

# **EEE225    Advanced Electrical Circuits and Electromagnetics**

## **Lecture 5 Electric and Magnetic Components (Resistor, Capacitors and Inductors)**

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Room EE322

# Content

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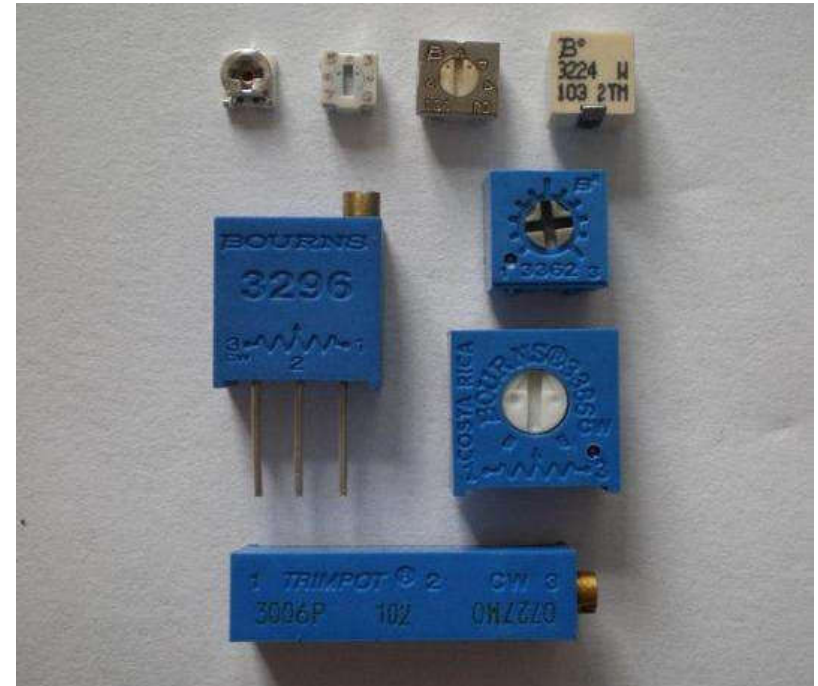
- 1. Resistor
- 2. Capacitor
  - Calculation
  - Dielectric material filled
  - Energy storage
  - I-V relationship
- 3. Inductor
  - Similar to capacitors
- 4. Case study
  - Parallel plates
  - Coaxial cable

# 1. Resistor

- A resistor is a passive two-terminal electrical component that implements electrical resistance as a circuit element.



Fixed resistors



Variable resistors

# 1. Resistance

## Ohm's Law

- The resistance of a conductor of length  $dl$  can be obtained by

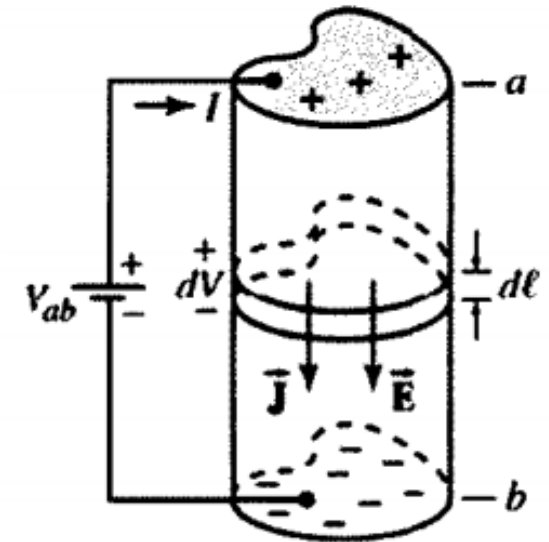
$$dR = \frac{dV}{I} = \frac{-\vec{E} \cdot d\vec{l}}{\iint_S \vec{J} \cdot d\vec{s}}$$

- If we assume that the potential at end  $a$  of the conductor is higher than that at end  $b$ .
- The total resistance of the conductor is:

$$R = \int_b^a \frac{-\vec{E} \cdot d\vec{l}}{\iint_S \vec{J} \cdot d\vec{s}}$$

- This is a general equation to determine the resistance of a conducting medium whose conductivity changes in the direction of the current. For homogeneous medium having constant  $\sigma$ , it reduces to:

$$R = \int_b^a \frac{-\vec{E} \cdot d\vec{l}}{\iint_S \vec{J} \cdot d\vec{s}} = \frac{-\int_b^a \vec{E} \cdot d\vec{l}}{\iint_S \vec{J} \cdot d\vec{s}} = \frac{V_{ab}}{I}$$



# 1. Resistance

How to calculate?

- Simplified model: A potential difference of  $V_0$  is maintained across the two ends of a conducting wire of length  $l$ . If  $A$  is the cross-sectional area of the wire, obtain an expression for the resistance of the wire.

- Assume the potential difference between the two ends of the conductor is  $V_0$ , the electric field holds:

$$V_0 = - \int_b^a \vec{E} \cdot d\vec{l} = El \Rightarrow E = \frac{V_0}{l}$$

- If  $\sigma$  is the conductivity of the conducting material, the current density at any cross section of the wire is:

$$J = \sigma E = \frac{\sigma V_0}{l}$$

- The current through the wire is:

$$I = \iint_S \vec{J} \cdot d\vec{s} = JA = \frac{\sigma V_0 A}{l} = \frac{V_0}{R}$$

- So the resistance of the piece of the conducting material is

$$R = \frac{V_0}{I} = \frac{l}{\sigma A}$$

# 1. Resistance

## Example 1

- A long, round wire of radius  $a$  and conductivity  $\sigma$  is coated with a material of conductivity  $0.1\sigma$ .
  - a) What must be the thickness of the coating so that the resistance per unit length of the uncoated wire is reduced by 50%?
  - b) Assuming a total current  $I$  in the coated wire, find  $\mathbf{J}$  and  $\mathbf{E}$  in both the core and the coating material.

## 2. Capacitor

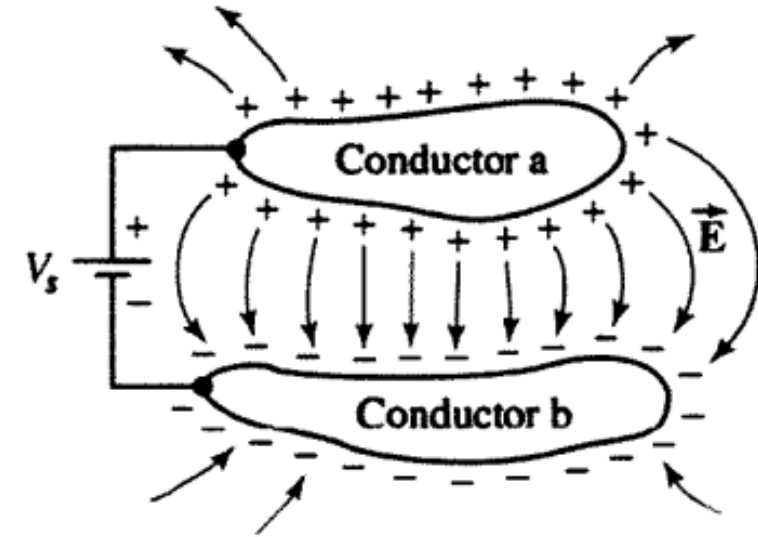
- A capacitor is also a two-terminal passive device which stores electric charge.



## 2. Capacitor

- Capacitor

- A capacitor is a device which stores electric charge.
- Its basic configuration is two conductors carrying equal but opposite charges

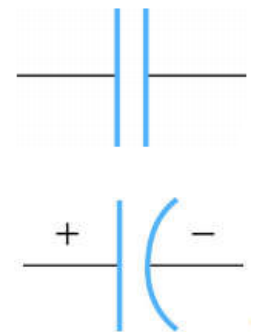


- Capacitance

- measures the capability of energy storage in electrical devices.
- the amount of charge  $Q$  stored in a capacitor is linearly proportional to the electric potential difference  $V$  between the two conductors:

$$\frac{Q}{V} = \text{constant} = C$$

- Unit: 1 F (farad) = 1 C / V (coulomb/volt)

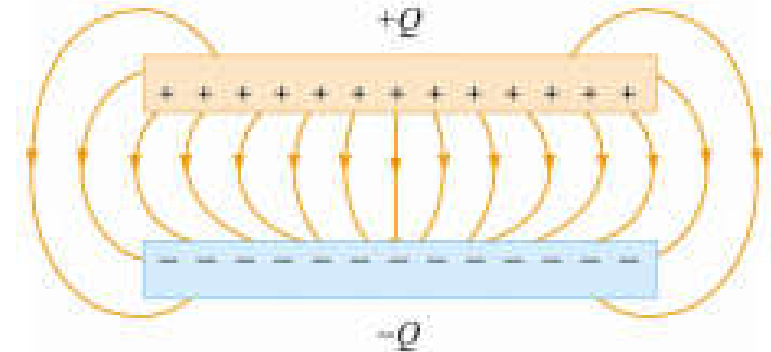
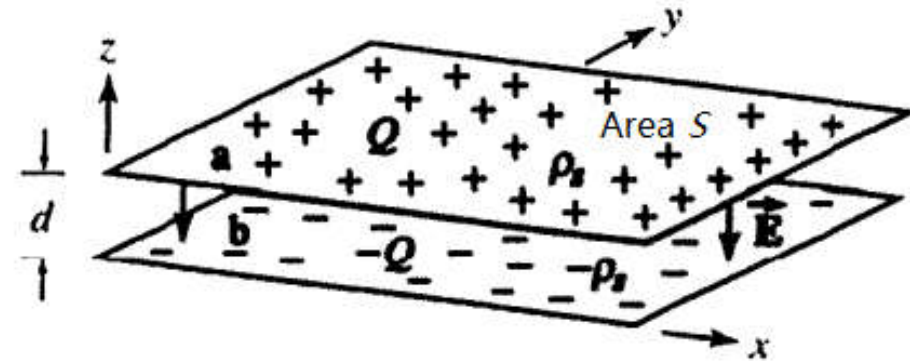




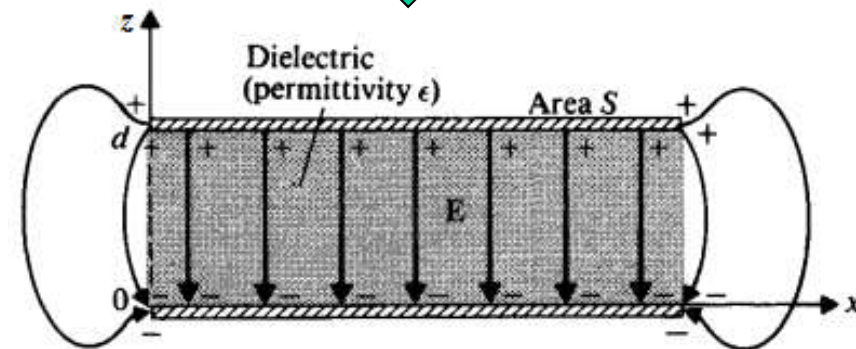
# 2.1 Capacitor Examples

## Parallel Plates

- Two parallel conducting plates, each of area  $S$ , and separated by a distance  $d$ , form a parallel-plate capacitor. The total charge on the top plate is  $+Q$  and that on the other plate is  $-Q$ .
  - What is its capacitance?
- Solution:
  - Edge effects:** The electric field lines at the edge of the plates are not straight lines, and the field is not contained entirely between the plates.
  - Fringing fields:** The non-uniform fields near the edge.



$d \ll \sqrt{S}$



# 2.1 Capacitor Examples

## Parallel Plates

- Solution:

- The surface charge density is:

$$\rho_s = Q/S$$

- Based on Gauss's Law:

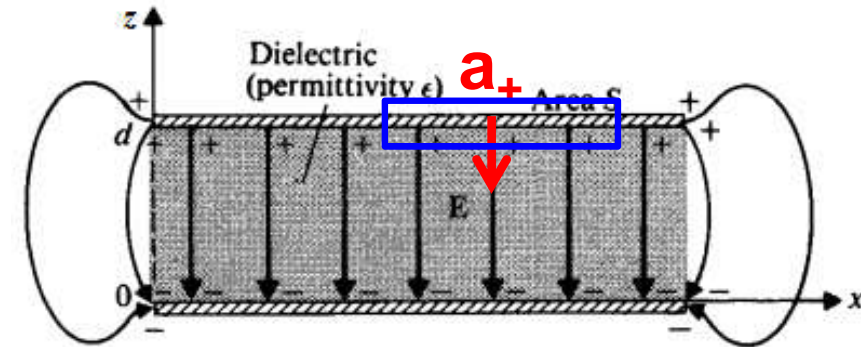
$$\oiint_{S'} \vec{E} \cdot d\vec{S} = \frac{q_{enc}}{\epsilon_0} \xrightarrow{S' \text{ is a unit area}} \vec{E} = -\hat{a}_z \frac{\rho_s}{\epsilon_0} = -\hat{a}_z \frac{Q}{\epsilon_0 S}$$

- The potential  $V$  is:

$$V = - \int_{z=0}^{z=d} \vec{E} \cdot d\vec{l} = - \int_0^d \left( -\hat{a}_z \frac{Q}{\epsilon_0 S} \right) \cdot (\hat{a}_z dz) = \frac{Q}{\epsilon_0 S} d$$

- Therefore, the capacitance of a parallel – plate is:

$$C = \frac{Q}{V} = \frac{\epsilon_0 S}{d}$$



# 2.1 Capacitor Examples

## Spherical capacitor

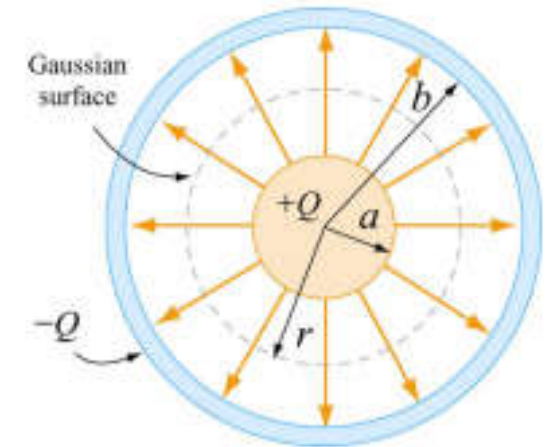
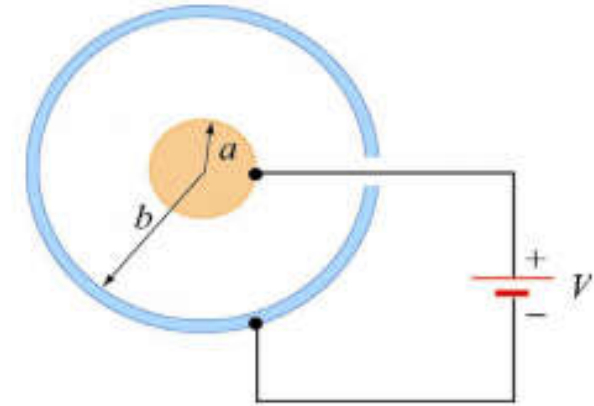
- A spherical capacitor:
  - two concentric spherical shells of radii  $a$  and  $b$
  - The inner shell has a charge  $+Q$  uniformly distributed over its surface, and the outer shell an equal but opposite charge  $-Q$ .
- Solution:
  - The electric field is non-vanishing only in the region  $a < r < b$ . Using Gauss's law, we obtain

$$E_r = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

- Therefore, the potential difference between the two conducting shells is

$$\Delta V = V_b - V_a = -\int_a^b E_r dr = -\frac{Q}{4\pi\epsilon_0} \int_a^b \frac{dr}{r^2} = -\frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{b} \right) = -\frac{Q}{4\pi\epsilon_0} \left( \frac{b-a}{ab} \right)$$

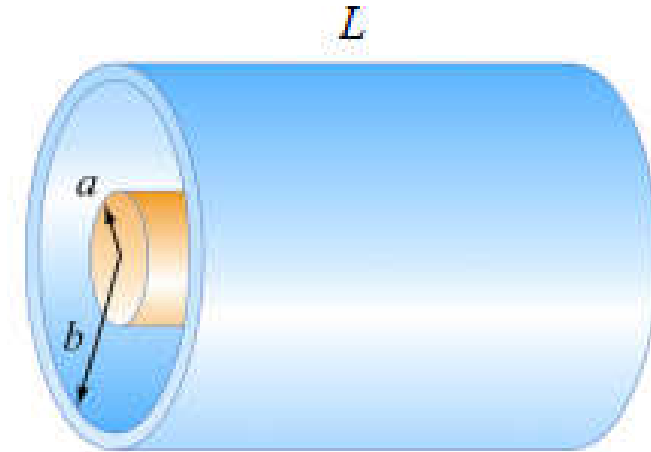
$$\longrightarrow C = \frac{Q}{|\Delta V|} = 4\pi\epsilon_0 \left( \frac{ab}{b-a} \right)$$



## 2.1 Capacitor Examples

### Cylindrical capacitor

- A cylindrical conductor
  - Inner radius  $a$  surrounded by a coaxial cylindrical shell of inner radius  $b$ . Filled with dielectrics with  $\epsilon$ . The length of both cylinders is  $L$ .
  - The capacitor is charged so that the inner cylinder has charge  $+Q$  while the outer shell has a charge  $-Q$ .

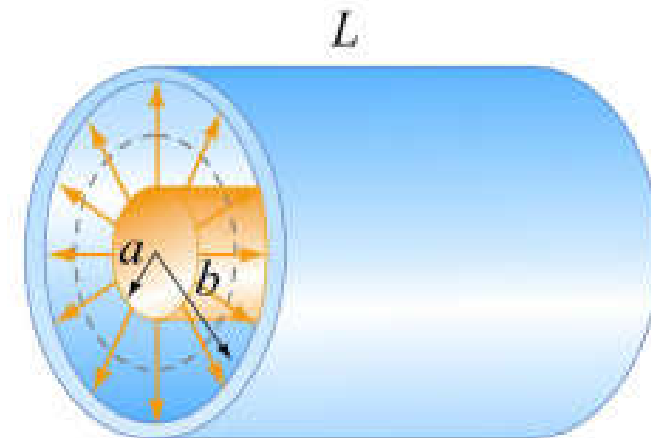


- Solution:
  - Assume that  $L$  is much larger than  $b-a$ , the separation of the cylinders, so that edge effects can be neglected.
  - $E$  can be calculated by:

$$\mathbf{E} = \mathbf{a}_r E_r = \mathbf{a}_r \frac{Q}{2\pi\epsilon L r}.$$

- So potential  $V$  is:

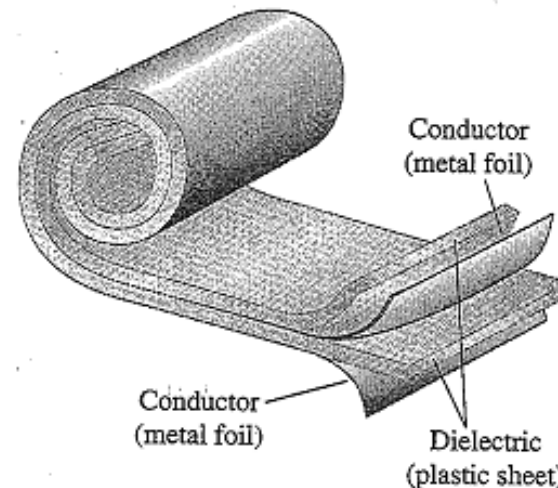
$$V_{ab} = -\int_{r=b}^{r=a} \mathbf{E} \cdot d\boldsymbol{\ell} = -\int_b^a \left( \mathbf{a}_r \frac{Q}{2\pi\epsilon L r} \right) \cdot (\mathbf{a}_r dr) \quad \longrightarrow \quad C = \frac{Q}{V_{ab}} = \frac{2\pi\epsilon L}{\ln\left(\frac{b}{a}\right)}.$$
$$= \frac{Q}{2\pi\epsilon L} \ln\left(\frac{b}{a}\right).$$



## 2.2 Capacitor with dielectrics

- Most capacitors have an insulating material, such as paper or plastic, between their conducting plates.
- Reasons:
  - To maintain a physical separation of the plates;
  - Increase the maximum possible potential difference between the conducting plates;
  - Capacitance increases when the space between the conductors is filled with dielectrics.

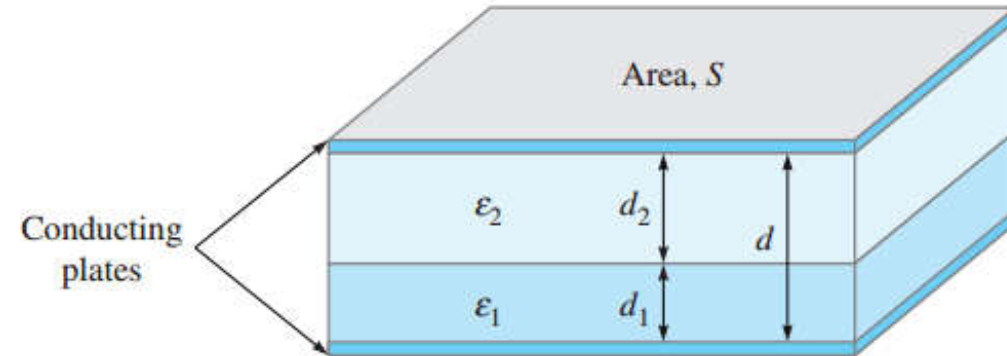
$$C = \frac{Q}{V} = \frac{\epsilon S}{d}$$



## 2.2 Capacitor with dielectrics

### Example 1

- A parallel-plate capacitor containing two dielectrics with the dielectric interface parallel to the plates.
- What is its capacitance?
- Solution 1:
  - It can be considered as two serially connected parallel-plate capacitors.
  - So the total capacitance is:  $C = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2}}$
  - where  $C_1 = \epsilon_1 S / d_1$        $C_2 = \epsilon_2 S / d_2$



This is the correct result, but let's try to obtain it using less intuition and a more basic approach (from the definition).



## 2.2 Capacitor with dielectrics

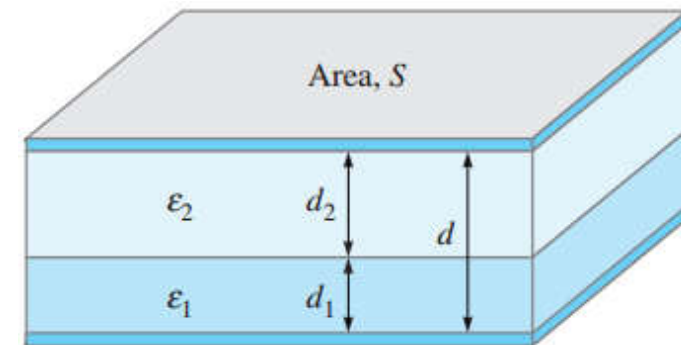
Example 1 cont.

- Solution 2:
  - Suppose we assume a potential difference  $V_0$  between the plates. The electric field intensities in the two regions,  $E_2$  and  $E_1$ , are both uniform, and  $V = E_1 d_1 + E_2 d_2$
  - At the dielectric interface,  $\mathbf{E}$  is normal to the interface, and our boundary condition tells us that  $D_1 = D_2$ , or  $\epsilon_1 E_1 = \epsilon_2 E_2$
  - The surface charge density

$$\rho_{s1} = D_1 = \epsilon_1 E_1 = \epsilon_2 E_2 = \rho_{s2} = \rho_s$$

- So we have:

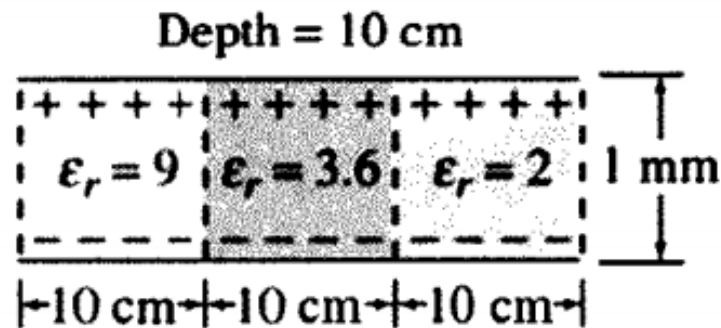
$$C = \frac{Q}{V} = \frac{\rho_s S}{V} = \frac{1}{\frac{d_1}{\epsilon_1 S} + \frac{d_2}{\epsilon_2 S}} = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2}}$$



## 2.2 Capacitor with dielectrics

### Example 2

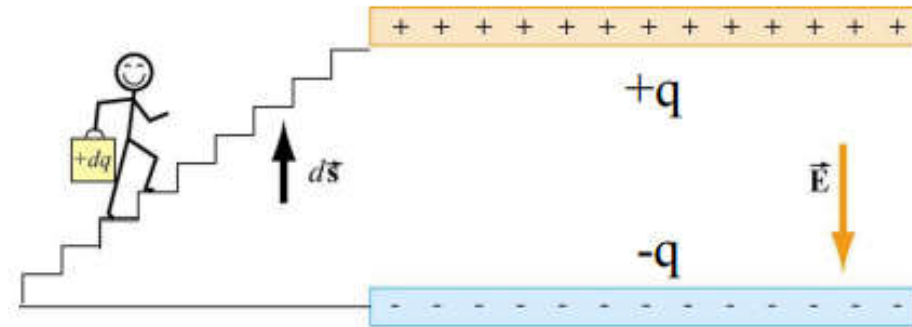
- A parallel-plate capacitor with three dielectric media is shown below. What is the total capacitance of the system?





## 2.3 Energy stored in a capacitor

- ✓ 1. Capacitor starts uncharged.
- ✓ 2. Carry  $+dq$  from bottom to top.  
Now top has charge  $q = +dq$ , bottom  $-dq$
- ✓ 3. Repeat
- ✓ 4. Finish when top has charge  $q = +Q$ , bottom  $-Q$ .

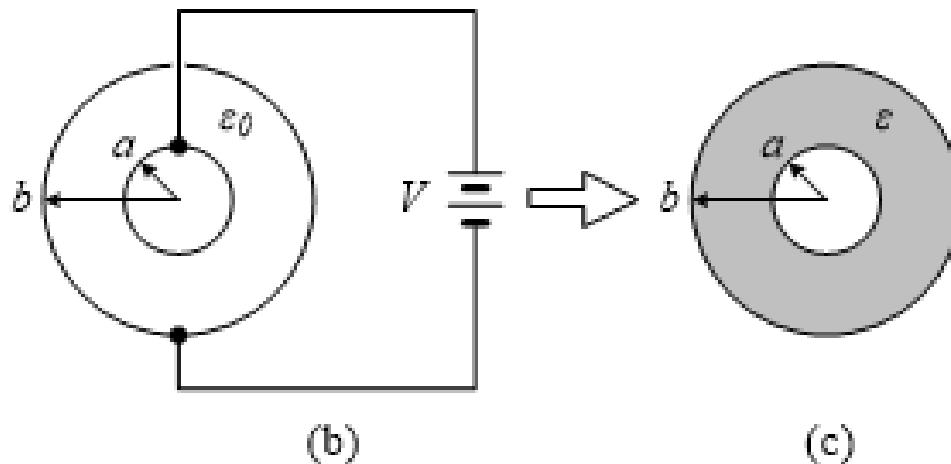


- At some point top plate has  $+q$ , potential difference is:  $\Delta\varphi = q/C$
- Work done to lift  $dq$  from the bottom to top is:  $dW = dq\Delta\varphi = qdq/C$
- So work done to move  $Q$  from bottom to top is:  
$$W = \int dW = \frac{1}{C} \int_0^Q q dq = \frac{1}{C} \frac{Q^2}{2}$$
- After charging, in  $\Delta\varphi = q/C$ ,  $\Delta\varphi$  is  $V$ , and  $q$  is  $Q$ , the total energy stored is:  
$$W = \frac{1}{C} \frac{Q^2}{2} = \frac{1}{C} \frac{(CV)^2}{2} = \frac{1}{2} CV^2$$

## 2.3 Energy stored in a capacitor

### Example 3

- An air-filled spherical capacitor with conductor radii  $a = 3$  cm and  $b = 15$  cm is connected to a source of voltage  $V = 15$  kV as shown in Figure (b).
- After an electrostatic state is established, the source is disconnected. The capacitor is then filled with a liquid dielectric of dielectric constant  $\epsilon_r = 2$  as shown in Figure (c).
- Determine the energy stored between the electrodes of the capacitor.



## 2.4 Current – Voltage Relationship

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- Start from the known relationship:

$$C = \frac{Q}{V} \Rightarrow V = \frac{Q}{C}$$

- In a time-dependent scenario:

$$V(t) = \frac{Q(t)}{C} = \frac{1}{C} \int_{t_0}^t I(\tau) d\tau + V(t_0)$$

- Taking the derivative of this and multiplying by  $C$ , get:

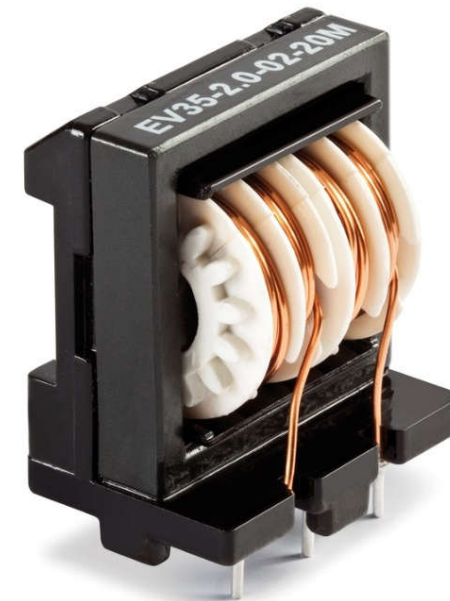
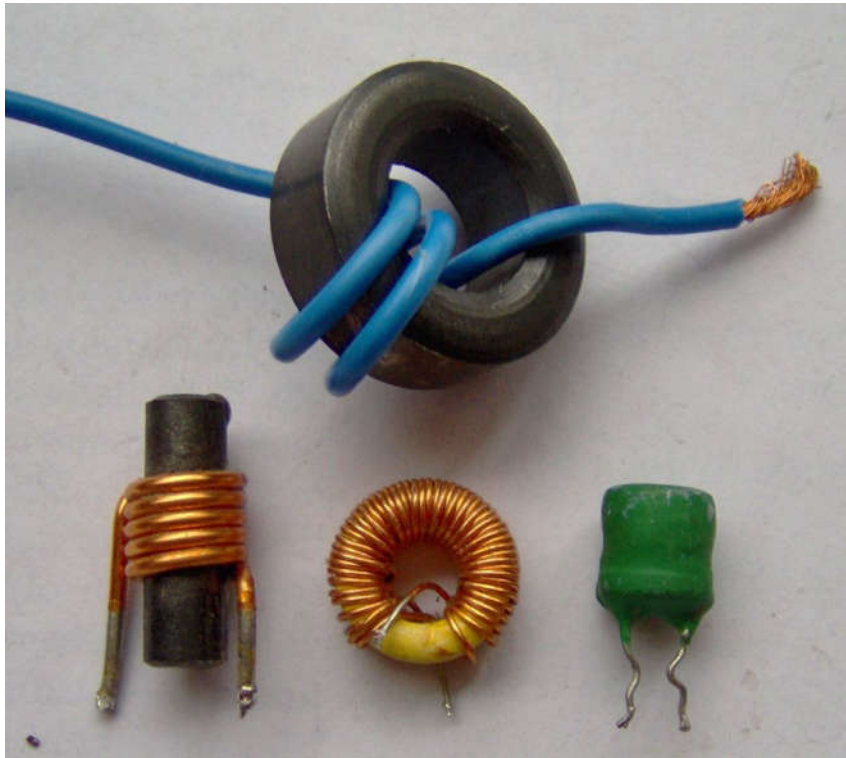
$$I(t) = \frac{dQ(t)}{dt} = C \frac{dV(t)}{dt}$$

- which means “the voltage on the capacitor is always continuous”;
- Also points out that the current “flows” through the capacitor is proportional to the capacitance and the changing rate of the voltage on the capacitor.



### 3. Inductors

- An inductor, also called a coil, is a passive two-terminal electrical component that stores energy in a magnetic field when electric current flows through it.



# 3.1 Self-inductance

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- Flux linkage:
  - Consider a toroid of  $N$  turns in which a current  $I$  produces a total flux  $\Phi$ . We assume first that this flux links or encircles each of the  $N$  turns, and we also see that each of the  $N$  turns links the total flux  $\Phi$ . The flux linkage  $N\Phi$  is defined as the product of the number of turns  $N$  and the flux  $\Phi$  linking each of them.
  - For a single turn coil, the flux linkage is equal to the total flux.
- Inductance(or self-inductance) is defined as the ratio of the total flux linkages to the current which they link

$$L = \frac{N\Phi}{I}$$

- This definition is for linear materials (not ferromagnetic materials)
- SI unit for inductance: H (Henry) = 1 Weber-turn per ampere.



# 3.1 Self-inductance

## Example 1

- Calculate the inductance per meter length of a coaxial cable of inner radius  $a$  and outer radius  $b$
- Solution:

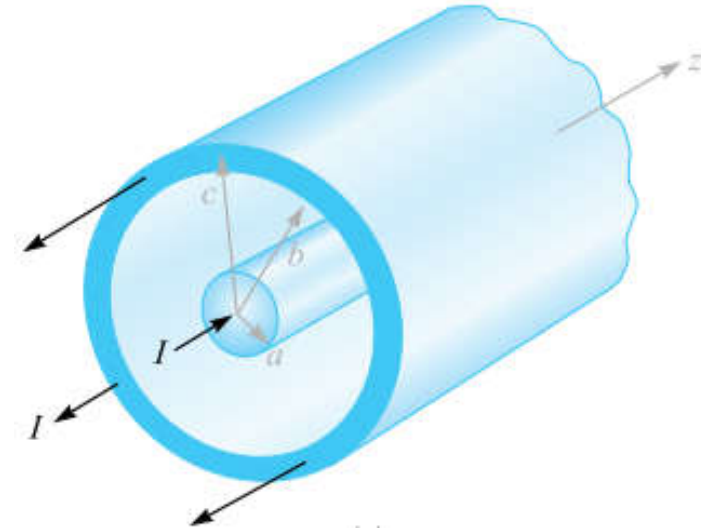
- The magnetic field intensity should be

$$H_\phi = \frac{I}{2\pi\rho} \quad (a < \rho < b) \quad \longrightarrow \quad \mathbf{B} = \mu_0 \mathbf{H} = \frac{\mu_0 I}{2\pi\rho} \mathbf{a}_\phi$$

- The magnetic flux contained between the conductors in a length  $d$  is

$$\Phi = \int_S \mathbf{B} \cdot d\mathbf{S} = \int_0^d \int_a^b \frac{\mu_0 I}{2\pi\rho} \mathbf{a}_\phi \cdot d\rho dz \mathbf{a}_\phi = \frac{\mu_0 I d}{2\pi} \ln \frac{b}{a}$$

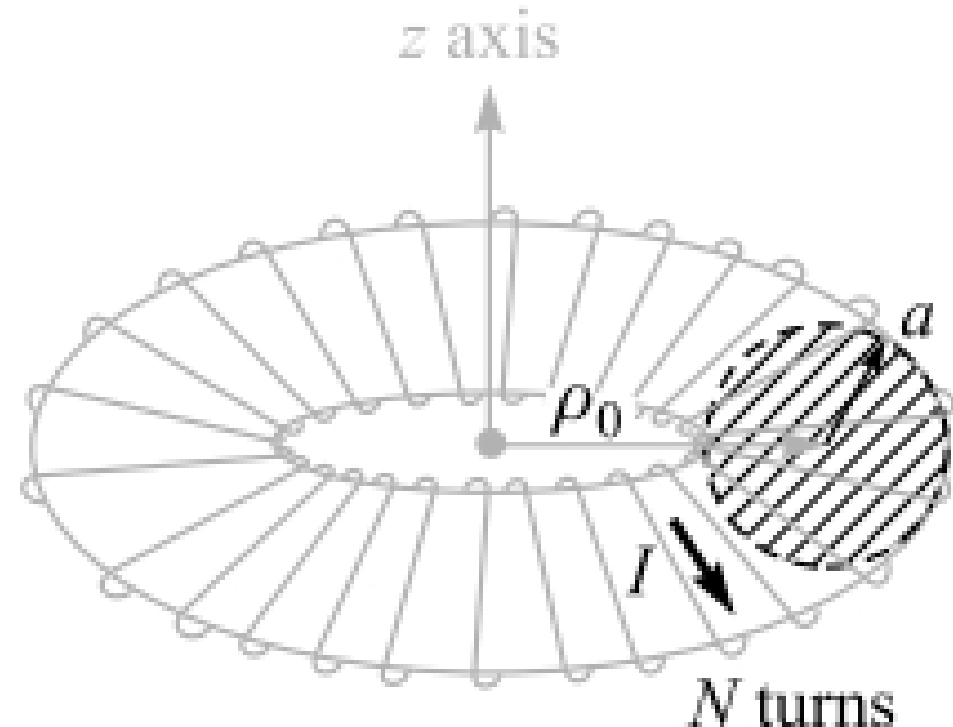
- So the inductance rapidly for a length  $d$  is  $L = \frac{\mu_0 d}{2\pi} \ln \frac{b}{a} \text{ H}$
- or, on a per-meter basis  $L = \frac{\mu_0}{2\pi} \ln \frac{b}{a} \text{ H/m}$



## 3.1 Self-inductance

### Example 2

- A toroid with a cross-section radius of  $a$ , consisting of  $N$  closely wound turns of a wire that carries a current  $I$  is shown below.
  - Determine the magnetic field intensity at points within the toroid;
  - Determine the inductance of the toroid.

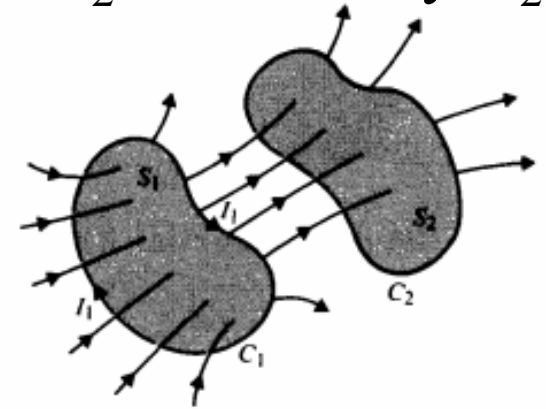




## 3.2 Mutual inductance

- Consider two neighboring closed loops  $C_1$  and  $C_2$  bounding surfaces  $S_1$  and  $S_2$ . If a current  $I_1$  flows in  $C_1$ , a magnetic field  $B_1$  will be created and its flux will pass through the surface  $S_2$  bounded by  $C_2$ .
- The mutual flux  $\Phi_{12}$  is defined as:
$$\Phi_{12} = \int_{S_2} \mathbf{B}_1 \cdot d\mathbf{s}_2 = L_{12}I_1$$
  - For the Biot-Savart law, we know that  $B_1$  is proportional to  $I_1$ , hence  $\Phi_{12}$  is also proportional to  $I_1$ , the proportionality constant  $L_{12}$  is called the mutual inductance.
- Neumann formula: the mutual inductance can be calculated:

$$L_{12} = \frac{\mu_0}{4\pi} \oint_{C_1} \oint_{C_2} \frac{d\mathbf{l}_1 \cdot d\mathbf{l}_2}{R}$$

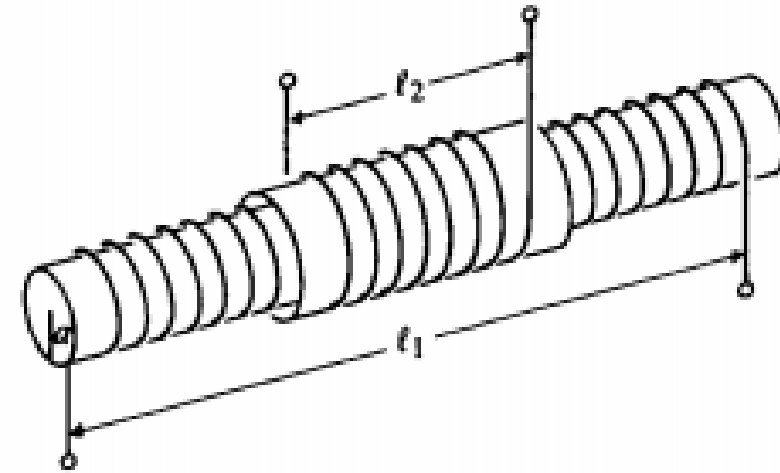




## 3.2 Mutual inductance

Example

- Two coils of  $N_1$  and  $N_2$  turns are wound concentrically on a straight cylindrical core of radius  $a$  and permeability  $\mu$ . The windings have lengths  $l_1$  and  $l_2$ , respectively.
- Find the mutual inductance between the coils.



## 3.3 Energy in a magnetic field

- In an electric field

- The energy density:

$$w_e = \frac{1}{2} \vec{\mathbf{D}} \cdot \vec{\mathbf{E}}$$

- The total electric energy stored in a medium:

$$W_e = \frac{1}{2} \int_v \vec{\mathbf{D}} \cdot \vec{\mathbf{E}} dv$$

- In a capacitor

$$W_C = \frac{1}{2} CV^2$$

- In a magnetic field

- The energy density:

$$w_m = \frac{1}{2} \vec{\mathbf{B}} \cdot \vec{\mathbf{H}}$$

$$= \frac{1}{2} \mu H^2 = \frac{1}{2\mu} B^2$$

- The total electric energy stored in a medium:

$$W_m = \int_v w_m dv$$

- In an inductor:

$$W_L = \frac{1}{2} LI^2$$



## 3.3 Energy in a magnetic field

Example

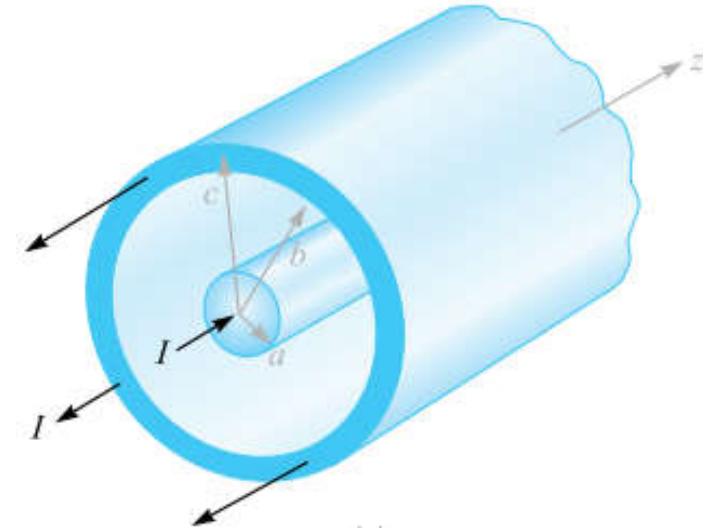
- Calculate the energy stored in a unit-length coaxial cable of inner radius  $a$  and outer radius  $b$
- Solution:
  - The magnetic field intensity should be

$$H_\phi = \frac{I}{2\pi\rho} \quad (a < \rho < b)$$

- So the total energy stored in this coaxial cable is

$$W_m = \int_v \frac{1}{2} \mu_0 H^2 dv = \frac{\mu_0}{2} \int_0^1 dz \int_0^{2\pi} \int_a^b \left( \frac{I}{2\pi\rho} \right)^2 \rho d\varphi d\rho = \frac{\mu_0 I^2}{4\pi} \ln \left( \frac{b}{a} \right) \quad (W)$$

- which agrees with the result calculated from  $W_L = \frac{1}{2} LI^2$



## 3.4 Current – voltage relationship

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- Any change in the current through an inductor creates a changing flux, inducing a voltage across the inductor. By Faraday's law of induction, the voltage induced by any change in magnetic flux through the circuit is given by:

$$\mathcal{E} = -\frac{d\Phi_B}{dt}$$

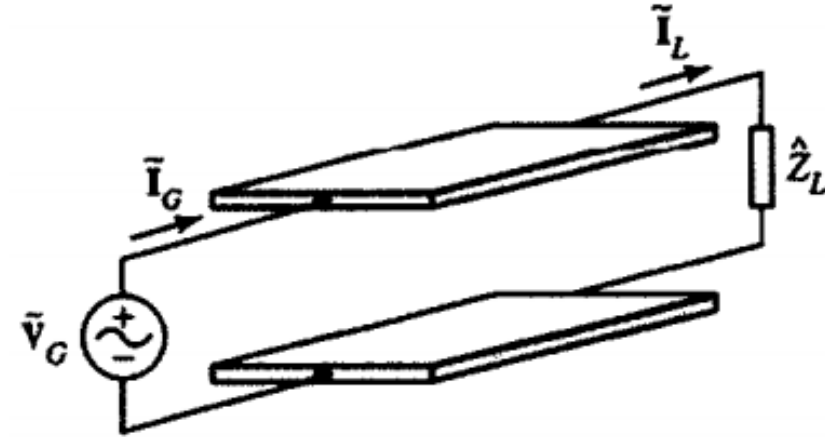
- And we know that  $\Phi_B = LI$ , substitute into the equation above, get:

$$\mathcal{E} = -\frac{d(LI)}{dt} = -L\frac{dI}{dt}$$

- which means “the current on the inductor is always continuous”;
- Also points out that the voltage “induced” on the inductor is proportional to the inductance and the changing rate of the current flowing through the inductor.

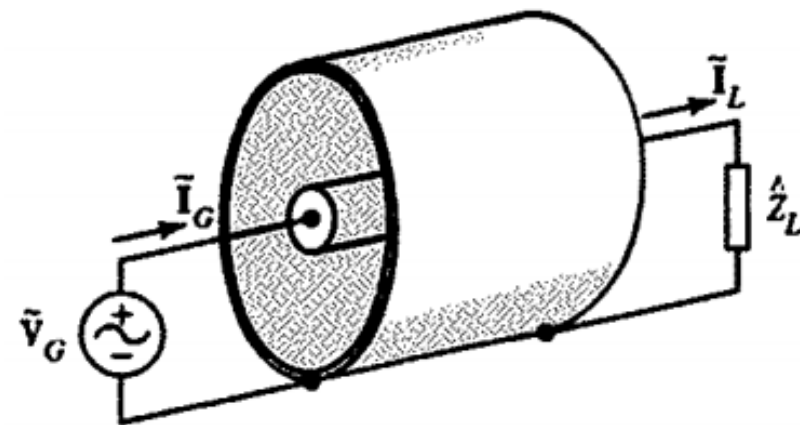
## 4.1 Case Study – Parallel plates

- For a pair of parallel plates, determine the resistance, capacitance, inductance and the leakage resistance (also called admittance).



## 4.2 Case Study – Coaxial cable

- For a piece of coaxial cable, determine the resistance, capacitance, inductance and the leakage resistance (also called admittance).



- Electrical circuit
  - Review the electric circuital properties of R, L and C
  - Review the AC circuit evaluation
  - Introduce “frequency response”