## Upwind discretization

Consider the model problem:

$$\frac{\partial u}{\partial t} + a \frac{\partial u}{\partial x} = 0$$
 (a is the convective velocity)

Use an explicit first-order Euler for discretization of the time derivative

$$u_i^{n+1} = u_i^n - \Delta t \ a(u_x)_i^n$$

It is known that a central-difference discretization of  $(u_x)_i^n$  gives an unstable scheme.

then use a backward discretization when a > 0

$$u_i^{n+1} = u_i^n - a \frac{\Delta t}{\Delta x} (u_i^n - u_{i-1}^n).$$

This equation is stable for  $0 \le a \frac{\Delta t}{\Delta x} \le 1$ .

$$if \ a < 0$$
 
$$\underbrace{ \begin{array}{c} \text{propagation} \\ \hline i-1 & i & i+1 \end{array} }$$

then use a forward discretization

$$u_i^{n+1} = u_i^n - a \frac{\Delta t}{\Delta x} (u_{i+1}^n - u_i^n).$$

This equation is stable for  $-1 \le a \frac{\Delta t}{\Delta x} \le 0$ .

None of these discretizations are stable for both negative and positive wave propagation or convection velocity.

Let us define

$$a^{+} = \max(a, 0) = \frac{1}{2}(a + |a|),$$
  
 $a^{-} = \min(a, 0) = \frac{1}{2}(a - |a|).$ 

These yield

if 
$$a > 0$$
 then  $a^{-} = 0$ ,  
if  $a < 0$  then  $a^{+} = 0$ .

Then write

$$u_i^{n+1} = u_i^n - \frac{\Delta t}{\Delta x} \left[ a^+(u_i^n - u_{i-1}^n) + a^-(u_{i+1}^n - u_i^n) \right].$$

This is first-order accurate upwind scheme for scaler form of Euler equations.

## Flux splitting

## Scalar problem

Split the convetive flux  $u\phi$  as

$$u\phi = f^+ + f^-,$$

where

$$f^{+} = \frac{1}{2}(u + |u|)\phi,$$
  
$$f^{-} = \frac{1}{2}(u - |u|)\phi.$$

If u > 0 then  $f^{+} = u\phi$ ,  $f^{-} = 0$ 

If 
$$u < 0$$
 then  $f^+ = 0$ ,  $f^- = u\phi$ 

Then, one can write

$$[(u\phi)_x]_i = (f_x^+ + f_x^-)_i \approx \frac{f_i^+ - f_{i-1}^+}{\Delta x} + \frac{f_{i+1}^+ - f_i^+}{\Delta x}$$
 (first-order accurate).

## Euler equation

1D case: 
$$\frac{\partial \underline{U}}{\partial t} + \frac{\partial \underline{F}}{\partial x} = 0.$$

$$U = \begin{pmatrix} \rho \\ \rho u \\ E_t \end{pmatrix}, \qquad \tilde{F} = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ (E_t + p)u \end{pmatrix}.$$

Can be written as  $\frac{\partial \underline{U}}{\partial t} + \underset{\approx}{A} \frac{\partial \underline{U}}{\partial x} = 0$ , where  $\underset{\approx}{A} = \frac{\partial \underline{F}}{\partial \underline{U}}$ .

One can write F = AU.

Apply eigendecomposition on  $\underset{\approx}{A} \Rightarrow \underset{\approx}{A} = \underset{\approx}{T} \underset{\approx}{\Lambda} T^{-1}$ , where

 $T = \mathbb{Z}$  is a matrix with eigenvectors of  $A = \mathbb{Z}$  as columns.

 $\stackrel{\wedge}{\mathbb{A}}$  is a matrix with eigenvalues of  $\stackrel{\wedge}{\mathbb{A}}$  on diagonal.

Decompose  $\underset{\sim}{\wedge}$  into  $\underset{\sim}{\wedge}^+ + \underset{\sim}{\wedge}^-$ , where

 $\mathring{\mathbb{A}}^+$  contains only positive elements (associated with right-running characteristics)

 $\mathop{\mathbb{A}^{-}}$  contains only negative elements (associated with left-running characteristics)

Here, the eigenvalues of matrix  $\underset{\approx}{A}$  are  $(u+c,\ u,\ u-c)$  which yield

$$\Lambda^{+} = \begin{bmatrix} u+c & & \\ & u & \\ & & 0 \end{bmatrix}, \quad \Lambda^{-} = \begin{bmatrix} 0 & & \\ & 0 & \\ & u-c \end{bmatrix}$$

Then, we have 
$$A = \mathbb{Z} A^+ \mathbb{Z}^{-1} + \mathbb{Z} A^- \mathbb{Z}^{-1} = A^+ + A^-$$
$$\Rightarrow \mathbb{Z} = A^+ \mathbb{Z} + A^- \mathbb{Z} = \mathbb{Z}^+ + \mathbb{Z}^-.$$

Now we can write

$$\frac{\partial \underline{U}}{\partial t} + \frac{\partial \underline{F}^+}{\partial x} + \frac{\partial \underline{F}^-}{\partial x} = 0.$$

 $\mathcal{E}^+$  and  $\mathcal{E}^-$  are associated with wave propagation in positive and negative directions, respectively. The discretized equations are then

$$U_i^{n+1} = U_i^n - \frac{\Delta t}{\Delta x} \left[ (E_i^+ - E_{i-1}^+) + (E_{i+1}^- - E_i^-) \right].$$

This figure illustrates the idea behind the flux-splitting scheme discussed above.

