

Homework 4 Iterative Methods

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a)

The answer is this problem is well posed. The reason is compatibility condition in Neumann conditions is required forcing f and the boundary condition need to fulfill specific compatibility conditions. Physically, it expresses that flux through the boundary needs to be compensated with a source term inside the domain. For this question, the $dp/dn=0$, so the integral of force should also be 0. For this homework, I calculated sum of f , the ans is $1.7e-15$, which is nearly equal to 0, so it fulfill the compatibility condition.

b)

For this question, the steps are below:

- 1: Discretized the Laplace equation in all boundaries.
- 2: Insert the discretization equation in Gauss-Seidel method
- 3: Implemented to Successive Over-Relaxation method.

Then got the answer.

for $i=1$

$$p_{1,1}^{(m+1)} = (1-\omega)p_{1,1}^{(m)} + \frac{\omega}{1+\beta^2} [p_{2,1}^{(m)} + \beta^2 p_{1,2}^{(m)} - h_x^2 f_{1,1}]$$

$$p_{1,j}^{(m+1)} = (1-\omega)p_{1,j}^{(m)} + \frac{\omega}{2(1+\beta^2)-1} [p_{2,j}^{(m)} + \beta^2 p_{1,j+1}^{(m)} + \beta^2 p_{1,j-1}^{(m+1)} - h_x^2 f_{1,j}] (j=2, N_y-1)$$

$$p_{1,N_y}^{(m+1)} = (1-\omega)p_{1,N_y}^{(m)} + \frac{\omega}{1+\beta^2} [p_{2,N_y}^{(m)} + \beta^2 p_{1,N_y-1}^{(m+1)} - h_x^2 f_{1,N_y}]$$

For $i=N_x$

$$p_{N_x,1}^{(m+1)} = (1-\omega)p_{N_x,1}^{(m)} + \frac{\omega}{1+\beta^2} [p_{N_x-1,1}^{(m+1)} + \beta^2 p_{N_x,2}^{(m)} - h_x^2 f_{N_x,1}]$$

$$p_{N_x,j}^{(m+1)} = (1-\omega)p_{N_x,j}^{(m)} + \frac{\omega}{2(1+\beta^2)-1} [p_{N_x-1,j}^{(m+1)} + \beta^2 p_{N_x,j+1}^{(m)} + \beta^2 p_{N_x,j-1}^{(m+1)} - h_x^2 f_{N_x,j}] (j=2, N_y-1)$$

$$p_{N_x,N_y}^{(m+1)} = (1-\omega)p_{N_x,N_y}^{(m)} + \frac{\omega}{1+\beta^2} [p_{N_x-1,N_y}^{(m+1)} + \beta^2 p_{N_x,N_y-1}^{(m+1)} - h_x^2 f_{N_x,N_y}]$$

For $j=1$

$$p_{i,1}^{(m+1)} = (1-\omega)p_{i,1}^{(m)} + \frac{\omega}{1+\beta^2} [p_{2,1}^{(m)} + \beta^2 p_{i,2}^{(m)} - h_y^2 f_{i,1}]$$

$$p_{i,1}^{(m+1)} = (1-\omega)p_{i,1}^{(m)} + \frac{\omega}{2(1+\beta^2)-\beta^2} [p_{i-1,1}^{(m+1)} + p_{i+1,1}^{(m)} + \beta^2 p_{i,2}^{(m)} - h_y^2 f_{i,1}] (i=2, N_x-1)$$

$$p_{N_x,1}^{(m+1)} = (1-\omega)p_{N_x,1}^{(m)} + \frac{\omega}{1+\beta^2} [p_{N_x-1,1}^{(m+1)} + \beta^2 p_{N_x,2}^{(m)} - h_y^2 f_{N_x,1}]$$

$j=N_y$

$$p_{1,N_y}^{(m+1)} = (1-\omega)p_{1,N_y}^{(m)} + \frac{\omega}{1+\beta^2} \left[p_{2,N_y}^{(m)} + \beta^2 p_{1,N_y-1}^{(m+1)} - h_x^2 f_{1,N_y} \right]$$

$$p_{i,N_y}^{(m+1)} = (1-\omega)p_{i,N_y}^{(m)} + \frac{\omega}{2(1+\beta^2)-\beta^2} \left[p_{i+1,N_y}^{(m)} + p_{i-1,N_y}^{(m+1)} + \beta^2 p_{i,N_y-1}^{(m+1)} - h_x^2 f_{i,N_y} \right] (i = 2, N_x - 1)$$

$$p_{N_x,N_y}^{(m+1)} = (1-\omega)p_{N_x,N_y}^{(m)} + \frac{\omega}{1+\beta^2} \left[p_{N_x-1,N_y}^{(m+1)} + \beta^2 p_{N_x,N_y-1}^{(m+1)} - h_x^2 f_{N_x,N_y} \right]$$

c)

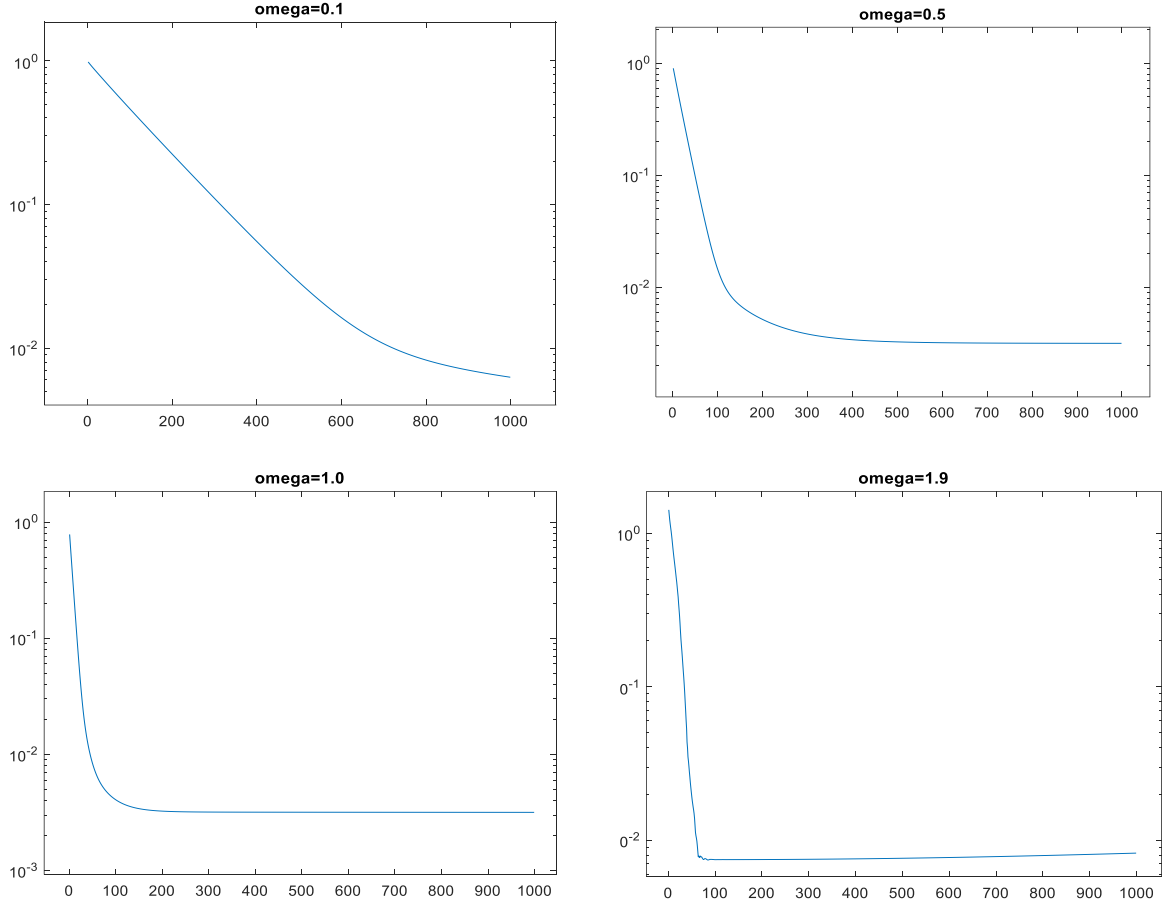


Figure 1 Relative residual with interaction number m

As shown in figure 1, the relative residual are all decreased. While for the $\omega=1.9$, the convergence is the fastest and $\omega=0.1$ is the slowest to convergence. When ω is 1.0, it is the GS method, which is slower than SOR method when ω larger than 1. So, the SOR could accelerate convergence.

d)

As shown in figure 2, is the different $N_x=N_y=N$, from 20 to 100. As can be seen, the coarse mesh, $N=20$ could accelerate convergence. While the fine mesh, $N=100$ induce slow convergence, the reason is the error reduction is related to h^2 , large N cause small h so the convergence is slow.

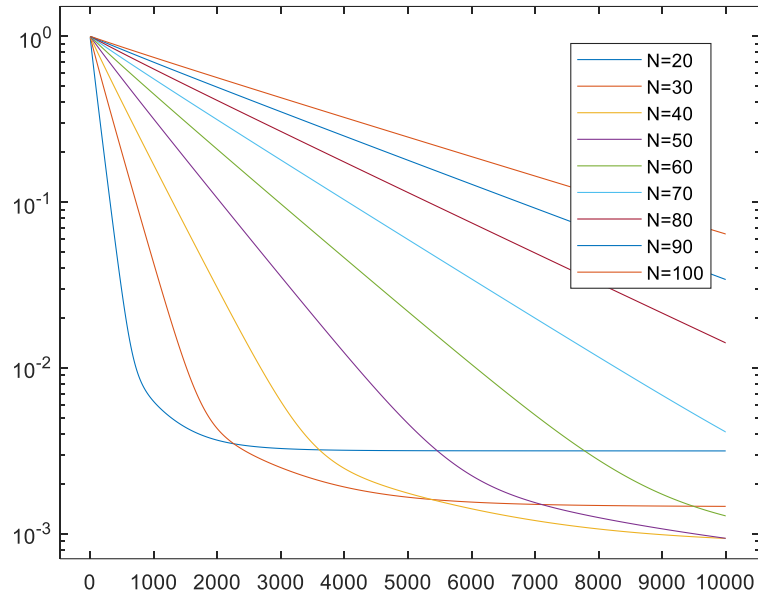


Figure 2 Different N_x and N_y induced relative residual

e)

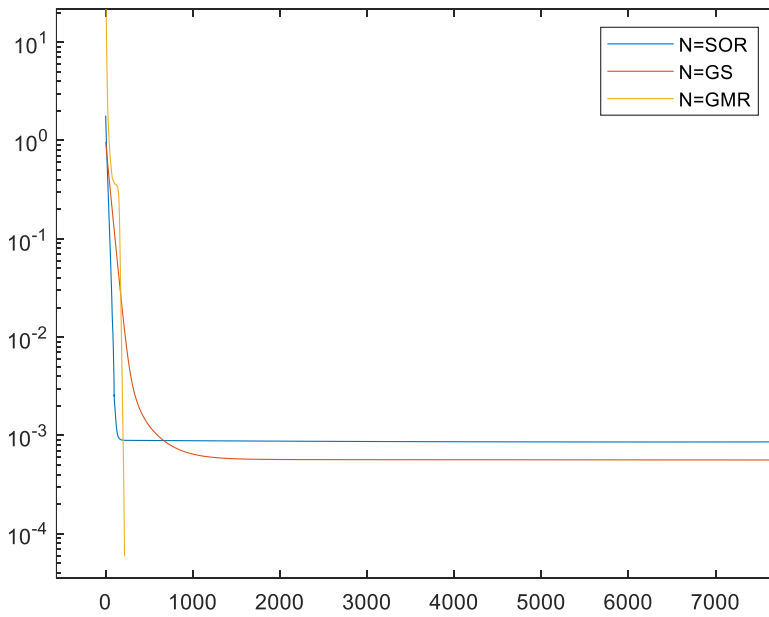


Figure 3 Comparisons of three different iteration methods

As shown in figure 3, the comparisons of GS, SOR and GMR. It could be seen that when ω is 1.9, the SOR is the fastest method to convergence while the relative residual is higher than GMR. The GMR is also fast, nearly the same with SOR, but it has a very high accuracy, the relative residual is smallest. The GS is the slowest one to convergence, also has a high relative residual.