1 QUANTUM MECHANIC

1.1 Wave Function

· Schrodinger function

$$i\hbarrac{\partial\Psi}{\partial t}=-rac{\hbar^2}{2m}rac{\partial^2\Psi}{\partial x^2}+V\Psi$$

ullet Actually, V is called potential function and is independent of time t and it reference to the potential energy of particle

According to the definition of potential function

$$oldsymbol{E} = -
abla arphi(oldsymbol{r})$$

 $m{E}$ is a electrical field function

· hence Schrodinger function can be discribe as

$$\Psi = \psi(x)\phi(t)$$

 $\psi(x)$ is the position-dependent portion of wave function and $\phi(t)$ is the time-dependent portion of wave function.

hence ,we can write the wave function as

$$i\hbarrac{\partial(\psi(x)\phi(t))}{\partial t}=-rac{\hbar^2}{2m}rac{\partial^2(\psi(x)\phi(t))}{\partial x^2}+\psi(x)\phi(t)V(x)$$

$$i\hbar\psi(x)rac{\partial\phi(t)}{\partial t}=-rac{\hbar^2}{2m}\phi(t)rac{\partial^2\psi(x)}{\partial x^2}+\psi(x)\phi(t)V(x)$$

$$i\hbarrac{\partial\phi(t)}{\phi(t)\partial t}=-rac{\hbar^2}{2m}rac{\partial^2\psi(x)}{\psi(x)\partial x^2}+V(x)$$

• For the left portion of the function is only related to t and the right portion is only related to x, hence both two parts are constant. Denote this constant by η , then

$$\eta=i\hbarrac{\partial\phi(t)}{\phi(t)\partial t}=-rac{\hbar^2}{2m}rac{\partial^2\psi(x)}{\psi(x)\partial x^2}+V(x)$$

$$\phi(t) = C e^{-jrac{\eta t}{\hbar}}$$

And according to eular therom we can get,

$$\phi(t) = C(cos(rac{\eta t}{\hbar}) - jsin(rac{\eta t}{\hbar}))$$

It's a sinusoidal wave, with $\omega=rac{\eta}{\hbar}$,then we can get ,

$$T = \frac{2\pi}{\omega} = \frac{2\pi\hbar}{\eta} = \frac{h}{\eta}$$

$$\nu = \frac{1}{T} = \frac{\eta}{h}$$

$$\eta = h \nu = E$$

 η refrence to the energy of the sinusoidal wave.

hence,

$$E=\eta=-rac{\hbar^2}{2m}rac{\partial^2\psi(x)}{\psi(x)\partial x^2}+V(x)$$

$$rac{2m}{\hbar^2}(E-V(x))\psi(x)+rac{\partial^2\psi(x)}{\partial x^2}=0$$

1.2 Physical Meaning of Wave Function

• $|\Psi(x,t)|^2$ refers to the probability density of finding the particle.

$$\phi(t) = Ce^{-j\omega t}$$

$$|\Psi(x,t)|^2 = \Psi \Psi^* = |\psi(x)|^2 C e^{-j\omega t} C e^{j\omega t} = C^2 |\psi(x)|^2$$

We can incorporate C into $\psi(x)$, then

$$|\Psi(x,t)|^2=|\psi(x)|^2$$

1.3 Boundry Conditions

· The function need to be normalized, for it is a discription of probability density, hence

$$\int_{-\infty}^{\infty}\left|\psi(x)
ight|^{2}\!dx=1$$

- ullet If the energy E and the V(X) is both finite ,then there will be two boundry conditions:
 - 1. $\psi(x)$ must be finite , single-valued , continous
 - 2. $\frac{\partial \psi(x)}{\partial x}$ must be finite , single-valued , continous

1.4 Application of Wave Function

1.4.1 Electrons in free space

ullet When there is no force acting on the particle , so V(x) must be a constant .

we have

$$E = h\nu = E_k + V(x)$$

$$E_k = \frac{h^2}{2m}$$

$$E - V(x) = \frac{h^2}{2m}$$

 $ullet E_k \geq 0$ in free space , hence $E \geq V(x)$ Assume V(x) = 0 ,

$$rac{2m}{\hbar^2}E\psi(x)+rac{\partial^2\psi(x)}{\partial x^2}=0$$

The solution of linear differential equation with constant coefficients of the second order is as follows:

if the characteristic function has

1. 2 different real root

$$y = C_1 e^{y_1 x} + C_2 e^{y_2 x}$$

2. 2 same real root

$$y = (C_1 + C_2 x)e^{y_1 x}$$

3. 2 conjugate complex root

$$y=e^{lpha x}(C_1 cos(eta x)+C_2 sin(eta x))$$

$$y_1 = e^{\alpha x} cos(\beta x)$$
 $y_2 = e^{\alpha x} sin(\beta x)$

• hence, the solution of the partial differential equation is:

$$\psi(x) = C_1 exp[rac{\sqrt{2mE}}{\hbar}ix] + C_2 exp[-rac{\sqrt{2mE}}{\hbar}ix]$$

Assume,

$$k=rac{\sqrt{2mE}}{\hbar}$$

then,

$$\psi(x) = C_1 exp[ikx] + C_2 exp[-ikx]$$

ullet hence , recall the time-dependent function $\phi(t)$, there is

$$\Psi(x,t) = exp[-jwt](C_1exp[ikx] + C_2exp[-ikx])$$

We can assume that the particle motion direction is +x then there is $C_2=0\,$

ullet The relationship between k and λ can be derived as follows :

$$k=rac{\sqrt{2mE}}{\hbar}$$

And because V(x)=0 , hence $E=rac{1}{2}mv^2$

$$E = \frac{1}{2}mv^2 = \frac{p^2}{2m}$$

$$k=rac{p}{\hbar}$$

And according to

$$p = \frac{h}{\lambda}$$

$$k = \frac{2\pi}{\lambda}$$

And ultimately , we get the result of Ψ

$$\Psi(x) = C_1 exp[j(kx-\omega t)]$$

$$\Psi\Psi^* = C_1 C_1^*$$

 $\Psi\Psi^*$ is independent to positio , when the particle's momentum is well-defined

1.4.2 The Infinite potential well

· In infinite potential well, there is

$$V(oldsymbol{x}) = \left\{egin{array}{ccc} 0 & 0 < r < a \ \infty & x > a & x < 0 \end{array}
ight.$$

When $x>a \qquad x<0$, if E is finite then , $\psi(x)=0$ When 0< x< a , there is

$$rac{2m}{\hbar^2}E\psi(x)+rac{\partial^2\psi(x)}{\partial x^2}=0$$

we define:

$$rac{2mE}{\hbar^2}=k^2 \qquad k=rac{\sqrt{2mE}}{\hbar}$$

the solution is like:

$$\psi(x) = Asin(kx) + Bcos(kx)$$

• According to boundry condition , when x=a and x=0 , $\psi(x)=0$ hence B=0 and Asin(ka)=0 , hence $ka=n\pi$. For $k\neq 0$ $a\neq 0$, $kx\neq 0$, hence $n\pi\neq 0$ $n\neq 0$ And when n<0 , for $Asin(-n\pi)=-Asin(n\pi)$ the negative ones can be conbined into A

$$\psi(x) = Asin(kx) = Asin(rac{n\pi}{a}x) \quad (n=1,2,3,...)$$

Normalization:

$$|\Psi(x,t)|^2=|\psi(x)|^2$$

$$\int_{-\infty}^{+\infty} |\Psi(x,t)|^2 = \int_{-\infty}^{+\infty} |\psi(x)|^2 = \int_0^a A^2 sin^2 (rac{n\pi}{a}x) = rac{aA^2}{2} = 1$$

hence,

$$A = \sqrt{\frac{2}{a}}$$

$$\psi(x)=\sqrt{rac{2}{a}}sin(rac{n\pi}{a}x) \quad (n=1,2,3,...)$$

1.4.3 The Step potential function

• The potential function is:

$$V(x) = egin{cases} 0 & x < 0 \ V_0 & x > 0 \end{cases}$$

- The incident flux and reflected flux are neither 0 ,
- hence in the region x < 0

$$rac{2m}{\hbar^2}E\psi(x)+rac{\partial^2\psi(x)}{\partial x^2}=0$$

the solution is like:

$$\psi_1(x) = A_1 e^{ik_1 x} + B_1 e^{-ik_1 x}$$

And
$$k_1=rac{\sqrt{2mE}}{\hbar}$$

• The first term is a traveling wave in direction of +xThe second term is a traveling wave in direction of -xSo , the first term refers to incident wave and the second term refers to the reflect wave.

- In th region of x>0 , V(x)
eq 0 , Assume V(x)>E , then

$$rac{2m}{\hbar^2}(E-V(x))\psi(x)+rac{\partial^2\psi(x)}{\partial x^2}=0$$

$$rac{2m}{\hbar^2}(V(x)-E)\psi(x)=rac{\partial^2\psi(x)}{\partial x^2}$$

$$k_2=rac{\sqrt{2m(V(x)-E)}}{\hbar}$$

$$V(x)=V_0 \qquad k_2=rac{\sqrt{2m(V_0-E)}}{\hbar}$$

Then there will be:

$$\psi_2(x) = A_2 e^{k_2 x} + B_2 e^{-k_2 x}$$

According to the boundry condition of finite $A_2=0$

$$\psi_2(x)=B_2e^{-k_2x}$$

According to the boundry condition

$$\psi_1(0) = \psi_2(0) = A_1 + B_1 = B_2$$

According to my comprehention , the variation of perticle at the two sides of x=0 is equal then ,

$$rac{\partial \psi_1(0)}{\partial x} = rac{\partial \psi_2(0)}{\partial x}$$

$$ik_1A_1 - ik_1B_1 = -B_2k_2$$

Then there will be:

$$B_1 = \frac{ik_1 + k_2}{ik_1 - k_2} A_1$$

$$B_2=\frac{2ik_1}{ik_1-k_2}A_1$$