

Homework 5

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November 8, 2017

1. Suppose there is a flat piece of glass whose front surface is at $x = 0$ and whose back surface is at $x = l$. Suppose as well that its index of refraction depends only on the height: $n^2(y) = 1 + (n_0^2 - 1)e^{-|ay|}$. ($n(0) = n_n$, $\lim_{|y| \rightarrow \infty} n(y) = 1$)
 - (a) Write down the function that needs to be extremized to use Fermat's principle. You may restrict yourself to paths that have positive y in the glass.

From Fermat's principle, we want the time to be extremized.

$$\begin{aligned}
 v &= \frac{l}{t} \\
 t &= \frac{l}{v} \\
 &= \int_{y_1}^{y_2} \frac{\sqrt{dx^2 + dy^2} n(x, y)}{c} \\
 &= \int_{y_1}^{y_2} \frac{\sqrt{\frac{dx^2}{dy^2} + 1} n(x, y)}{c} dy \\
 &= \int_{y_1}^{y_2} \frac{1}{c} \sqrt{\dot{x}^2 + 1} n(x, y) dy
 \end{aligned}$$

Thus, we have our function waiting to be extremized:

$$L = \frac{1}{c} \sqrt{\dot{x}^2 + 1} n(x, y)$$

When we are outside the glass, i.e., if $x \in (-\infty, 0) \cup (l, +\infty)$, $n(x, y) = 1$. When we are inside the glass, i.e., if $x \in (0, l)$, $n(x, y) = n(y) = \sqrt{1 + (n_0^2 - 1)e^{-|ay|}}$.

- (b) Find the general form of the function that describes the path $x(y)$ of a light beam through the glass that starts at $x = 0$, $y = 0$ by extremizing the time the light takes. Be careful to find the integrating constants: express it in terms of the angle the light makes to the surface at $x = 0$

We just plug the function in to the Lagrangian's equation:

$$\frac{d}{dy} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0$$

Notice we have no x in L , this gives us $\frac{\partial L}{\partial x} = 0$. Thus we have:

$$\begin{aligned}
 \frac{d}{dy} \frac{\partial L}{\partial \dot{x}} &= 0 \\
 \frac{d}{dy} \frac{\partial}{\partial \dot{x}} \frac{1}{c} \sqrt{\dot{x}^2 + 1} n(x, y) &= 0 \\
 \frac{1}{c} \frac{d}{dy} \frac{\dot{x}}{\sqrt{\dot{x}^2 + 1}} n(x, y) &= 0
 \end{aligned}$$

This means $\frac{\dot{x}}{\sqrt{\dot{x}^2+1}}n(x,y)$ is a constant:

$$\frac{\dot{x}}{\sqrt{\dot{x}^2+1}}n(x,y) = K$$

Which gives us

$$\begin{aligned}\frac{\dot{x}}{\sqrt{\dot{x}^2+1}}n(x,y) &= K \\ \dot{x}n(x,y) &= K\sqrt{\dot{x}^2+1} \\ \dot{x}^2n(x,y)^2 &= K^2\dot{x}^2 + K^2 \\ \dot{x}^2(n^2 - K^2) &= K^2 \\ \dot{x} &= \frac{K}{\sqrt{n^2 - K^2}}\end{aligned}$$

When we are in the glass and $y > 0$, we have

$$\begin{aligned}\frac{dx}{dy} &= \frac{K}{\sqrt{n^2 - K^2}} \\ \frac{dx}{dy} &= \frac{K}{\sqrt{1 + (n_0^2 - 1)e^{-ay} - K^2}} \\ x &= \int \frac{K}{\sqrt{1 + (n_0^2 - 1)e^{-ay} - K^2}} dy\end{aligned}$$

This give us result:

$$\begin{aligned}& -\frac{2K \tan^{-1}\left(\frac{\sqrt{n^2-K^2}}{\sqrt{K^2-1}}\right)}{a\sqrt{K^2-1}} + J, \text{ if } K^2 > 1 \\ & -\frac{K \ln\left(\frac{|\sqrt{n^2-K^2}-\sqrt{1-K^2}|}{\sqrt{n^2-K^2}+\sqrt{1-K^2}}\right)}{a\sqrt{1-K^2}} + J, \text{ if } K^2 < 1\end{aligned}$$

From above we find

$$\frac{dx}{dy} = \cot \theta = \frac{K}{\sqrt{n^2 - K^2}}$$

Thus,

$$\tan \theta = \frac{1}{\cot \theta} = \frac{\sqrt{n^2(x,y) - K^2}}{K}$$

where θ is the angle while the light is inside the glass.

In this problem, we are finding the solution on $x(y) = x(0) = 0$.

$$\tan \theta = \frac{\sqrt{n^2 - K^2}}{K}$$

This give us

$$K^2 = \frac{n_0^2}{1 + \tan^2 \theta} = n^2 \cos^2 \theta$$

This give us the $K = n_0 \cos \theta$ at $x = y = 0$. Which is the constent we want to find depend on our angle in initial condiction.

Now we want to find the initial condiction for J . For $x = y = 0$, $n = n_0$, we have:

$$J = \frac{2n_0 \cos \theta \tan^{-1} \left(\frac{\sqrt{n_0^2 - n_0^2 \cos^2 \theta}}{\sqrt{n_0^2 \cos^2 \theta - 1}} \right)}{a \sqrt{n_0^2 \cos^2 \theta - 1}}, \text{ if } K^2 > 1$$

$$J = \frac{n_0 \cos \theta \ln \left(\frac{|\sqrt{n_0^2 - K^2} - \sqrt{1 - n_0^2 \cos^2 \theta}|}{\sqrt{n_0^2 - n_0^2 \cos^2 \theta} + \sqrt{1 - n_0^2 \cos^2 \theta}} \right)}{a \sqrt{1 - n_0^2 \cos^2 \theta}}, \text{ if } K^2 < 1$$

- (c) Consider two points, one at $x = -d, y = h$ and the second at $x = d + l, y = h$. Calculate the time it would take for light to go between the two points along a straight line, taking into account the time delay in glass. Now consider a path between the two points that goes a little way higher a point $h + \delta$ to the glass, then straight across the glass, and then back down to the other point, and calculate the time this would take. Is this longer or shorter than the time the straight line takes. (You can try an example such as $d = a = h = l = 1, n_0 = 2$ and plot the result as a function of δ or give an analytic argument)

For our first path:

$$t = \frac{2d}{c} + \frac{l}{c} \sqrt{1 + (n_0^2 - 1)e^{-ay}}$$

For our second path:

$$t = \frac{2\sqrt{d^2 + \delta^2}}{c} + \frac{l}{c} \sqrt{1 + (n_0^2 - 1)e^{-a(y+\delta)}}$$

Notice $\frac{2\sqrt{d^2 + \delta^2}}{c} \approx \frac{2d}{c}$, while $\frac{l}{c} \sqrt{1 + (n_0^2 - 1)e^{-a(y+\delta)}} < \frac{l}{c} \sqrt{1 + (n_0^2 - 1)e^{-ay}}$. This means the higher path takes less time. It suprise me a little bit.

2. Consider a particle of mass M that is constrained to move on a cone defined by $z^2 = A^2(x^2 + y^2)$ (A is a constant) with no external force other than the constraint on the particle. ($z > 0$)

- (a) Write down the Lagrangian in Cartesian coordinates with Lagrange multiplier that imposes the constraints.

$$L = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - \lambda(A^2(x^2 + y^2) - z^2)$$

- (b) Find the Euler-Lagrange equations in Cartesian coordinates.

The Euler-Lagrange equations is given as

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} = \frac{\partial L}{\partial q}$$

For us, we have:

$$\begin{aligned}
 x : \quad & \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} = \frac{\partial L}{\partial x} \\
 & \frac{d}{dt} m\dot{x} = -2\lambda A^2 x \\
 & m\ddot{x} = -2\lambda A^2 x \\
 y : \quad & \frac{d}{dt} \frac{\partial L}{\partial \dot{y}} = \frac{\partial L}{\partial y} \\
 & \frac{d}{dt} m\dot{y} = -2\lambda A^2 y \\
 & m\ddot{y} = -2\lambda A^2 y \\
 z : \quad & \frac{d}{dt} \frac{\partial L}{\partial \dot{z}} = \frac{\partial L}{\partial z} \\
 & \frac{d}{dt} m\dot{z} = -2\lambda z \\
 & m\ddot{z} = -2\lambda z
 \end{aligned}$$

That is we have following equations:

$$\begin{aligned}
 m\ddot{x} &= -2\lambda A^2 x \\
 m\ddot{y} &= -2\lambda A^2 y \\
 m\ddot{z} &= -2\lambda z
 \end{aligned}$$

- (c) Rewrite the Lagrangian in cylindrical polar coordinates. Solve the constraint for z and substitute the solution into the Lagrangian to get Lagrangian $L(R, \dot{R}, \dot{\phi})$

For cylindrical, we have

$$dl^2 = dr^2 + r^2 d\phi^2 + dz^2 \Rightarrow \vec{v}^2 = \dot{r}^2 + r^2 \dot{\phi}^2 + \dot{z}^2$$

Thus we have Lagrangian as:

$$L = \frac{1}{2} m(\dot{r}^2 + r^2 \dot{\phi}^2 + \dot{z}^2) - \lambda(A^2 r^2 - z^2)$$

Now we substitute $z^2 = A^2(x^2 + y^2) = A^2 r^2$, $\dot{z}^2 = A^2 \dot{r}^2$

$$L = \frac{1}{2} m(\dot{r}^2 + r^2 \dot{\phi}^2 + A^2 \dot{r}^2) = \frac{1}{2} m((1 + A^2)\dot{r}^2 + r^2 \dot{\phi}^2)$$

- (d) Find the Euler-Lagrange equations that follow from L that you found.

$$\begin{aligned}
 r : \quad & \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \frac{\partial L}{\partial r} \\
 & m(1 + A^2)\ddot{r} = mr\dot{\phi}^2 \\
 \phi : \quad & \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial \phi} \\
 & mr^2 \dot{\phi} = L_z
 \end{aligned}$$

- (e) Find the effective potential for r in terms of L_z , the z-component of the angular momentum.
From above we find

$$mr^2 \dot{\phi} = L_z$$

this give us

$$\dot{\phi} = \frac{L_z}{mr^2}$$

which give us

$$\begin{aligned} m(1 + A^2)\ddot{r} &= \frac{L_z^2}{mr^3} \\ m\ddot{r} &= \frac{L_z^2}{(1 + A^2)mr^3} = F \\ U_{\text{eff}} &= - \int \frac{L_z^2}{(1 + A^2)mr^3} dr \\ &= \frac{L_z^2}{(1 + A^2)m} \frac{1}{2r^2} \end{aligned}$$

3. Consider the same system as above, but now turn on a gravitational force.

- (a) Write the Lagrangian in cylindrical polar coordinates. Solve the constraint for z and substitute the solution into the Lagrangian to get a Lagrangian $L(R, \dot{R}, \dot{\phi})$

$$L = \frac{1}{2}m((1 + A^2)\dot{r}^2 + r^2\dot{\phi}^2) - mgAr$$

- (b) Find the Euler-Lagrange equations that follow from L you find:

$$\begin{aligned} r : \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} &= \frac{\partial L}{\partial r} \\ m(1 + A^2)\ddot{r} &= mr\dot{\phi}^2 - mgA \\ \phi : \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} &= \frac{\partial L}{\partial \phi} \\ mr^2\dot{\phi} &= L_z \end{aligned}$$

- (c) Find the effective potential for R in terms of L_z , the z -component of the angular momentum. Sketch the effective potential ($z > 0$)

From above we find:

$$\begin{aligned} m(1 + A^2)\ddot{r} &= mr\dot{\phi}^2 - mgA \\ m\ddot{r} &= \frac{L_z^2 - gAm^2r^3}{mr^3(1 + A^2)} = F_r \\ U_{\text{eff}} &= - \int \frac{L_z^2 - gAm^2r^3}{mr^3(1 + A^2)} dr \\ &= \frac{2Agm^2r^3 + L^2}{2(A^2 + 1)mr^2} \end{aligned}$$

- (d) For a given L_z , what is the equilibrium value of r ?

For equilibrium value, we want the force is zero:

$$\frac{L_z^2 - gAm^2r^3}{mr^3(1 + A^2)} = F_r = 0$$

the solution is given as

$$r = \left(\frac{L_z^2}{Agm^2} \right)^{\frac{1}{3}}$$

- (e) Is the equilibrium stable or unstable?

$$-\frac{\partial}{\partial r} \frac{L_z^2 - gAm^2r^3}{mr^3(1+A^2)} = \frac{3L_z^2}{(A^2+1)mr^4} > 0$$

This means we have an stable equilibrium.

4. Consider a bead of mass m constrained to move along a helical wire described in cylindrical coordinates by the equations $r = r_0$, and $z = \beta\phi$. Let gravity act on the bead.

- (a) Write down the Lagrangian - since the constraints are so simple, there is no need to use Lagrange multipliers, and you can simply use z as your variable.

$$L = \frac{1}{2}m(\dot{r}^2 + r^2\dot{\phi}^2 + \dot{z}^2) - mgz$$

Plug the constraints:

$$L = \frac{1}{2}m(r_0^2\dot{\phi}^2 + \beta^2\dot{\phi}^2) - mg\beta\phi$$

- (b) Write down the Euler-Lagrange equation. Find the general solution for arbitrary initial z_0 and \dot{z}_0 . How does this differ from a freely falling bead?

$$\begin{aligned} r : \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} &= \frac{\partial L}{\partial r} \\ 0 &= 0 \\ \phi : \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} &= \frac{\partial L}{\partial \phi} \\ \frac{d}{dt} (r_0^2 m \dot{\phi} + m \beta^2 \dot{\phi}) &= -mg\beta \\ \ddot{\phi} &= -\frac{mg\beta}{m(r_0^2 + \beta^2)} \end{aligned}$$

Now we want to find $\dot{\phi}$ and ϕ

$$\dot{\phi} = -\frac{mg\beta}{m(r_0^2 + \beta^2)}t + \dot{\phi}_0$$

and

$$\phi = -\frac{mg\beta}{m(r_0^2 + \beta^2)} \frac{t^2}{2} + \dot{\phi}_0 t + \phi_0$$

now we know that $z = \beta\phi$

$$\begin{aligned} \dot{z} &= -\beta \frac{mg\beta}{m(r_0^2 + \beta^2)}t + \dot{z}_0 \\ z &= -\beta \frac{mg\beta}{m(r_0^2 + \beta^2)} \frac{t^2}{2} + \dot{z}_0 t + z_0 \end{aligned}$$

Now we see that for a free falling

$$\dot{z} = -mgt + \dot{z}_0$$