

理论力学 A 第二次作业答案

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Problem 1. 以下假设 $\vec{a} = a_i \vec{e}_i, \vec{b} = b_i \vec{e}_i, \vec{c} = c_i \vec{e}_i$, 其中 \vec{e}_i 为直角坐标下的单位基矢量:

$$\begin{aligned}\vec{a} \times (\vec{b} \times \vec{c}) &= \vec{a} \times (\vec{e}_i \epsilon_{ijk} b_j c_k) \\ &= \epsilon_{lmi} \vec{e}_l a_m \epsilon_{ijk} b_j c_k \\ &= \delta_{jl} \delta_{km} \vec{e}_l a_m b_j c_k - \delta_{jm} \delta_{kl} \vec{e}_l a_m b_j c_k \\ &= (\vec{a} \cdot \vec{c}) \vec{b} - (\vec{a} \cdot \vec{b}) \vec{c}\end{aligned}$$

$$\begin{aligned}(\vec{a} \times \vec{b}) \cdot (\vec{c} \times \vec{d}) &= (\epsilon_{ijk} \vec{e}_i a_j b_k) \cdot (\epsilon_{lmn} \vec{e}_l c_m d_n) \\ &= \epsilon_{ijk} \epsilon_{lmn} a_j b_k c_m d_n \\ &= \delta_{jm} \delta_{kn} a_j b_k c_m d_n - \delta_{jn} \delta_{km} a_j b_k c_m d_n \\ &= (\vec{a} \cdot \vec{c})(\vec{b} \cdot \vec{d}) - (\vec{a} \cdot \vec{d})(\vec{b} \cdot \vec{c})\end{aligned}$$

$$\begin{aligned}&\vec{a} \times (\vec{b} \times \vec{c}) + \vec{b} \times (\vec{c} \times \vec{a}) + \vec{c} \times (\vec{a} \times \vec{b}) \\ &= (\vec{a} \cdot \vec{c}) \vec{b} - (\vec{a} \cdot \vec{b}) \vec{c} + (\vec{b} \cdot \vec{a}) \vec{c} - (\vec{b} \cdot \vec{c}) \vec{a} + (\vec{c} \cdot \vec{b}) \vec{a} - (\vec{c} \cdot \vec{a}) \vec{b} \\ &= 0\end{aligned}$$

Problem 2. (1) 根据: $C_i = \frac{1}{2} \epsilon_{ijk} T_{jk}$, 可知:

$$\begin{aligned}\epsilon_{ijk} C_k &= \frac{1}{2} \epsilon_{ijk} \epsilon_{kmn} T_{mn} \\ &= \frac{1}{2} (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) T_{mn} \\ &= \frac{1}{2} (T_{ij} - T_{ji}) \\ &= T_{ij}\end{aligned}$$

(2)	点乘	矢量	轴矢量	叉乘	矢量	轴矢量
	矢量	标量	赝标量	矢量	赝矢量	矢量
	轴矢量	赝标量	标量	轴矢量	矢量	赝矢量

Problem 3. (1) 根据转动公式的分量形式：

$$\begin{aligned} x'_i &= x_i \cos \theta + n_i n_j (1 - \cos \theta) x_j + \epsilon_{ijk} n_k x_j \sin \theta \\ &= [\delta_{ij} \cos \theta + n_i n_j (1 - \cos \theta) - \epsilon_{ijk} n_k \sin \theta] x_j \end{aligned}$$

又 $x'_i = R_{ij} x_j$, 得：

$$R_{ij} = \delta_{ij} \cos \theta + n_i n_j (1 - \cos \theta) - \epsilon_{ijk} n_k \sin \theta$$

据此可得：

$$R = I \cos \theta + \begin{pmatrix} n_1^2 & n_1 n_2 & n_1 n_3 \\ n_1 n_2 & n_2^2 & n_2 n_3 \\ n_1 n_3 & n_2 n_3 & n_3^2 \end{pmatrix} (1 - \cos \theta) + \begin{pmatrix} 0 & -n_3 & n_2 \\ n_3 & 0 & -n_1 \\ -n_2 & n_1 & 0 \end{pmatrix} \sin \theta$$

(2)

$$\hat{n} = \frac{1}{\sqrt{3}}(1, 1, 1), R = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} (R^T \text{ 亦可})$$

(3)

$$\begin{aligned} R_{ik} R_{jk} &= [\delta_{ik} \cos \theta + n_i n_k (1 - \cos \theta) - \epsilon_{ikl} n_l \sin \theta] [\delta_{jk} \cos \theta + n_j n_k (1 - \cos \theta) - \epsilon_{jkm} n_m \sin \theta] \\ &= \delta_{ik} \delta_{jk} \cos^2 \theta + n_i n_j n_k n_k (1 - \cos \theta)^2 + \epsilon_{ikl} \epsilon_{jkm} n_l n_m \sin^2 \theta \\ &\quad + \delta_{ik} n_j n_k (1 - \cos \theta) \cos \theta - \delta_{ik} \epsilon_{jkm} n_m \cos \theta \sin \theta - \epsilon_{ikl} n_j n_l n_k (1 - \cos \theta) \sin \theta \\ &\quad + \delta_{jk} n_i n_k (1 - \cos \theta) \cos \theta - \delta_{jk} \epsilon_{ikm} n_m \cos \theta \sin \theta - \epsilon_{jkm} n_m n_i n_k (1 - \cos \theta) \sin \theta \end{aligned}$$

由于 $n_i n_j$ 为对称张量, ϵ_{ijk} 为全反对称张量, 则根据上次作业的结论可知, 形如 $\epsilon_{ijk} n_i n_j$ 的项事实上对求和没有贡献. 再虑及 $n_k n_k = \hat{n}^2 = 1$, 那么：

$$\begin{aligned} R_{ik} R_{jk} &= \delta_{ij} \cos^2 \theta + n_i n_j (1 - \cos \theta)^2 + (\delta_{ij} \delta_{ml} - \delta_{im} \delta_{jl}) n_m n_l \sin^2 \theta \\ &\quad + 2n_i n_j (1 - \cos \theta) \cos \theta - (\epsilon_{ijm} + \epsilon_{jim}) n_m \cos \theta \sin \theta \\ &= \delta_{ij} (\cos^2 \theta + \sin^2 \theta) + n_i n_j [(1 - \cos \theta)^2 - \sin^2 \theta + 2(1 - \cos \theta) \cos \theta] \\ &= \delta_{ij} \end{aligned}$$

Problem 4. 设 $\overleftrightarrow{T} = T_{ij} \hat{x}_i \hat{x}_j = T'_{ij} \hat{x}'_i \hat{x}'_j$, 其中 \hat{x}_i, \hat{x}'_i 分别为原参考系与转动参考系中的单位矢量, 由于在转动参考系中, 转动参考系的直角坐标基矢自然是常矢量, 有：

$$\left(\frac{d\overleftrightarrow{T}}{dt} \right)_{rot} = \left(\frac{d}{dt} (T'_{ij} \hat{x}'_i \hat{x}'_j) \right)_{rot} = \left(\frac{dT'_{ij}}{dt} \hat{x}'_i \hat{x}'_j \right)_{rot} = \frac{dT'_{ij}}{dt} \hat{x}'_i \hat{x}'_j$$

而：

$$\begin{aligned}
 \frac{d\overleftrightarrow{T}}{dt} &= \frac{d}{dt}(T'_{ij}\hat{x}'_i\hat{x}'_j) \\
 &= \frac{dT'_{ij}}{dt}\hat{x}'_i\hat{x}'_j + T'_{ij}\frac{d\hat{x}'_i}{dt}\hat{x}'_j + T'_{ij}\hat{x}'_i\frac{d\hat{x}'_j}{dt} \\
 &= \left(\frac{d\overleftrightarrow{T}}{dt}\right)_{rot} + T'_{ij}\vec{\omega} \times \hat{x}'_i\hat{x}'_j + T'_{ij}\hat{x}'_i\vec{\omega} \times \hat{x}'_j \\
 &= \left(\frac{d\overleftrightarrow{T}}{dt}\right)_{rot} + T'_{ij}\vec{\omega} \times \hat{x}'_i\hat{x}'_j - T'_{ij}\hat{x}'_i\hat{x}'_j \times \vec{\omega} \\
 &= \left(\frac{d\overleftrightarrow{T}}{dt}\right)_{rot} + \vec{\omega} \times \overleftrightarrow{T} - \overleftrightarrow{T} \times \vec{\omega}
 \end{aligned}$$

Problem 5. (1) 据柱坐标定义：

$$\begin{cases} x = s \cos \phi = \sqrt{\xi \eta} \cos \phi \\ y = s \sin \phi = \sqrt{\xi \eta} \sin \phi \\ z = z = \frac{1}{2}(\xi - \eta) \end{cases}$$

在 $x-z$ 平面内, 有 $y=0$, 不妨 $\phi=0$, 即：

$$\begin{cases} x = \sqrt{\xi \eta} \\ z = \frac{1}{2}(\xi - \eta) \end{cases}$$

η 曲线(消去 η):

$$z = \frac{1}{2}\left(\xi - \frac{x^2}{\xi}\right)$$

ξ 曲线(消去 ξ):

$$z = \frac{1}{2}\left(\frac{x^2}{\eta} - \eta\right)$$

显见曲线均为抛物线.

(2)

$$\vec{r} = x\hat{x} + y\hat{y} + z\hat{z} = \sqrt{\xi \eta} \cos \phi \hat{x} + \sqrt{\xi \eta} \sin \phi \hat{y} + \frac{1}{2}(\xi - \eta)\hat{z}$$

则有：

$$\begin{aligned}
 \left|\frac{\partial \vec{r}}{\partial \xi}\right| \hat{\xi} &= \frac{\partial \vec{r}}{\partial \xi} = \frac{1}{2}\sqrt{\frac{\eta}{\xi}}(\cos \phi \hat{x} + \sin \phi \hat{y}) + \frac{1}{2}\hat{z} \\
 \left|\frac{\partial \vec{r}}{\partial \eta}\right| \hat{\eta} &= \frac{\partial \vec{r}}{\partial \eta} = \frac{1}{2}\sqrt{\frac{\xi}{\eta}}(\cos \phi \hat{x} + \sin \phi \hat{y}) - \frac{1}{2}\hat{z} \\
 \left|\frac{\partial \vec{r}}{\partial \phi}\right| \hat{\phi} &= \frac{\partial \vec{r}}{\partial \phi} = \sqrt{\xi \eta}(-\sin \phi \hat{x} + \cos \phi \hat{y})
 \end{aligned}$$

可得： $\hat{\xi} \cdot \hat{\eta} = \hat{\xi} \cdot \hat{\phi} = 0$

(3) 由正交性可得:

$$\begin{aligned} v^2 &= \left| \frac{\partial \vec{r}}{\partial \xi} \right|^2 \dot{\xi}^2 + \left| \frac{\partial \vec{r}}{\partial \eta} \right|^2 \dot{\eta}^2 + \left| \frac{\partial \vec{r}}{\partial \phi} \right|^2 \dot{\phi}^2 \\ &= \frac{\xi + \eta}{4} \left(\frac{\dot{\xi}^2}{\xi} + \frac{\dot{\eta}^2}{\eta} \right) + \xi \eta \dot{\phi}^2 \end{aligned}$$

Problem 6. (1) 类似上题, 可得:

η 曲线 (消去 η):

$$\frac{\xi^2}{(\xi^2 - 1)\sigma^2} + \frac{z^2}{\xi^2 \sigma^2} = 1$$

可见曲线为椭圆.

ξ 曲线 (消去 ξ):

$$\frac{\xi^2}{(1 - \eta^2)\sigma^2} - \frac{z^2}{\eta^2 \sigma^2} = -1$$

可见曲线为双曲线.

(2)

$$\vec{r} = s\hat{s} + z\hat{z} = \sigma\sqrt{(\xi^2 - 1)(1 - \eta^2)}\hat{s} + \sigma\xi\eta\hat{z}$$

注意到 \hat{s} 仅明显地依赖于 ϕ , \hat{z} 是常量, 有:

$$\frac{\partial \hat{s}}{\partial \phi} = \dot{\hat{s}} / \dot{\phi} = \frac{\vec{\omega} \times \hat{s}}{\dot{\phi}} = \hat{\phi}$$

式中 $\vec{\omega} = \frac{d\phi}{dt}\hat{z}$, 所以:

$$\begin{aligned} \left| \frac{\partial \vec{r}}{\partial \xi} \right| \dot{\xi} &= \frac{\partial \vec{r}}{\partial \xi} = \sigma\xi \sqrt{\frac{1 - \eta^2}{\xi^2 - 1}} \hat{s} + \sigma\eta\hat{z} \\ \left| \frac{\partial \vec{r}}{\partial \eta} \right| \dot{\eta} &= \frac{\partial \vec{r}}{\partial \eta} = -\sigma\eta \sqrt{\frac{\xi^2 - 1}{1 - \eta^2}} \hat{s} + \sigma\xi\hat{z} \\ \left| \frac{\partial \vec{r}}{\partial \phi} \right| \dot{\phi} &= \frac{\partial \vec{r}}{\partial \phi} = \sigma\sqrt{(\xi^2 - 1)(1 - \eta^2)} \hat{\phi} \end{aligned}$$

可得: $\hat{\xi} \cdot \hat{\eta} = \hat{\xi} \cdot \hat{\phi} = 0$

(3) 类似地, 有:

$$\begin{aligned} v^2 &= \left| \frac{\partial \vec{r}}{\partial \xi} \right|^2 \dot{\xi}^2 + \left| \frac{\partial \vec{r}}{\partial \eta} \right|^2 \dot{\eta}^2 + \left| \frac{\partial \vec{r}}{\partial \phi} \right|^2 \dot{\phi}^2 \\ &= \sigma^2 \left[(\xi^2 - \eta^2) \left(\frac{\dot{\xi}^2}{\xi^2 - 1} + \frac{\dot{\eta}^2}{\eta^2 - 1} \right) + (\xi^2 - 1)(1 - \eta^2) \dot{\phi}^2 \right] \end{aligned}$$