Notes on ESPRIT me21321ods for large array processing

December 7, 2023

1 Introduction

In this pa123123per, we would like to evaluate the performance of the estimation of Direction of Arrival (DoA) method ESPRIT [?], the idea of which is described in the following section.

2 System Model

For a Unitary Linear Array (ULA) of size N, we consider the following model for the signal received at time t = 1, ..., T

$$\mathbf{x}(t) = \sum_{\ell=1}^{K} \mathbf{a}(\theta_{\ell}) s_{\ell}(t) + \mathbf{z}(t) \in \mathbb{C}^{N}$$
 (1)

with $\mathbf{a}(\theta_\ell) \in \mathbb{C}^N$ the steering vector of source s_ℓ at angle of arrival θ_ℓ , its j-th entry given by \mathbf{b}

$$[\mathbf{a}(\theta_{\ell})]_{j} = \frac{1}{\sqrt{N}} e^{i\frac{2\pi d}{\lambda_{0}}(j-1)\sin(\theta_{\ell})} \equiv \frac{1}{\sqrt{N}} e^{i\omega(j-1)\sin(\theta_{\ell})}, \quad \omega \equiv \frac{2\pi d}{\lambda_{0}}$$
(2)

where there is in total k signal sources $\{s_\ell\}_{\ell=1}^K$, at angle $\{\theta_\ell\}_{\ell=1}^K$ for some $K \ll \min(N,T)$, as well as some independent Gaussian noise $\mathbf{z}(t) \overset{i.i.d.}{\sim} \mathcal{CN}(\mathbf{0},\mathbf{I}_N)$ for all t.

The above signal model can be rewritten in matrix model by cascading the total *T* observations as

$$X = AS + Z \tag{3}$$

with $\mathbf{X} = [\mathbf{x}(1), \dots, \mathbf{x}(T)] \in \mathbb{C}^{N \times T}$, $\mathbf{A} = [\mathbf{a}(\theta_1), \dots, \mathbf{a}(\theta_k)] \in \mathbb{C}^{N \times K}$, $\mathbf{S} = [\mathbf{s}(1), \dots, \mathbf{s}(T)] \in \mathbb{C}^{K \times T}$, column vector $\mathbf{s}(t) = [s_1(t), \dots, s_k(t)]^\mathsf{T} \in \mathbb{C}^K$ and $\mathbf{Z} = [\mathbf{z}(1), \dots, \mathbf{z}(T)] \in \mathbb{C}^{N \times T}$ a standard (circular) Gaussian random matrix.

3 ESPRIT method for DoA estimation

In this paper, we would like to evaluate the performance of the estimation of Direction of Arrival (DoA) method ESPRIT [?], the idea of which is described as follows.

¹The normalization by \sqrt{N} here is for notational convenience so that $\mathbf{a}(\theta_{\ell})$ is of unit norm, note that this is equivalent to a *rescaling* of the source signal s_{ℓ} .

ESPRIT: intuition. Note that the *population* covariance of received signal

$$\mathbf{C} = \mathbb{E}[\mathbf{x}(t)\mathbf{x}^{\mathsf{H}}(t)] = \mathbb{E}[(\mathbf{A}\mathbf{s}(t) + \mathbf{z}(t))(\mathbf{A}\mathbf{s}(t) + \mathbf{z}(t))^{\mathsf{H}}]$$
$$= \mathbf{A}\mathbb{E}[\mathbf{s}(t)\mathbf{s}^{\mathsf{H}}(t)]\mathbf{A}^{\mathsf{H}} + \mathbb{E}[\mathbf{z}(t)\mathbf{z}^{\mathsf{H}}(z)]$$
$$= \mathbf{A}\mathbf{P}(t)\mathbf{A}^{\mathsf{H}} + \sigma^{2}\mathbf{I}_{N}$$

where we used the fact that $\mathbf{z}(t)$ is independent of the signal $\mathbf{s}(t)$ and denote the signal power $\mathbf{P}(t) \equiv \mathbb{E}[\mathbf{s}(t)\mathbf{s}^{\mathsf{H}}(t)]$. Then, for diagonal $\mathbf{P}(t) = \mathrm{diag}\{p_{\ell}(t)\}_{\ell=1}^{k}$ (which implies uncorrelated signal in the Gaussian case), one has

$$\mathbf{C} = \mathbb{E}[\mathbf{x}(t)\mathbf{x}^{\mathsf{H}}(t)] = \sum_{k=1}^{K} p_k(t)\mathbf{a}(\theta_k)\mathbf{a}^{\mathsf{H}}(\theta_k) + \sigma^2 \mathbf{I}_N = \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}} + \sigma^2 \mathbf{I}_N, \tag{4}$$

so that the top subspace of *population* covariance is expected to obtain structure information about the subspace spanned by the steering vectors $\mathbf{a}(\theta_k)$. *If the sample covariance* $\hat{\mathbf{C}}$ *is a good "proxy" of the population* \mathbf{C} in the sense that, e.g.,

$$\|\hat{\mathbf{C}} - \mathbf{C}\| \to 0 \tag{5}$$

in spectral norm, then, one has, by Davis-Kahan theorem that

$$\|\hat{\mathbf{U}}_S - \mathbf{U}_S\|_F \to 0,\tag{6}$$

(in fact, this holds for each individual eigenvector).

On the other hand, using the rotational invariance of the matrix **A**, we have, for two selection matrices $\mathbf{J}_1, \mathbf{J}_2 \in \mathbb{R}^{n \times N}$ that selection n among the in total N rows of \mathbf{X} , with "distance" Δ , that

$$\mathbf{J}_{1}\mathbf{A}\operatorname{diag}\left\{e^{\imath\omega\Delta\cdot\sin(\theta_{\ell})}\right\}_{k=1}^{K}=\mathbf{J}_{2}\mathbf{A}$$
(7)

with

$$\mathbf{J}_{1} = \begin{bmatrix} \mathbf{e}_{k}^{\mathsf{T}} \\ \dots \\ \mathbf{e}_{n+k-1}^{\mathsf{T}} \end{bmatrix} \in \mathbb{R}^{n \times N}, \quad \mathbf{J}_{2} = \begin{bmatrix} \mathbf{e}_{k+\Delta}^{\mathsf{T}} \\ \dots \\ \mathbf{e}_{n+k+\Delta-1}^{\mathsf{T}} \end{bmatrix} \in \mathbb{R}^{n \times N}$$
(8)

for \mathbf{e}_k the canonical vector of \mathbb{R}^N with $[\mathbf{e}_k]_i = \delta_{ij}$. We take, without loss of generality, k = 1 here so that

$$\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1} = \begin{bmatrix} \mathbf{I}_{n} & \mathbf{0}_{N-n} \\ \mathbf{0}_{N-n} & \mathbf{0}_{N-n} \end{bmatrix}, \quad \mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2} = \begin{bmatrix} \mathbf{0}_{n \times \Delta} & \mathbf{I}_{n} & \mathbf{0}_{n \times (N-n-\Delta)} \\ \mathbf{0} & \mathbf{0} & \mathbf{0} \end{bmatrix}$$
(9)

so that with $\mathbf{U}_S = \mathbf{A}\mathbf{Q}^{-1}$ for some invertible $\mathbf{Q} \in \mathbb{C}^{K \times K}$ and $\hat{\mathbf{U}}_S \simeq \mathbf{U}_S$, we obtain

$$\mathbf{Q}\operatorname{diag}\left\{e^{i\omega\Delta\cdot\sin(\theta_{\ell})}\right\}_{\ell=1}^{k}\mathbf{Q}^{-1} = (\mathbf{J}_{1}\mathbf{U}_{S})^{\dagger}\mathbf{J}_{2}\mathbf{U}_{S} \simeq (\mathbf{J}_{1}\hat{\mathbf{U}}_{S})^{\dagger}\mathbf{J}_{2}\hat{\mathbf{U}}_{S}. \tag{10}$$

ESPRIT: algorithm.

- 1. define two selection matrices $J_1, J_2 \in \mathbb{R}^{n \times N}$ that selection n among the in total N rows of X, with "distance" Δ , for instance, J_1 from row i to i + n, and J_2 from $i + \Delta$ to $i + \Delta + n$;
- 2. compute the sample covariance matrix $\hat{\mathbf{C}} = \frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}} = \frac{1}{T}\sum_{t=1}^{T}\mathbf{x}(t)\mathbf{x}^{\mathsf{H}}(t)$, and denote $\hat{\mathbf{U}}_{S}$ the top subspace composed of the eigenvectors associated to the largest-k eigenvalues, or the so-called *signal subspace*;

- 3. compute $\mathbf{\Phi} = (\mathbf{J}_1 \hat{\mathbf{U}}_S)^{\dagger} \mathbf{J}_2 \hat{\mathbf{U}}_S \in \mathbb{C}^{K \times K}$, where we denote \mathbf{A}^{\dagger} the Moore–Penrose pseudoinverse of \mathbf{A} , note that the resulting matrix $\mathbf{\Phi}$ is, in general, non-Hermitian;
- 4. the estimate of the angles of $\hat{\theta}_k$ are given by

$$\hat{\theta}_k = \arcsin(\arg(\lambda_k(\mathbf{\Phi}))/\omega/\Delta),\tag{11}$$

with λ_k the *k*th (complex) eigenvalues of Φ .

4 Characterization of ESPRIT method for large linear array

It follows from (??) that, for selection matrix J_1 such that J_1U_S has linearly independent columns so that the inverse $(\mathbf{U}_S^H \mathbf{J}_1^H \mathbf{J}_1 \mathbf{U}_S)^{-1}$ is well defined, we have

$$\operatorname{diag}\{e^{\imath\omega\Delta\cdot\sin(\theta_\ell)}\}_{k=1}^K=(\mathbf{A}^\mathsf{H}\mathbf{J}_1^\mathsf{H}\mathbf{J}_1\mathbf{A})^{-1}\mathbf{A}^\mathsf{H}\mathbf{J}_1^\mathsf{H}\mathbf{J}_2\mathbf{A}$$

where we assume that the selection matrix It thus suffices to evaluate the two terms $\mathbf{A}^H \mathbf{J}_1^H \mathbf{J}_1 \mathbf{A}$ and $\mathbf{A}^H \mathbf{J}_1^H \mathbf{J}_2 \mathbf{A}$ so as to retrieve the DoA θ_k as desired.

We positive ourselves under the following large array scenario.

Assumption 1 (Large array). *As* $T \rightarrow \infty$, *we have that*

$$0 < \liminf_{T} N/T < \limsup_{T} N/T < \infty, \quad 0 < \liminf_{T} n/N < \limsup_{T} n/N < 1.$$
 (12)

Assumption 2 (Widely spaced DoA). *All DoA angles* $\theta_1, \ldots, \theta_K$ *are* fixed as $N, T \to \infty$.

The *widely spaced* DoA scenario as **??** practically arises, e.g., when the DoA have an angular separation much larger than a beamwidth [?], by considering the case of all DoAs $\theta_1, \ldots, \theta_K$ are *fixed* with respect to *N* large. In this case, we have in particular that

$$[\mathbf{A}^{\mathsf{H}}\mathbf{A}]_{ij} = \mathbf{a}(\theta_i)^{\mathsf{H}}\mathbf{a}(\theta_j) = \frac{1}{N} \sum_{\ell=1}^{N} e^{-i\omega(\ell-1)(\sin(\theta_j) - \sin(\theta_i))}$$

$$= \begin{cases} 1 & \text{for } i = j \\ \frac{1}{N} \frac{1 - e^{-i\omega(\ell-1)N(\sin(\theta_j) - \sin(\theta_i))}}{1 - e^{-i\omega(\ell-1)(\sin(\theta_j) - \sin(\theta_i))}} = O(N^{-1}) & \text{for } i \neq j. \end{cases}$$

$$[\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\mathbf{A}]_{ij} = \mathbf{a}(\theta_{i})^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\mathbf{a}(\theta_{j}) = \begin{cases} \frac{n}{N}, & \text{for } i = j; \\ O(N^{-1}), & \text{for } i \neq j \end{cases}$$
(13)

as well as

$$[\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{2}\mathbf{A}]_{ij} = \mathbf{a}(\theta_{i})^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\mathbf{a}(\theta_{j}) = \begin{cases} \frac{n}{N}e^{-i\omega\Delta\sin(\theta_{i})}, & \text{for } i = j; \\ O(N^{-1}), & \text{for } i \neq j. \end{cases}$$
(14)

As such, under ?? and ??, we have, in matrix form that

$$\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\mathbf{A} = \begin{bmatrix} \mathbf{a}(\theta_{1})^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\mathbf{a}(\theta_{1}) & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & \mathbf{a}(\theta_{K})^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\mathbf{a}(\theta_{K}) \end{bmatrix} + O_{\|\cdot\|}(N^{-1}), \quad (15)$$

and similarly

$$\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{2}\mathbf{A} = \begin{bmatrix} \mathbf{a}(\theta_{1})^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{2}\mathbf{a}(\theta_{1}) & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & \mathbf{a}(\theta_{K})^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{2}\mathbf{a}(\theta_{K}) \end{bmatrix} + O_{\|\cdot\|}(N^{-1}).$$
(16)

As a consequence, we get

$$\mathbf{Q}^{-1}\operatorname{diag}\{e^{\imath\omega\Delta\cdot\sin(\theta_k)}\}_{k=1}^K\mathbf{Q} = (\mathbf{Q}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{J}_1^{\mathsf{H}}\mathbf{J}_1\mathbf{A}\mathbf{Q})^{-1}\mathbf{Q}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{J}_1^{\mathsf{H}}\mathbf{J}_2\mathbf{A}\mathbf{Q} = (\mathbf{U}_S^{\mathsf{H}}\mathbf{J}_1^{\mathsf{H}}\mathbf{J}_1\mathbf{U}_S)^{-1}\mathbf{U}_S^{\mathsf{H}}\mathbf{J}_1^{\mathsf{H}}\mathbf{J}_2\mathbf{U}_S,$$
(17)

holds for some invertible matrix $\mathbf{Q} \in \mathbb{C}^{K \times K}$ such that

$$\mathbf{U}_{S} = \mathbf{AQ}.\tag{18}$$

[Zhenyu: The above claim to clarify!]

The above result illustrates two cases, one is when N tends to infinity, the elements on the off-diagonal are 0 and the division of the elements on the diagonal is the angle we want

$$\mathbf{Q} \operatorname{diag} \{ e^{i\omega\Delta \cdot \sin(\theta_{\ell})} \}_{\ell=1}^{k} \mathbf{Q}^{-1} = (\mathbf{J}_{1}\mathbf{U}_{S})^{\dagger} \mathbf{J}_{2}\mathbf{U}_{S}.$$

$$= (\operatorname{diag}(\mathbf{U}_{S}^{H} \mathbf{J}_{1}^{H} \mathbf{J}_{1}\mathbf{U}_{S})) \setminus \operatorname{diag}(\mathbf{U}_{S}^{H} \mathbf{J}_{1}^{H} \mathbf{J}_{2}\mathbf{U}_{S})$$

In practice we cannot correctly estimate the u_i , and we tend to approximate u_i by \hat{u}_i , which leads to an error.

We need to find two functions $u_i^H \mathbf{J}_1^H \mathbf{J}_1 u_i = f_1(\hat{u}_i^H \mathbf{J}_1 \mathbf{J}_1^H \hat{u}_i), u_i^H \mathbf{J}_1^H \mathbf{J}_2 u_i = f_2(\hat{u}_i^H \mathbf{J}_1^H \mathbf{J}_2 \hat{u}_i),$ We need to approximate the real eigenvector by these two functions

If we can perfectly estimate the true eigenvector,

$$(diag(\mathbf{A}^H \mathbf{J}_1^H \mathbf{J}_1 \mathbf{A})) \setminus diag(\mathbf{A}^H \mathbf{J}_1^H \mathbf{J}_2 \mathbf{A}) = diag\{e^{i\omega\Delta \cdot \sin(\theta_\ell)}\}_{\ell=1}^k$$

5 Simplified model with single DoA

In the case of k = 1 with angle θ , we have the following simplified model

$$\mathbf{X} = \mathbf{a}(\theta)[s(1), \dots, s(T)] + \mathbf{Z} \equiv \mathbf{a}(\theta)\mathbf{s}^{\mathsf{H}} + \mathbf{Z}$$
(19)

so that

$$\frac{1}{T}XX^{\mathsf{H}} = \frac{1}{T}ZZ^{\mathsf{H}} + \begin{bmatrix} \mathbf{a} & \frac{\mathbf{Z}\mathbf{s}}{T} \end{bmatrix} \begin{bmatrix} \rho & 1\\ 1 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{a}^{\mathsf{H}}\\ \frac{(\mathbf{Z}\mathbf{s})^{\mathsf{H}}}{T} \end{bmatrix} \equiv \frac{1}{T}ZZ^{\mathsf{H}} + \mathbf{V}\Lambda\mathbf{V}^{\mathsf{H}}$$
(20)

where we denote $\rho = \frac{1}{T} \|\mathbf{s}\|^2$ the (limit of the normalized) signal strength and $\mathbf{V} = \begin{bmatrix} \mathbf{a} & \frac{\mathbf{Z}\mathbf{s}}{T} \end{bmatrix} \in \mathbb{C}^{N \times 2}$.

As a consequence,

$$\left(\frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}} - z\mathbf{I}_{N}\right)^{-1} = \left(\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}} + \mathbf{V}\boldsymbol{\Lambda}\mathbf{V}^{\mathsf{H}} - z\mathbf{I}_{N}\right)^{-1}$$
$$= \mathbf{Q} - \mathbf{Q}\mathbf{V}\boldsymbol{\Lambda}(\mathbf{I}_{2} + \mathbf{V}^{\mathsf{H}}\mathbf{Q}\mathbf{V}\boldsymbol{\Lambda})^{-1}\mathbf{V}^{\mathsf{H}}\mathbf{Q}$$

where we denoted $\mathbf{Q}(z) = \mathbf{Q} = (\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}} - z\mathbf{I}_N)^{-1}$ and used Woodbury matrix identity. Now, since

$$\mathbf{V}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{V} = \begin{bmatrix} \mathbf{a}^{\mathsf{H}} \\ \frac{(\mathbf{Z}\mathbf{s})^{\mathsf{H}}}{T} \end{bmatrix} \mathbf{Q}(z) \begin{bmatrix} \mathbf{a} & \frac{\mathbf{Z}\mathbf{s}}{T} \end{bmatrix} = \begin{bmatrix} \mathbf{a}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{a} & 0 \\ 0 & \frac{1}{T}\mathbf{s}^{\mathsf{H}}\frac{1}{T}\mathbf{Z}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{Z}\mathbf{s} \end{bmatrix} + o(1)$$
(21)

with

$$\frac{1}{T}\mathbf{s}^{\mathsf{H}}\frac{1}{T}\mathbf{Z}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{Z}\mathbf{s} = \frac{1}{T}\mathbf{s}^{\mathsf{H}}\tilde{\mathbf{Q}}(z)\frac{1}{T}\mathbf{Z}^{\mathsf{H}}\mathbf{Z}\mathbf{s} = \frac{1}{T}\mathbf{s}^{\mathsf{H}}(\mathbf{I}_{T} + z\tilde{\mathbf{Q}}(z))\mathbf{s} = \rho + \frac{z}{T}\mathbf{s}^{\mathsf{H}}\tilde{\mathbf{Q}}(z)\mathbf{s}$$
(22)

for co-resolvent $\tilde{\mathbf{Q}}(z) = \left(\frac{1}{T}\mathbf{Z}^{\mathsf{H}}\mathbf{Z} - z\mathbf{I}_{T}\right)^{-1}$.

$$\mathbf{Q}(z) \leftrightarrow \mathbf{\bar{Q}}(z) = m(z)\mathbf{I}_N = \left(\frac{1}{1 + cm(z)} - z\right)^{-1}\mathbf{I}_N, \quad \mathbf{\tilde{Q}}(z) = -\left(\frac{1}{zm(z)} + 1\right)\mathbf{I}_T \quad (23)$$

for $c = \lim N/T$ and m(z) the unique solution of the Marčenko-Pastur equation

$$zcm^{2}(z) - (1 - c - z)m(z) + 1 = 0, (24)$$

Therefore

$$(\mathbf{I}_2 + \mathbf{V}^\mathsf{H} \mathbf{Q} \mathbf{V} \mathbf{\Lambda})^{-1} = \begin{bmatrix} 1 + \rho m(z) & m(z) \\ \rho \left(1 - z - \frac{1}{m(z)} \right) & 1 \end{bmatrix}^{-1} + o(1)$$
 (25)

and

$$\mathbf{V}\boldsymbol{\Lambda}(\mathbf{I}_{2}+\mathbf{V}^{\mathsf{H}}\mathbf{Q}\mathbf{V}\boldsymbol{\Lambda})^{-1}\mathbf{V}^{\mathsf{H}} = \begin{bmatrix} \mathbf{a} & \frac{\mathbf{Z}\mathbf{s}}{T} \end{bmatrix} \frac{1}{1+\rho+\rho z m(z)} \begin{bmatrix} \rho z + \frac{\rho}{m(z)} & \mathsf{H} \\ \mathsf{H} & \mathsf{H} \end{bmatrix} \begin{bmatrix} \mathbf{a}^{\mathsf{H}} \\ \frac{(\mathbf{Z}\mathbf{s})^{\mathsf{H}}}{T} \end{bmatrix} + o_{\|\cdot\|}(1)$$
(26)

so that it suffices to evaluate the following expectations:

- 1. $\mathbb{E}[\mathbf{Q}(z)\mathbf{a}\mathbf{a}^{\mathsf{H}}\mathbf{Q}(z)] = m^{2}(z)\mathbf{a}\mathbf{a}^{\mathsf{H}} + o_{\|\cdot\|}(1);$
- 2. $\frac{1}{T}\mathbb{E}[\mathbf{Q}(z)\mathbf{a}\mathbf{s}^{\mathsf{H}}\mathbf{Z}^{\mathsf{T}}\mathbf{Q}(z)] = o_{\|\cdot\|}(1)$ and it Hermitian transpose.

This thus allows to conclude that

$$\left(\frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}} - z\mathbf{I}_{N}\right)^{-1} \leftrightarrow m(z)\mathbf{I}_{N} - \frac{\rho m(z)(zm(z) + 1)}{1 + \rho(zm(z) + 1)}\mathbf{a}\mathbf{a}^{\mathsf{H}}.$$
 (27)

5.1 Random matrix analysis

In the case of single DoA, our object of interest is the following *complex* random variable

$$Z = \frac{\hat{\mathbf{u}}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \hat{\mathbf{u}}}{\hat{\mathbf{u}}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \hat{\mathbf{u}}},\tag{28}$$

with $\hat{\mathbf{u}} \in \mathbb{C}^N$ the dominant eigenvector of the sample covariance matrix $\hat{\mathbf{C}} = \frac{1}{T}\mathbf{X}\mathbf{X}^H$. Note that $\|\mathbf{J}_i^T\mathbf{J}_i\| = O(1)$ for $i, j \in \{1, 2\}$, we have

According to (??), we need to evaluate the random variable of the form $\mathbf{y}_1^H \mathbf{Q}(z) \mathbf{y}_2$ for $\mathbf{y}_1, \mathbf{y}_2 \in \mathbb{C}^N$ of bounded Euclidean norm. To this end, we introduce the following deterministic equivalent result.

Theorem 1 (Deterministic Equivalent).

$$\left(\frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}} - z\mathbf{I}_{N}\right)^{-1} \leftrightarrow m(z)\mathbf{I}_{N} - \frac{\rho m(z)(zm(z) + 1)}{1 + \rho(zm(z) + 1)}\mathbf{a}\mathbf{a}^{\mathsf{H}}.\tag{29}$$

As a consequence, we have, in the notation of (??),

$$\mathbf{e}_{j}^{\mathsf{H}}(\frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}} - z\mathbf{I}_{N})^{-1}\mathbf{e}_{j} \leftrightarrow m(z) - \frac{\rho m(z)(zm(z) + 1)}{1 + \rho(zm(z) + 1)}(\mathbf{e}_{j}^{\mathsf{H}}\mathbf{a})^{2}$$
(30)

and

$$\mathbf{e}_{j+\Delta}^{\mathsf{H}}(\frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}} - z\mathbf{I}_{N})^{-1}\mathbf{e}_{j} \leftrightarrow m(z)\delta_{\Delta=0} - \frac{\rho m(z)(zm(z)+1)}{1+\rho(zm(z)+1)}(\mathbf{e}_{j+\Delta}^{\mathsf{H}}\mathbf{a})(\mathbf{a}^{\mathsf{H}}\mathbf{e}_{j})$$
(31)

It follows from Cauchy's integral formula and residue theorem that, for Γ_S circling around the isolated eigenvalue, we have

$$\begin{split} &-\frac{1}{2\pi \iota} \oint_{\Gamma_S} \mathbf{e}_j^\mathsf{H} (\frac{1}{T} \mathbf{X} \mathbf{X}^\mathsf{H} - z \mathbf{I}_N)^{-1} \mathbf{e}_j \, dz \simeq \frac{1}{2\pi \iota} \oint_{\Gamma_S} \frac{m(z)}{1 + \rho(z m(z) + 1)} dz \cdot (\mathbf{e}_j^\mathsf{H} \mathbf{a})^2 \\ &= -\lim_{z \to \lambda_S} (z - \lambda_S) \frac{m(z)}{1 + \rho(z m(z) + 1)} \cdot (\mathbf{e}_j^\mathsf{H} \mathbf{a})^2 = -\frac{m(\lambda_S)}{\rho(m(\lambda_S) + \lambda_S m'(\lambda_S))} \cdot (\mathbf{e}_j^\mathsf{H} \mathbf{a})^2 \end{split}$$

with $\lambda_S \equiv 1 + \rho + c \frac{1+\rho}{\rho}$ the (asymptotic) position of the isolated eigenvalue and $m'(z) = \frac{m^2(z)}{1 - \frac{cm^2(z)}{(1+cm(z))^2}}$ obtained by differentiating the Marčenko-Pastur equation, where we used the

fact that Γ_S does not contain any pole of m(z).

And similarly that

$$-\frac{1}{2\pi i} \oint_{\Gamma_S} \mathbf{e}_{j+\Delta}^{\mathsf{H}} (\frac{1}{T} \mathbf{X} \mathbf{X}^{\mathsf{H}} - z \mathbf{I}_N)^{-1} \mathbf{e}_j \, dz \simeq -\frac{m(\lambda_S)}{\rho(m(\lambda_S) + \lambda_S m'(\lambda_S))} \cdot \mathbf{e}_{j+\Delta}^{\mathsf{H}} \mathbf{a} \cdot \mathbf{a}^{\mathsf{H}} \mathbf{e}_j \tag{32}$$

To-do list:

1. simulate the different between $\hat{\mathbf{u}}^H \mathbf{M} \hat{\mathbf{u}}$ and its deterministic equivalent, for spiked covariance model and big matrix \mathbf{M} with large rank, Note that two things happen at the same time: i) $\hat{\mathbf{u}}^H \mathbf{e} \mathbf{e}^H \hat{\mathbf{u}} = O(1)$ and ii) $\hat{\mathbf{u}}^H \mathbf{I}_N \hat{\mathbf{u}} = 1 = O(1)$.

6 Proof of ESPIRIT in the case of single DoA

To investigate the performance of ESPIRIT method, one needs to evaluate the statistic of the dominant eigenvector $\hat{\mathbf{u}}_S$ of the spiked random matrix model $\frac{1}{T}\mathbf{X}\mathbf{X}^H$, for $\mathbf{X} = \mathbf{a}\mathbf{s}^H + \mathbf{Z}$ with $\mathbf{a} \in \mathbb{C}^N$ such that $\|\mathbf{a}\| = 1$, $\mathbf{s} \in \mathbb{C}^T$ a standard (circular) Gaussian random vector, and $\mathbf{Z} \in \mathbb{C}^{N \times T}$ a standard (circular) Gaussian random matrix.

Denote $(\hat{\lambda}, \hat{\mathbf{u}})$ the pair of largest eigenvalue-eigenvector pair of $\frac{1}{T}\mathbf{X}\mathbf{X}^H$, and thus satisfies

$$\frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{T}}\hat{\mathbf{u}} = \hat{\lambda}\hat{\mathbf{u}} = \frac{\|\mathbf{s}\|^2}{T}\mathbf{a}\mathbf{a}^{\mathsf{H}}\hat{\mathbf{u}} + \frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}}\hat{\mathbf{u}} + \frac{1}{T}\left(\mathbf{a}\mathbf{s}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}} + \mathbf{Z}\mathbf{s}\mathbf{a}^{\mathsf{H}}\right)\hat{\mathbf{u}}.$$
 (33)

Denote $\mathbf{Q}(z) = (\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}} - z\mathbf{I}_N)$, for $z \in \mathbb{C}$ not an eigenvalue of $\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{T}}$ (which is known to have eigenvalues lying within the MP support as $N, T \to \infty$), we obtain

$$\begin{split} \mathbf{0} &= \frac{\|\mathbf{s}\|^2}{T} \mathbf{a} \mathbf{a}^\mathsf{H} \hat{\mathbf{u}} + \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^\mathsf{H} - \hat{\lambda} \mathbf{I}_N\right) \hat{\mathbf{u}} + \frac{1}{T} \left(\mathbf{a} \mathbf{s}^\mathsf{H} \mathbf{Z}^\mathsf{H} + \mathbf{Z} \mathbf{s} \mathbf{a}^\mathsf{H}\right) \hat{\mathbf{u}} \\ &\Leftrightarrow -\hat{\mathbf{u}} = \frac{\|\mathbf{s}\|^2}{T} \mathbf{a}^\mathsf{H} \hat{\mathbf{u}} \cdot \mathbf{Q}(\hat{\lambda}) \mathbf{a} + \frac{1}{T} \mathbf{s}^\mathsf{H} \mathbf{Z}^\mathsf{H} \hat{\mathbf{u}} \cdot \mathbf{Q}(\hat{\lambda}) \mathbf{a} + \mathbf{a}^\mathsf{H} \hat{\mathbf{u}} \cdot \frac{1}{T} \mathbf{Q}(\hat{\lambda}) \mathbf{Z} \mathbf{s} \\ &\Rightarrow \sqrt{N} [\hat{\mathbf{u}}]_i = \sqrt{N} \mathbf{e}_i^\mathsf{T} \hat{\mathbf{u}} = -\frac{\|\mathbf{s}\|^2}{T} \mathbf{a}^\mathsf{H} \hat{\mathbf{u}} \cdot \sqrt{N} \mathbf{e}_i^\mathsf{T} \mathbf{Q}(\hat{\lambda}) \mathbf{a} - \frac{1}{T} \mathbf{s}^\mathsf{H} \mathbf{Z}^\mathsf{H} \hat{\mathbf{u}} \cdot \sqrt{N} \mathbf{e}_i^\mathsf{T} \mathbf{Q}(\hat{\lambda}) \mathbf{a} - \mathbf{a}^\mathsf{H} \hat{\mathbf{u}} \cdot \frac{\sqrt{N}}{T} \mathbf{e}_i^\mathsf{T} \mathbf{Q}(\hat{\lambda}) \mathbf{Z} \mathbf{s} \end{split}$$

Note that till now no asymptotic approximation has been performed, we have only used linear algebraic results.

Proof to-do list:

- (i) establish the asymptotic *complex* limit of $\mathbf{a}^{\mathsf{H}}\hat{\mathbf{u}} = ? + o(1)$; and
- (ii) establish the asymptotic *complex* limit of $\sqrt{N}\mathbf{e}_{i}^{\mathsf{T}}\mathbf{Q}(\hat{\lambda})\mathbf{a} = ? + o(1)$; and
- (iii) show that $\frac{1}{T}\mathbf{s}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}}\hat{\mathbf{u}} \to 0$ almost surely (this, together with item (ii), allows us to asymptotic discard the term $\frac{1}{T}\mathbf{s}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}}\hat{\mathbf{u}} \cdot \sqrt{N}\mathbf{e}_{i}^{\mathsf{T}}\mathbf{Q}(\hat{\lambda})\mathbf{a}$);

this allows us to conclude that the i-th entry of $\hat{\mathbf{u}}$ satisfies

$$\sqrt{N}[\hat{\mathbf{u}}]_{i} = -\mathbf{a}^{\mathsf{H}} \hat{\mathbf{u}} \left(\underbrace{\frac{\|\mathbf{s}\|^{2}}{T} \sqrt{N} \mathbf{e}_{i}^{\mathsf{T}} \mathbf{Q}(\hat{\lambda}) \mathbf{a}}_{\text{deterministic } O(1) + o(1)} + \underbrace{\frac{\sqrt{N}}{T} \mathbf{e}_{i}^{\mathsf{T}} \mathbf{Q}(\hat{\lambda}) \mathbf{Z} \mathbf{s}}_{\text{Gaussian fluctuation } O(1)} + o(1).$$
(34)

This further leads to

$$N[\hat{\mathbf{u}}]_i^2 = +o(1),$$

$$N[\hat{\mathbf{u}}]_i[\hat{\mathbf{u}}]_j = +o(1).$$

We thus obtain

$$(\mathbf{J}_{1}\hat{\mathbf{u}})^{\dagger}\mathbf{J}_{2}\hat{\mathbf{u}} = \frac{\hat{\mathbf{u}}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{2}\hat{\mathbf{u}}}{\hat{\mathbf{u}}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{H}}\mathbf{J}_{1}\hat{\mathbf{u}}} = \frac{\sum_{j=i}^{i+n} \overline{[\hat{\mathbf{u}}]_{j}} [\hat{\mathbf{u}}]_{j+\Delta}}{\sum_{j=i}^{i+n} [\hat{\mathbf{u}}]_{j}^{2}}$$
(35)

7 Alternative proof of ESPIRIT in the case of single DoA: random signal case

In the case of (proper) complex Gaussian signal $\mathbf{s} \sim \mathcal{CN}(\mathbf{0}, \rho^2 \mathbf{I}_T)$ with signal strength ρ^2 , we have that the observation matrix $\mathbf{X} \in \mathbb{C}^{N \times}$ is equivalently given by

$$\mathbf{X} = \left(\mathbf{I}_N + \rho^2 \mathbf{a} \mathbf{a}^\mathsf{H}\right)^{\frac{1}{2}} \mathbf{Z},\tag{36}$$

for standard complex Gaussian $\mathbf{Z} \in \mathbb{C}^{N \times T}$.

Let us first consider the form $\hat{\mathbf{u}}^H \mathbf{J}_1^T \mathbf{J}_2 \hat{\mathbf{u}}$, with $\hat{\mathbf{u}}$ the dominant eigenvector of the SCM $\hat{\mathbf{C}} = \frac{1}{T} \mathbf{X} \mathbf{X}^H$. We have, for Γ circling around the isolated eigenvalue of $\hat{\mathbf{C}}$, that

$$\begin{split} \hat{\mathbf{u}}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \hat{\mathbf{u}} &= \sum_{i=1}^{n} \hat{\mathbf{u}}^{\mathsf{H}} \mathbf{e}_{i} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \hat{\mathbf{u}} \\ &= -\frac{1}{2\pi \iota} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\hat{\mathbf{C}} - z \mathbf{I}_{N} \right)^{-1} \mathbf{e}_{i} \, dz \\ &= -\frac{1}{2\pi \iota} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^{\mathsf{H}} - z \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-1} \right)^{-1} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi \iota} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^{\mathsf{H}} - z \mathbf{I}_{N} + z \frac{\rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}}}{1 + \rho^{2}} \right)^{-1} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi \iota} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \left(\mathbf{Q}(z) - \frac{\rho^{2}}{1 + \rho^{2}} \frac{z \mathbf{Q}(z) \mathbf{a} \mathbf{a}^{\mathsf{H}} \mathbf{Q}(z)}{1 + \frac{\rho^{2}}{1 + \rho^{2}} \cdot z \mathbf{a}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{a}} \right) \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= \frac{1}{2\pi \iota} \frac{\rho^{2}}{1 + \rho^{2}} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \frac{z \mathbf{Q}(z) \mathbf{a} \mathbf{a}^{\mathsf{H}} \mathbf{Q}(z)}{1 + \frac{\rho^{2}}{1 + \rho^{2}} \cdot z \mathbf{a}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{a}} \left(\mathbf{I}_{N} + \rho^{2} \mathbf{a} \mathbf{a}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \end{split}$$

with the resolvent

$$\mathbf{Q}(z) \equiv \left(\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}} - z\mathbf{I}_{N}\right)^{-1}.$$
 (37)

This leads to

$$\begin{split} \hat{\mathbf{u}}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\hat{\mathbf{u}} &= \frac{1}{2\pi\imath}\frac{\rho^{2}}{1+\rho^{2}}\sum_{i=1}^{n}\oint_{\Gamma}\frac{z\mathbf{a}^{\mathsf{H}}\mathbf{Q}(z)\left(\mathbf{I}_{N}+\rho^{2}\mathbf{a}\mathbf{a}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{e}_{i}\mathbf{e}_{i+\Delta}^{\mathsf{T}}\left(\mathbf{I}_{N}+\rho^{2}\mathbf{a}\mathbf{a}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{Q}(z)\mathbf{a}}{1+\frac{\rho^{2}}{1+\rho^{2}}\cdot z\mathbf{a}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{a}}\,dz,\\ &= \frac{1}{2\pi\imath}\frac{\rho^{2}}{1+\rho^{2}}\oint_{\Gamma}\frac{z\mathbf{a}^{\mathsf{H}}\mathbf{Q}(z)\left(\mathbf{I}_{N}+\rho^{2}\mathbf{a}\mathbf{a}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\left(\mathbf{I}_{N}+\rho^{2}\mathbf{a}\mathbf{a}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{Q}(z)\mathbf{a}}{1+\frac{\rho^{2}}{1+\rho^{2}}\cdot z\mathbf{a}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{a}}\,dz, \end{split}$$

Note that the only pole in this case is $\lambda \in \mathbb{R}$ such that

$$z\mathbf{a}^{\mathsf{H}}\mathbf{Q}(\lambda)\mathbf{a} = -\frac{1+\rho^2}{\rho^2}.$$
 (38)

Since $\mathbf{Q}(z) \leftrightarrow \mathbf{\bar{Q}}(z) = m(z)\mathbf{I}_N$, we have

$$h(\lambda) \equiv \lambda m(\lambda) = -1 - \rho^{-2},\tag{39}$$

with

$$\lambda = 1 + \rho^2 + c \frac{1 + \rho^2}{\rho^2},\tag{40}$$

and therefore the following first-order result

$$\begin{split} &\frac{1}{2\pi\imath}\frac{\rho^2}{1+\rho^2}\oint_{\Gamma}\mathbf{e}_{i+\Delta}^\mathsf{T}\left(\mathbf{I}_N+\rho^2\mathbf{a}\mathbf{a}^\mathsf{H}\right)^{-\frac{1}{2}}\frac{z\mathbf{Q}(z)\mathbf{a}\mathbf{a}^\mathsf{H}\mathbf{Q}(z)}{1+\frac{\rho^2}{1+\rho^2}\cdot z\mathbf{a}^\mathsf{H}\mathbf{Q}(z)\mathbf{a}}\left(\mathbf{I}_N+\rho^2\mathbf{a}\mathbf{a}^\mathsf{H}\right)^{-\frac{1}{2}}\mathbf{e}_i\,dz\\ &\simeq \frac{1}{2\pi\imath}\frac{\rho^2}{1+\rho^2}\oint_{\Gamma}\frac{zm^2(z)\,dz}{1+\rho^2+\rho^2zm(z)}\mathbf{e}_{i+\Delta}^\mathsf{T}\mathbf{a}\mathbf{a}^\mathsf{H}\mathbf{e}_i+O(N^{-1/2})\\ &=-\mathrm{Res}\left(\frac{zm^2(z)\,dz}{1+\rho^2+\rho^2zm(z)}\right)\frac{\rho^2}{1+\rho^2}\mathbf{e}_{i+\Delta}^\mathsf{T}\mathbf{a}\mathbf{a}^\mathsf{H}\mathbf{e}_i+O(N^{-1/2})\\ &=\frac{m(\lambda)(1+h(\lambda))}{h'(\lambda)}\mathbf{e}_{i+\Delta}^\mathsf{T}\mathbf{a}\mathbf{a}^\mathsf{H}\mathbf{e}_i+O(N^{-1/2}). \end{split}$$

Since m(z) is the solution to

$$zcm^{2}(z) - (1 - c - z)m(z) + 1 = 0, (41)$$

we have

$$m'(z) = \frac{m^2(z)}{1 - \frac{cm^2(z)}{(1 + cm(z))^2}},$$
(42)

so that

$$h'(z) = m(z) + zm'(z).$$
 (43)

Further note that

$$\sum_{i=1}^{n} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \mathbf{a} \mathbf{a}^{\mathsf{H}} \mathbf{e}_{i} = \mathbf{a}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{a}. \tag{44}$$

8 Proof of ESPIRIT in the case of multiple DoAs: random signal case

In this case, we have

$$\mathbf{X} = (\mathbf{I}_N + \mathbf{A}\mathbf{P}\mathbf{A}^\mathsf{H})^{\frac{1}{2}}\mathbf{Z}, \quad \mathbf{A}\mathbf{P}\mathbf{A}^\mathsf{H} = \mathbf{U}\boldsymbol{\Lambda}\mathbf{U}^\mathsf{H}, \quad \mathbf{U} \in \mathbb{C}^{N \times K}, \quad \boldsymbol{\Lambda} = \operatorname{diag}\{\rho_i\}_{i=1}^K,$$
 (45)

for standard complex Gaussian $\mathbf{Z} \in \mathbb{C}^{N \times T}$.

Let us first consider the diagonal entries of the form $\hat{\mathbf{u}}_k^H \mathbf{J}_1^T \mathbf{J}_2 \hat{\mathbf{u}}_k$, with $\hat{\mathbf{u}}_k$ the kth dominant eigenvector of the SCM $\hat{\mathbf{C}} = \frac{1}{T} \mathbf{X} \mathbf{X}^H$. We have, for Γ circling around the isolated eigenvalue of $\hat{\mathbf{C}}$, that

$$\begin{split} \hat{\mathbf{u}}_{k}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \hat{\mathbf{u}}_{k} &= \sum_{i=1}^{n} \hat{\mathbf{u}}_{k}^{\mathsf{H}} \mathbf{e}_{i} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \hat{\mathbf{u}}_{k} \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\hat{\mathbf{C}} - z \mathbf{I}_{N} \right)^{-1} \mathbf{e}_{i} \, dz \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^{\mathsf{H}} - z \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-1} \right)^{-1} \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^{\mathsf{H}} - z \mathbf{I}_{N} + z \mathbf{A} (\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} \mathbf{A}^{\mathsf{H}} \right)^{-1} \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \left(\mathbf{Q}(z) - z \mathbf{Q}(z) \mathbf{A} \left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A} + z \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{A} \right)^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \right) \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \cdot z \mathbf{Q}(z) \mathbf{A} \left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A} + z \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{A} \right)^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \int_{\Gamma} z \operatorname{tr} \left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A} + z \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{A} \right)^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \int_{\Gamma} z \operatorname{tr} \left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A} + z \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{A} \right)^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \int_{\Gamma} z \operatorname{tr} \left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A} + z \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{A} \right)^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^{-\frac{1}{2}} \mathbf{e}_{i} \, dz, \\ &= -\frac{1}{2\pi i} \int_{\Gamma} z \operatorname{tr} \left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}} \mathbf{A} + z \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{A} \right)^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z) \left(\mathbf{I}_{N} + \mathbf{A} \mathbf{P} \mathbf{A}^{\mathsf{H}} \right)^$$

with the resolvent

$$\mathbf{Q}(z) \equiv \left(\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}} - z\mathbf{I}_{N}\right)^{-1}.\tag{46}$$

We have the following first and second deterministic equivalent results.

$$\mathbf{Q}(z) \leftrightarrow \bar{\mathbf{Q}}(z) = m(z)\mathbf{I}_N, \quad \mathbf{Q}(z)\mathbf{B}\mathbf{Q}(z) \leftrightarrow \bar{\mathbf{Q}}(z)\mathbf{B}\bar{\mathbf{Q}}(z) + \frac{1}{T}\operatorname{tr}(\mathbf{B})\frac{m'(z)m^2(z)}{(1+cm(z))^2}\mathbf{I}_N, \quad (47)$$

for any deterministic matrix $\mathbf{B} \in \mathbb{C}^{N \times N}$ of bounded operator norm. As such,

$$z\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{A} = zm(z)\mathbf{A}^{\mathsf{H}}\mathbf{A} + o_{\parallel \cdot \parallel}(1), \tag{48}$$

and

$$\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z) \left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}} \mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2} \left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}} \mathbf{Q}(z)\mathbf{A}$$

$$= m^{2}(z)\mathbf{A}^{\mathsf{H}} \left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}} \mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2} \left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}} \mathbf{A}$$

$$+ \frac{1}{T} \operatorname{tr} \left(\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2} \left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-1}\right) \frac{m'(z)m^{2}(z)}{(1 + cm(z))^{2}} \mathbf{A}^{\mathsf{H}}\mathbf{A} + o_{\parallel \cdot \parallel}(1),$$

[Zhenyu: Note that we have in general $\frac{1}{T}\operatorname{tr}\left(\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\left(\mathbf{I}_{N}+\mathbf{APA}^{\mathsf{H}}\right)^{-1}\right)=o(1)$ and $\mathbf{A}^{\mathsf{H}}\mathbf{A}=\mathbf{I}_{K}+o_{\|\cdot\|}(1)$. This leads to]

For the non-diagonal entries with $k \neq \ell$, we have instead

$$\begin{split} &|\mathring{\mathbf{u}}_{k}^{\mathsf{H}}J_{1}^{\mathsf{T}}J_{2}\mathring{\mathbf{u}}_{\ell}|^{2} = \mathring{\mathbf{u}}_{k}^{\mathsf{H}}J_{1}^{\mathsf{T}}J_{2}\mathring{\mathbf{u}}_{\ell}\mathring{\mathbf{u}}_{\ell}^{\mathsf{H}}J_{2}^{\mathsf{T}}J_{1}\mathring{\mathbf{u}}_{k} \\ &= -\frac{1}{2\pi\iota}\oint_{\Gamma_{1}}\mathring{\mathbf{u}}_{k}^{\mathsf{H}}J_{1}^{\mathsf{T}}J_{2}(\mathring{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}J_{2}^{\mathsf{T}}J_{1}\mathring{\mathbf{u}}_{k}\,dz_{1} \\ &= -\frac{1}{2\pi\iota}\oint_{\Gamma_{1}}\sum_{i=1}^{n}\mathring{\mathbf{u}}_{k}^{\mathsf{H}}e_{i}\mathbf{e}_{i+\Delta}^{\mathsf{T}}(\mathring{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}J_{2}^{\mathsf{T}}J_{1}\mathring{\mathbf{u}}_{k}\,dz_{1} \\ &= -\frac{1}{2\pi\iota}\oint_{\Gamma_{1}}\int_{\Gamma_{1}}\sum_{i=1}^{n}\mathbf{e}_{i+\Delta}^{\mathsf{T}}(\mathring{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}J_{2}^{\mathsf{T}}J_{1}\mathring{\mathbf{u}}_{k}\mathring{\mathbf{u}}_{k}^{\mathsf{H}}e_{i}\,dz_{1} \\ &= -\frac{1}{4\pi^{2}}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\sum_{i=1}^{n}\mathbf{e}_{i+\Delta}^{\mathsf{T}}(\mathring{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}J_{2}^{\mathsf{T}}J_{1}(\mathring{\mathbf{C}} - z_{2}\mathbf{I}_{N})^{-1}\mathbf{e}_{i}\,dz_{1}\,dz_{2} \\ &= -\frac{1}{4\pi^{2}}\sum_{i=1}^{n}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\mathbf{e}_{i+\Delta}^{\mathsf{T}}\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}\left(\frac{1}{T}\mathbf{Z}\mathbf{Z}^{\mathsf{H}} - z_{1}\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-1}\right)^{-1}\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{e}_{i}\,dz_{1}\,dz_{2} \\ &= -\frac{1}{4\pi^{2}}\sum_{i=1}^{n}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\mathbf{e}_{i+\Delta}^{\mathsf{T}}\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}z_{1}\mathbf{Q}(z_{1})\mathbf{A}\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}}\mathbf{A} + z_{1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{A}\right)^{-1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{1})\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}z_{2}\mathbf{Q}(z_{2})\mathbf{A}\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}}\mathbf{A} + z_{2}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{A}\right)^{-1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{I}_{2}^{\mathsf{T}}\mathbf{I}_{1} \\ &= -\frac{1}{4\pi^{2}}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}z_{1}z_{2}\operatorname{tr}\left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}}\mathbf{A} + z_{1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{A}\right)^{-1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{e}_{i}\,dz_{1}dz_{2} \\ &= -\frac{1}{4\pi^{2}}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}z_{1}z_{2}\operatorname{tr}\left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}}\mathbf{A} + z_{1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{A}\right)^{-1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)^{-\frac{1}{2}}\mathbf{e}_{i}\,dz_{1}\,dz_{2} \\ &= -\frac{1}{4\pi^{2}}\frac{1}{2}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}z_{1}z_{2}\operatorname{tr}\left(\left(\mathbf{P}^{-1} + \mathbf{A}^{\mathsf{H}}\mathbf{A} + z_{1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{A}\right)^{-1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\left(\mathbf{I}_{N} + \mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\right)$$

with

9 Proof of ESPIRIT in the case of multiple DoAs: deterministic signal case

In this case, we have

$$X = AS + Z, \tag{49}$$

with $\mathbf{X} = [\mathbf{x}(1), \dots, \mathbf{x}(T)] \in \mathbb{C}^{N \times T}$, $\mathbf{A} = [\mathbf{a}(\theta_1), \dots, \mathbf{a}(\theta_k)] \in \mathbb{C}^{N \times K}$, $\mathbf{S}^=[\mathbf{s}(1), \dots, \mathbf{s}(T)] \in \mathbb{C}^{K \times T}$, column vector $\mathbf{s}(t) = [s_1(t), \dots, s_k(t)]^\mathsf{T} \in \mathbb{C}^K$ and $\mathbf{Z} = [\mathbf{z}(1), \dots, \mathbf{z}(T)] \in \mathbb{C}^{N \times T}$ a standard (circular) Gaussian random matrix.

Let us first consider the form $\hat{\mathbf{u}}_k^\mathsf{H} \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \hat{\mathbf{u}}_k$, with $\hat{\mathbf{u}}$ the kth dominant eigenvector of the

SCM $\hat{\mathbf{C}} = \frac{1}{T}\mathbf{X}\mathbf{X}^{\mathsf{H}}$. We have, for Γ circling around the kth isolated eigenvalue of $\hat{\mathbf{C}}$, that

$$\begin{split} \hat{\mathbf{u}}_{k}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \hat{\mathbf{u}}_{k} &= \sum_{i=1}^{n} \hat{\mathbf{u}}_{k}^{\mathsf{H}} \mathbf{e}_{i} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \hat{\mathbf{u}}_{k} \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} (\hat{\mathbf{C}} - z \mathbf{I}_{N})^{-1} \mathbf{e}_{i} dz \\ &= -\frac{1}{2\pi i} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^{\mathsf{H}} - z \mathbf{I}_{N} + \mathbf{U} \Lambda \mathbf{U}^{\mathsf{H}} \right)^{-1} \mathbf{e}_{i} dz, \end{split}$$

with X = AS + Z and that

$$\mathbf{U} = \begin{bmatrix} \mathbf{A} & \frac{1}{T} \mathbf{Z} \mathbf{S}^{\mathsf{H}} \end{bmatrix} \in \mathbb{C}^{N \times 2K}, \quad \mathbf{\Lambda} = \begin{bmatrix} \mathbf{P} & \mathbf{I}_K \\ \mathbf{I}_K & \mathbf{0}_K \end{bmatrix} \in \mathbb{R}^{2K \times 2K}, \mathbf{P} = \lim \frac{1}{T} \mathbf{S} \mathbf{S}^{\mathsf{H}}. \tag{50}$$

As such, we get, by Woodbury identity that

$$\begin{split} \hat{\mathbf{u}}_{k}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \hat{\mathbf{u}}_{k} &= -\frac{1}{2\pi \iota} \sum_{i=1}^{n} \oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \left(\mathbf{Q}(z) - \mathbf{Q}(z) \mathbf{U} \left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{U} \right)^{-1} \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \right) \mathbf{e}_{i} \, dz, \\ &= \frac{1}{2\pi \iota} \sum_{i=1}^{n} \operatorname{tr} \left(\oint_{\Gamma} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \mathbf{Q}(z) \mathbf{U} \left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{U} \right)^{-1} \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{e}_{i} \, dz \right), \\ &= \frac{1}{2\pi \iota} \operatorname{tr} \left(\oint_{\Gamma} \left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{U} \right)^{-1} \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{Q}(z) \mathbf{U} \, dz \right), \end{split}$$

with

$$\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}} \mathbf{Q}(z) \mathbf{U} \right)^{-1} = \begin{bmatrix} (z + \frac{1}{m(z)}) (\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A}))^{-1} & (\mathbf{I}_{K} + (zm(z) + 1) \mathbf{P} \mathbf{A}^{\mathsf{H}} \mathbf{A}) \\ (\mathbf{I}_{K} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A} \mathbf{P}) & -m(z) ((\mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} + (zm(z) + 1) \mathbf{P})^{-1} \end{bmatrix}$$

$$\mathbf{Q}(z) \equiv \left(\frac{1}{T} \mathbf{Z} \mathbf{Z}^{\mathsf{H}} - z \mathbf{I}_{N} \right)^{-1} ,$$

$$(52)$$

the resolvent.

As such, we have

$$\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\mathbf{U} = \begin{bmatrix} \mathbf{A}^{\mathsf{H}} \\ \frac{1}{T}\mathbf{S}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}} \end{bmatrix}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\begin{bmatrix} \mathbf{A} & \frac{1}{T}\mathbf{Z}\mathbf{S} \end{bmatrix} \simeq \begin{bmatrix} \mathbf{A}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\mathbf{A} & \mathbf{0}_{K} \\ \mathbf{0}_{K} & \frac{1}{T}\mathbf{S}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\frac{1}{T}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\mathbf{J}_{2$$

with

$$\frac{1}{T}\mathbf{S}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\frac{1}{T}\mathbf{Z}\mathbf{S} = [\mathsf{Zhenyu}: \frac{1}{\mathsf{T}}\mathbf{S}^{\mathsf{H}}\tilde{\mathbf{Q}}(z)\frac{1}{\mathsf{T}}\mathbf{Z}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Z}\tilde{\mathbf{Q}}(z)\mathbf{S}] \tag{54}$$

for the co-resolvent

$$\tilde{\mathbf{Q}}(z) \equiv \left(\frac{1}{T}\mathbf{Z}^{\mathsf{H}}\mathbf{Z} - z\mathbf{I}_{T}\right)^{-1}.$$
 (55)

For the non-diagonal entries with $k \neq \ell$, we have instead

$$\begin{split} &|\hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\hat{\mathbf{u}}_{\ell}|^{2} &= \hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\hat{\mathbf{u}}_{\ell}\hat{\mathbf{u}}_{\ell}^{\mathsf{H}}\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\hat{\mathbf{u}}_{k} \\ &= -\frac{1}{2\pi\iota} \oint_{\Gamma_{1}} \hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}(\hat{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\hat{\mathbf{u}}_{k} dz_{1} \\ &= -\frac{1}{2\pi\iota} \oint_{\Gamma_{1}} \sum_{i=1}^{n} \hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{e}_{i}\mathbf{e}_{i+\Delta}^{\mathsf{T}}(\hat{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\hat{\mathbf{u}}_{k} dz_{1} \\ &= -\frac{1}{2\pi\iota} \oint_{\Gamma_{1}} \oint_{\Gamma_{1}} \sum_{i=1}^{n} \mathbf{e}_{i+\Delta}^{\mathsf{T}}(\hat{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\hat{\mathbf{u}}_{k}\hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{e}_{i} dz_{1} \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \sum_{i=1}^{n} \mathbf{e}_{i+\Delta}^{\mathsf{T}}(\hat{\mathbf{C}} - z_{1}\mathbf{I}_{N})^{-1}\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}(\hat{\mathbf{C}} - z_{2}\mathbf{I}_{N})^{-1}\mathbf{e}_{i} dz_{1} dz_{2} \\ &= -\frac{1}{4\pi^{2}} \sum_{i=1}^{n} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \mathbf{e}_{i+\Delta}^{\mathsf{T}}\left(\mathbf{Q}(z_{1}) - \mathbf{Q}(z_{1})\mathbf{U}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\right)\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\left(\mathbf{Q}(z_{2}) - \mathbf{Q}(z_{2})\mathbf{U}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{2})\mathbf{U}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{2})\mathbf{U}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{2})\mathbf{U}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}$$

with

$$\begin{split} \hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\hat{\mathbf{u}}_{k} &= -\frac{1}{2\pi\imath}\sum_{i=1}^{n}\oint_{\Gamma}\mathbf{e}_{i+\Delta}^{\mathsf{T}}\left(\mathbf{Q}(z)-\mathbf{Q}(z)\mathbf{U}\left(\boldsymbol{\Lambda}^{-1}+\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\right)\mathbf{e}_{i}\,dz,\\ &= \frac{1}{2\pi\imath}\sum_{i=1}^{n}\operatorname{tr}\left(\oint_{\Gamma}\mathbf{e}_{i+\Delta}^{\mathsf{T}}\mathbf{Q}(z)\mathbf{U}\left(\boldsymbol{\Lambda}^{-1}+\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{e}_{i}\,dz\right),\\ &= \frac{1}{2\pi\imath}\operatorname{tr}\left(\oint_{\Gamma}\left(\boldsymbol{\Lambda}^{-1}+\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\mathbf{U}\,dz\right),\\ &\left(\boldsymbol{\Lambda}^{-1}+\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{U}\right)^{-1}=\begin{bmatrix}(z+\frac{1}{m(z)})(\mathbf{P}^{-1}+(zm(z)+1)\mathbf{A}^{\mathsf{H}}\mathbf{A})^{-1} & (\mathbf{I}_{K}+(zm(z)+1)\mathbf{P}\mathbf{A}^{\mathsf{H}}\mathbf{A})^{-1}\\ &(\mathbf{I}_{K}+(zm(z)+1)\mathbf{A}^{\mathsf{H}}\mathbf{A}\mathbf{P})^{-1} & \frac{-m(z)\mathbf{P}^{-1}}{zm(z)+1}(\mathbf{I}_{K}-(\mathbf{I}_{K}+(zm(z)+1)\mathbf{P}\mathbf{A}^{\mathsf{H}}\mathbf{A})^{-1}\\ &\text{(56)}\\ \operatorname{set}\mathbf{H}_{1}=(z+\frac{1}{m(z)})(\mathbf{P}^{-1}+(zm(z)+1)\mathbf{A}^{\mathsf{H}}\mathbf{A})^{-1},\\ \mathbf{H}_{2}=\frac{-m(z)\mathbf{P}^{-1}}{zm(z)+1}(\mathbf{I}_{K}-(\mathbf{I}_{K}+(zm(z)+1)\mathbf{P}\mathbf{A}^{\mathsf{H}}\mathbf{A})^{-1}),\\ \operatorname{then}\\ &\left(\boldsymbol{\Lambda}^{-1}+\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z)\mathbf{U}=\begin{bmatrix}\mathbf{H}_{1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}\mathbf{A}\\ \mathbf{0}_{K} & \mathbf{H}_{2}\frac{1}{\pi}\mathbf{S}^{\mathsf{H}}\mathbf{Z}^{\mathsf{H}}\mathbf{Q}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}\frac{1}{\pi}\mathbf{Z}\mathbf{S}\end{bmatrix} \end{split}$$

then

$$\hat{\mathbf{u}}_k^\mathsf{H} \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \hat{\mathbf{u}}_k = \frac{1}{2\pi \imath} \oint_{\Gamma} \operatorname{tr} \left(\mathsf{H}_1 \mathbf{A}^\mathsf{H} \mathbf{Q} \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \mathbf{Q} \mathbf{A} \right) + \operatorname{tr} \left(\mathsf{H}_2 \frac{1}{T} \mathbf{S}^\mathsf{H} \mathbf{Z}^\mathsf{H} \mathbf{Q} \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \mathbf{Q} \frac{1}{T} \mathbf{Z} \mathbf{S} \right) \, dz.$$

where

$$\begin{split} &\frac{1}{2\pi\iota} \oint_{\Gamma} \operatorname{tr} \left(\mathsf{H}_{1} \mathbf{A}^{\mathsf{H}} \mathbf{Q} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{Q} \mathbf{A} \right) dz = \frac{1}{2\pi\iota} \oint_{\Gamma} \operatorname{tr} \left((z + \frac{1}{m(z)}) (\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{Q} \mathbf{A} \right) dz \\ &= \frac{1}{2\pi\iota} \oint_{\Gamma} (zm^{2}(z) + m(z)) \operatorname{tr} \left((\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{A} \right) dz \\ &= \frac{1}{2\pi\iota} \sum_{i=1}^{l+n-1} \oint_{\Gamma} (zm^{2}(z) + m(z)) \mathbf{e}_{i+\Delta}^{\mathsf{T}} \mathbf{A} (\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{e}_{i} dz \\ &= \frac{1}{2\pi\iota} \sum_{i=1}^{l+n-1} \sum_{j=1}^{K} \oint_{\Gamma} \frac{zm^{2}(z) + m(z)}{zm(z) + 1 + l_{j}^{-1}} \mathbf{e}_{i+\Delta}^{\mathsf{T}} \mathbf{A} \mathbf{u}_{j} \mathbf{u}_{j}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{e}_{i} dz \\ &= \lim_{z \to \bar{\lambda}_{k}} \frac{(z - \bar{\lambda}_{k}) (zm^{2}(z) + m(z))}{zm(z) + 1 + l_{k}^{-1}} \operatorname{tr} \left(\mathbf{A} \mathbf{u}_{k} \mathbf{u}_{k}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \right) \\ &= \frac{1 - cl_{k}^{-2}}{1 + cl_{k}^{-1}} \mathbf{u}_{k}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{A} \mathbf{u}_{k} \end{split}$$

where we consider the eigendecomposition of $PA^{H}A = UVU^{H}$.

$$\frac{1}{2\pi \iota} \oint_{\Gamma} \operatorname{tr}\left(\mathsf{H}_{2} \frac{1}{T} \mathbf{S}^{\mathsf{H}} \mathbf{Z}^{\mathsf{H}} \mathbf{Q} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{Q} \frac{1}{T} \mathbf{Z} \mathbf{S}\right) dz = 0?$$

so we have $\hat{\mathbf{u}}_k^H \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \hat{\mathbf{u}}_k = -\frac{1-cl_k^{-2}}{1+cl_k^{-1}} \mathbf{u}_k^H \mathbf{A}^H \mathbf{J}_1^T \mathbf{J}_2 \mathbf{A} \mathbf{u}_k + o(1)$. For the non-diagonal entries with $k \neq \ell$, we have instead

$$\begin{split} &|\hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\hat{\mathbf{u}}_{\ell}|^{2} = -\frac{1}{4\pi^{2}}\operatorname{tr}\left(\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{2})\mathbf{U}\left(\mathbf{\Lambda} + \mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{U}\right)^{-1}\mathbf{U}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{J}_{1}^{\mathsf{T}}\\ &= -\frac{1}{4\pi^{2}}\operatorname{tr}\left(\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\mathbf{H}_{1}(z_{1})\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{2})\mathbf{A}\mathbf{H}_{1}(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{Q}(z_{1})\mathbf{A}\,dz_{1}\,dz_{2}\right)\\ &= -\frac{1}{4\pi^{2}}\sum_{i=l}^{l+n-1}\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\sum_{n=1}^{K}\sum_{m=1}^{K}\sum_{m=1}^{K}\frac{zm^{2}(z_{1})+m(z_{1})}{zm(z_{1})+1+l_{m}^{-1}}\cdot\frac{zm^{2}(z_{2})+m(z_{2})}{zm(z_{2})+1+l_{n}^{-1}}\mathbf{e}_{i}^{T}\mathbf{A}\mathbf{u}_{n}\mathbf{u}_{n}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{T}\mathbf{J}_{2}\mathbf{A}\mathbf{u}_{m}\mathbf{u}_{m}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{e}_{i+\Delta}\\ &=\frac{1-cl_{k}^{-2}}{1+cl_{k}^{-1}}\cdot\frac{1-cl_{l}^{-2}}{1+cl_{l}^{-1}}\mathbf{u}_{l}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{T}\mathbf{J}_{2}\mathbf{A}\mathbf{u}_{k}\mathbf{u}_{k}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{J}_{2}^{T}\mathbf{J}_{1}\mathbf{A}\mathbf{u}_{l}. \end{split}$$

The conclusion about the diagonal elements and the nondiagonal elements are the same as $\mathbf{\Phi}_2$ in the case of random signals.

$$\begin{split} \hat{\mathbf{u}}_k^\mathsf{H} \mathbf{J}_1^\mathsf{T} \mathbf{J}_1 \hat{\mathbf{u}}_k &= -\frac{1}{2\pi \iota} \sum_{i=1}^n \oint_\Gamma \mathbf{e}_{i+\Delta}^\mathsf{T} \left(\mathbf{Q}(z) - \mathbf{Q}(z) \mathbf{U} \left(\mathbf{\Lambda}^{-1} + \mathbf{U}^\mathsf{H} \mathbf{Q}(z) \mathbf{U} \right)^{-1} \mathbf{U}^\mathsf{H} \mathbf{Q}(z) \right) \mathbf{e}_i \, dz, \\ &= \frac{1}{2\pi \iota} \sum_{i=1}^n \operatorname{tr} \left(\oint_\Gamma \mathbf{e}_{i+\Delta}^\mathsf{T} \mathbf{Q}(z) \mathbf{U} \left(\mathbf{\Lambda}^{-1} + \mathbf{U}^\mathsf{H} \mathbf{Q}(z) \mathbf{U} \right)^{-1} \mathbf{U}^\mathsf{H} \mathbf{Q}(z) \mathbf{e}_i \, dz \right), \\ &= \frac{1}{2\pi \iota} \operatorname{tr} \left(\oint_\Gamma \left(\mathbf{\Lambda}^{-1} + \mathbf{U}^\mathsf{H} \mathbf{Q}(z) \mathbf{U} \right)^{-1} \mathbf{U}^\mathsf{H} \mathbf{Q}(z) \mathbf{J}_1^\mathsf{T} \mathbf{J}_1 \mathbf{Q}(z) \mathbf{U} \, dz \right), \end{split}$$

Similarly,

$$\begin{split} &\frac{1}{2\pi \iota} \oint_{\Gamma} \operatorname{tr} \left(\mathsf{H}_{1} \mathbf{A}^{\mathsf{H}} \mathbf{Q} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{Q} \mathbf{A} \right) dz = \frac{1}{2\pi \iota} \oint_{\Gamma} \operatorname{tr} \left((z + \frac{1}{m(z)}) (\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} \mathbf{A}^{\mathsf{H}} \mathbf{Q} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{Q} \mathbf{A} \right) dz \\ &= \frac{1}{2\pi \iota} \oint_{\Gamma} \operatorname{tr} \left((z + \frac{1}{m(z)}) (\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} [m^{2}(z) \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{A} + \frac{n}{T} \cdot \frac{m'(z)m^{2}(z)}{(1 + cm(z))^{2}} \mathbf{I}_{K}] \right) dz \\ &= \frac{1 - cl_{k}^{-2}}{1 + cl_{k}^{-1}} \mathbf{u}_{k}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{A} \mathbf{u}_{k} + \frac{1}{2\pi \iota} \oint_{\Gamma} \operatorname{tr} \left(\frac{n}{T} \frac{(zm(z) + 1)m'(z)m(z)}{(1 + cm(z))^{2}} (\mathbf{P}^{-1} + (zm(z) + 1) \mathbf{A}^{\mathsf{H}} \mathbf{A})^{-1} \right) dz \\ &= \frac{1 - cl_{k}^{-2}}{1 + cl_{k}^{-1}} \mathbf{u}_{k}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{A} \mathbf{u}_{k} + \lim_{z \to \bar{\lambda}_{k}} \oint_{\Gamma} \sum_{i=1}^{K} \frac{n}{T} \frac{(zm(z) + 1)m'(z)m(z)}{(1 + cm(z))^{2}} \cdot \frac{\mathbf{u}_{i}^{\mathsf{H}} \mathbf{u}_{i}}{zm(z) + 1 + l_{i}^{-1}} dz \\ &= \frac{1 - cl_{k}^{-2}}{1 + cl_{k}^{-1}} \mathbf{u}_{k}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{A} \mathbf{u}_{k} + \lim_{z \to \bar{\lambda}_{k}} \frac{n}{T} \frac{(z - \bar{\lambda}_{k})(zm(z) + 1)m'(z)m(z)}{(1 + cm(z))^{2}(zm(z) + 1 + l_{k}^{-1})} \\ &= \frac{1 - cl_{k}^{-2}}{1 + cl_{k}^{-1}} \mathbf{u}_{k}^{\mathsf{H}} \mathbf{A}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{A} \mathbf{u}_{k} + \frac{n}{T} \frac{1}{l_{k}^{2} - cl_{k}} + o(1)(?????) \end{split}$$

Similarly for off-diagonal entries:

$$\begin{split} |\hat{\mathbf{u}}_{k}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\hat{\mathbf{u}}_{\ell}|^{2} &= -\frac{1}{4\pi^{2}}\operatorname{tr}\left(\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\mathsf{H}_{1}(z_{1})\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{1})\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{2})\mathbf{A}\mathsf{H}_{1}(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z_{2})\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{Q}(z_{1})\mathbf{A}\,dz_{1}\,dz_{2}\right) \\ &= -\frac{1}{4\pi^{2}}\operatorname{tr}\left(\oint_{\Gamma_{1}}\oint_{\Gamma_{2}}\mathsf{H}_{1}(z_{1})(m(z_{1})m(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A} + \eta(z_{1},z_{2})\frac{n}{T}\mathbf{I}_{K})\mathsf{H}_{1}(z_{2})(m(z_{1})m(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A} + \eta(z_{1},z_{2})\frac{n}{T}\mathbf{I}_{K})\mathsf{H}_{1}(z_{2})(m(z_{1})m(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A} + \eta(z_{1},z_{2})\frac{n}{T}\mathbf{I}_{K})\mathsf{H}_{2}(z_{1})(m(z_{1})m(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A} + \eta(z_{1},z_{2})\frac{n}{T}\mathbf{I}_{K})\mathsf{H}_{2}(z_{1})(m(z_{1})m(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A} + \eta(z_{1},z_{2})\frac{n}{T}\mathbf{I}_{K})\mathsf{H}_{2}(z_{1})(m(z_{1})m(z_{2})\mathbf{J}_{1}^{\mathsf{H}}$$

10 Proof of ESPIRIT in the case of multiple DoAs: uncorrelated signal case

$$\mathbf{T}_{1} = \mathbf{P}^{-1} + (1 + zm(z))\mathbf{I}_{K} + o_{\|.\|}(1), \quad \mathbf{T}_{2} = m^{2}(z)\mathbf{A}^{\mathsf{H}}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{A} + o_{\|.\|}(1).$$
 (57)

Considering the spectral decomposition $\mathbf{P} = \mathbf{U}L\mathbf{U}^H$ with $L = \text{diag}\{l_i, \dots, l_k\}$ and $\mathbf{U} = [\mathbf{u}_1, \dots, \mathbf{u}_k]$. Then we have

$$\begin{aligned} (\mathbf{T}_1)^{-1} &= \frac{1}{1 + zm(z)} \mathbf{I}_K - \frac{1}{(1 + zm(z))^2} \mathbf{U} (\mathbf{I}_K + \frac{1}{1 + zm(z)} \mathbf{L}^{-1})^{-1} \mathbf{U}^H \\ &= \frac{1}{1 + zm(z)} \mathbf{I}_K - \sum_{i=1}^K \frac{1}{(1 + zm(z))^2} \mathbf{u}_i \frac{(1 + zm(z))l_i^{-1}}{1 + zm(z) + l_i^{-1}} \mathbf{u}_i^H \end{aligned}$$

Then

$$\begin{split} [\mathbf{\Phi}_{2}]_{11} &= -\frac{1}{2\pi \iota} \oint_{\Gamma_{1}} z \operatorname{tr}(\mathbf{T}_{1}^{-1}\mathbf{T}_{2}) dz \\ &= \frac{1}{2\pi \iota} \oint_{\Gamma_{1}} \sum_{i=1}^{K} \frac{zm^{2}(z)}{1 + zm(z)} \frac{l_{i}^{-1}}{1 + zm(z)l_{i}^{-1}} \operatorname{tr}(\mathbf{u}_{i}\mathbf{A}^{\mathsf{H}}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{A}\mathbf{u}_{i}^{H}) dz \\ &= \lim_{z \to \bar{\lambda}_{1}} \frac{(z - \bar{\lambda}_{1})zm^{2}(z)}{1 + zm(z)} \frac{l_{1}^{-1}}{1 + zm(z) + l_{1}^{-1}} \mathbf{u}_{1}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{A}\mathbf{u}_{1} \\ &= \lim_{z \to \bar{\lambda}_{1}} \frac{(z - \bar{\lambda}_{1})zm^{2}(z)}{1 + zm(z) + l_{1}^{-1}} \mathbf{u}_{1}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{A}\mathbf{u}_{1} + o(1) \\ &= \frac{1 - cl_{1}^{-2}}{1 + cl_{1}^{-1}} \cdot \mathbf{u}_{1}^{H}\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{A}\mathbf{u}_{1} + o(1) \end{split}$$

The conclusion is similar to the previous formula, except that $\mathbf{a}(\theta_1)$ is replaced by eigenvector $\mathbf{u}_i^H \mathbf{a}(\theta_1)^H$.

$$\mathbf{T}_{2}(z_{1}, z_{2}) = m(z_{1})m(z_{2})\mathbf{A}^{\mathsf{H}}\mathbf{B}\mathbf{A} + o_{\|\cdot\|}(1)$$

where $\mathbf{B} = \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \mathbf{C}^{-\frac{1}{2}} \in \mathbb{C}^{N \times N}$. Similarly,

$$\begin{split} &|[\boldsymbol{\Phi}_{2}]_{12}|^{2} = -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} z_{1}z_{2} \operatorname{tr}(\mathbf{T}_{2}^{\mathsf{H}}(z_{1}, z_{2})\mathbf{T}_{1}^{-1}(z_{1})\mathbf{T}_{2}(z_{1}, z_{2})\mathbf{T}_{1}^{-1}(z_{2})) \, dz_{1} \, dz_{2}, \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \sum_{i=1}^{K} \sum_{j=1}^{K} \frac{z_{1}m^{2}(z_{1})l_{i}^{-1}z_{2}m^{2}(z_{2})l_{j}^{-1} \operatorname{tr}(\mathbf{u}_{j}\mathbf{u}_{j}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{B}^{\mathsf{H}}\mathbf{A}\mathbf{u}_{i}\mathbf{u}_{i}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{B}\mathbf{A}) \, dz_{1}dz_{2}}{(1+z_{1}m(z_{1}))(1+z_{1}m(z_{1})+l_{i}^{-1})(1+z_{2}m(z_{2}))(1+z_{2}m(z_{2})+l_{j}^{-1})} \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \frac{z_{1}m^{2}(z_{1})l_{1}^{-1}z_{2}m^{2}(z_{2})l_{2}^{-1}\mathbf{u}_{2}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{B}^{\mathsf{H}}\mathbf{A}\mathbf{u}_{1}\mathbf{u}_{1}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{B}\mathbf{A}\mathbf{u}_{2} \, dz_{1}dz_{2}}{(1+z_{1}m(z_{1}))(1+z_{1}m(z_{1})+l_{1}^{-1})(1+z_{2}m(z_{2}))(1+z_{2}m(z_{2})+l_{2}^{-1})} \\ &= \frac{1-cl_{1}^{-2}}{1+cl_{1}^{-1}} \cdot \frac{1-cl_{2}^{-2}}{1+cl_{2}^{-1}}\mathbf{u}_{2}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{J}_{2}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A}\mathbf{u}_{1}\mathbf{u}_{1}^{\mathsf{H}}\mathbf{A}^{\mathsf{H}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{A}\mathbf{u}_{2}. \end{split}$$

where we use the fact that $\mathbf{u}_2^H \mathbf{A}^H \mathbf{B}^H \mathbf{A} \mathbf{u}_1 \mathbf{u}_1^H \mathbf{A}^H \mathbf{A} \mathbf{u}_2 = \bar{\lambda}_1 \bar{\lambda}_2 \mathbf{u}_2^H \mathbf{A}^H \mathbf{J}_2^T \mathbf{J}_1 \mathbf{A} \mathbf{u}_1 \mathbf{u}_1^H \mathbf{A}^H \mathbf{J}_1^T \mathbf{J}_2 \mathbf{A} \mathbf{u}_2$ in the last line.

$$\mathbf{T}_{3} = m^{2}(z)\mathbf{A}^{\mathsf{H}}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{C}^{-\frac{1}{2}}\mathbf{A} + \frac{n}{T}\frac{m^{4}(z)\mathbf{I}_{K}}{(1+cm(z))^{2}-cm^{2}(z)} + o_{\|\cdot\|}(1)$$

As such, we have

$$\begin{split} &[\mathbf{\Phi}_{1}]_{11} = -\frac{1}{2\pi \iota} \oint_{\Gamma_{1}} z \operatorname{tr}(\mathbf{T}_{1}^{-1}\mathbf{T}_{3}) dz \\ &= \frac{1}{2\pi \iota} \oint_{\Gamma_{1}} \sum_{i=1}^{K} \frac{zm^{2}(z)}{1 + zm(z)} \frac{l_{i}^{-1}}{1 + zm(z)l_{i}^{-1}} \operatorname{tr}(\mathbf{u}_{i}\mathbf{u}_{i}^{H}\mathbf{A}^{H}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{C}^{-\frac{1}{2}}\mathbf{A}) \\ &+ z \operatorname{tr}(\sum_{i=1}^{K} \frac{zm^{2}(z)}{1 + zm(z)} \frac{l_{i}^{-1}}{1 + zm(z)l_{i}^{-1}} \frac{n}{T} \frac{m^{4}(z)}{(1 + cm(z))^{2} - cm^{2}(z)} \cdot \mathbf{u}_{i}\mathbf{u}_{i}^{H}) dz \\ &= \frac{1 - cl_{1}^{-2}}{1 + cl_{1}^{-1}} \cdot \mathbf{u}_{1}^{H}\mathbf{A}^{H}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{1}\mathbf{A}\mathbf{u}_{1} + \frac{n}{T} \frac{1 + \rho_{1}^{-1}}{c + \rho_{1}} + o(1) \\ &= \frac{1 - cl_{1}^{-2}}{1 + cl_{1}^{-1}} \cdot \frac{n}{N} + \frac{n}{T} \frac{1 + \rho_{1}^{-1}}{c + \rho_{1}} + o(1). \end{split}$$

Similarly for off-diagonal entries of Φ_1 as

$$\begin{split} &|[\boldsymbol{\Phi}_{1}]_{12}|^{2} = -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} z_{1}z_{2} \operatorname{tr}(\mathbf{T}_{4}(z_{2}, z_{1})\mathbf{T}_{1}^{-1}(z_{1})\mathbf{T}_{4}(z_{1}, z_{2})\mathbf{T}_{1}^{-1}(z_{2})) \, dz_{1} \, dz_{2}, \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \operatorname{tr}(\sum_{i=1}^{K} \sum_{j=1}^{K} \frac{z_{1}m^{2}(z_{1})l_{i}^{-1}z_{2}m^{2}(z_{2})l_{j}^{-1} \operatorname{tr}(\mathbf{T}_{4}(z_{2}, z_{1})u_{j}u_{j}^{H}\mathbf{T}_{4}(z_{1}, z_{2})u_{i}u_{i}^{H}) \quad dz_{1}dz_{2}}{(1 + z_{1}m(z_{1}))(1 + z_{1}m(z_{1}) + l_{i}^{-1})(1 + z_{2}m(z_{2}))(1 + z_{2}m(z_{2}) + l_{j}^{-1})} \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \frac{z_{1}m^{2}(z_{1})l_{1}^{-1}z_{2}m^{2}(z_{2})l_{2}^{-1}\mathbf{u}_{2}^{H}\mathbf{A}^{H}\mathbf{H}^{H}\mathbf{A}\mathbf{u}_{1}\mathbf{u}_{1}^{H}\mathbf{A}^{H}\mathbf{H}\mathbf{A}\mathbf{u}_{2} \quad dz_{1}dz_{2}}{(1 + z_{1}m(z_{1}))(1 + z_{1}m(z_{1}) + l_{1}^{-1})(1 + z_{2}m(z_{2}))(1 + z_{2}m(z_{2}) + l_{2}^{-1})} \\ &= \frac{1 - cl_{1}^{-2}}{1 + cl_{1}^{-1}} \cdot \frac{1 - cl_{2}^{-2}}{1 + cl_{2}^{-1}}\mathbf{u}_{2}^{H}\mathbf{A}^{H}\mathbf{J}_{1}^{T}\mathbf{J}_{1}\mathbf{A}\mathbf{u}_{1}\mathbf{u}_{1}^{H}\mathbf{A}^{H}\mathbf{J}_{1}^{T}\mathbf{J}_{1}\mathbf{A}\mathbf{u}_{2} + o(1) \\ &= o(1) \end{split}$$

with

$$\begin{split} \mathbf{T}_{4}(z_{1},z_{2}) &= \mathbf{A}^{\mathsf{H}} \mathbf{Q}(z_{1}) \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{C}^{-\frac{1}{2}} \mathbf{Q}(z_{2}) \mathbf{A} \\ &= m(z_{1}) m(z_{2}) \mathbf{A}^{\mathsf{H}} \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{C}^{-\frac{1}{2}} \mathbf{A} + \frac{n}{T} \eta(z_{1},z_{2}) \mathbf{I}_{K} + o_{\|\cdot\|}(1) = \mathbf{T}_{4}(z_{2},z_{1}), \end{split}$$

and $\mathbf{H} = \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_1^\mathsf{T} \mathbf{J}_1 \mathbf{C}^{-\frac{1}{2}} \in \mathbb{C}^{N \times N}$. The last line is because $\mathbf{u}_2^H \mathbf{A}^H \mathbf{J}_1^T \mathbf{J}_1 \mathbf{A} \mathbf{u}_1 = \mathbf{u}_2^H \frac{n}{N} \mathbf{I}_K \mathbf{u}_1 = 0$. So we can draw the conclusion that:

$$\bar{\mathbf{\Phi}}_{1} = \operatorname{diag} \left\{ \frac{n}{N} \frac{1 - cl_{k}^{-2}}{1 + cl_{k}^{-1}} + \frac{n}{N} \frac{1 + l_{k}^{-1}}{1 + l_{k}/c} \right\}_{k=1}^{K},$$
 (58)

$$[\bar{\mathbf{\Phi}}_2]_{kk} = \frac{1 - cl_1^{-2}}{1 + cl_1^{-1}} \cdot \mathbf{u}_k^H \mathbf{A}^H \mathbf{J}_1^\mathsf{T} \mathbf{J}_2 \mathbf{A} \mathbf{u}_k$$
 (59)

However, we can only get the modulus of the off-diagonal elements of $\bar{\Phi}_2$.

$$|[\mathbf{\Phi}_2]_{mn}|^2 = \frac{1 - cl_m^{-2}}{1 + cl_m^{-1}} \cdot \frac{1 - cl_n^{-2}}{1 + cl_n^{-1}} \mathbf{u}_m^H \mathbf{A}^H \mathbf{J}_2^T \mathbf{J}_1 \mathbf{A} \mathbf{u}_n \mathbf{u}_n^H \mathbf{A}^H \mathbf{J}_1^T \mathbf{J}_2 \mathbf{A} \mathbf{u}_m.$$

11 Closely DoA

Settings: $\mathbf{a}(\theta) = \frac{1}{\sqrt{N}} [1, e^{i\theta}, \dots, e^{i(N-1)\theta}]^{\mathsf{T}}.$

We set that $\theta_2 = \theta_1 + \frac{\alpha}{T}$, where *T* is the number of time slots, then:

$$\begin{split} \mathbf{a}^{\mathsf{H}}(\theta_1)\mathbf{a}(\theta_2) &= \frac{1}{N} \sum_{n=0}^{N-1} e^{\imath n\alpha/T} = \frac{1}{N} \frac{1 - e^{\imath N\alpha/T}}{1 - e^{\imath \alpha/T}} = \lim_{N,T \to \infty} \frac{\frac{1}{T}(1 - e^{\imath N\alpha/T})}{\frac{N}{T}(1 - e^{\imath \alpha/T})} \\ &= \lim_{N,T \to \infty} \frac{-\frac{1}{T^2}(1 - e^{\imath \alpha c})}{\frac{1}{T^2}\imath \alpha c e^{\imath \frac{\alpha}{T}}} = \imath \frac{1 - e^{\imath c\alpha}}{c\alpha} \end{split}$$

where the *j*-th entry of $\mathbf{a}(\theta)$ is given by

$$[\mathbf{a}(\theta_{\ell})]_{j} = \frac{1}{\sqrt{N}} e^{i\frac{2\pi d}{\lambda_{0}}(j-1)\theta_{\ell}} \equiv \frac{1}{\sqrt{N}} e^{i\omega(j-1)\theta_{\ell}}, \quad \omega \equiv \frac{2\pi d}{\lambda_{0}}$$
(60)

So we have:

$$\mathbf{A}^{\mathsf{H}}\mathbf{A} \to \begin{bmatrix} 1 & e^{i\alpha c/2}sinc(\alpha c/2) \\ e^{-i\alpha c/2}sinc(\alpha c/2) & 1 \end{bmatrix}$$

with its eigenvalue decomposition $\mathbf{V}L\mathbf{V}^H$ with $\Sigma = \mathrm{diag}\{1 + \mathrm{sinc}\frac{\alpha c}{2}, 1 - \mathrm{sinc}\frac{\alpha c}{2}\}$ and $\mathbf{U} = [\mathbf{v}_1, \mathbf{v}_2]$, where $\mathbf{v}_1 = \frac{1}{\sqrt{2}}[e^{\frac{i\alpha c}{2}}, 1]^T$ and $\mathbf{v}_2 = \frac{1}{\sqrt{2}}[1, -e^{-\frac{i\alpha c}{2}}]^T$. In this part, we have:

$$T_1 = P^{-1} + A^H A + z A^H Q(z) A = P^{-1} + (1_z m(z)) A^H A,$$
 (61)

and

$$\mathbf{T}_1^{-1} = \mathbf{P} - (1 + zm(z))\mathbf{P}\mathbf{A}^{\mathsf{H}}(\mathbf{I}_N + (1 + zm(z))\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}})^{-1}\mathbf{A}\mathbf{P}.$$

Then, we have

$$\begin{split} [\mathbf{\Phi}_{2}]_{kk} &= \frac{1}{2\pi \iota} \oint_{\Gamma_{k}} z \operatorname{tr}(\mathbf{T}_{1}^{-1}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{C}^{\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{\frac{1}{2}}\mathbf{Q}(z)\mathbf{A}) dz \\ &= -\frac{1}{2\pi \iota} \oint_{\Gamma_{k}} z \operatorname{tr}((1+zm(z))\mathbf{P}\mathbf{A}^{\mathsf{H}}(\mathbf{I}_{N}+(1+zm(z))\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}})^{-1}\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\mathbf{Q}(z)\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{Q}(z)\mathbf{A}) dz \\ &= -\frac{1}{2\pi \iota} \oint_{\Gamma_{k}} z \operatorname{tr}(\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}(\mathbf{I}_{N}-(\mathbf{I}_{N}+(1+zm(z))\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}})^{-1})\mathbf{Q}(z)\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{Q}(z)) dz \\ &= \frac{1}{2\pi \iota} \oint_{\Gamma_{k}} z \operatorname{tr}(\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}(\mathbf{I}_{N}+(1+zm(z))\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}})^{-1}\mathbf{Q}(z)\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{Q}(z)) dz \end{split}$$

we define the spectral decomposition $\mathbf{APA}^{\mathsf{H}} = \mathbf{U}L\mathbf{U}^{H}$ with $L = \mathrm{diag}\{\rho_{i}, \cdots, \rho_{N}\}$ and $\mathbf{U} = [\mathbf{u}_{1}, \cdots, \mathbf{u}_{N}]$. Then:

$$\begin{split} [\mathbf{\Phi}_{2}]_{kk} &= \frac{1}{2\pi i} \oint_{\Gamma_{k}} \sum_{i=1}^{N} \frac{z}{1 + (1 + zm(z)\rho_{i})} \operatorname{tr}(\mathbf{A}\mathbf{P}\mathbf{A}^{\mathsf{H}}\mathbf{u}_{i}\mathbf{u}_{i}^{H}\mathbf{Q}(z)\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{Q}(z)) dz \\ &= \frac{1}{2\pi i} \oint_{\Gamma_{k}} \sum_{i=1}^{N} \frac{z\rho_{k}}{1 + (1 + zm(z)\rho_{i})} \mathbf{u}_{k}^{H}\mathbf{Q}(z)\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{Q}(z)\mathbf{u}_{k} dz \\ &= \frac{1}{2\pi i} \oint_{\Gamma_{k}} \frac{zm^{2}(z)}{\rho_{i}^{-1} + 1 + zm(z)} \mathbf{u}_{i}^{H}\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{C}^{-\frac{1}{2}}\mathbf{u}_{i} dz \\ &= \frac{1 - c\rho_{k}^{-2}}{1 + c\rho_{k}^{-1}} \cdot \mathbf{u}_{k}^{H}\mathbf{J}_{1}^{\mathsf{T}}\mathbf{J}_{2}\mathbf{u}_{k} + o(1) \end{split}$$

the conclusion is the same as $[\Phi_2]_{kk}$ in the case of widely DoA. Similarly for $[\mathbf{\Phi}_1]_{kk}$ as:

$$\begin{split} & [\boldsymbol{\Phi}_1]_{kk} = \frac{1}{2\pi \imath} \oint_{\Gamma_k} \frac{z}{\rho_k^{-1} + 1 + zm(z)} \mathbf{u}_k^H \mathbf{Q}(z) \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_1^\mathsf{T} \mathbf{J}_1 \mathbf{C}^{-\frac{1}{2}} \mathbf{Q}(z) \mathbf{u}_k \, dz \\ & = \frac{1 - c\rho_k^{-2}}{1 + c\rho_k^{-1}} \cdot \mathbf{u}_K^H \mathbf{J}_1^\mathsf{T} \mathbf{J}_1 \mathbf{u}_k + \frac{1}{2\pi \imath} \oint_{\Gamma_k} \frac{z}{\rho_k^{-1} + 1 + zm(z)} \frac{n}{T} \frac{m'(z)m^2(z)}{(1 + cm(z))^2} \, dz \\ & = \frac{1 - c\rho_k^{-2}}{1 + c\rho_k^{-1}} \cdot \mathbf{u}_K^H \mathbf{J}_1^\mathsf{T} \mathbf{J}_1 \mathbf{u}_k + \frac{n}{T} \frac{1 + \rho_k^{-1}}{c + \rho_k} + o(1). \end{split}$$

where \mathbf{u}_k and ρ_k represent the k-th eigenvetor and eigenvalue of \mathbf{APA}^H , respectively. The extra bias term is the same as the previous conclusions.

For the off-diagonal term:

$$\begin{split} |[\mathbf{\Phi}_{2}]_{12}|^{2} &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} z_{1}z_{2} \operatorname{tr}(\mathbf{T}_{2}^{\mathsf{H}}(z_{1}, z_{2}) \mathbf{T}_{1}^{-1}(z_{1}) \mathbf{T}_{2}(z_{1}, z_{2}) \mathbf{T}_{1}^{-1}(z_{2})) \, dz_{1} \, dz_{2}, \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \sum_{i=1}^{K} \sum_{j=1}^{K} \frac{z_{1}m^{2}(z_{1})z_{2}m^{2}(z_{2}) \mathbf{u}_{j}^{H} \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_{2}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{C}^{-\frac{1}{2}} \mathbf{u}_{j} \mathbf{u}_{i}^{H} \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{C}^{-\frac{1}{2}} \mathbf{u}_{i} \quad dz_{1} dz_{2}}{(1 + z_{1}m(z_{1}) + \rho_{i}^{-1})(1 + z_{2}m(z_{2}) + \rho_{j}^{-1})} \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \frac{z_{1}m^{2}(z_{1})l_{1}^{-1}z_{2}m^{2}(z_{2})l_{2}^{-1} \mathbf{u}_{1}^{H} \mathbf{J}_{2}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{u}_{2} \mathbf{u}_{2}^{H} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{u}_{1} \quad dz_{1} dz_{2}}{(1 + z_{1}m(z_{1}) + \rho_{1}^{-1})(1 + z_{2}m(z_{2}) + \rho_{2}^{-1})} \\ &= \frac{1 - c\rho_{1}^{-2}}{1 + c\rho_{1}^{-1}} \cdot \frac{1 - c\rho_{2}^{-2}}{1 + c\rho_{2}^{-1}} \mathbf{u}_{1}^{H} \mathbf{J}_{2}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{u}_{2} \mathbf{u}_{2}^{H} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{u}_{1}. \end{split}$$

Similarly for $[\Phi_1]_{12}$ as:

$$\begin{split} |[\mathbf{\Phi}_{1}]_{12}|^{2} &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} z_{1}z_{2} \operatorname{tr}(\mathbf{T}_{3}^{\mathsf{H}}(z_{1}, z_{2}) \mathbf{T}_{1}^{-1}(z_{1}) \mathbf{T}_{3}(z_{1}, z_{2}) \mathbf{T}_{1}^{-1}(z_{2})) dz_{1} dz_{2}, \\ &= -\frac{1}{4\pi^{2}} \oint_{\Gamma_{1}} \oint_{\Gamma_{2}} \sum_{i=1}^{K} \sum_{j=1}^{K} \frac{z_{1}m^{2}(z_{1})z_{2}m^{2}(z_{2}) \mathbf{u}_{j}^{\mathsf{H}} \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_{2}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{C}^{-\frac{1}{2}} \mathbf{u}_{j} \mathbf{u}_{i}^{\mathsf{H}} \mathbf{C}^{-\frac{1}{2}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{2} \mathbf{C}^{-\frac{1}{2}} \mathbf{u}_{i} dz_{1} dz_{2} \\ &= \frac{1 - c\rho_{1}^{-2}}{1 + c\rho_{1}^{-1}} \cdot \frac{1 - c\rho_{2}^{-2}}{1 + c\rho_{2}^{-1}} \mathbf{u}_{1}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{u}_{2} \mathbf{u}_{2}^{\mathsf{H}} \mathbf{J}_{1}^{\mathsf{T}} \mathbf{J}_{1} \mathbf{u}_{1}. \end{split}$$

where $\mathbf{T}_3^\mathsf{H}(z_1,z_2) = m(z_1)m(z_2)\mathbf{C}^{-\frac{1}{2}}\mathbf{J}_1^\mathsf{T}\mathbf{J}_1\mathbf{C}^{-\frac{1}{2}} + \frac{n}{T}\eta(z_1,z_2)\mathbf{I}_N$ For the special case $\mathbf{P} = \mathbf{I}_K$, setting that \mathbf{x} is the eigenvector of $\mathbf{A}^\mathsf{H}\mathbf{A}\mathbf{P}$, then $(\mathbf{A}\mathbf{P})\mathbf{x}$ is the eigenvector of APA^H . So that \mathbf{u}_1 and \mathbf{u}_2 have the following form:

$$\mathbf{u}_1 = \mathbf{A}\mathbf{x}_1 = \frac{1}{\sqrt{2}} [e^{\frac{\imath \alpha c}{2}} \mathbf{a}(\theta_1) + \mathbf{a}(\theta_2)]^\mathsf{T}$$
$$\mathbf{u}_2 = \mathbf{A}\mathbf{x}_2 = \frac{1}{\sqrt{2}} [\mathbf{a}(\theta_1) - e^{-\frac{\imath \alpha c}{2}} \mathbf{a}(\theta_2)]^\mathsf{T}$$

Then $\mathbf{u}_1 \mathbf{J}_1^T \mathbf{J}_2 \mathbf{u}_1$ has the form as:

$$\mathbf{u}_{1}\mathbf{J}_{1}^{T}\mathbf{J}_{2}\mathbf{u}_{1} = \frac{1}{2}\left[\frac{n}{N}e^{\imath\Delta\theta_{1}} + \frac{n}{N}e^{\imath\Delta\theta_{2}} + \frac{1 - e^{\imath\alpha cn/N}}{-\alpha c}e^{\frac{\imath\alpha c}{2}}\frac{n}{N}e^{-\imath\Delta\theta_{2}} + \frac{1 - e^{\imath\alpha cn/N}}{-\imath\alpha c}e^{-\frac{\imath\alpha c}{2}}\frac{n}{N}e^{\imath\Delta\theta_{2}}\right]$$

12 Widely spaced DoA(P is not diagonal)

Signal Mode1231:

$$\mathbf{X} = (\mathbf{A}\mathbf{P}\mathbf{A}^{H} + \mathbf{I}_{k})^{\frac{1}{2}}\mathbf{Z}$$

$$\mathbf{C} = (\mathbf{A}\mathbf{P}\mathbf{A}^{H} + \mathbf{I}_{k})$$

$$P = BLB^{H}, B = \begin{pmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{pmatrix}, B^{H}B = BB^{H} = I, \mathbf{U}_{s} = \mathbf{A}\mathbf{B}$$
(62)

Q is

We aim to get the eigenvalue of this $\Phi_1^{-1}\Phi_2$, (denoted as $\lambda_1...\lambda_k$)

$$\overline{\Phi}_1 = \frac{n}{N} I_k, \Phi_2 = \begin{pmatrix} \hat{u}_1^H J_1^H J_2 \hat{u}_1 & \hat{u}_1^H J_1^H J_2 \hat{u}_2 \\ \hat{u}_2^H J_1^H J_2 \hat{u}_1 & \hat{u}_2^H J_1^H J_2 \hat{u}_2 \end{pmatrix}$$
(63)

We do not have direct access to the deterministic equivalent for Φ_2 referring to $\overline{\Phi}_2$, but an asymptotic eigenvalue equivalence (denoted as $\overline{\lambda}_1...\overline{\lambda}_k$) can be obtained.

$$Eig(\Phi_1^{-1}\Phi_2) = Eig(\overline{\Phi}_1^{-1}\Phi_2)$$

$$= \frac{n}{N}Eig(\Phi_2)$$
(64)

The above equation is the solution of the following equation

$$\begin{split} \det(\Phi_{2} - \lambda I_{2}) &= (\hat{u}_{1}^{H} J_{1}^{H} J_{2} \hat{u}_{1} - \lambda)(\hat{u}_{2}^{H} J_{1}^{H} J_{2} \hat{u}_{2} - \lambda) - \hat{u}_{1}^{H} J_{1}^{H} J_{2} \hat{u}_{2} \hat{u}_{2}^{H} J_{1}^{H} J_{2} \hat{u}_{1} \\ &= \lambda^{2} - \lambda(\hat{u}_{1}^{H} J_{1}^{H} J_{2} \hat{u}_{1} + \hat{u}_{2}^{H} J_{1}^{H} J_{2} \hat{u}_{2}) + \hat{u}_{1}^{H} J_{1}^{H} J_{2} \hat{u}_{1} \hat{u}_{2}^{H} J_{1}^{H} J_{2} \hat{u}_{2} - \hat{u}_{1}^{H} J_{1}^{H} J_{2} \hat{u}_{2} \hat{u}_{2}^{H} J_{1}^{H} J_{2} \hat{u}_{1} \\ &= \overline{\lambda}^{2} - \overline{\lambda}(g_{1} u_{1}^{H} J_{1}^{H} J_{2} u_{1} + g_{2} u_{2}^{H} J_{1}^{H} J_{2} u_{2}) + g_{1} g_{2} u_{1}^{H} J_{1}^{H} J_{2} u_{1} u_{2}^{H} J_{1}^{H} J_{2} u_{2} - g_{1} g_{2} u_{1}^{H} J_{1}^{H} J_{2} u_{2} u_{2}^{H} J_{1}^{H} J_{2} u_{1} \\ &= \overline{\lambda}^{2} - \overline{\lambda}(g_{1} u_{1}^{H} J_{1}^{H} J_{2} u_{1} + g_{2} u_{2}^{H} J_{1}^{H} J_{2} u_{2}) + g_{1} g_{2} Det(\mathbf{U}_{s}^{H} \mathbf{J}_{1}^{H} \mathbf{J}_{2} \mathbf{U}_{s}) \\ &= \overline{\lambda}^{2} - \overline{\lambda}(g_{1} (b_{11}^{2} \frac{n}{N} e^{i\theta_{1}} + b_{12}^{2} \frac{n}{N} e^{i\theta_{1}}) + g_{2} (b_{21}^{2} \frac{n}{N} e^{i\theta_{1}} + b_{22}^{2} \frac{n}{N} e^{i\theta_{2}})) + g_{1} g_{2} Det(\mathbf{A}^{H} \mathbf{J}_{1}^{H} \mathbf{J}_{2} \mathbf{A}) \\ &= \overline{\lambda}^{2} - \overline{\lambda}((g_{1} b_{11}^{2} + g_{2} b_{21}^{2}) \frac{n}{N} e^{i\theta_{1}} + (g_{1} b_{12}^{2} + g_{2} b_{22}^{2}) \frac{n}{N} e^{i\theta_{2}})) + g_{1} g_{2} (\frac{n}{N})^{2} e^{i\theta_{1}} e^{i\theta_{2}} \end{split}$$

$$\begin{split} \alpha_1 &= (g_1 b_{11}^2 + g_2 b_{21}^2) \frac{n}{N} e^{i\theta_1} + (g_1 b_{12}^2 + g_2 b_{22}^2) \frac{n}{N} e^{i\theta_2}) \quad g_i = (1 - c l_i^{-2}) / (1 + c l_i^{-1}) \\ \alpha_2 &= g_1 g_2 (\frac{n}{N})^2 e^{i\theta_1} e^{i\theta_2} \\ \overline{\lambda}_1 &= \frac{\alpha_1 + \sqrt{\Delta}}{2} \\ \overline{\lambda}_2 &= \frac{\alpha_1 - \sqrt{\Delta}}{2} \\ \Delta &= \alpha_1^2 - 4\alpha_2 \end{split}$$

$$\overline{\lambda}_{1} = \frac{(g_{1}b_{11}^{2} + g_{2}b_{21}^{2})\frac{n}{N}e^{i\theta_{1}} + (g_{1}b_{12}^{2} + g_{2}b_{22}^{2})\frac{n}{N}e^{i\theta_{2}}) + \sqrt{((g_{1}b_{11}^{2} + g_{2}b_{21}^{2})\frac{n}{N}e^{i\theta_{1}} + (g_{1}b_{12}^{2} + g_{2}b_{22}^{2})\frac{n}{N}e^{i\theta_{2}}))^{2} - 4*}}{\text{I guess the two roots of the equation are } (g_{1}b_{11}^{2} + g_{2}b_{21}^{2})\frac{n}{N}e^{i\theta_{1}}, (g_{1}b_{21}^{2} + g_{2}b_{22}^{2})\frac{n}{N}e^{i\theta_{2}}}$$

I guess the two roots of the equation are $(g_1b_{11}^2+g_2b_{21}^2)\frac{n}{N}e^{i\theta_1}$, $(g_1b_{21}^2+g_2b_{22}^2)\frac{n}{N}e^{i\theta_2}$ Specially, P is diagonal, $B = \mathbf{I}_k$, $b_{11} = b_{22} = 1$, $b_{12} = b_{21} = 0$, $\overline{\lambda}_1 = g_1\frac{n}{N}e^{i\theta_1}$, $\overline{\lambda}_2 = g_2\frac{n}{N}e^{i\theta_2}$ (consistent with the results of our paper)

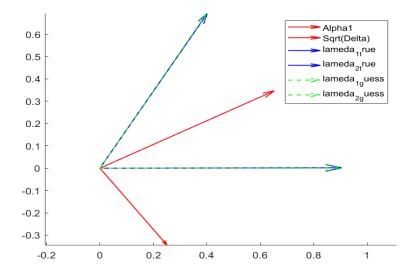


Figure 1: The caption

$$g_1g_2 = (g_1b_{11}^2 + g_2b_{21}^2)(g_1b_{21}^2 + g_2b_{22}^2) - b_{11}^2b_{12}^2(g_1 - g_2)^2$$
(67)

If we can prove that the second item $(b_{11}^2b_{12}^2(g_1-g_2)^2)$ asymptotically $(N,T\to\infty)$ converges to 0 (**Note**: the item is equal to 0 absolutely when P is diagonal, due to $b_{12}=0$).

$$\overline{\lambda}_{1} = (g_{1}b_{11}^{2} + g_{2}b_{21}^{2})\frac{n}{N}e^{i\theta_{1}}$$

$$\overline{\lambda}_{2} = (g_{1}b_{21}^{2} + g_{2}b_{22}^{2})\frac{n}{N}e^{i\theta_{2}}$$
(68)

All we need do is to prove the rate of the item2 is slowly than the first ite12312321321435342m $\lambda_1=1$