
PHYSICS 252 EXPERIMENT NO. 11

THE DYNAMICS OF ELEMENTARY PARTICLES

Introduction

Elementary particles are everywhere around us. Apart from the standard matter particles proton, neutron and electron, hundreds of other particles have been found, produced in cosmic ray interactions in the atmosphere or by accelerators. Hundreds of charged particles traverse our bodies per second, and some will damage our DNA, one of the reasons for the existence of a sophisticated DNA repair mechanism in the cell.

The data for this experiment is in the form of a bubble chamber photograph which shows bubble tracks made by elementary particles as they traverse liquid hydrogen. In the experiment under study, a beam of low-energy negative pions (π^- beam) hits a hydrogen (p for proton) target in the form of a container with liquid hydrogen normally kept just below its boiling point ($T=20$ K). As the pions enter the detector a piston slightly decompresses the liquid so it becomes "super-critical" and starts boiling (i.e. bubble formation) first at the ionization (protons and electrons) trails left by the charged particles traversing the liquid.

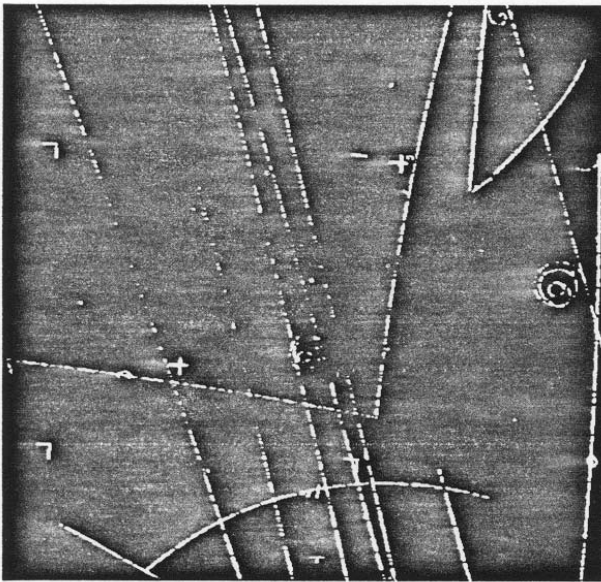


Figure 1: Photograph of the interaction between a high-energy π^- -meson from the Berkeley Bevatron accelerator and a proton in a liquid hydrogen bubble chamber, which produces two neutral particles Λ^0 and K^0 , which are short-

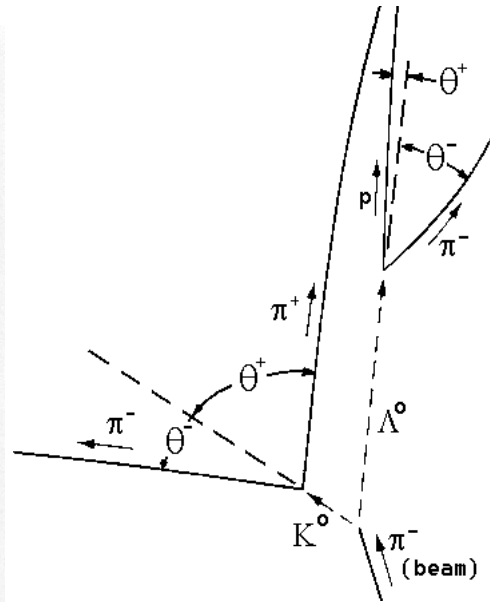


Figure 2: illustration of the interaction, and identification of bubble trails and variables to be measured in the *photograph* in

lived and decay into charged particles a bit further.

Figure 3.

The reaction photographed in Figure 1 shows the production of a pair of neutral particles (that do *not* leave a ionized trail in their wake), which after a short while decay into pairs of charged particles:

$$\pi^- + p \rightarrow \Lambda^0 + K^0 ,$$

where the neutrals decay as follows:

$$\Lambda^0 \rightarrow p + \pi^- , \quad K^0 \rightarrow \pi^+ + \pi^- .$$

In this experiment, we know the masses of the proton ($m_p = 938.3 \text{ MeV}/c^2$) and the pions ($m_{\pi^+} = m_{\pi^-} = 139.4 \text{ MeV}/c^2$) precisely, and will determine the masses of the Λ^0 and the K^0 , also in these mass energy units.

Momentum measurement

In order to "reconstruct" the interaction completely, one uses the conservation laws of (relativistic) momentum and energy, plus the knowledge of the initial pion beam parameters (mass and momentum). In order to measure momenta of the produced charged particles, the bubble chamber is located inside a magnet that bends the charged particles in helical paths. The 1.5 T magnetic field is directed up out of the photograph. The momentum p of each particle is directly proportional to the radius of curvature r , which in turn can be calculated from a measurement of the "chord length" l and sagitta s as:

$$r = [l^2/(8s)] + [s/2] .$$

Note, that the above is strictly true only if all momenta are perfectly in the plane of the photograph; in actual experiments stereo photographs of the interaction are taken to be able to reconstruct the interaction in all three dimensions. The interaction in this photograph was specially selected for its planarity.

In the reproduced photograph the actual radius of curvature R of the track in the bubble chamber is multiplied by the magnification factor m , $r = mR$. For the reproduction in Figure 3 $m = \text{height of photograph (in mm)} / 173 \text{ mm}$.

The momentum p of the particles is proportional to their radius of curvature R in the chamber. To derive this relationship for relativistic particles we begin with Newton's law in the form:

$$\mathbf{F} = d\mathbf{p}/dt = e \mathbf{v} \times \mathbf{B} \quad (\text{Lorentz' force}).$$

Here the momentum (\mathbf{p}) is the relativistic momentum $m\mathbf{v}\gamma$, where the relativistic γ -factor is defined in the usual way $\gamma = [\sqrt{1-v^2/c^2}]^{-1}$. Thus, because the *speed* v is constant:

$$\mathbf{F} = d\mathbf{p}/dt = d(m\mathbf{v}\gamma)/dt = m\gamma d\mathbf{v}/dt = m\gamma (v^2/R)(-\mathbf{r}) = e v B (-\mathbf{r}) ,$$

where \mathbf{r} is the *unit* vector in the radial direction. Division by v on both sides of the last equality finally yields:

$$m\gamma v / R = p / R = e B ,$$

identical to the non-relativistic result! In atomic units we find:

$$p c \text{ (in eV)} = c R B , \quad \text{thus} \quad p \text{ (in MeV/c)} = 2.998 \times 10^8 R B \times 10^{-6} = 300 R(\text{in m}) B(\text{in T}) \quad (1)$$

Measurement of angles

Draw straight lines from the point of primary interaction to the points where the Λ^0 and the K^0 decay. Extend the lines beyond the decay vertices. Draw tangents to the four decay product tracks at the two vertices. (Take care drawing these tangents, as doing it carelessly is a source of large errors.) Use a protractor to measure the angles of the decay product tracks relative to the parent directions (use Fig. 3 for measurements and Fig. 2 for definitions).

Analysis

The laws of relativistic dynamics relevant to this calculation are written below. We use the subscripts zero, plus, and minus to refer to the charges of the decaying particles and the decay products.

$$p_+ \sin \theta_+ = p_- \sin \theta_- \quad (2)$$

$$p_0 = p_+ \cos \theta_+ + p_- \cos \theta_- \quad (3)$$

$$E_0 = E_+ + E_-, \quad \text{where} \quad E_+ = \sqrt{(p_+^2 c^2 + m_+^2 c^4)}, \quad \text{and} \quad E_- = \sqrt{(p_-^2 c^2 + m_-^2 c^4)}$$

$$m_0 c^2 = \sqrt{(E_0^2 - p_0^2 c^2)}$$

Note that there is a redundancy here. That is, if p_+ , p_- , θ_+ , and θ_- are all known, equation (2) is not needed to find m_0 . In our two-dimensional case we have two equations (2 and 3), and only one unknown quantity m_0 , and the system is over-determined. This is fortunate, because sometimes (as here) one of the four measured quantities will have a large experimental error. When this is the case, it is usually advantageous to use only three of the variables and to use equation (2) to *calculate* the fourth. Alternatively, one may use the over-determination to "fit" the m_0 , and to determine it more precisely.

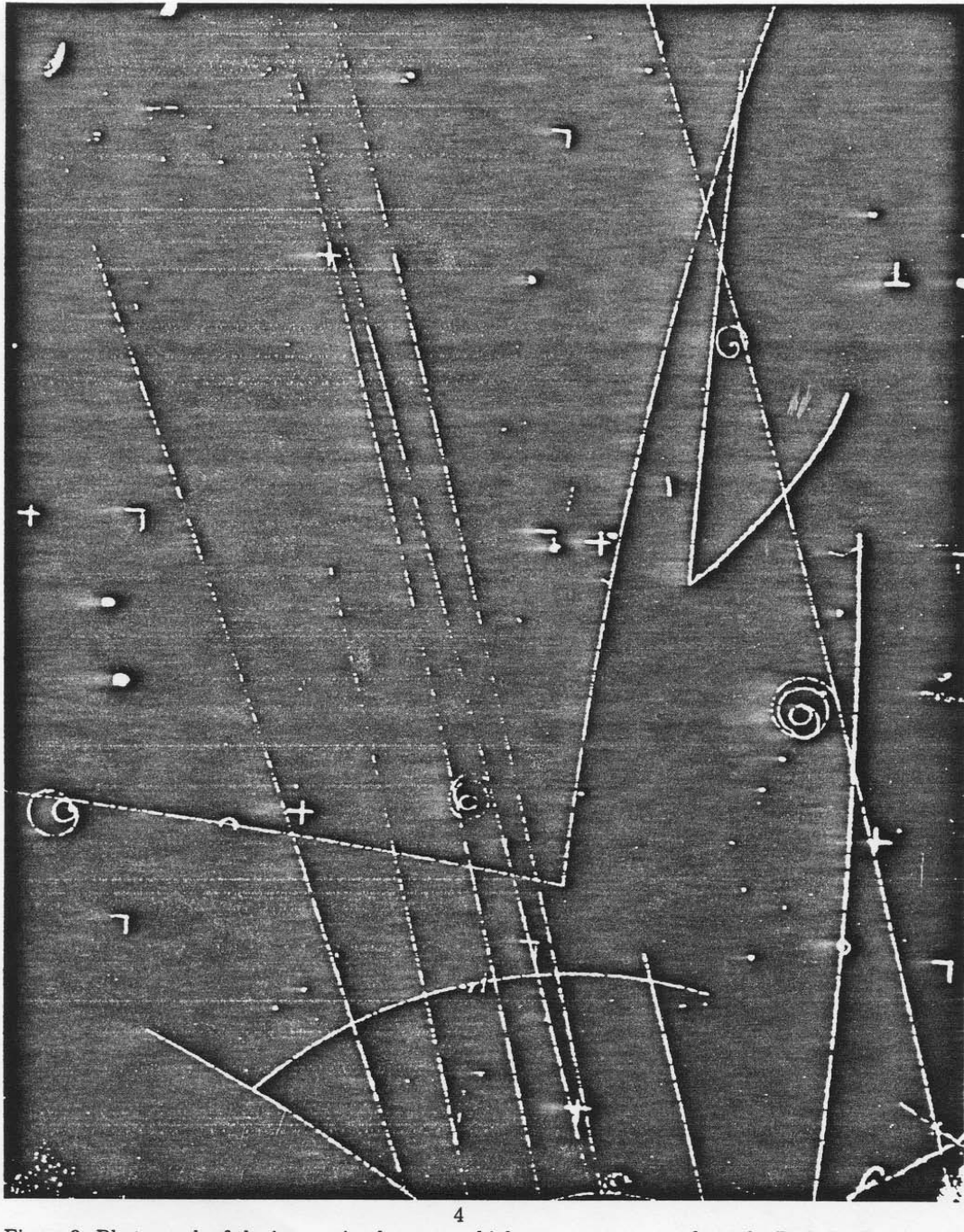
1. Focus your attention on the K^0 decay. Measure three of the quantities r_+ , r_- , θ_+ , and θ_- . Omit the one which you believe would introduce the largest experimental error if used to determine m_K .
2. Use the magnification factor m to calculate the actual radii R and equation (1) to calculate the momenta (in MeV/c) of one or both pions.
3. Use the equations above to determine the rest mass (in MeV/c²) of the K^0 .
4. Carefully estimate the error in your result from the errors in the measured quantities.
5. Now turn your attention to the Λ^0 : the proton track is too straight to be well measured in curvature. Also, θ_+ is small, and the value of m_Λ is quite sensitive to this measurement. Assume that $\theta_+ = (0.32 \pm 0.05)^\circ$ (check this with your protractor). Measure r_- and θ_- .
6. Calculate m_Λ and its error the same way as for the K^0 .
7. Finally, compare your values with the accepted mass values (the world average), and discuss.

Bonus - 25% extra credit

Calculate the momenta for both neutral particles, and hence find their lifetimes, both in the laboratory, and in their own rest-frames. Compare the latter with the accepted values.

Figure 3 (Next Page): Photograph of the interaction between a high-energy π^- -meson from the Berkeley Bevatron accelerator and a proton in a liquid hydrogen bubble chamber. The interaction produces two neutral particles Λ^0 and K^0 , which are short-lived and decay into charged particles a bit further. The photo covers an area (H×W) of 173 mm × 138 mm of the bubble chamber. In this enlargement m = actual height (in mm) of the photograph / 173 mm.

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Figure 3: Photograph of the interaction between a high-energy π^- -meson from the Berkeley Bevatron and a proton in a hydrogen bubble chamber and produces two neutral particles (Λ^0 and K^0 , which are short-lived and decay into charged particles a bit further. In this enlargement $m = 1.26$.