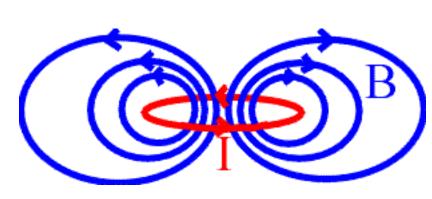
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MAGNETIC FIELD LINE INDUCED BY A CIRCULAR COIL AND A MAGNETIC DIPOLE



See also this article from the union of the physicists.

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Since the potential vector of the magnetic field created by a current with intensity I flowing through the curve (Γ_0) is

 $\overrightarrow{A}(M) = \frac{\mu_0 I}{4\pi} \int_{\Gamma_0} \frac{d\overrightarrow{M}_0}{M_0 M}, \text{ the magnetic field is } \overrightarrow{B}(M) = \overrightarrow{\text{rot}} \left(\overrightarrow{A}(M)\right) = \frac{\mu_0 I}{4\pi} \int_{\Gamma_0} d\overrightarrow{M}_0 \wedge \frac{\overrightarrow{M}_0 M}{(M_0 M)^3}, \text{ and the field lines are given by the}$

differential equation: $\overrightarrow{B}(M) \land d\overrightarrow{M} = \overrightarrow{0}$.

For a circular coil (Γ_0) : $\begin{cases} x_0 = a \cos t_0 \\ y_0 = 0 \\ z_0 = a \sin t_0 \end{cases}$ and a point M, we get the coordinates of the potential vector:

$$\begin{cases} A_{x} = \frac{\mu_{0} I a}{4\pi} \int_{-\pi}^{\pi} \frac{-\sin t}{\sqrt{(x - a\cos t)^{2} + y^{2} + (z - a\sin t)^{2}}} dt \\ A_{y} = 0 \\ A_{z} = \frac{\mu_{0} I a}{4\pi} \int_{-\pi}^{\pi} \frac{\cos t}{\sqrt{(x - a\cos t)^{2} + y^{2} + (z - a\sin t)^{2}}} dt \end{cases}$$

and the coordinates of the magnetic field:

$$\begin{cases} B_{x} = \frac{\partial A_{z}}{\partial y} = \frac{\mu_{0} Ia}{4\pi} \int_{-\pi}^{\pi} \frac{-y \cos t}{\left((x - a \cos t)^{2} + y^{2} + (z - a \sin t)^{2}\right)^{3/2}} dt \\ B_{y} = -\frac{\partial A_{z}}{\partial x} + \frac{\partial A_{x}}{\partial z} = \frac{\mu_{0} Ia}{4\pi} \int_{-\pi}^{\pi} \frac{(x - a \cos t) \cos t + (z - a \sin t) \sin t}{\left((x - a \cos t)^{2} + y^{2} + (z - a \sin t)^{2}\right)^{3/2}} dt \\ B_{z} = -\frac{\partial A_{x}}{\partial y} = \frac{\mu_{0} Ia}{4\pi} \int_{-\pi}^{\pi} \frac{y \sin t}{\left((x - a \cos t)^{2} + y^{2} + (z - a \sin t)^{2}\right)^{3/2}} dt \end{cases}$$

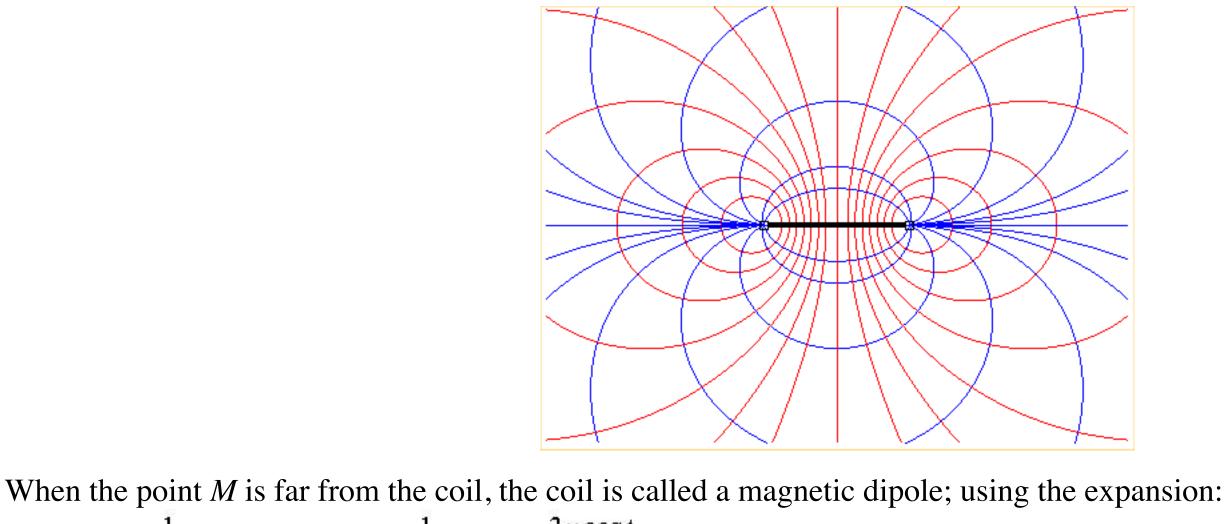
which, in xOy, reduces to:

$$\begin{cases} B_{x} = \frac{\mu_{0}Ia}{2\pi} \int_{0}^{\pi} \frac{-y\cos t}{\left(x^{2} + y^{2} + a^{2} - 2ax\cos t\right)^{3/2}} dt \\ B_{y} = \frac{\mu_{0}Ia}{2\pi} \int_{0}^{\pi} \frac{x\cos t - a}{\left(x^{2} + y^{2} + a^{2} - 2ax\cos t\right)^{3/2}} dt \text{ (Bessel integrals)} \\ B_{z} = 0 \end{cases}$$

hence the differential equation of the field lines: $y' = \frac{B_x}{B_y}$

and that of the orthogonal lines: $y' = -\frac{B_y}{R}$

The curves studied above are the magnetic field lines created by a continuous current flowing through the circular coil with radius a centred on O in the plane xOz, and their orthogonal trajectories:

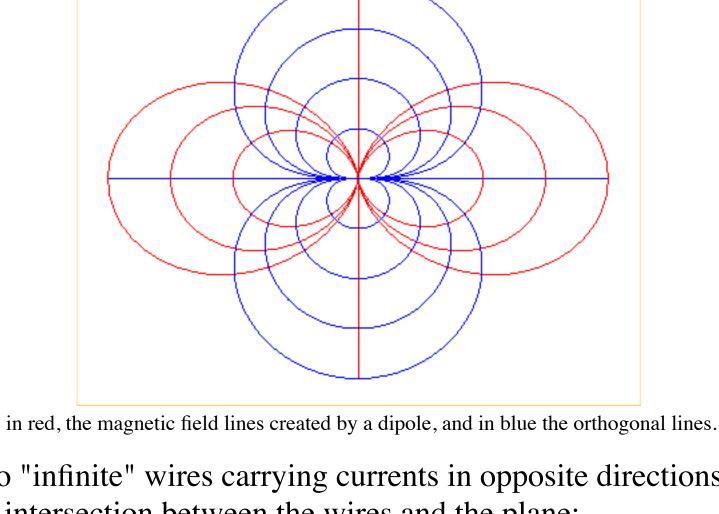


 $\frac{1}{\left(x^2+y^2+a^2-2ax\cos t\right)^{3/2}} = \frac{1}{\left(x^2+y^2\right)^{3/2}} + \frac{3x\cos t}{\left(x^2+y^2\right)^{5/2}}a + \dots, \text{ we get}$

$$(x^{2} + y^{2} + a^{2} - 2ax \cos t)^{-1} \quad (x^{2} + y^{2})^{-1} \quad (x^{2} + y^{2})^{-1}$$
the approximate magnetic field:
$$\begin{cases}
B_{x} \approx \widetilde{B}_{x} = \frac{\mu_{0}Ia^{2}}{4} \frac{-3xy}{\left(x^{2} + y^{2}\right)^{5/2}} \\
B_{y} \approx \widetilde{B}_{y} = \frac{\mu_{0}Ia^{2}}{4} \frac{x^{2} - 2y^{2}}{\left(x^{2} + y^{2}\right)^{5/2}}
\end{cases}$$
See another proof at romain.bel.free.fr/agregation/Lecons/LP27.doc

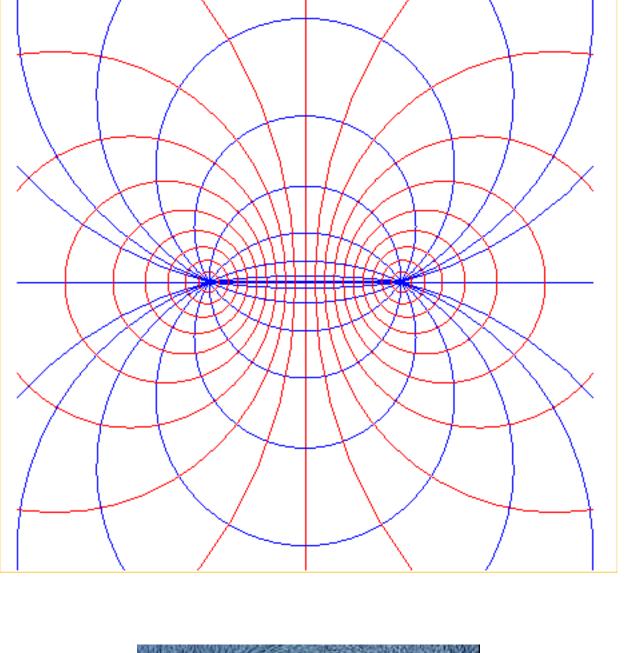
The field lines are then double eggs: $\rho = a\cos^2\theta = \frac{a}{2}\cos 2\theta + \frac{a}{2}$, and the orthogonal lines, curves of the dipole, with polar

equation $\rho^2 = a^2 \sin \theta$.



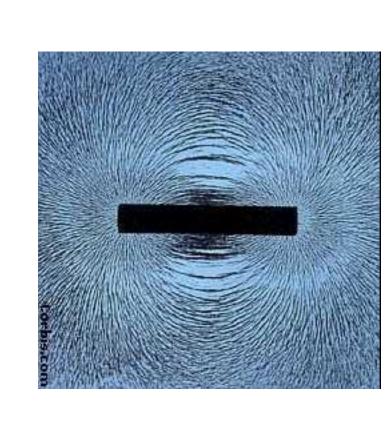
Remark: if the coil is replaced by two "infinite" wires carrying currents in opposite directions, the field lines form an intersecting

pencil of circles with base points the intersection between the wires and the plane:



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See other field lines <u>here</u>



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