

理论力学 A

2024 年秋季学期第四次习题课

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内容摘要

① 知识点整理

- 中心力与散射
- Hamilton 力学
- 刚体

中心力与散射

中心力问题：

- 角动量守恒： $l = mr^2\dot{\theta} = \text{const.}$, 方向始终与 $\mathbf{r} \times \mathbf{v}$ 所在平面垂直;
- 径向运动与圆周运动稳定性分析：

$$E = \frac{1}{2}m\dot{r}^2 + V(r), \quad V(r) = U(r) + \frac{l^2}{2mr^2}, \quad \frac{\partial V}{\partial r} = 0, \quad \frac{\partial^2 V}{\partial r^2} > 0$$

- 轨道方程：

① 积分形式：

$$\theta = \pm \frac{l}{\sqrt{2m}} \int \frac{dr}{r^2 \sqrt{E - V(r)}}, \quad \Delta\theta = \int_{r_1}^{r_2} \frac{l^2 dr}{r^2 \sqrt{2m(E - U) - l^2/r^2}}$$

② 一阶方程：

$$\left(\frac{du}{d\theta}\right)^2 + u^2 = \frac{2m}{l^2}(E - U), \quad u = \frac{1}{r}$$

③ 二阶方程 (Binet 公式)

$$\frac{d^2 u}{d\theta^2} + u = -\frac{mF}{l^2 u^2}, \quad F = -\frac{\partial U}{\partial r}$$

- 中心力作用下轨道关于拱线对称, 周期运动 (轨道闭合) 的条件:

$$\Delta\theta = q\pi, \quad q \in \mathbb{Q}$$

散射问题:

- 散射角: $\Theta = \pi - 2\theta_b$ (排斥力), $\Theta = |2\theta_b - \pi - 2n\pi|$ (吸引力), θ_b :

$$\theta_b = \int_{r_{\min}}^{\infty} \frac{b dr}{r^2 \sqrt{1 - U/E - b^2/r^2}}$$

- 微分散射截面:

$$\frac{d\sigma}{d\Omega} = \frac{b}{\sin \theta} \left| \frac{db}{d\theta} \right| \quad (\text{轴对称情形})$$

Legendre 变换与 Hamilton 方程

- 利用 Legendre 变换可从 Lagrange 函数给出 Hamilton 函数, 需要满足 Hess 条件:

$$H(q, p, t) = p_k \dot{q}_k - L(q, \dot{q}, t), \quad \det \left(\frac{\partial^2 L}{\partial \dot{q}_i \partial \dot{q}_j} \right) \neq 0$$

- Hamilton 方程:

$$\boxed{\dot{q}_k = \frac{\partial H}{\partial p_k}, \quad \dot{p}_k = -\frac{\partial H}{\partial q_k}}$$

或写成正则变量的形式:

$$\dot{\xi} = \Omega \frac{\partial H}{\partial \xi} = [\xi, H]_{\xi}, \quad \Omega = \begin{pmatrix} 0 & \mathbf{1} \\ -\mathbf{1} & 0 \end{pmatrix}, \quad \xi = (q, p)$$

- Ω 矩阵的性质:

$$\Omega^T = \Omega^{-1} = -\Omega, \quad \Omega_{\alpha\beta} \Omega_{\gamma\beta} = \delta_{\alpha\gamma}$$

Poission 括号

- Poission 括号:

$$[A, B]_{\xi} = \frac{\partial A}{\partial \xi_{\alpha}} \Omega_{\alpha\beta} \frac{\partial B}{\partial \beta} = \frac{\partial A}{\partial q_k} \frac{\partial B}{\partial p_k} - \frac{\partial A}{\partial p_k} \frac{\partial B}{\partial q_k}$$

- 基本 Poission 括号:

$$[q_k, p_l] = \delta_{kl}, \quad [q_k, q_l] = [p_k, p_l] = 0$$

- Poission 括号的性质 (Poission 代数):

$$\begin{aligned} [A, B] &= -[B, A], & [A, [B, C]] + [B, [C, A]] + [C, [A, B]] &= 0 \\ [A, BC] &= [A, B]C + B[A, C], & [AB, C] &= A[B, C] + [A, C]B \end{aligned}$$

- 力学量的运动方程及 Taylor 展开:

$$\dot{f} = \frac{\partial f}{\partial t} + [f, H]_{\xi}, \quad f(t + \tau) = \exp(-\tau D_H) f, \quad D_H f = [H, f]$$

- 判断 Hamilton 体系 $\Leftrightarrow [\dot{\xi}_{\alpha}, \xi_{\beta}] + [\xi_{\alpha}, \dot{\xi}_{\beta}] = 0$

正则变换与哈密顿-雅可比理论

受限正则变换的条件:

- 基本 Poisson 括号的不变性: $\xi \mapsto \eta, [\eta_\alpha, \eta_\beta]_\xi = \Omega_{\alpha\beta}$
- 辛条件:

$$M_{\alpha\beta} = \frac{\partial \eta_\alpha}{\partial \xi_\beta}, \quad M^T \Omega M = \Omega$$

- 可积条件:

$$p_k dq_k - H = P_i dQ_i - K + dF(q, p, t)$$

$$\Rightarrow \frac{\partial F}{\partial q_k} = p_k - P_i \frac{\partial Q_i}{\partial q_k}$$

$$\frac{\partial F}{\partial p_k} = -P_i \frac{\partial Q_i}{\partial p_k}$$

$$K = H + P_i \frac{\partial Q_i}{\partial t} + \frac{\partial F}{\partial t}$$

正则变换的性质：

- hamilton 体系的演化可视为正则变换
- Liouville 体积定理： $\Gamma(t) = \Gamma(0)$
- Liouville 定理：

$$n = \frac{\Delta N}{\Delta \Gamma} = n_0, \quad \rho = \frac{n}{N} = n_0, \quad J_\alpha = \rho_\alpha \Omega_{\alpha\beta} \frac{\partial H}{\partial \beta},$$
$$\frac{\partial \rho}{\partial t} + \partial_\alpha J_\alpha = 0$$

正则变换的分类：

- 第一类正则变换： q, Q 独立, $F_1(q, Q, t) = F(q, p, t)$,
- 第二类正则变换： q, P 独立, $F_2(q, P, t) = F(q, p, t) + Q_k P_k$,

$$p_k = \frac{\partial F_2}{\partial q_k}, \quad Q_k = \frac{\partial F_2}{\partial P_k}, \quad K = H + \frac{\partial F_2}{\partial t}$$

- 第三类正则变换: p, Q 独立, $F_3(p, Q, t) = F(q, p, t) - q_k p_k$,
- 第四类正则变换: p, P 独立, $F_4(p, P, t) = F(q, p, t) - q_k p_k + Q_i P_i$.

Hamilton-Jacobi 理论:

- Hamilton-Jacobi 方程:

$$S(q, P, t) = F_2(q, P, t), \quad \boxed{-\frac{\partial S}{\partial t} = H(q, \frac{\partial S}{\partial q}, t)}$$

- 自治体系: $H = H(q, p), \quad S = T(t) + W(q)$

$$-\frac{\partial T}{\partial t} = H(q, \frac{\partial W}{\partial q}), \quad H(q, \frac{\partial W}{\partial q}) = P_1, \quad \Rightarrow T = -P_1 t + C$$

若 H 不显含某个 q_k , 则令对应的 p_k 为运动常数 P_2 , 有:

$$p_n = \frac{\partial S}{\partial q_n} = P_n = \frac{\partial W}{\partial q_n}, \Rightarrow W = P_n q_n$$

刚体运动学

- 欧拉运动学方程(只写了一种形式, 考试如要用大概率会给):

$$\begin{cases} \omega_1 = \dot{\varphi} \sin \theta \sin \psi + \dot{\theta} \cos \psi \\ \omega_2 = \dot{\varphi} \sin \theta \cos \psi - \dot{\theta} \sin \psi \\ \omega_3 = \dot{\varphi} \cos \theta + \dot{\psi} \end{cases}$$

- 惯量张量相关:

$$\mathbf{L} = \mathbf{J} \cdot \boldsymbol{\omega}, \quad \mathbf{J} = \int dm (r^2 \mathbf{1} - \mathbf{r} \mathbf{r})$$
$$T = \frac{1}{2} \boldsymbol{\omega} \cdot \mathbf{L} = \frac{1}{2} \boldsymbol{\omega} \cdot \mathbf{J} \cdot \boldsymbol{\omega} = \frac{1}{2} J_{ij} \omega_i \omega_j$$

- 正交轴定理: $J_{11} + J_{22} \geq J_{33}$, 平行轴定理: $\mathbf{J} = \mathbf{J}^* + J_C$
- 主轴与主转动惯量: 本质上是惯量张量矩阵本征值问题的解。

刚体动力学

- Euler 动力学方程:

$$\begin{cases} \tau_1 = J_1 \dot{\omega}_1 - (J_2 - J_3) \omega_2 \omega_3 \\ \tau_2 = J_2 \dot{\omega}_2 - (J_3 - J_1) \omega_3 \omega_1 \\ \tau_3 = J_3 \dot{\omega}_3 - (J_1 - J_2) \omega_1 \omega_2 \end{cases}$$

- Euler 陀螺: $\tau = 0, L^2, T$ 守恒
- Lagrange 陀螺:

$$\begin{aligned} L &= \frac{1}{2} J_{12} (\omega_1^2 + \omega_2^2) + \frac{1}{2} J_3 \omega_3^2 - mgl \cos \theta \\ &= \frac{1}{2} J_{12} (\dot{\theta}^2 + \dot{\varphi}^2 \sin^2 \theta) + \frac{1}{2} J_3 (\dot{\varphi} \cos \theta + \dot{\psi})^2 - mgl \cos \theta \end{aligned}$$

p_ψ, p_φ, E 守恒, 其中:

$$E = \frac{1}{2} J_{12} \dot{\theta}^2 + \frac{(p_\varphi - p_\psi \cos \theta)^2}{2 J_{12} \sin^2 \theta} + \frac{p_\psi^2}{2 J_3} + mgl \cos \theta$$