



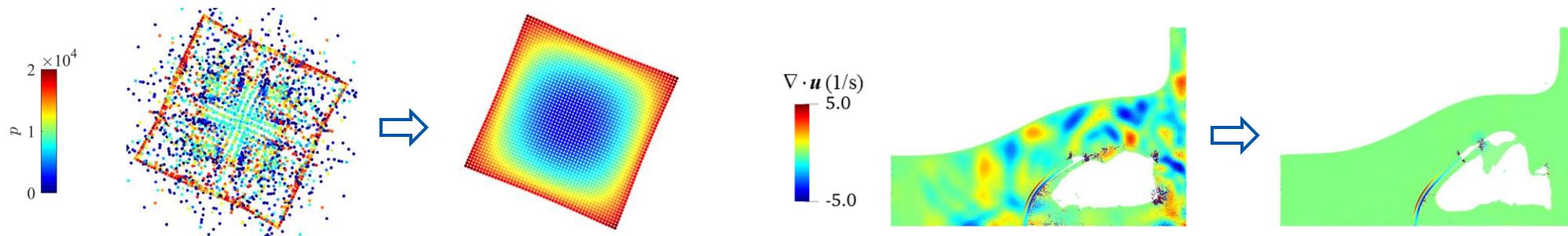
An enhanced SPH-based hydroelastic FSI solver with structural dynamic hourglass control

Yi Zhan¹, Min Luo¹, Abbas Khayyer²

¹ Ocean College, Zhejiang University

² Department of Civil and Earth Resources Engineering, Kyoto University

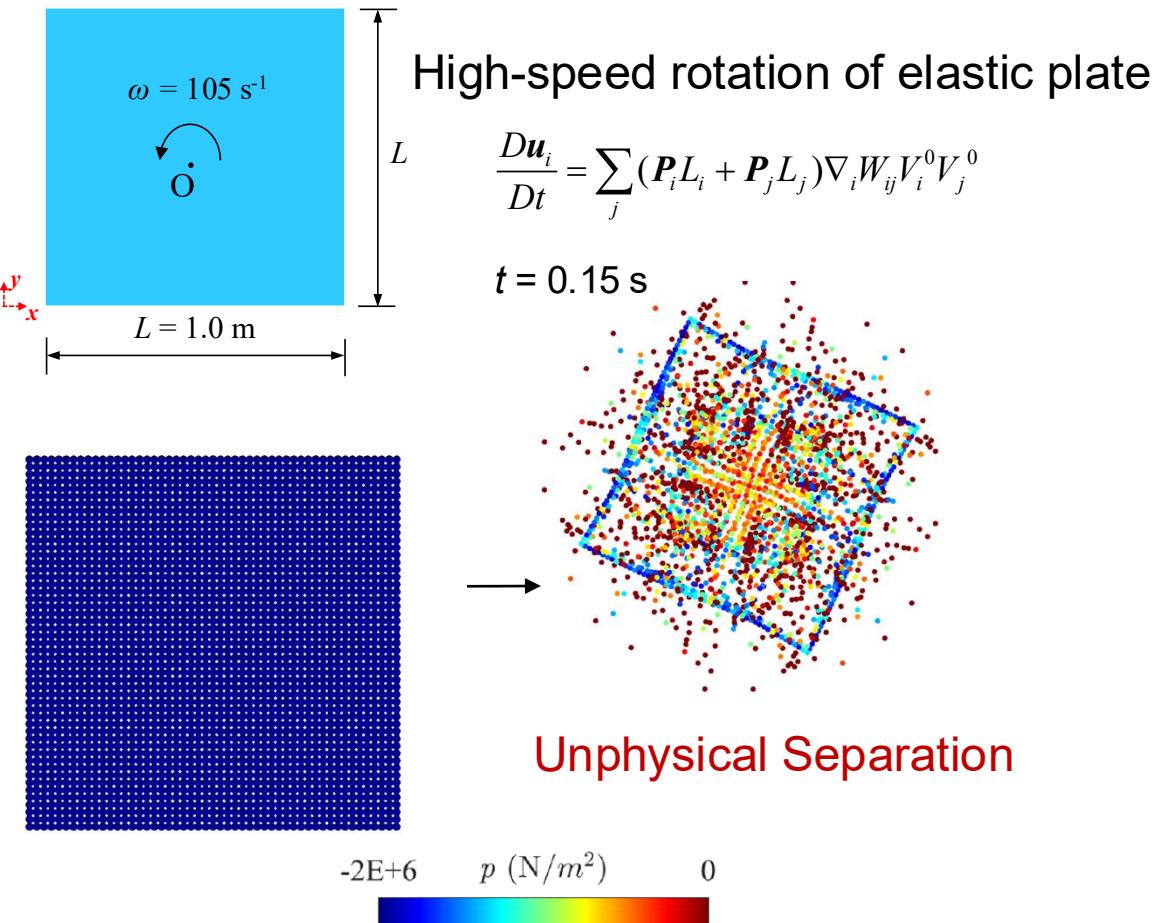
17th June 2025



Contents

- Motivation
- Enhanced FSI model
- Numerical validations

➤ Low solution accuracy

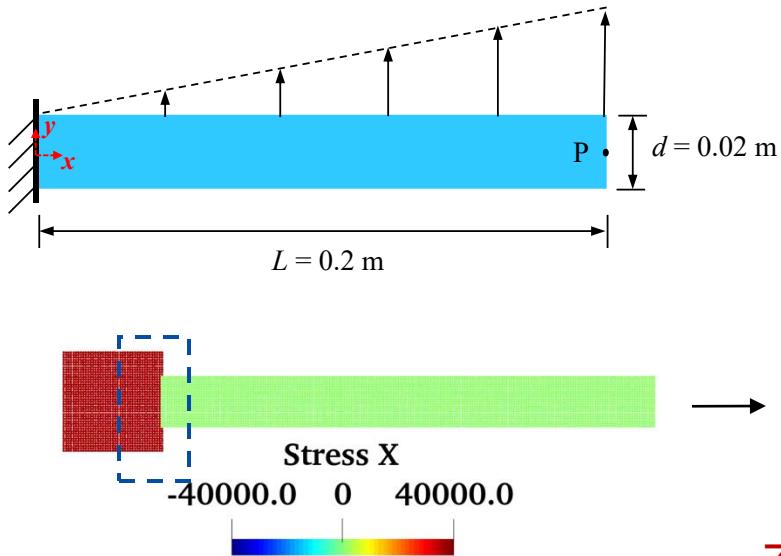


- Low solution accuracy leads to non-physical results
- Rank deficiency / Hourglass mode

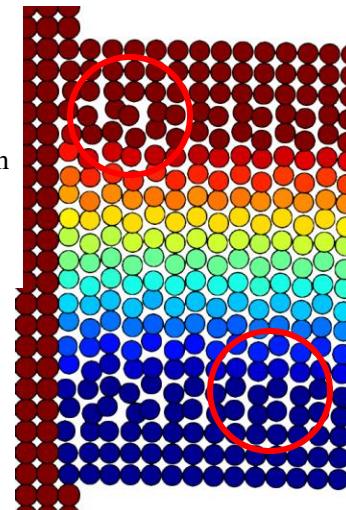
Simulated by Total Lagrangian SPH

➤ Serious hourglass modes

Free oscillation of a cantilever plate

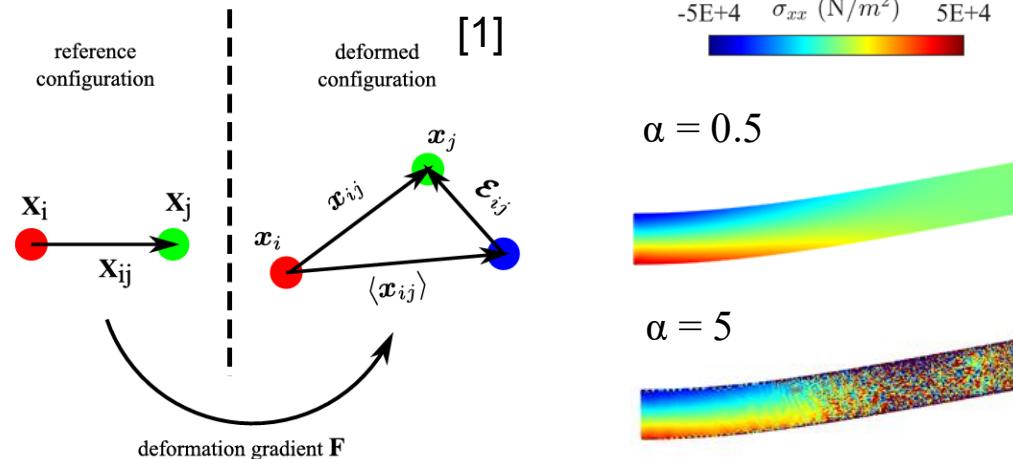


Unphysical clustering



Zigzag distribution

- Hourglass mode → unphysical particle distribution

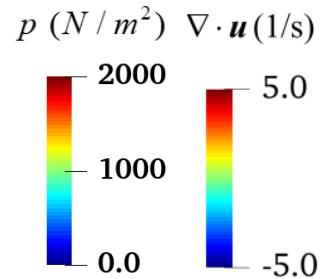


$$\mathbf{f}_i^{\text{HG}} = -\frac{1}{2} \alpha \sum_j \frac{V_{0i} V_{0j} W_{0ij}}{|\mathbf{r}_{ij}^0 \cdot \mathbf{r}_{ij}^0|} (E_i \delta_{ij}^i + E_j \delta_{ji}^j) \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|}$$

- Penalizes any local displacement not described by linear transformation of \mathbf{F}
- Parameter needs proper calibrations

[1] Ganzenmüller, Georg C. "An hourglass control algorithm for Lagrangian smooth particle hydrodynamics." *CMAME* (2015)

➤ Noises in fluid field



Dam-break with an elastic plate



Simulated by δ -SPH - TLSPH

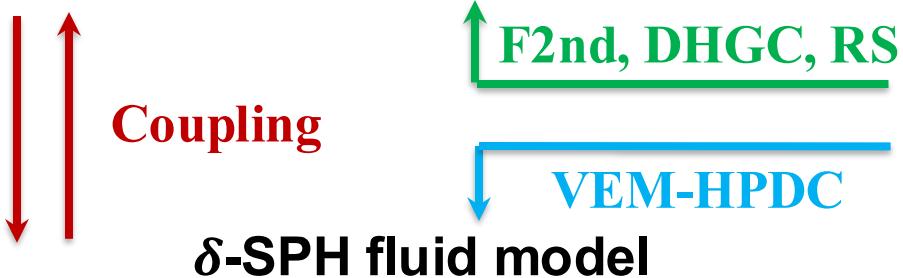
- Velocity divergence errors → pressure noises → FSI coupling → stress noises

Contents

- Motivation
- Enhanced FSI model
- Numerical validations

TLSPH structure model

$$\frac{D\mathbf{u}_i}{Dt} = \frac{1}{\rho_i^0} \sum_j \left(-\mathbf{P}_j \cdot \begin{bmatrix} \mathbf{A}_j^{0,1} \\ \mathbf{A}_j^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ji}^0 + \mathbf{P}_i \cdot \begin{bmatrix} \mathbf{A}_i^{0,1} \\ \mathbf{A}_i^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ij}^0 \right) \frac{\partial W_{ij}^0}{\partial r_{ij}^0} \frac{V_j^0}{|\mathbf{r}_{ij}^0|}$$



$$\frac{D\rho_i}{Dt} = -\rho_i \sum_j \mathbf{u}_{ij} \cdot \nabla_i W_{ij} V_j + D_i$$

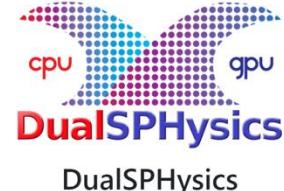
$$\frac{D\mathbf{u}_i}{Dt} = -\sum_j m_j \left(\frac{p_i + p_j}{\rho_i \rho_j} \right) \nabla_i W_{ij} + \sum_j m_j \left(\frac{4\nu \mathbf{r}_{ij} \cdot \nabla_i W_{ij}}{(\rho_i + \rho_j)(\mathbf{r}_{ij}^2 + \eta^2)} \right) \mathbf{u}_{ij}$$

$$p_i = c_f^2 (\rho_i - \rho_0)$$

Enhanced schemes

(1) F2nd (Second-order deformation gradient tensor \mathbf{F})

$$\mathbf{F}_i^S = \sum_j \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}^0|} \begin{bmatrix} \mathbf{A}_i^{0,1} \\ \mathbf{A}_i^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ij}^0 V_j^0 \frac{\partial W_{ij}^0}{\partial r_{ij}^0}$$



(2) DHGC (Dynamic Hourglass Control)

$$\mathbf{f}_i^{\text{HG}} = -\frac{1}{2} \sum_j \frac{\alpha_i^k + \alpha_j^k}{2} \frac{V_{0i} V_{0j} W_{0ij}}{|\mathbf{r}_{ij}^0 \cdot \mathbf{r}_{ij}^0|} (E_i \delta_{ij}^a + E_j \delta_{ji}^b) \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|} \quad \alpha_i^k = |\boldsymbol{\varepsilon}_i|^k / |\boldsymbol{\varepsilon}_i|^{k-N}$$

$$\boldsymbol{\varepsilon}_{ij}^i = \langle \mathbf{r}_{ij} \rangle^i - \mathbf{r}_{ij} = \mathbf{F}_i \mathbf{r}_{ij}^0 - \mathbf{r}_{ij}$$

(3) RS (Riemann SPH-based Stabilization term)

$$\frac{D\mathbf{u}_i}{Dt} = \frac{1}{\rho_i^0} \sum_j \left(-\mathbf{P}_j \cdot \begin{bmatrix} \mathbf{A}_j^{0,1} \\ \mathbf{A}_j^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ji}^0 + \mathbf{P}_i \cdot \begin{bmatrix} \mathbf{A}_i^{0,1} \\ \mathbf{A}_i^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ij}^0 \right) \frac{\partial W_{ij}^0}{\partial r_{ij}^0} \frac{V_j^0}{|\mathbf{r}_{ij}^0|} + \boldsymbol{\Pi}_i^{\text{RS}}$$

(4) VEM-HPDC (Velocity divergence Error Mitigating and Hyperbolic/Parabolic Divergence Cleaning schemes)

$$\frac{D\mathbf{u}_i}{Dt} = -2 \sum_j m_j \left(\frac{p_i + p_j}{\rho_i \rho_j} \right) \nabla_i W_{ij} + \mathbf{g} + \mathbf{a}_i^{\text{VEM}} - \nabla \psi_i$$

(1) **F2nd:** Improving accuracy of stress and strain computation

(2) **DHGC:** Suppressing zero-energy modes and dynamically adjusting hourglass control coefficient

(3) **RS:** Reducing high-frequency structure stress noises

(4) **VEM-HPDC:** Mitigating velocity divergence errors and fluid pressure noises

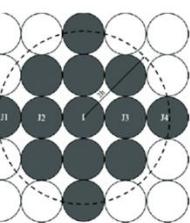
Traditional first-order discretization (F1st)



$$\frac{D\mathbf{u}_i}{Dt} = \sum_j (\mathbf{P}_i L_i + \mathbf{P}_j L_j) \nabla_i W_{ij} V_i^0 V_j^0$$

Second-order discretization (F2nd)

- 1st Order → 2nd Order, without reducing much computational efficiency
- The first three-dimensional 2nd order accuracy GPU-based structure solver



$$\mathbf{F} = \nabla_0 \mathbf{r}$$

Taylor series

$$\mathbf{M}_y^0 = \left[\frac{x_y^0}{|\mathbf{r}_y^0|}, \frac{y_y^0}{|\mathbf{r}_y^0|}, \frac{z_y^0}{|\mathbf{r}_y^0|}, \frac{1}{d_0} \frac{x_y^0 x_y^0}{|\mathbf{r}_y^0|}, \frac{1}{d_0} \frac{y_y^0 y_y^0}{|\mathbf{r}_y^0|}, \frac{1}{d_0} \frac{z_y^0 z_y^0}{|\mathbf{r}_y^0|}, \frac{1}{d_0} \frac{x_y^0 y_y^0}{|\mathbf{r}_y^0|}, \frac{1}{d_0} \frac{x_y^0 z_y^0}{|\mathbf{r}_y^0|}, \frac{1}{d_0} \frac{y_y^0 z_y^0}{|\mathbf{r}_y^0|} \right]^T$$

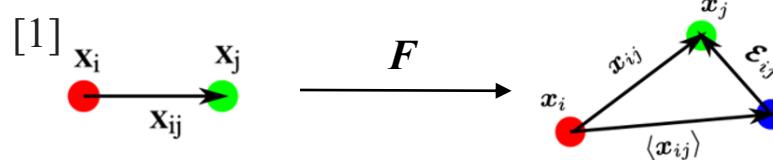
- Coefficient matrix \mathbf{A} is computed at the beginning

$$\frac{D\mathbf{u}_i}{Dt} = \frac{1}{\rho_i^0} \sum_j \left(-\mathbf{P}_j \cdot \begin{bmatrix} \mathbf{A}_j^{0,1} \\ \mathbf{A}_j^{0,2} \\ \mathbf{A}_j^{0,3} \end{bmatrix} \cdot \mathbf{M}_{ji}^0 + \mathbf{P}_i \cdot \begin{bmatrix} \mathbf{A}_i^{0,1} \\ \mathbf{A}_i^{0,2} \\ \mathbf{A}_i^{0,3} \end{bmatrix} \cdot \mathbf{M}_{ij}^0 \right) \frac{\partial W_{ij}^0}{\partial \dot{r}_{ij}^0} \frac{V_j^0}{|\mathbf{r}_{ij}^0|}$$

$$\begin{bmatrix} \mathbf{A}_i^{0,1} \\ \mathbf{A}_i^{0,2} \\ \vdots \\ \vdots \\ \mathbf{A}_i^{0,9} \end{bmatrix} = \begin{bmatrix} (M_{ij}^{0,1}, M_{ij}^{0,1}) & (M_{ij}^{0,1}, M_{ij}^{0,2}) & \dots & \dots & (M_{ij}^{0,1}, M_{ij}^{0,9}) \\ (M_{ij}^{0,2}, M_{ij}^{0,1}) & (M_{ij}^{0,2}, M_{ij}^{0,2}) & & & \vdots \\ \vdots & \vdots & & & \vdots \\ \vdots & \vdots & & & \vdots \\ (M_{ij}^{0,9}, M_{ij}^{0,1}) & \dots & \dots & \dots & (M_{ij}^{0,9}, M_{ij}^{0,9}) \end{bmatrix}^{-1}$$

[1] Khayyer et al. "An improved Riemann SPH-Hamiltonian SPH coupled solver for hydroelastic fluid-structure interactions." *EABE* (2024)

➤ Dynamic hourglass control (DHGC)



Traditional hourglass control scheme (HGC)

$$f_i^{\text{HG}} = -\frac{1}{2} \alpha \sum_j \frac{V_{0i} V_{0j} W_{0ij}}{|\mathbf{r}_{ij}^0 \cdot \mathbf{r}_{ij}^0|} (E_i \delta_{ij}^i + E_j \delta_{ji}^j) \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|} \quad \delta_{ij}^i = \frac{\epsilon_{ij}^i \cdot \mathbf{r}_{ij}}{r_{ij}} \text{ with } \epsilon_{ij}^i = \langle \mathbf{r}_{ij} \rangle^i - \mathbf{r}_{ij} = \mathbf{F}_i \mathbf{r}_{ij}^0 - \mathbf{r}_{ij}$$

- α is a constant and needs proper calibrations

Dynamic hourglass control scheme (DHGC)

$$f_i^{\text{HG}} = -\frac{1}{2} \sum_j \frac{\alpha_i^k + \alpha_j^k}{2} \frac{V_{0i} V_{0j} W_{0ij}}{|\mathbf{r}_{ij}^0 \cdot \mathbf{r}_{ij}^0|} (E_i \delta_{ij}^a + E_j \delta_{ji}^b) \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|} \quad \alpha_i^k = \frac{|\epsilon_i|^k}{|\epsilon_i|^{k-N}}$$

- Adaptively adjust based on error vector

➤ Riemann Stabilization (RS)

$$\frac{D\mathbf{u}_i}{Dt} = \frac{1}{\rho_i^0} \sum_j \left(-\mathbf{P}_j \cdot \begin{bmatrix} \mathbf{A}_j^{0,1} \\ \mathbf{A}_j^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ji}^0 + \mathbf{P}_i \cdot \begin{bmatrix} \mathbf{A}_i^{0,1} \\ \mathbf{A}_i^{0,2} \end{bmatrix} \cdot \mathbf{M}_{ij}^0 \right) \frac{\partial W_{ij}^0}{\partial r_{ij}^0} \frac{V_j^0}{|\mathbf{r}_{ij}^0|} + \boxed{\mathbf{\Pi}_i^{\text{RS}}}$$

$$\boxed{\mathbf{\Pi}_i^{\text{RS}}} = \sum_j \mathbf{P}_{ij}^{\text{RS}} \cdot \nabla_0 W_{ij} V_j^0, \mathbf{P}_{ij}^{\text{RS}} = \det(\mathbf{F}_i) \Pi_{ij}^{\text{RS}} \mathbf{F}_i^{-T}$$

$$\Pi_{ij}^{\text{RS}} = \frac{\beta c_0}{\rho_i} \frac{\rho_i + \rho_j}{2} \frac{\mathbf{u}_{ij} \cdot \mathbf{r}_{ij}}{|\mathbf{r}_{ij}|}; \beta = \max \left(0, \frac{u_L - u_R}{|u_L - u_R|} \right)$$

$$u_L = \mathbf{u}_i \cdot \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|}; u_R = \mathbf{u}_j \cdot \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij}|}$$

- Reduce noises in stress field
- No artificial parameters

- [1] Ganzenmüller, Georg C. "An hourglass control algorithm for Lagrangian smooth particle hydrodynamics." *CMAME* (2015)
[2] Khayyer et al. "An improved updated Lagrangian SPH method for structural modelling." *CPM* (2024)



➤ Velocity-divergence Error

Mitigating (VEM) [1]

$$p_a^{\text{VEM}} = c_s^2 d \rho_a = c_s^2 \Delta t \left(\frac{D \rho}{D t} \right)_a^{k-1} = -\rho_a c_s^2 \Delta t \langle \nabla \cdot \mathbf{u} \rangle_a^{k-1}$$

Pressure related to velocity divergence error



This error will be accumulated

$$\mathbf{a}_a^{\text{VEM}} = -\frac{1}{\rho_a} \sum_b F(p_a^{\text{VEM}}, p_b^{\text{VEM}}) \nabla_a W_{ab} V_b$$

$$F(p_a, p_b) = \begin{cases} p_b + p_a & (p_a \geq 0 \cup a \notin \Omega_{IN}) \\ p_b - p_a & (p_a < 0 \cap a \in \Omega_{IN}) \end{cases}$$

$$\frac{D \mathbf{u}_a}{D t} = -\sum_b m_b \left(\frac{F(p_a, p_b)}{\rho_a \rho_b} + \Pi_{ab} \right) \nabla_a W_{ab} + \mathbf{g}_a + \mathbf{a}_a^{\text{VEM}}$$

➤ Hyperbolic/Parabolic Divergence Cleaning (HPDC) [2]

$$\begin{cases} \left(\frac{D \mathbf{u}}{D t} \right)_\psi + \nabla \psi = 0 \\ \mathcal{D}(\psi) + \nabla \cdot \mathbf{u} = 0 \end{cases}$$

Hyperbolic term

$$\frac{\partial^2 (\nabla \cdot \mathbf{u})}{\partial t} - \boxed{c_h^2 \nabla^2 (\nabla \cdot \mathbf{u})} + \boxed{\frac{c_h^2}{c_p^2} \frac{\partial (\nabla \cdot \mathbf{u})}{\partial t}} = 0$$

$$\frac{D \mathbf{u}_a}{D t} = -2 \sum_b m_b \left(\frac{p^*}{\rho_a \rho_b} \right) \nabla_a W_{ab} + \mathbf{g} + \mathbf{a}_a^{\text{VEM}} \boxed{- \nabla \psi_a}$$

$\psi = 0$ for boundary particles

[1] Khayyer et al. "Enhanced resolution of the continuity equation in explicit weakly compressible SPH simulations of incompressible free-surface fluid flows." *AMM* (2023)

[2] Fourtakas et al. "Divergence cleaning for weakly compressible smoothed particle hydrodynamics." *C&F* (2025)

Contents

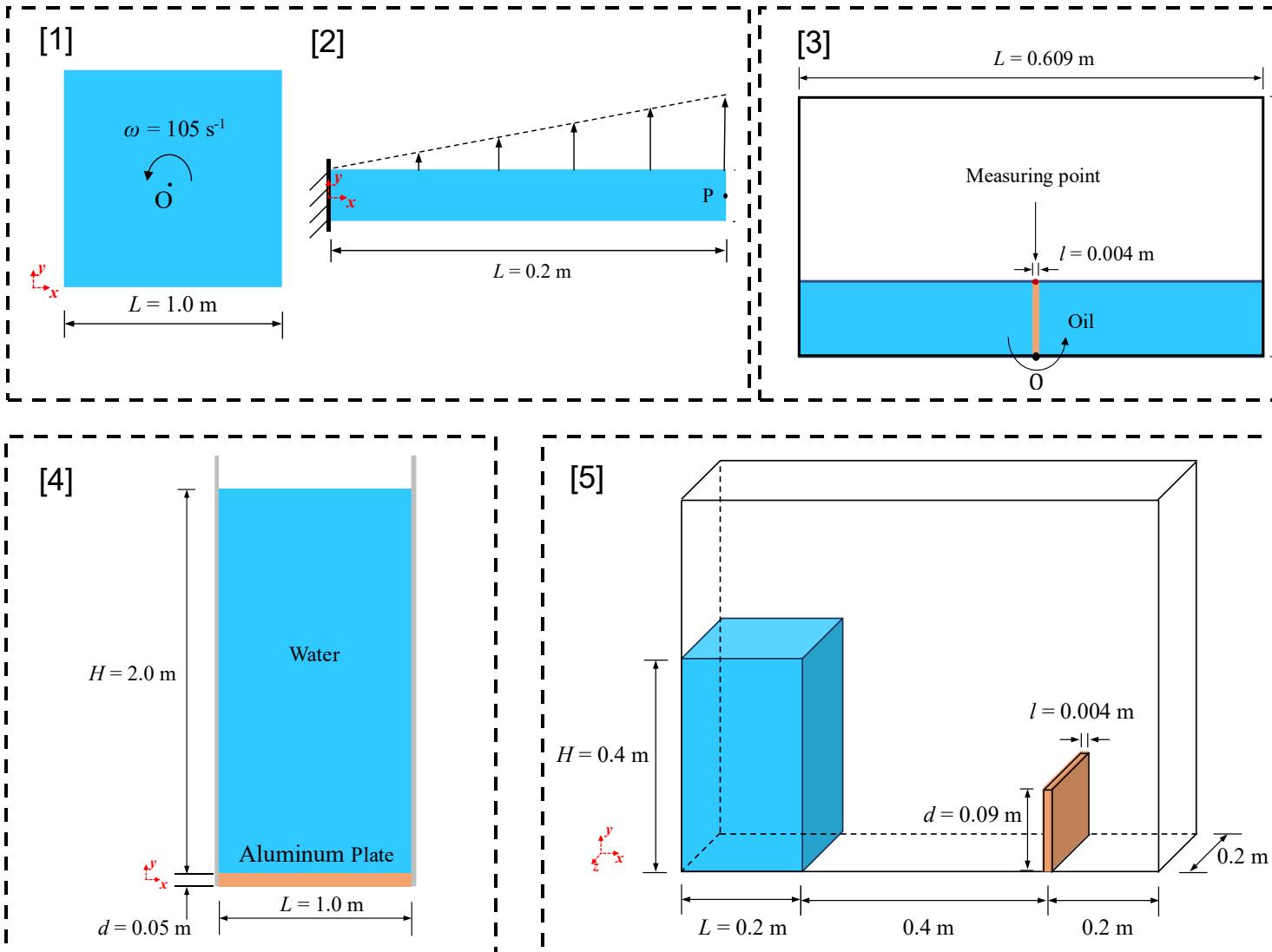
- Motivation
- Enhanced FSI model
- Numerical validations

➤ Five cases

- [1] High speed rotation of elastic square
- [2] Free oscillation of a cantilever plate
- [3] Liquid sloshing with an elastic plate
- [4] Hydrostatic water column
- [5] Dam break with an elastic plate

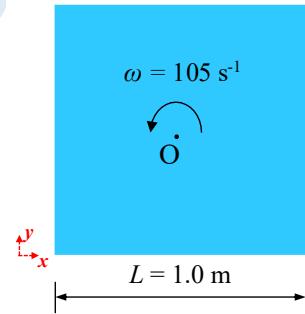
References

- [1] Lee, C. H, et al. "Development of a cell centred upwind finite volume algorithm for a new conservation law formulation in structural dynamics." *C&S* (2013)
- [2] Gray, James P., et al. "SPH elastic dynamics." *CMAME* (2001)
- [3] Idelsohn, S. R., et al. "Interaction between an elastic structure and free-surface flows: experimental versus numerical comparisons using the PFEM." *CM* (2008)
- [4] Fourey, G., et al. "An efficient FSI coupling strategy between smoothed particle hydrodynamics and finite element methods." *CPC* (2017)
- [5] Liao, et al. "Free surface flow impacting on an elastic structure: Experiment versus numerical simulation." *APOR* (2015)



3.1

High speed rotation of elastic square: Pressure field



$L = 1 \text{ m}, dp = 0.01 \text{ m}$

$E = 2 \text{ MPa}, \rho = 1000 \text{ kg/m}^3$

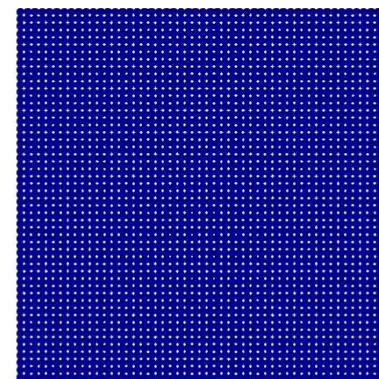
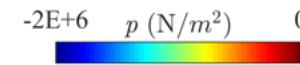
$\nu = 0.49, C_s = 242 \text{ m/s}$

Predictor-Corrector,

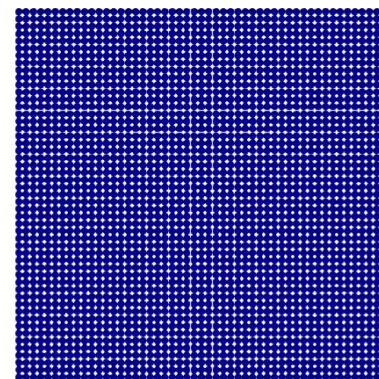
$C_{\text{CFL}} = 0.2$

5th Wendland C2, $h/dp = 2$

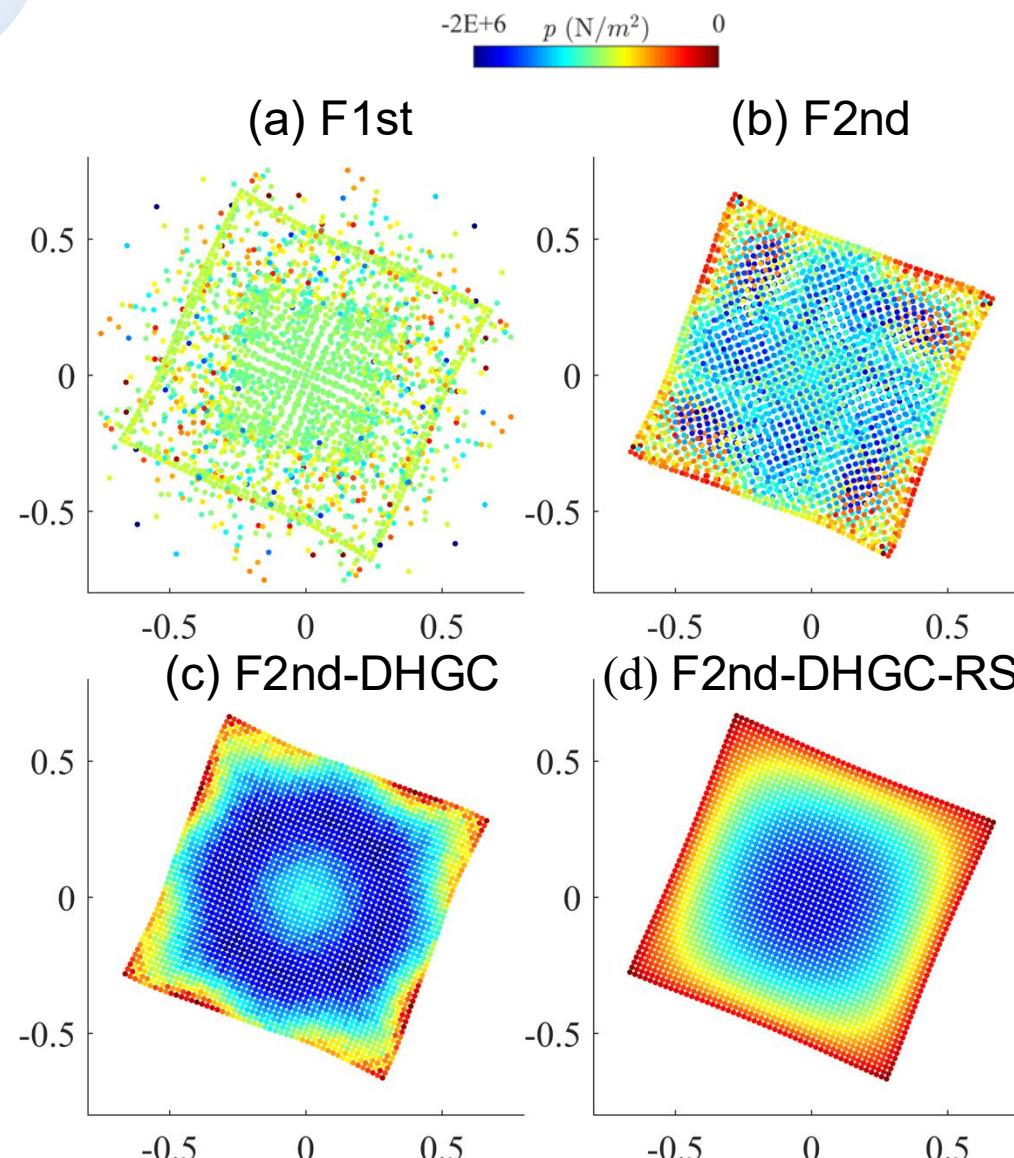
Original TLSPH model



Present model

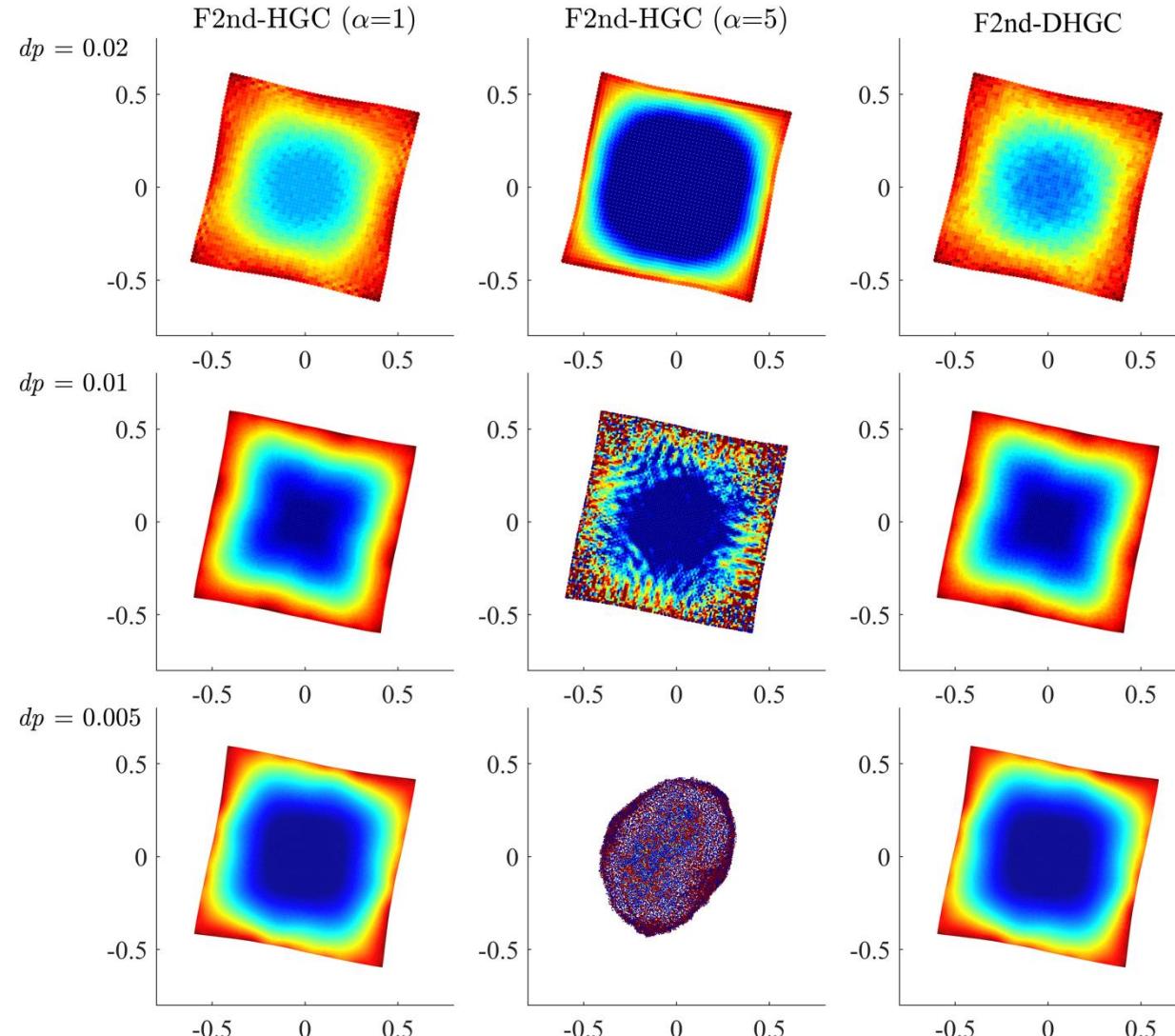


- Improves accuracy, reduces noises and effectively suppresses hourglass modes



- F2nd → improves the solution accuracy
- DHGC → mitigates hourglass modes
- RS → reduces stress noises

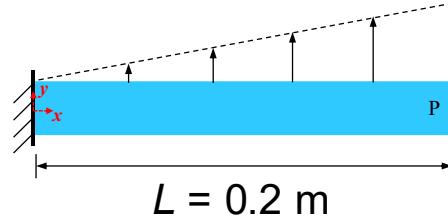
➤ Convergence analysis



- HGC: noticeable unphysical stress fluctuations ($\alpha = 5$) at $dp = 0.01, 0.005$
- DHGC: resolution-independent

3.2

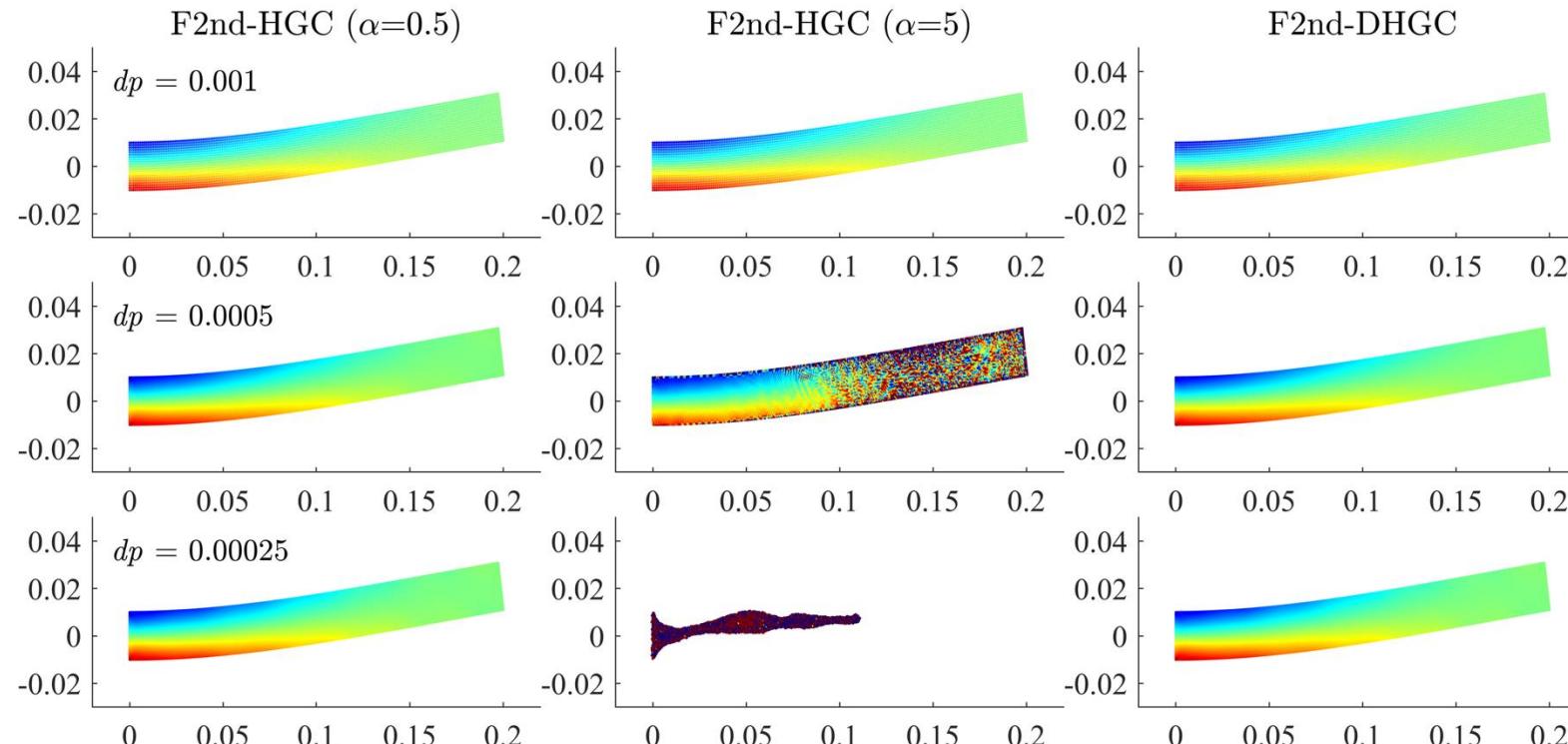
Free oscillation of a cantilever plate: Stress field



$$E = 2 \text{ MPa}, \rho = 1000 \text{ kg/m}^3, \nu = 0.3975$$

$$\begin{cases} u_y(x) = \xi c_s \frac{f(x)}{f(L)} \\ f(x) = (\cos kL + \cosh kL)(\cosh kx - \cos kx) + (\sin kL - \sinh kL)(\sinh kx - \sin kx) \end{cases}$$

$$-5E+4 \quad \sigma_{xx} (\text{N/m}^2) \quad 5E+4$$



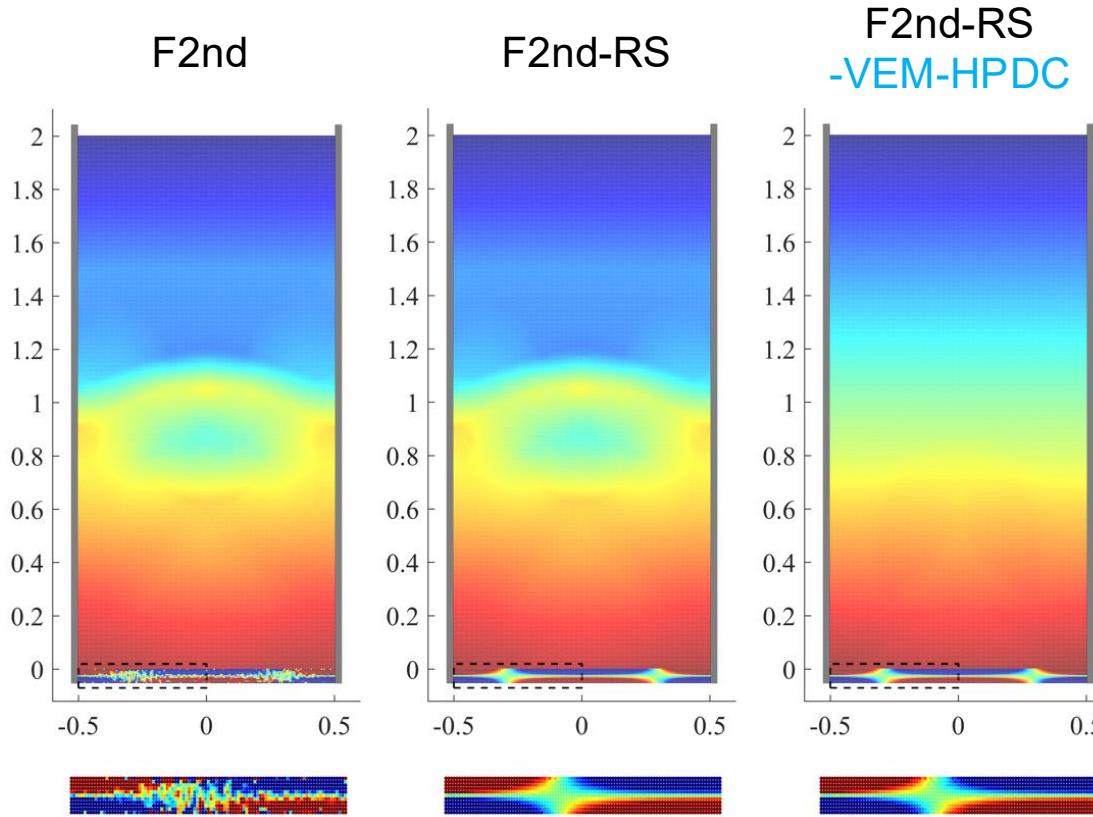
- HGC ($\alpha = 5$) \rightarrow excessive hourglass control \rightarrow over-stiffness
- DHGC \rightarrow Dynamically adjusted based on the errors

3.3

Hydrostatic water column & Sloshing

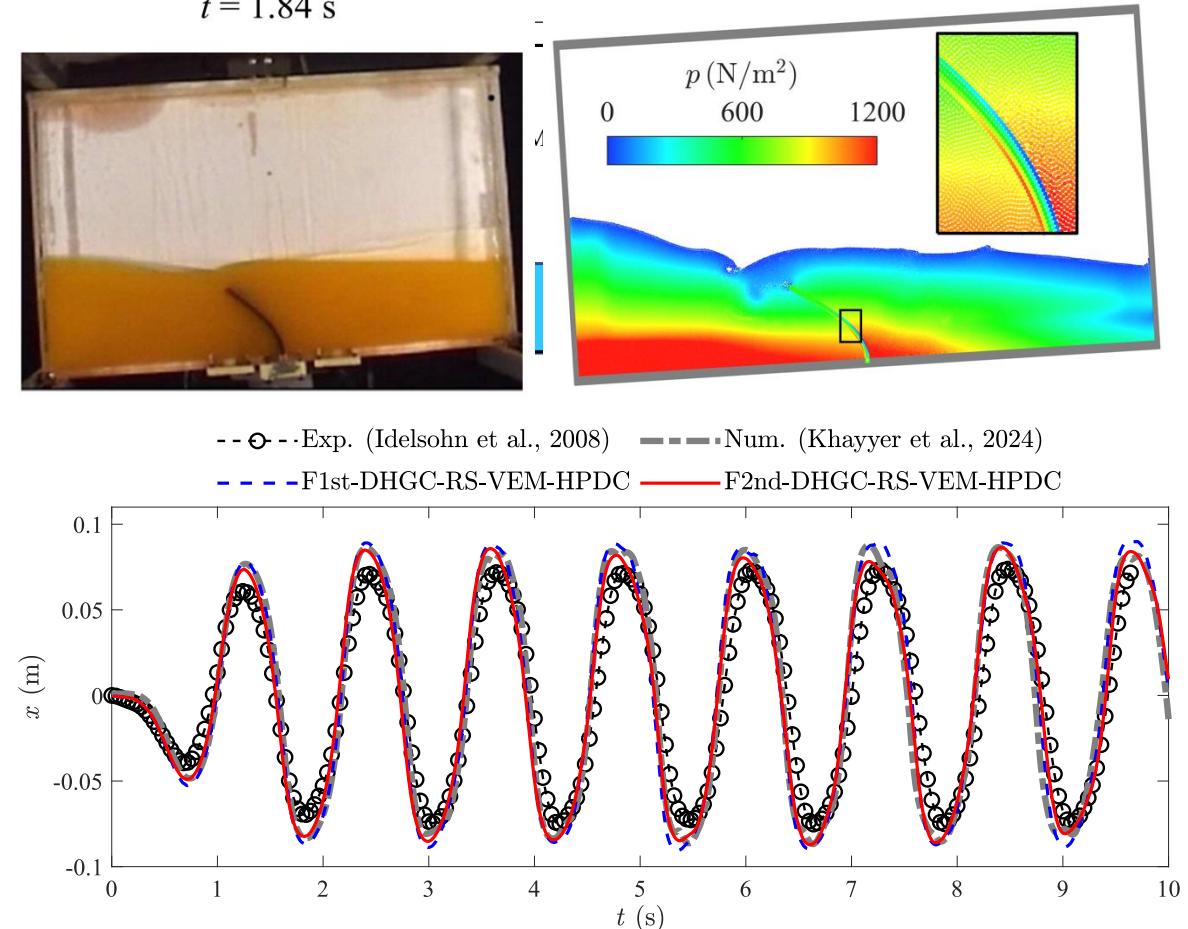


$$E = 67.5 \text{ Gpa}, \rho = 2700 \text{ kg/m}^3, \nu = 0.34$$



$$E = 6 \text{ Mpa}, \rho = 1100 \text{ kg/m}^3, \nu = 0.49$$

$$t = 1.84 \text{ s}$$

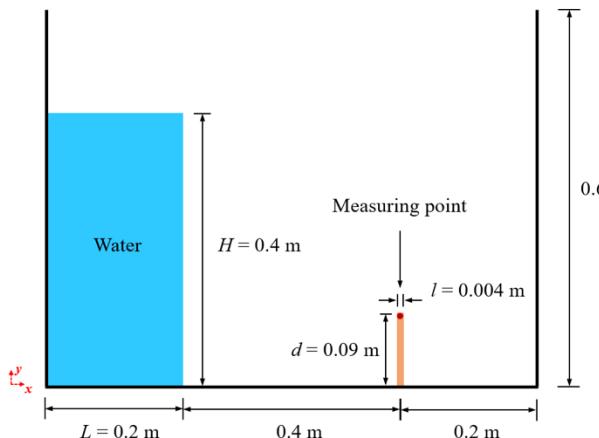


- Reduces noises in both fluid and structure fields

- The displacement history predicted by the present model is closer to the experimental result

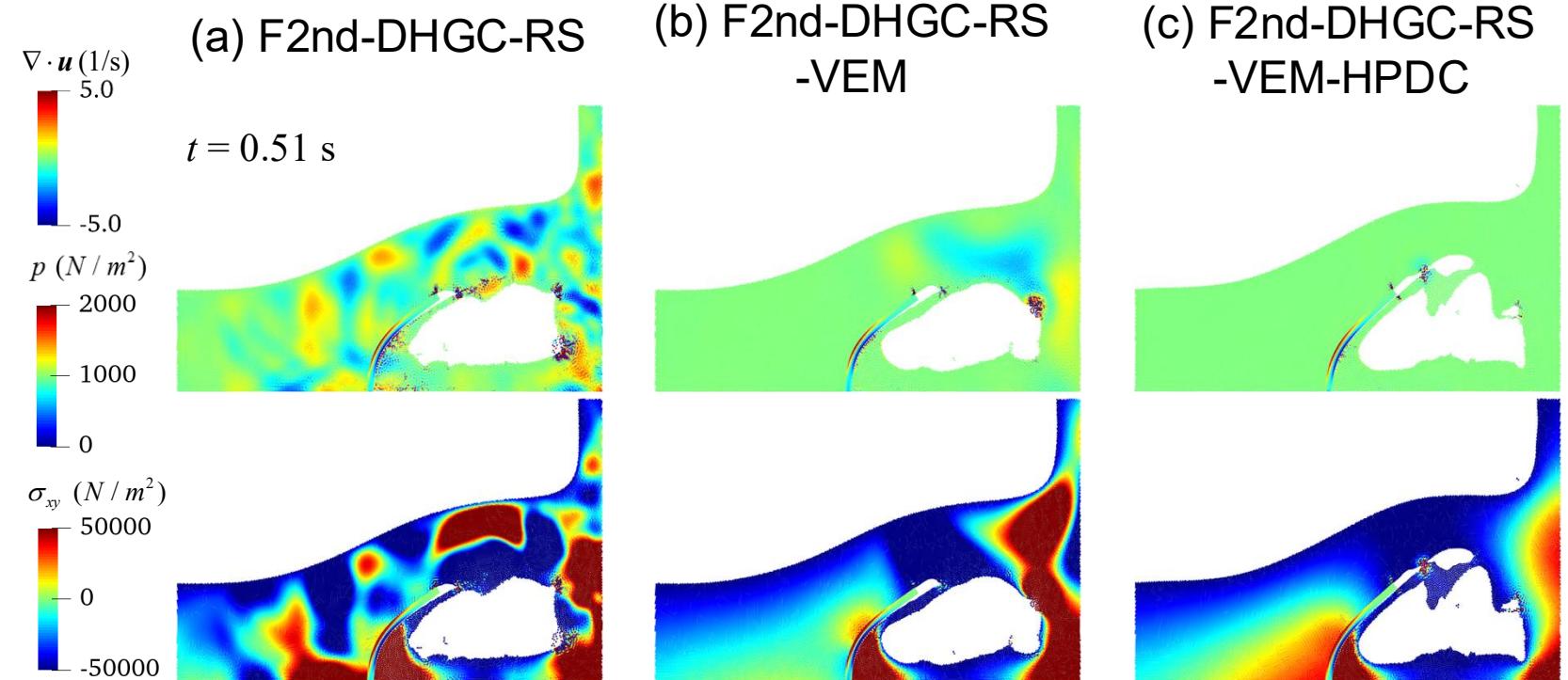
3.4

Dam break with elastic plate : Pressure and velocity divergence fields



$$E = 12 \text{ MPa}$$

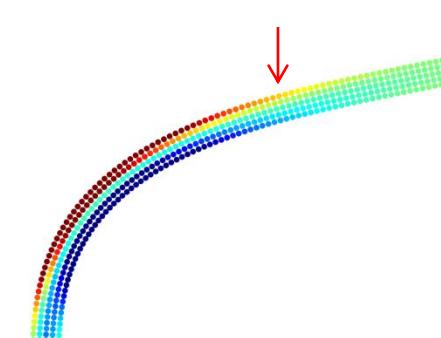
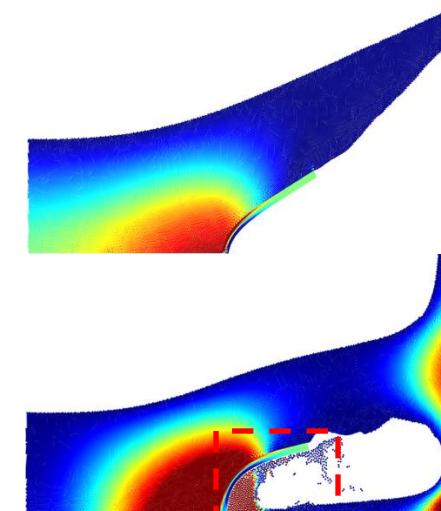
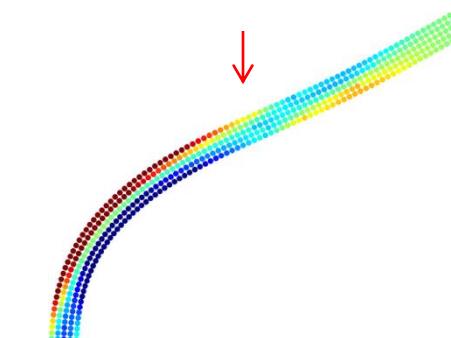
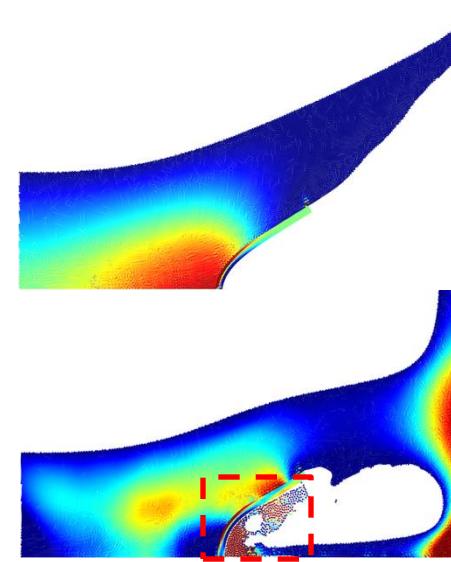
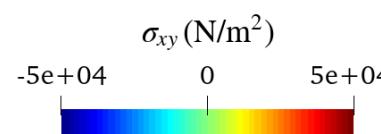
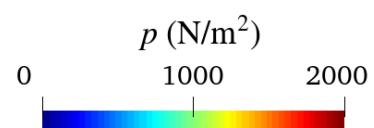
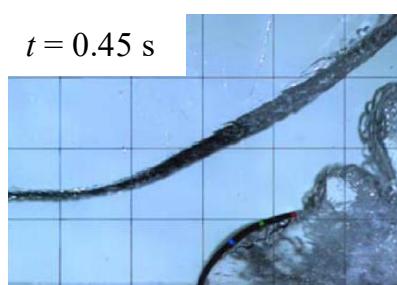
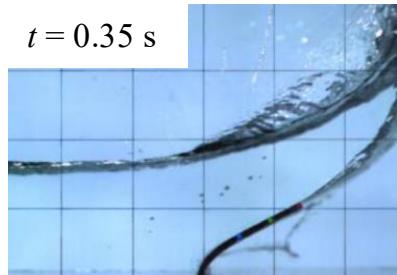
$$\rho = 1161.5 \text{ kg/m}^3 \quad \nu = 0.4$$



- VEM, HPDC → Reduces velocity divergence error and enhances pressure field

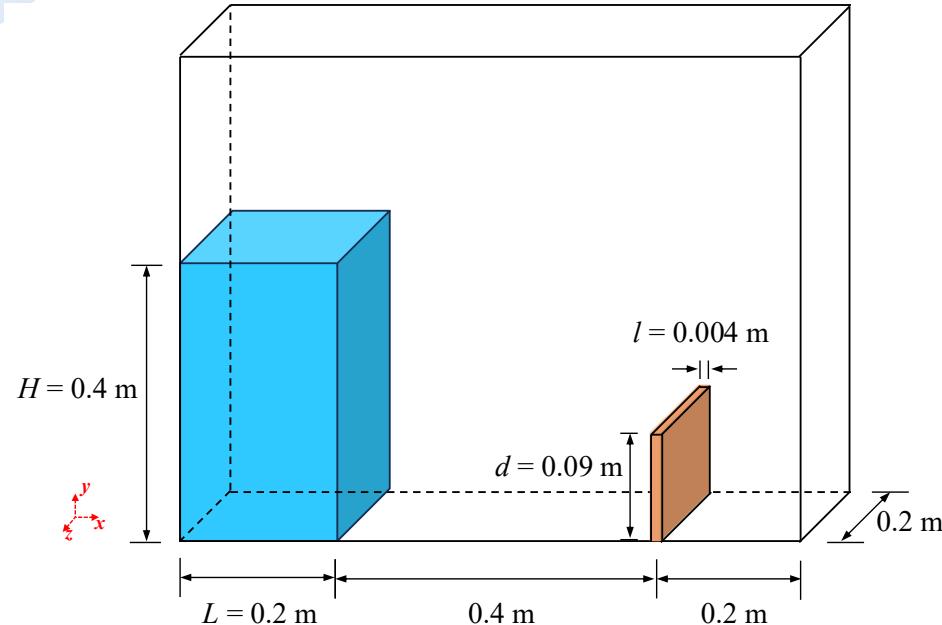
Exp. (Liao et al., 2015)

(a) F2nd-DHGC-RS

(b) F2nd-DHGC-RS
-VEM-HPDC

- Satisfactory agreement with the experimental photos
- Divergence cleaning → reduces fluid pressure noises → enhances the fluid-structure coupling → reduces stress noises

3.4 Dam break with elastic plate : 3D simulation

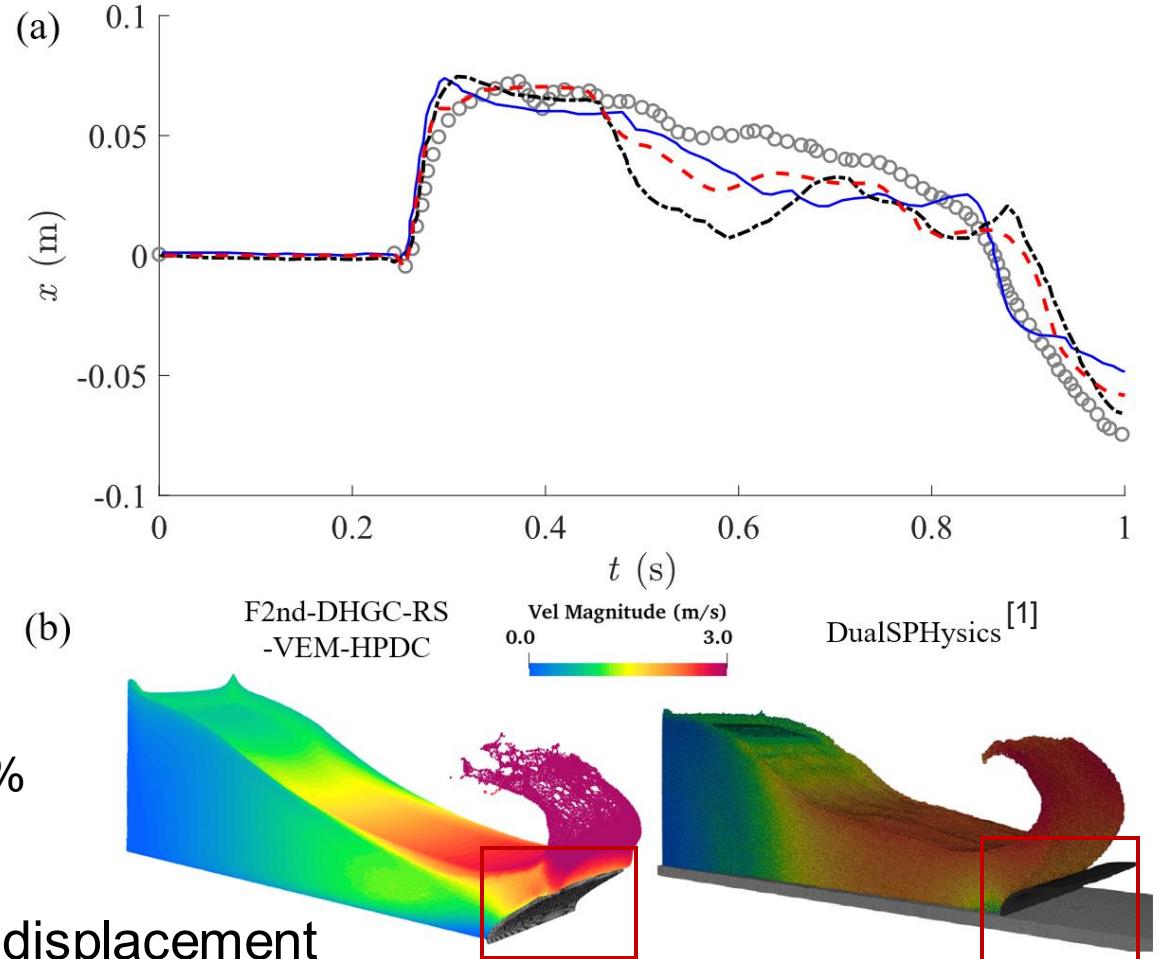


$dp = 0.001 \text{ m}$, $t = 1 \text{ s}$, $N \approx 24,000,000$, $\approx 60 \text{ hours}$

Computational time and memory increases $\approx 20\%$ and 25%

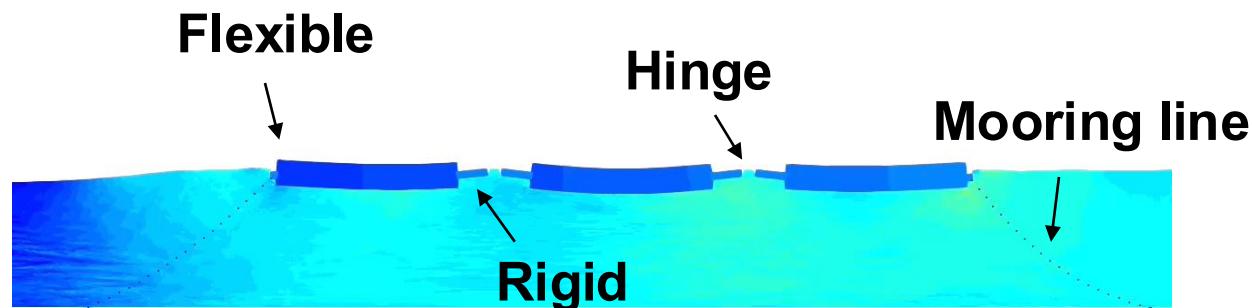
■ Improved accuracy in predicting deformation and displacement

- Exp. (Liao et al., 2015)
- Num. (Khayyer et al., 2021)
- Num. (O'Connor and Rogers, 2021)
- - - F2nd-DHGC-RS-VEM-HPDC



[1] O'Connor and Rogers. "A fluid–structure interaction model for free-surface flows and flexible structures using smoothed particle hydrodynamics on a GPU." *JFS* (2021)

- F2nd → Improve solution accuracy of TLSPH
- DHGC → Mitigate hourglass modes and being case- and resolution-independent
- RS → Reduce noises in structure stress field and being parameter free
- VEM, HPDC → Reduce noises in fluid pressure field, enhancing the fluid structure coupling
- On going/ future work → Engineering applications: Flexible floating arrays





Thank you for your attention



Welcome to ZJU @ Zhoushan