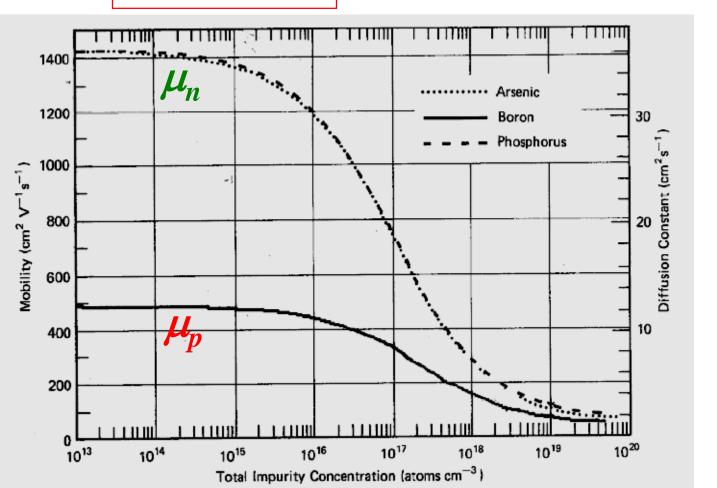


## **Drift Velocity and Carrier Mobility**

Mobile charge-carrier drift velocity is proportional to applied *E*-field:

$$/v/=\mu E$$

 $\mu$  is the **mobility** (Units: cm<sup>2</sup>/V·s)

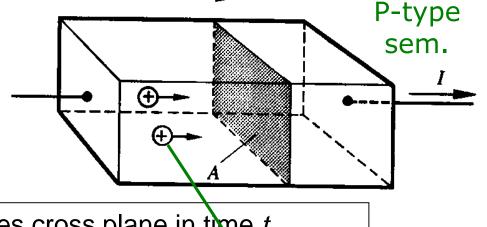


Note: Carrier mobility depends on *total* dopant concentration  $(N_D + N_A)$ !

#### **Drift Current**

 Drift current is proportional to the carrier velocity and carrier concentration:

- 1) *p*---hole density
- 2)  $q = 1.6 \times 10^{-19} C$
- --- One electron charge
- 3) Charges passing through 'A' per second
- --- The definition of current.



 $v_h t A = volume from which all holes cross plane in time t$ 

 $p v_h tA = \#$  of holes crossing plane in time t

 $q p v_h t A =$ charge crossing plane in time t

 $q p v_h A = \text{charge crossing plane per unit time} = \text{hole current}$ 

 $\rightarrow$  Hole current per unit area (i.e. current density)  $J_{p,drift} = q p v_h$ 

# Electrical Conductivity $\sigma$

Negatively charged electron
Direction of electron drift

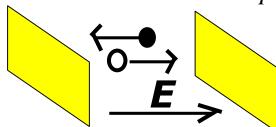
When an electric field is applied, current flows due to drift of mobile electrons and holes:

electron current density:

$$J_n = (-q)nv_e = qn\mu_n E$$

hole current density:

$$\boldsymbol{J}_p = (+q)p\boldsymbol{v}_h = qp\boldsymbol{\mu}_p\boldsymbol{E}$$



total current density:

$$J = J_n + J_p = (qn\mu_n + qp\mu_p)E$$

电导率

$$J = \sigma E$$

conductivity

$$\sigma \equiv qn\mu_n + qp\mu_p$$

Units:  $(\Omega \cdot cm)^{-1}$ 

# Electrical Resistivity $\rho$

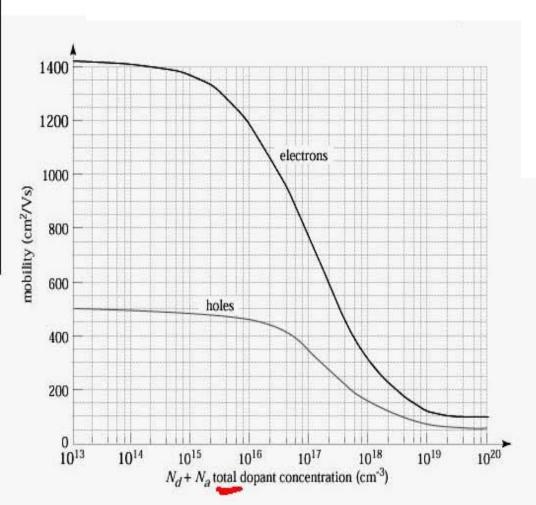
$$\rho \equiv \frac{1}{\sigma} = \frac{1}{qn\mu_n + qp\mu_p}$$

$$\rho \cong \frac{1}{qn\mu_n}$$
 for n-type mat'l

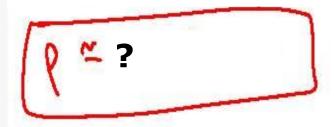
$$\rho \cong \frac{1}{qp\mu_p}$$
 for p-type mat'l

(Units: ohm•cm)

#### HW4



Estimate the resistivity of a Si sample doped with phosphorus to a concentration of 10<sup>15</sup> cm<sup>-3</sup> and boron to a concentration of 10<sup>17</sup> cm<sup>-3</sup>.



The electron mobility and hole mobility are 700 cm<sup>2</sup>/Vs and 350 cm<sup>2</sup>/Vs, respectively.

# **Example**

Consider a Si sample doped What is its resistivity?

#### Answer:

$$N_A = 10^{16} / \text{cm}^3$$
,  $N_D = 0$   $(N_A >> N_D \rightarrow \text{p-type})$ 

Total Impurity Concentration (atoms cm<sup>-3</sup>)

 $\rightarrow p \approx 10^{16} \text{/cm}^3$  and  $n \approx 10^4 \text{/cm}^3$ 

$$\rho = \frac{1}{qn\mu_n + qp\mu_p} \cong \frac{1}{qp\mu_p}$$
$$= \left[ (1.6 \times 10^{-19})(10^{16})(450) \right]^{-1} = 1.4 \,\Omega \cdot \text{cm}$$

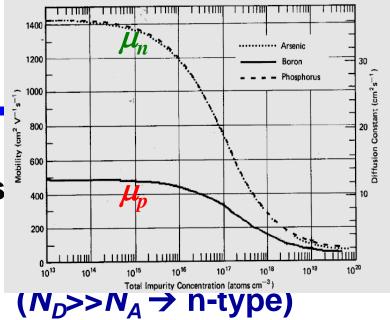
From  $\mu$  vs. ( $N_A + N_D$ ) plot

# Example (cont'd)

Consider the same Si sample, with 10<sup>17</sup>/cm<sup>3</sup> Arsenic. What is

#### <u>Answer:</u>

$$N_A = 10^{16} \text{/cm}^3$$
,  $N_D = 10^{17} \text{/cm}^3$  ( $N_D > N_A \rightarrow \text{n-type}$ )

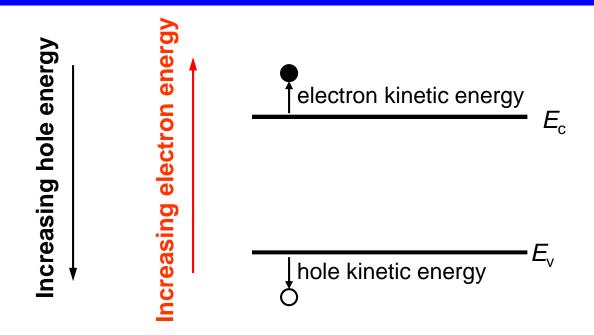


 $\rightarrow n \approx 9 \times 10^{16} / \text{cm}^3$  and  $p \approx 1.1 \times 10^3 / \text{cm}^3$ 

$$\rho = \frac{1}{qn\mu_n + qp\mu_p} \cong \frac{1}{qn\mu_n}$$
$$= \left[ (1.6 \times 10^{-19})(9 \times 10^{16})(700) \right]^{-1} = 0.10 \,\Omega \cdot \text{cm}$$

The sample is converted to n-type material by adding more donors than acceptors, and is said to be "compensated".

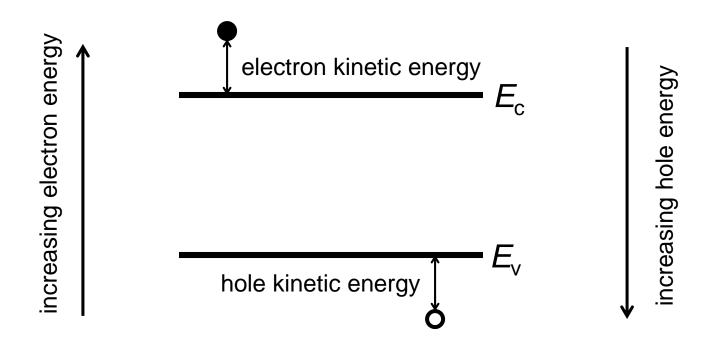
### **Electrons and Holes (Band Model)**



- Electrons and holes tend to seek lowestenergy positions
  - Electrons tend to fall
  - Holes tend to float up (like bubbles in water)

### Potential vs. Kinetic Energy

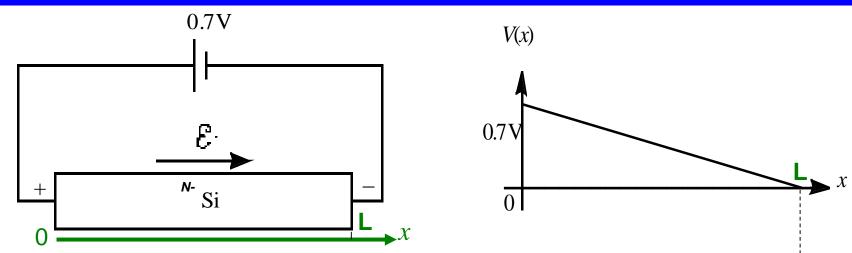
#### 势能和动能



#### $E_{\rm c}$ represents the electron potential energy:

$$P.E. = E_c - E_{reference}$$
 电子的势能

## Electrostatic Potential, V 电势或电位

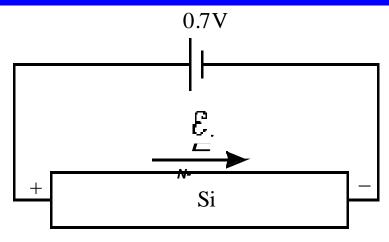


• The potential energy of a particle with charge -q is related to the electrostatic potential V(x):

$$P.E. = -qV$$
 
$$E_{\text{reference}} \text{ is } V = \frac{1}{q}(E_{\text{reference}} - E_{\text{c}})$$
 constant

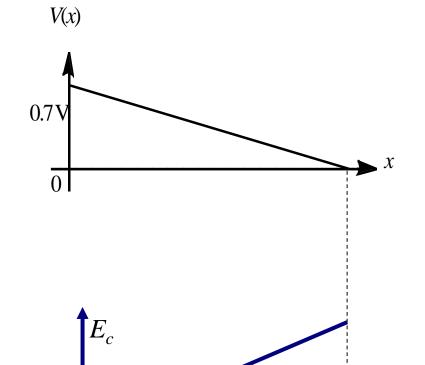
### Electric Field, $\varepsilon$

$$\mathcal{E} = -\frac{dV}{dx}$$



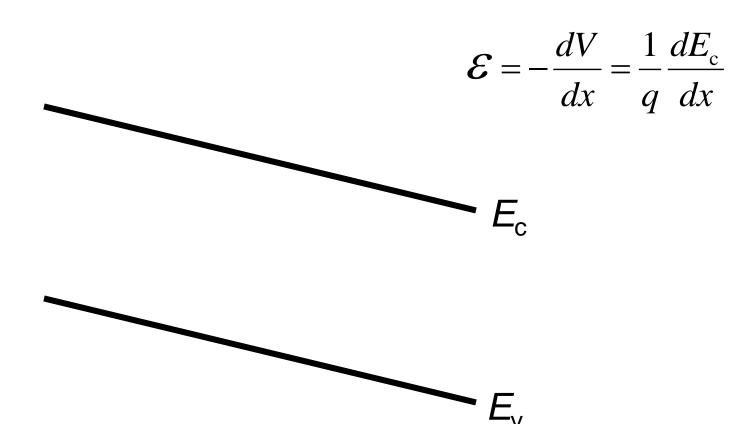
$$V = \frac{1}{q} (E_{\text{reference}} - E_{\text{c}})$$

$$\mathcal{E} = -\frac{dV}{dx} = \frac{1}{q} \frac{dE_{c}}{dx}$$



Variation of E<sub>c</sub> with position is called "band bending."

### **HW 5: Carrier Drift (Band Diagram Visualization)**



Q1: what is the direction of electric field?

Q2: what is the direction of carriers' drift?

## 2.6 Carrier drift and diffusion

Carrier scattering

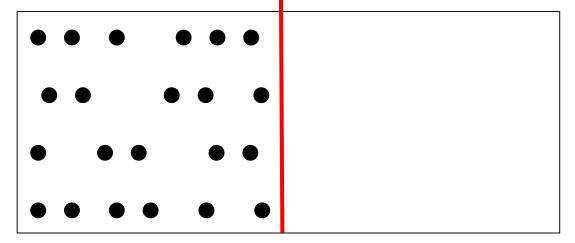
- Carrier drift:
  - Carrier mobility
  - Conductivity & Resistivity
  - Energy band model

载流子扩散

Carrier diffusion

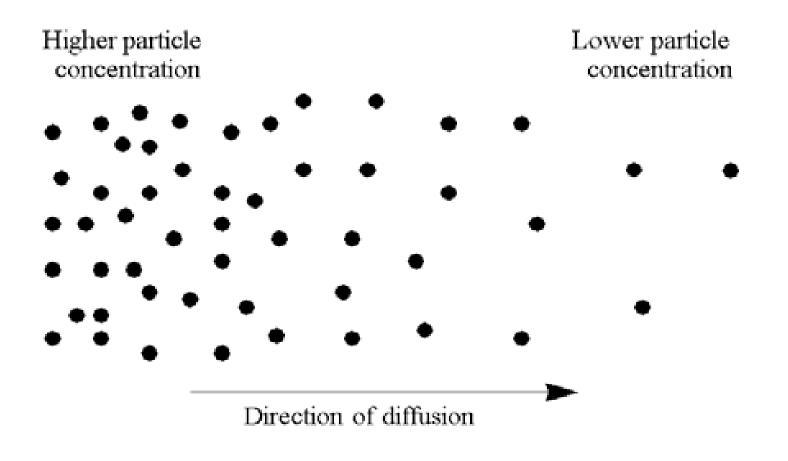
## **Diffusion**

- Diffusion occurs when there exists a concentration gradient
- In the figure below, imagine that we fill the left chamber with a gas at temperate T
- If we suddenly remove the divider, what happens?
- The gas will fill the entire volume of the new chamber.
- How does this occur?



## **Diffusion**

 Particles diffuse from higher concentration to lower concentration locations.



## **Carrier Diffusion**

- Due to thermally induced random motion, mobile particles tend to move from a region of high concentration to a region of low concentration.
  - Analogy: ink droplet in water

浓度梯度

- Current flow due to mobile charge diffusion is proportional to the carrier concentration gradient.
  - The proportionality constant is the *diffusion* constant.

Semiconductor Material

Injection of Carriers

$$J_p = -qD_p \frac{dp}{dx}$$

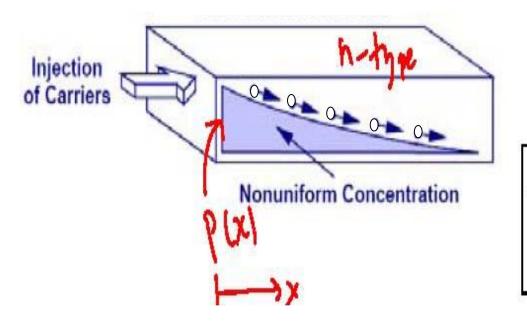
**Nonuniform Concentration** 

#### **Notation:**

 $D_p \equiv \text{hole diffusion constant (cm}^2/\text{s})$ 

 $D_n \equiv \text{electron diffusion constant (cm}^2/\text{s})$ 

$$J_p = -qD_p \frac{dp}{dx}$$



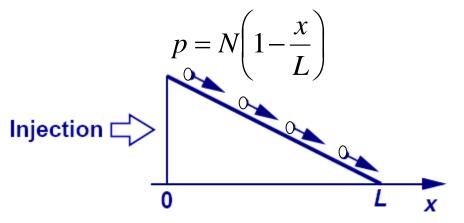
#### Notation:

 $D_p = \text{hole diffusion constant (cm}^2/\text{s})$ 

 $D_n = \text{electron diffusion constant (cm}^2/\text{s})$ 

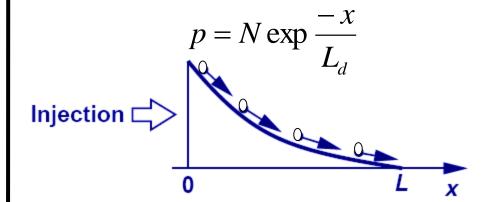
# **Diffusion Examples**

- Linear concentration profile
  - → constant diffusion current



$$J_{p,diff} = -qD_{p} \frac{d_{p}}{d_{p}}$$
$$= qD_{p} \frac{N}{L}$$

Non-linear concentration profile
 → varying diffusion current



$$J_{p,diff} = -qD_{p} \frac{dp}{dx}$$

$$= \frac{qD_{p}N}{L_{d}} \exp \frac{-x}{L_{d}}$$

## Total **Diffusion** Current

Due to the non-uniform distribution of carriers

$$J_n = qD_n \frac{dn}{dx}$$

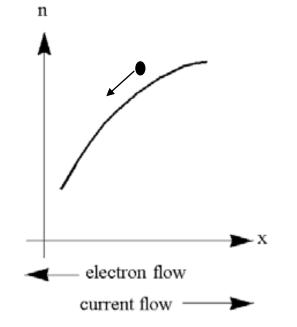
- $D_n$  --- e Diffusion constant.
- Driving force: thermal energy, not electric field
- dn/dx--- density gradient
- Total diffusion current

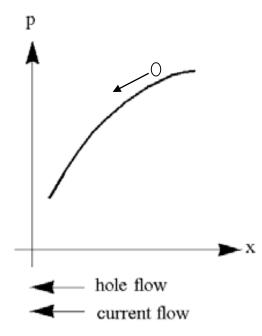
$$\rightarrow$$
  $J = J_n + J_p$ 

## Total **Diffusion** Current

Diffusion current within a semiconductor consists of hole and electron components:

$$egin{align} J_{p, diff} &= -qD_p \, rac{dp}{dx} & J_{n, diff} &= qD_n \, rac{dn}{dx} \ J_{tot, diff} &= q(D_n \, rac{dn}{dx} - D_p \, rac{dp}{dx}) \ \end{array}$$





## The total current

 The total current flowing in a semiconductor is the sum of drift current and diffusion current:

$$\left| \boldsymbol{J}_{tot} = \boldsymbol{J}_{p,dri\!f\!t} + \boldsymbol{J}_{n,dri\!f\!t} + \boldsymbol{J}_{p,di\!f\!f} + \boldsymbol{J}_{n,di\!f\!f} 
ight|$$

$$J_{p,drift} = qp\mu_{p}E, \qquad J_{n,drift} = qn\mu_{n}E$$
 
$$J_{p,diff} = -qD_{p}\frac{dp}{dx}, \quad J_{n,diff} = qD_{n}\frac{dn}{dx}$$

## The Einstein Relation

 The characteristic constants for drift and diffusion are related:

$$\frac{D}{\mu} = \frac{kT}{q} = 26 \text{ mV}$$
at  $T = 300 \text{ K}$ 

- Note that  $\frac{kT}{q} \cong 26 \mathrm{mV}$  at room temperature (300K)
  - > This is often referred to as the "thermal voltage".

### **Important Constants**

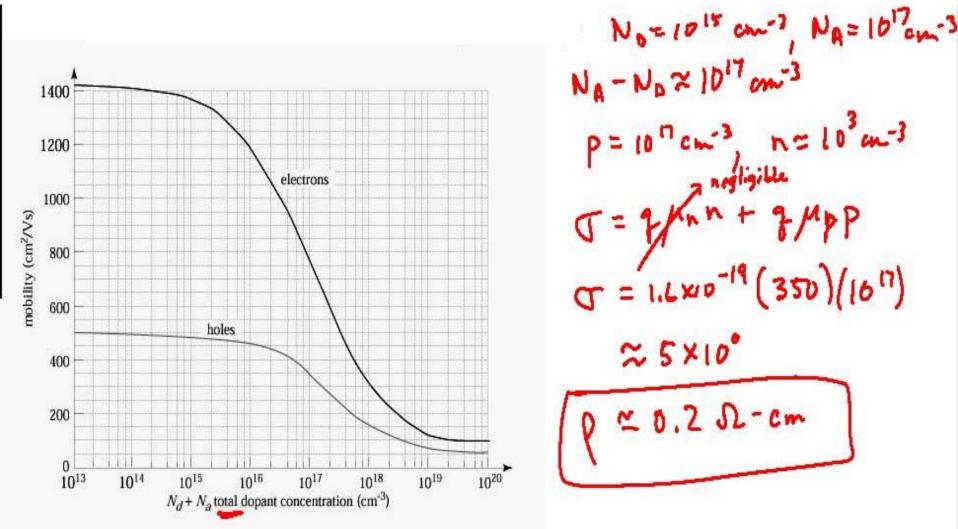
- Electronic charge,  $q = 1.6 \times 10^{-19}$  C
- Permittivity of free space,  $\varepsilon_0 = 8.854 \times 10^{-14}$  F/cm
- Boltzmann constant,  $k = 8.62 \times 10^{-5}$  eV/K
- Planck constant,  $h = 4.14 \times 10^{-15} \text{ eV} \cdot \text{s}$
- Free electron mass,  $m_0 = 9.1 \times 10^{-31}$  kg
- Thermal voltage kT/q = 26 mV, at T=300K

# HW3: Energy-band diagram

Question: Where is  $E_F$  for  $n = 10^{17}$  cm<sup>-3</sup>?

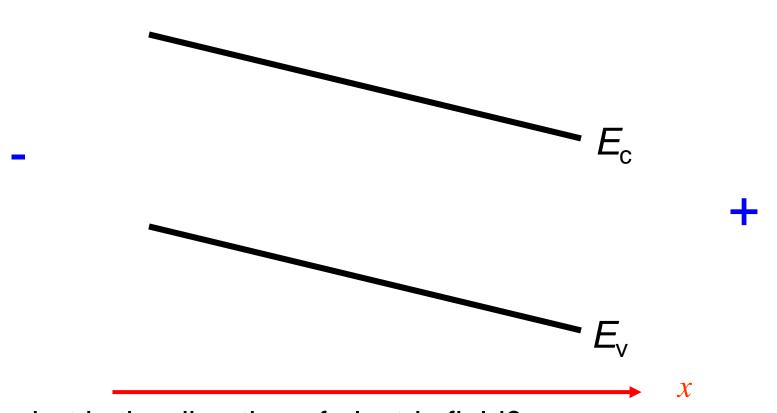
$$n = n_i \exp\left[\frac{(E_F - E_i)}{kT}\right]$$

#### HW4



The electron mobility and hole mobility are 700 cm<sup>2</sup>/Vs and 350 cm<sup>2</sup>/Vs, respectively.

#### **HW5: Carrier Drift (Band Diagram Visualization)**



Q1: what is the direction of electric field?

Q2: what is the direction of carriers' drift?