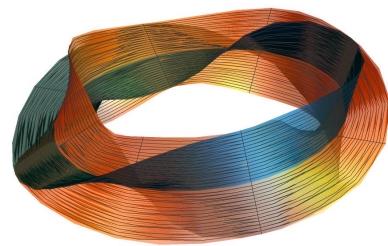


Topology and geometry of non-Hermitian Möbius bands and their deformations

Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program



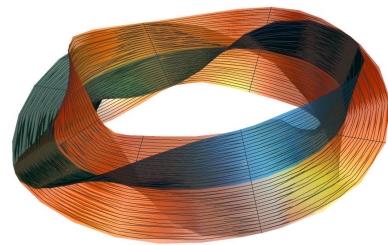
Yifei Zhu

Southern University of Science and Technology

2024.5.26

Topology and geometry of non-Hermitian Möbius bands and their deformations

Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program



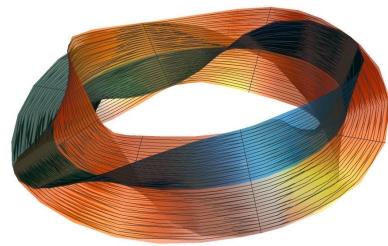
Yifei Zhu

Southern University of Science and Technology

2024.5.26

Topology and geometry of non-Hermitian Möbius bands and their deformations

Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program



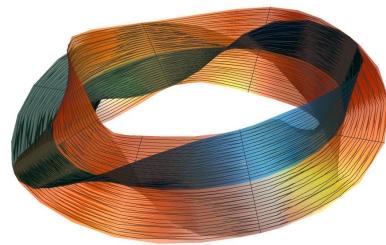
Yifei Zhu

Ongoing joint work with H. Jia, J. Hu, C. T. Chan

W. Yang, Z. Fang, C. Huang, et al.

Topology and geometry of non-Hermitian Möbius bands and their deformations

Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program



Yifei Zhu

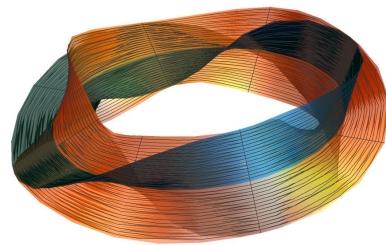
Physicists at HKUST

Ongoing joint work with H. Jia, J. Hu, C. T. Chan

W. Yang, Z. Fang, C. Huang, et al.

Topology and geometry of non-Hermitian Möbius bands and their deformations

Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program



Yifei Zhu

Physicists at HKUST

Ongoing joint work with H. Jia, J. Hu, C. T. Chan

W. Yang, Z. Fang, C. Huang, et al.

Math students at SUSTech

Topology and geometry of non-Hermitian Möbius bands and their deformations

**Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program**

Disclaimers: 1. Physics + math modeling, *not* mathematical physics

Yifei Zhu

Ongoing joint work with H. Jia, J. Hu, C. T. Chan
W. Yang, Z. Fang, C. Huang, et al.

Topology and geometry of non-Hermitian Möbius bands and their deformations

**Explicit examples of rank-2 and rank-3 Higgs bundles
in the contexts of
quantum materials and geometric Langlands program**

Disclaimers: 1. Physics + math modeling, *not* mathematical physics
2. Some parts completed, more in progress

Yifei Zhu

Ongoing joint work with H. Jia, J. Hu, C. T. Chan
W. Yang, Z. Fang, C. Huang, et al.

Motivations: Quantum materials and their math modeling

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials**

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ... physical properties*

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ... physical properties* at the *macroscopic* level

*Holography, optical devices,
absorption devices, ...*

Motivations: Quantum materials and their math modeling

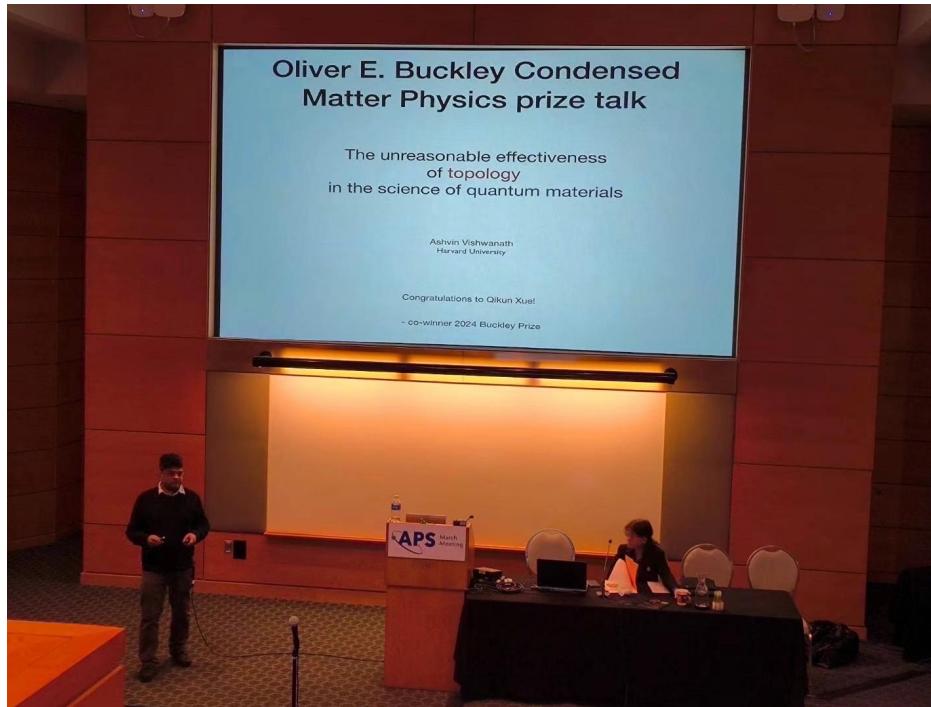
As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the **macroscopic** level that arise from the interactions of their electrons at the **microscopic** level

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause **unique and unexpected behaviors**.

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.



Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.

- *The unreasonable effectiveness of topology in the science of quantum materials*

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.

- *The unreasonable effectiveness of topology in the science of quantum materials*, Ashvin Vishwanath of Harvard delivering the Buckley prize talk at this year's APS March Meeting in Minneapolis

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.

- *The unreasonable effectiveness of topology in the science of quantum materials*, Ashvin Vishwanath of Harvard delivering the Buckley prize talk at this year's APS March Meeting in Minneapolis (薛其坤 of Tsinghua and SUSTech was the co-winner of the prize)

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.

- *The unreasonable effectiveness of topology in the science of quantum materials*, Ashvin Vishwanath of Harvard delivering the Buckley prize talk at this year's APS March Meeting in Minneapolis (薛其坤 of Tsinghua and SUSTech was the co-winner of the prize)
- U.S. Department of Energy, Office of Science. *Basic research needs for quantum materials: Research to discover, harness, and exploit exotic electronic properties* (brochure), 2016.

Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.

- *The unreasonable effectiveness of topology in the science of quantum materials*, Ashvin Vishwanath of Harvard delivering the Buckley prize talk at this year's APS March Meeting in Minneapolis (薛其坤 of Tsinghua and SUSTech was the co-winner of the prize)
- U.S. Department of Energy, Office of Science. *Basic research needs for quantum materials: Research to discover, harness, and exploit exotic electronic properties* (brochure), 2016.
- 方忠 等, “拓扑电子材料计算预测”, 2023年度国家自然科学奖一等奖

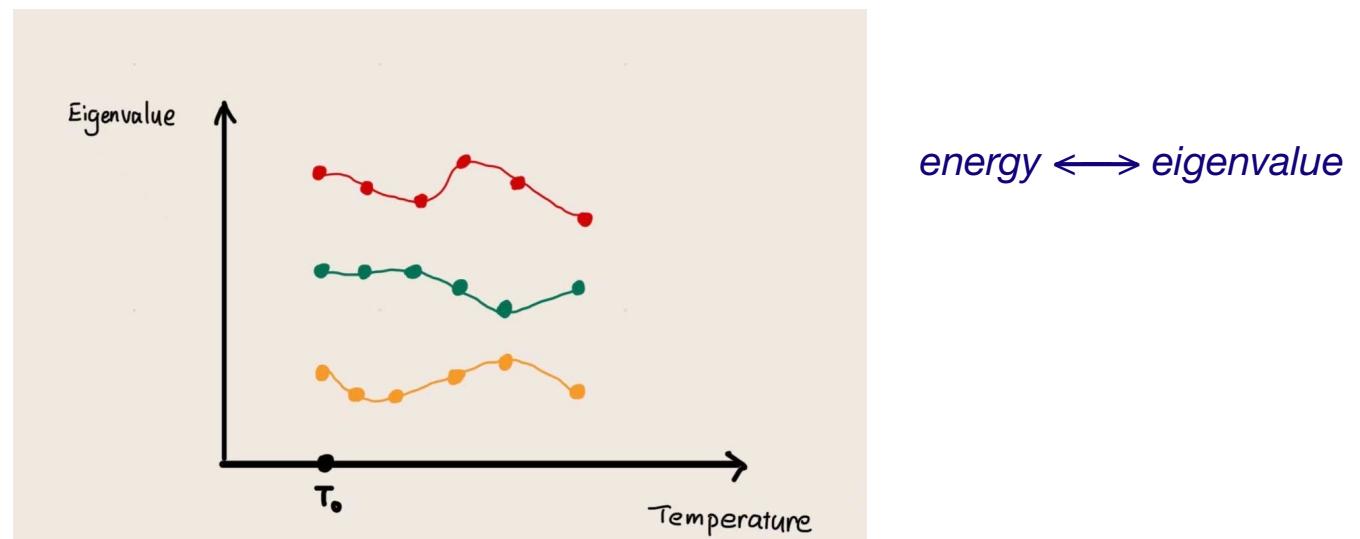
Motivations: Quantum materials and their math modeling

As an umbrella term in condensed matter physics, **quantum materials** refer to solids with *exotic/exceptional/not found in nature ...* physical properties at the *macroscopic* level that arise from the interactions of their electrons at the *microscopic* level, beginning at atomic and subatomic scales where the extraordinary effects of quantum mechanics cause unique and unexpected behaviors.

- *The unreasonable effectiveness of topology in the science of quantum materials*, Ashvin Vishwanath of Harvard delivering the Buckley prize talk at this year's APS March Meeting in Minneapolis (薛其坤 of Tsinghua and SUSTech was the co-winner of the prize)
- U.S. Department of Energy, Office of Science. *Basic research needs for quantum materials: Research to discover, harness, and exploit exotic electronic properties* (brochure), 2016.
- 方忠 等, “**拓扑电子材料计算预测**”, 2023年度国家自然科学奖一等奖
- 第一届魅丽数学与交叉应用会议“**数学与生物医药、数学与先进材料**”, 2024年5月, 苏州

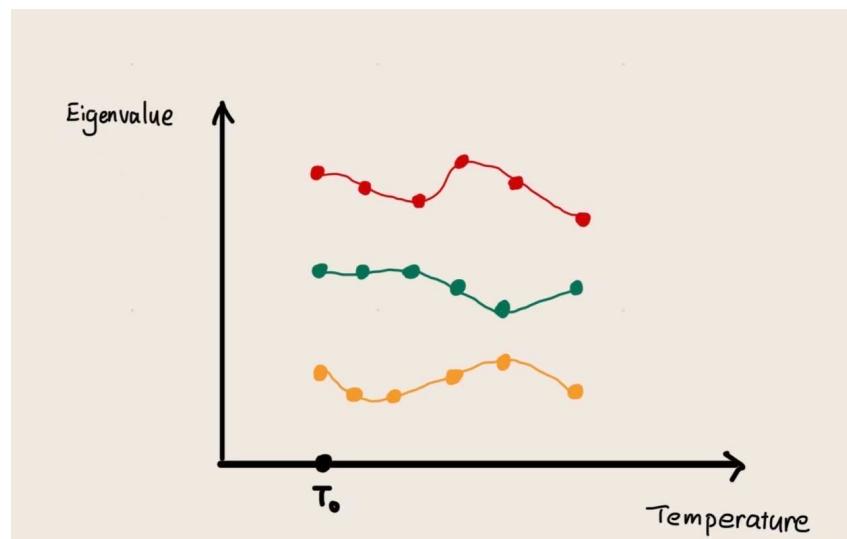
Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy *band structures* therein



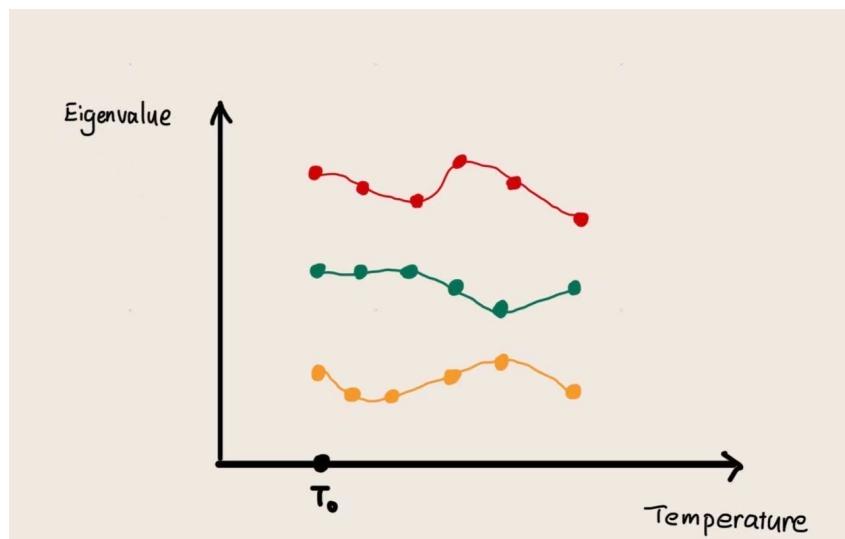
Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians*



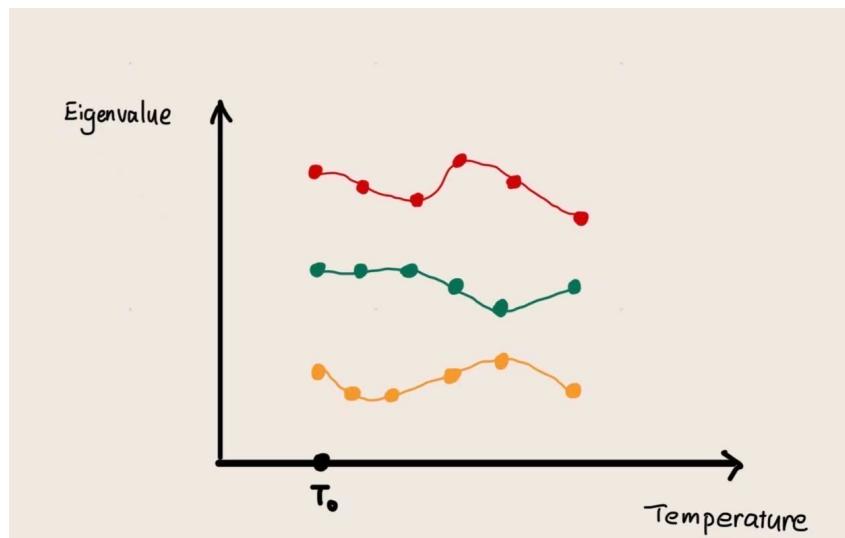
Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems



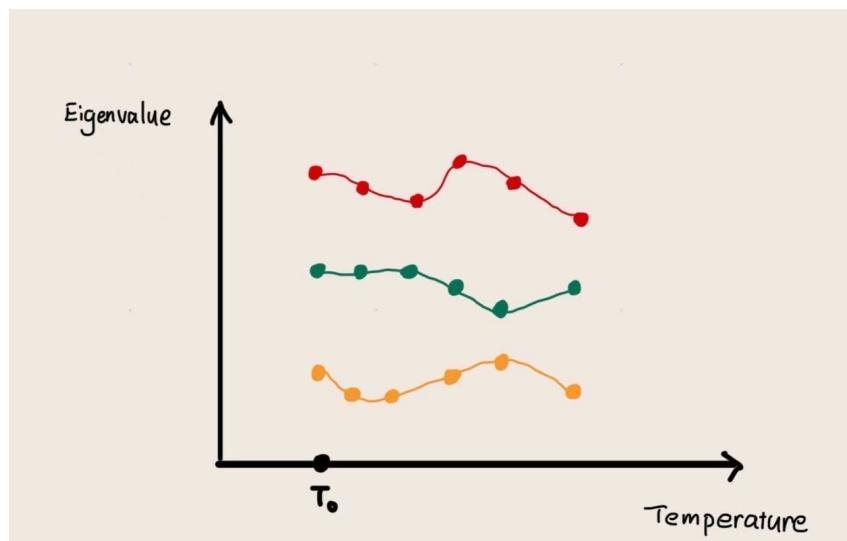
Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries]



Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed **symmetries**]

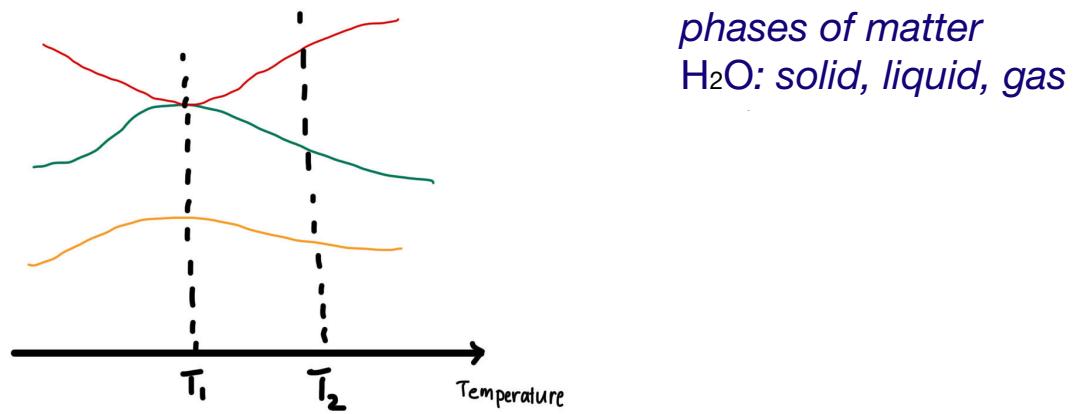


*Hermitian vs.
non-Hermitian*

*real eigenvalues
(observable
energies) vs.
eigenvalues with
imaginary part
(counts for
energy exchange
with surrounding
environment or
other systems)*

Motivations: Quantum materials and their math modeling

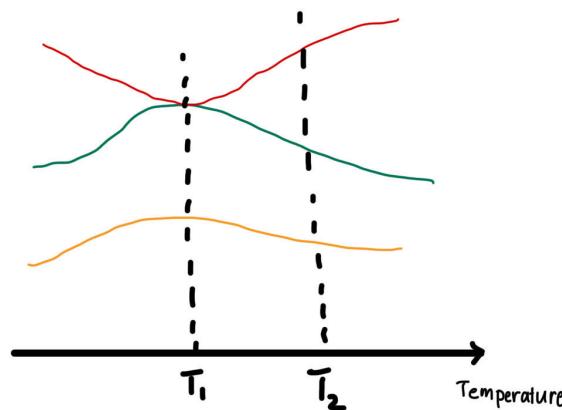
Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, **singularity/degeneracy** in the relevant moduli spaces



T_1 : singular points (points where eigenvalues degenerate)
 $H(T_1)$: gapless Hamiltonian
 $H(T_2)$: gapped Hamiltonian

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, **singularity/degeneracy** in the relevant moduli spaces, against which fine-tuning a system leads to **exceptional properties** of solid materials.



T_1 : singular points (points where eigenvalues degenerate)
 $H(T_1)$: gapless Hamiltonian
 $H(T_2)$: gapped Hamiltonian

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by **experimental** realization, engineering, ...

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by experimental realization, engineering, ... (though there is approach the other way around).

Guzman et al., Model-free characterization of topological edge and corner states in mechanical networks, PNAS 2024.

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy *band structures* therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed **symmetries**] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by experimental realization, engineering, ... (though there is approach the other way around).

In collaboration with physicists, our preliminary works explored the intriguing topological structures arising from certain novel ***non-Hermitian*** systems

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by experimental realization, engineering, ... (though there is approach the other way around).

In collaboration with physicists, our preliminary works explored the intriguing topological structures arising from certain novel *non-Hermitian* systems, whose moduli spaces have **stratified, non-isolated** singularities,

- Hongwei Jia, Ruo-Yang Zhang, Jing Hu, Yixin Xiao, Shuang Zhang, **Yifei Zhu**, and C. T. Chan. *Topological classification for intersection singularities* of exceptional surfaces in pseudo-Hermitian systems. **Communication Physics**, 6:293, 2023.
- Jing Hu, Ruo-Yang Zhang, Yixiao Wang, Xiaoping Ouyang, **Yifei Zhu**, Hongwei Jia, and Che Ting Chan. *Non-Hermitian swallowtail catastrophe revealing transitions among diverse topological singularities*. **Nature Physics**, 19:1098–1103, 2023.

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by **experimental realization, engineering, ...** (though there is approach the other way around).

In collaboration with physicists, our preliminary works explored the intriguing topological structures arising from certain novel *non-Hermitian* systems, whose moduli spaces have stratified, non-isolated singularities, as well as their **circuit realizations**

- Hongwei Jia, Ruo-Yang Zhang, Jing Hu, Yixin Xiao, Shuang Zhang, **Yifei Zhu**, and C. T. Chan. *Topological classification for intersection singularities of exceptional surfaces in pseudo-Hermitian systems.* **Communication Physics**, 6:293, 2023.
- Jing Hu, Ruo-Yang Zhang, Yixiao Wang, Xiaoping Ouyang, **Yifei Zhu**, Hongwei Jia, and Che Ting Chan. *Non-Hermitian swallowtail catastrophe revealing transitions among diverse topological singularities.* **Nature Physics**, 19:1098–1103, 2023.

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by experimental realization, engineering, ... (though there is approach the other way around).

In collaboration with physicists, our preliminary works explored the intriguing topological structures arising from certain novel *non-Hermitian* systems, whose moduli spaces have stratified, non-isolated singularities, as well as their circuit realizations and **unconventional physical consequences**.

- Hongwei Jia, Jing Hu, Ruo-Yang Zhang, Yixin Xiao, Dongyang Wang, Mudi Wang, Shaojie Ma, Xiaoping Ouyang, **Yifei Zhu**, and C. T. Chan. *Anomalous bulk-edge correspondence intrinsically beyond line-gap topology in non-Hermitian swallowtail gapless phase*, 2024. **Preprint**.
- Hongwei Jia, Jing Hu, Ruo-Yang Zhang, Yixin Xiao, Dongyang Wang, Mudi Wang, Shaojie Ma, Xiaoping Ouyang, **Yifei Zhu**, and C. T. Chan. *Unconventional topological edge states* beyond the paradigms of line-gap topology, 2024. **Preprint**.

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by experimental realization, engineering, ... (though there is approach the other way around).

In collaboration with physicists, our preliminary works explored the intriguing topological structures arising from certain novel *non-Hermitian* systems, whose moduli spaces have stratified, non-isolated singularities, as well as their circuit realizations and unconventional physical consequences. However, the mathematical modeling was rather **ad hoc** and the topological classifications remain **incomplete**.

Motivations: Quantum materials and their math modeling

Mathematical modeling of electronic energy band structures therein concerns topological/homotopical classification of *Hamiltonians* [= quantum mechanical systems = (families of) matrices with prescribed symmetries] and, in particular, singularity/degeneracy in the relevant moduli spaces, against which fine-tuning a system leads to exceptional properties of solid materials. This mathematical modeling is then followed by experimental realization, engineering, ... (though there is approach the other way around).

In collaboration with physicists, our preliminary works explored the intriguing topological structures arising from certain novel *non-Hermitian* systems, whose moduli spaces have stratified, non-isolated singularities, as well as their circuit realizations and unconventional physical consequences. However, the mathematical modeling was rather **ad hoc** and the topological classifications remain **incomplete**.

Thanks to Hopf bundles and Higgs bundles as **eigenbundles**, we now have a **conceptually more systematic**, visibly more intuitive understanding of the topic.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the **real** matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a **variant of Hermitian symmetry** such that $\eta H \eta^{-1} = \overline{H^t}$

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a **variant of Hermitian symmetry** such that $\eta H \eta^{-1} = \overline{H^t}$ where

$$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{is a Riemannian metric form.}$$

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$

has a double root if and only if $f_2 = \pm f_3$.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

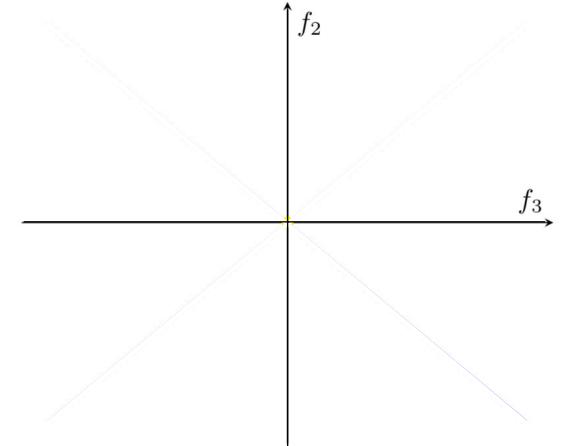
It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$

has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:



Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

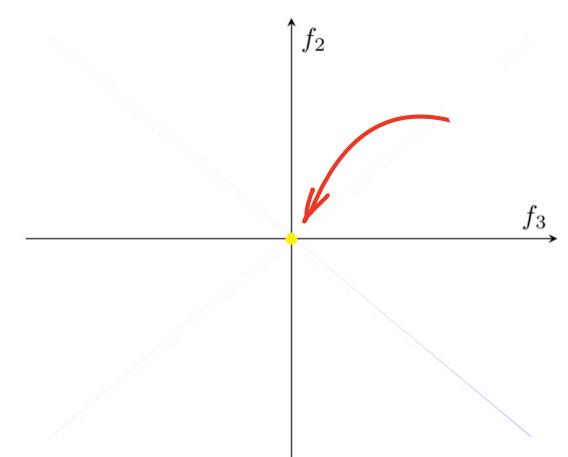
$$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{is a Riemannian metric form.}$$

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$

has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

0. Over $\{(0, 0)\}$



Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

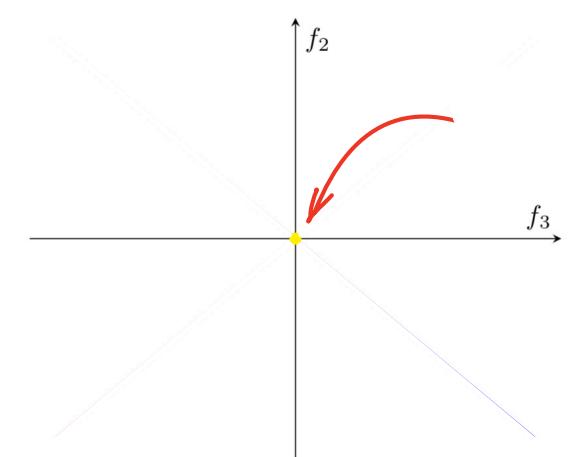
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

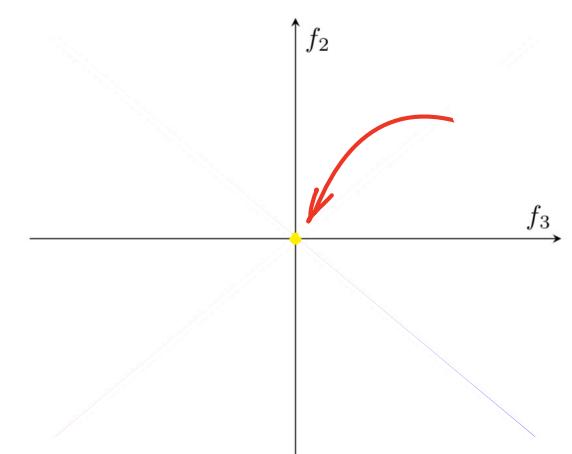
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{is a Riemannian metric form.}$$

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

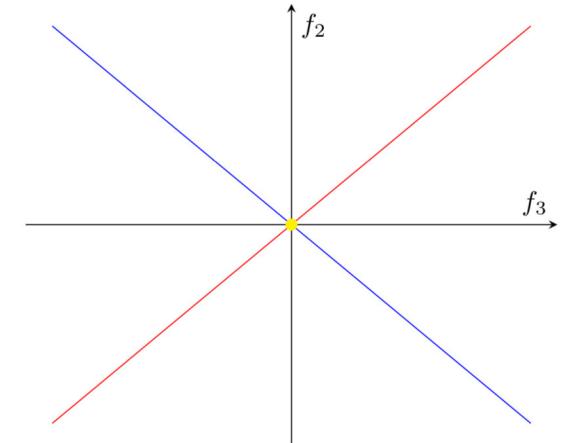
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{is a Riemannian metric form.}$$

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

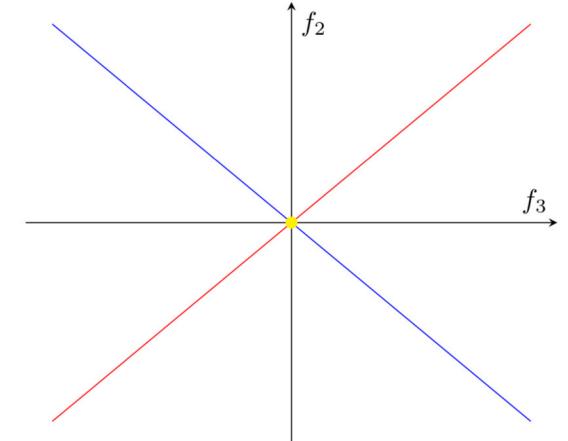
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{is a Riemannian metric form.}$$

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

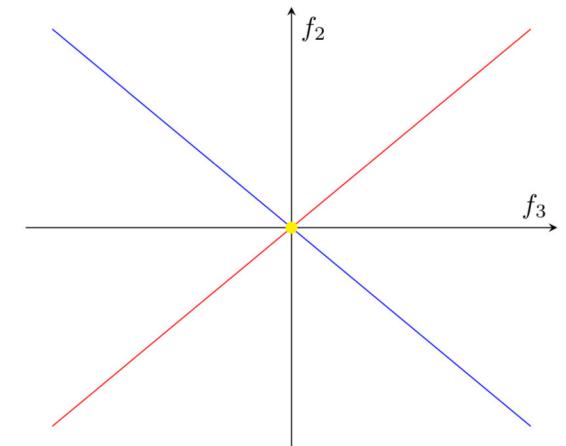
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix} \quad \text{is a Riemannian metric form.}$$

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

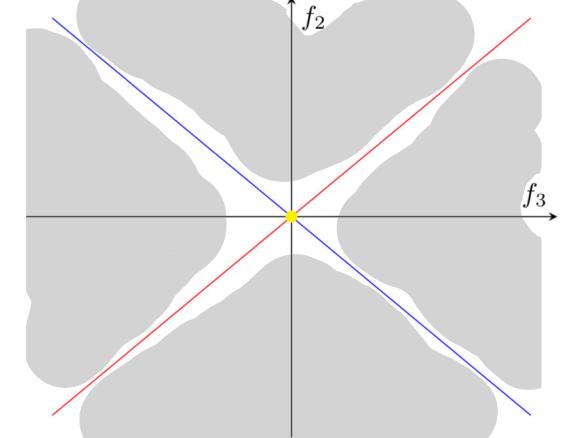
$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$

has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

2. Over $\{f_2 \neq \pm f_3\}$



Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

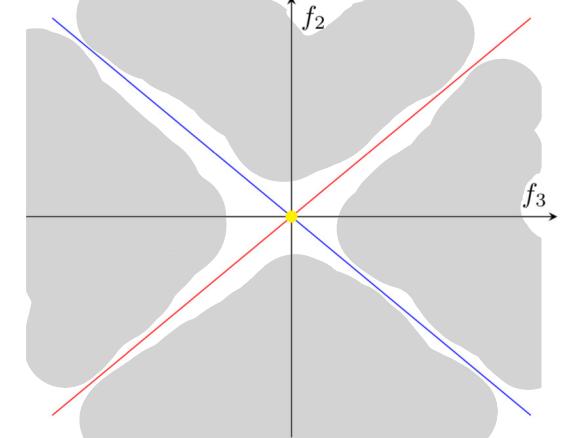
$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$

has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

2. Over $\{f_2 \neq \pm f_3\}$, H has 2 **distinct** eigenvalues.



Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

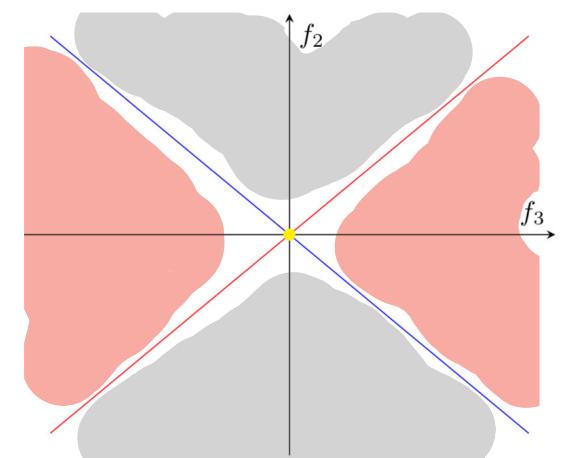
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are **real**.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

With motivations from topological classifications of non-Hermitian gapless quantum mechanical systems and their applications to materials science, let us consider the real matrix (a Hamiltonian)

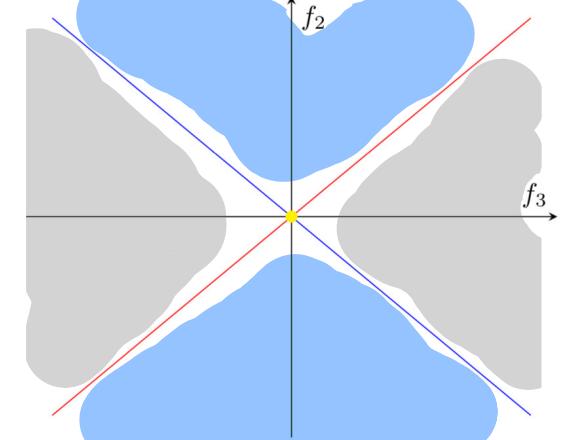
$$H = H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

It satisfies a variant of Hermitian symmetry such that $\eta H \eta^{-1} = \overline{H^t}$ where

$\eta = \begin{bmatrix} -1 & 0 \\ 0 & 1 \end{bmatrix}$ is a Riemannian metric form.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_2 \\ -f_2 & -f_3 - \omega \end{vmatrix} = \omega^2 + f_2^2 - f_3^2$$



has a double root if and only if $f_2 = \pm f_3$. As a parameter space for this Hamiltonian H , the $f_2 f_3$ -plane becomes a **stratified space**:

2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are real. When $|f_2| > |f_3|$, the eigenvectors are **not real**.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this **stratified space**

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “**3-band** systems” is also of particular interest to us.

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

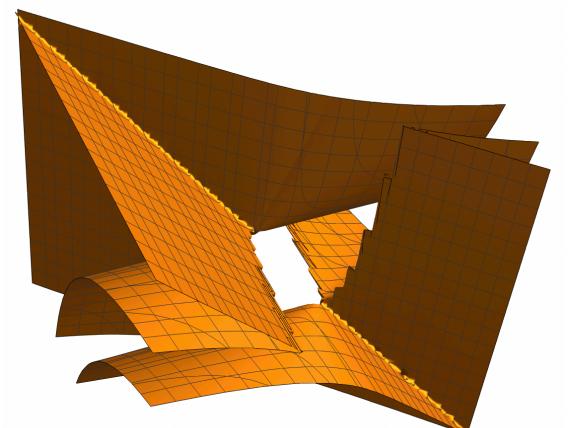
Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:

The equation for this surface is a non-homogeneous real polynomial in f_1, f_2, f_3 of degree 6.



Swallowtail couple sw2

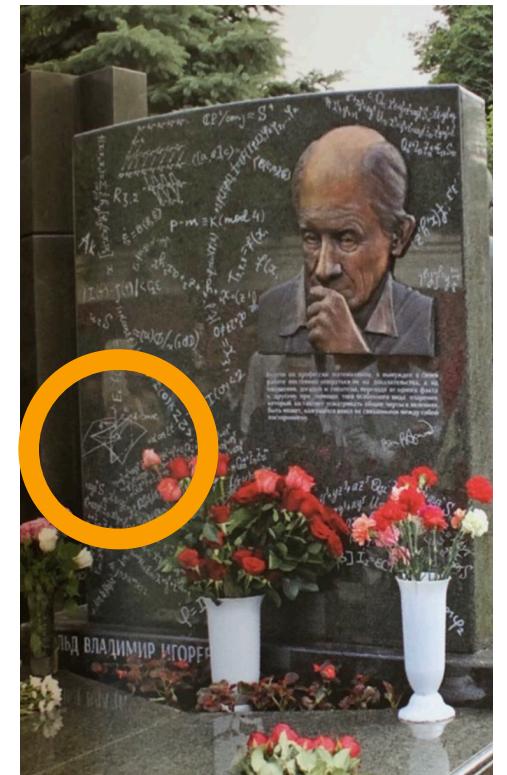
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



V. I. Arnold's tombstone at the Novodevichy Cemetery in Moscow

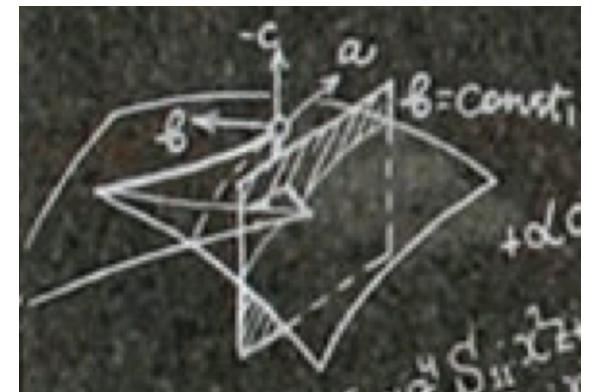
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



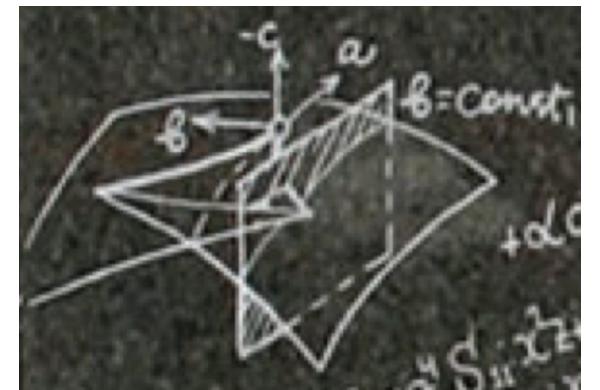
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



A **local** model for moduli spaces of 3-band Hamiltonians

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

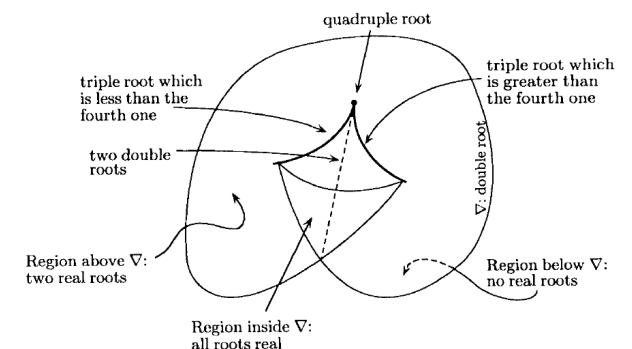
$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:

Arnold, Braids of algebraic functions and the cohomology of swallowtails, 1968.

Homological stability of braid groups

*Portrait from Gelfand, Kapranov, Zelevinsky,
Discriminants, resultants, and multidimensional determinants.*



The space of polynomials $x^4 + ax^2 + bx + c$

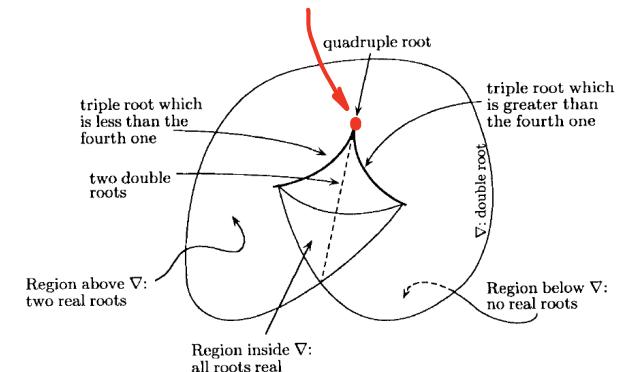
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



The space of polynomials $x^4 + ax^2 + bx + c$

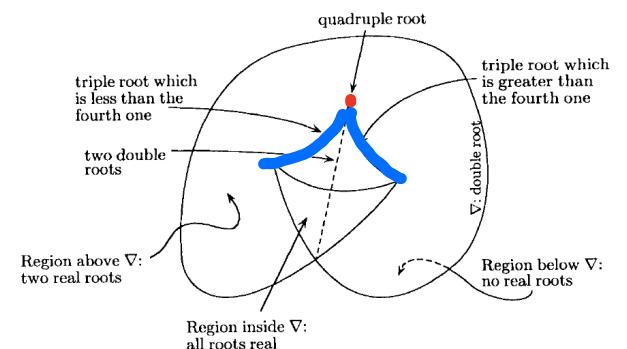
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



The space of polynomials $x^4 + ax^2 + bx + c$

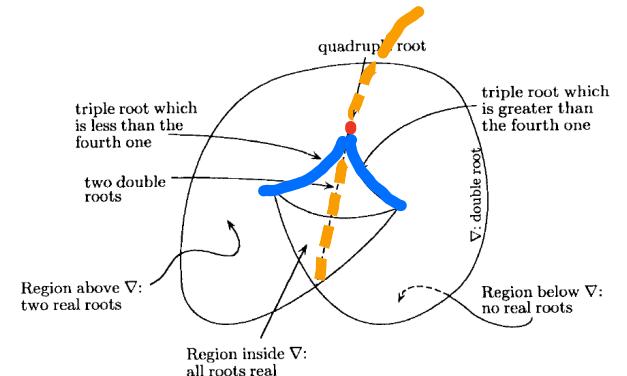
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



The space of polynomials $x^4 + ax^2 + bx + c$

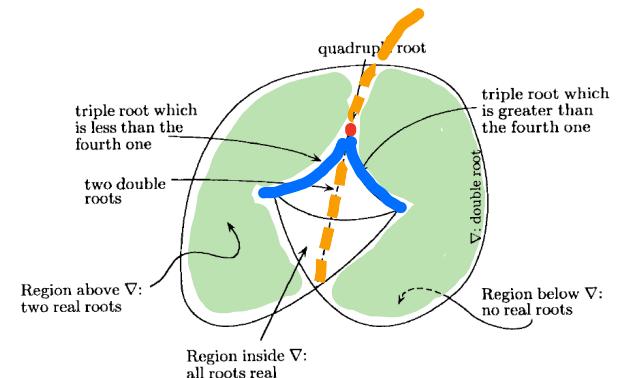
Mathematical set-up: Eigenframe rotation of non-Hermitian systems

Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:



The space of polynomials $x^4 + ax^2 + bx + c$

Mathematical set-up: Eigenframe rotation of non-Hermitian systems

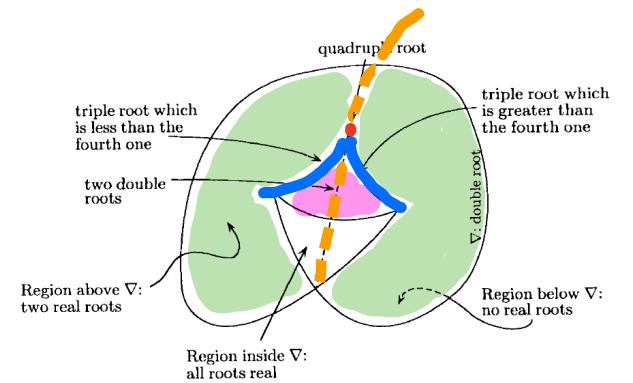
Question. We would like to classify, up to “intersection” homotopy, the loops in this stratified space, according to how the eigenvectors evolve along each loop and the resulting monodromy. What would be computable algebraic invariants for such a “stratified vector bundle?”

Another example of non-Hermitian “3-band systems” is also of particular interest to us. One such Hamiltonian takes the form

$$H_3 = \begin{bmatrix} 1 - f_1 - f_2 & f_1 & f_2 \\ -f_1 & f_1 - f_3 & f_3 \\ -f_2 & f_3 & f_2 - f_3 \end{bmatrix}$$

Governing eigenvalues with multiplicity, the discriminant surface of its characteristic polynomial is a pair of *swallowtails* in the $f_1 f_2 f_3$ -space:

Again, we aim to find computable algebraic invariants that systematically classify the evolutions of eigenvectors along loops in such stratified parameter spaces, including when they cross the discriminant surface resulting in degeneracies of various sorts.



The space of polynomials $x^4 + ax^2 + bx + c$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is **isolated**, yet has profound physical implications already known to Arnold.

*Remarks on eigenvalues and eigenvectors of Hermitian matrices,
Berry phase, adiabatic connections and quantum Hall effect, 1995.*

Also: Polymathematics, 2000.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is isolated, yet has profound physical implications already known to Arnold. Let us consider the real Hamiltonian

$$H(f_1, f_3) = \begin{bmatrix} f_3 & f_1 \\ f_1 & -f_3 \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is isolated, yet has profound physical implications already known to Arnold. Let us consider the real Hamiltonian

$$H(f_1, f_3) = \begin{bmatrix} f_3 & f_1 \\ f_1 & -f_3 \end{bmatrix}$$

It represents all symmetric 2×2 matrices **spectrally**, since any $\begin{bmatrix} a & b \\ b & c \end{bmatrix}$ has the same eigenvalues and eigenvectors as $\begin{bmatrix} (a-c)/2 & b \\ b & (c-a)/2 \end{bmatrix}$.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is isolated, yet has profound physical implications already known to Arnold. Let us consider the real Hamiltonian

$$H(f_1, f_3) = \begin{bmatrix} f_3 & f_1 \\ f_1 & -f_3 \end{bmatrix}$$

It represents all symmetric 2×2 matrices spectrally, since any $\begin{bmatrix} a & b \\ b & c \end{bmatrix}$ has the same eigenvalues and eigenvectors as $\begin{bmatrix} (a-c)/2 & b \\ b & (c-a)/2 \end{bmatrix}$.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_1 \\ f_1 & -f_3 - \omega \end{vmatrix} = \omega^2 - f_1^2 - f_3^2$$

has a double root if and only if $f_1 = f_3 = 0$.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is isolated, yet has profound physical implications already known to Arnold. Let us consider the real Hamiltonian

$$H(f_1, f_3) = \begin{bmatrix} f_3 & f_1 \\ f_1 & -f_3 \end{bmatrix}$$

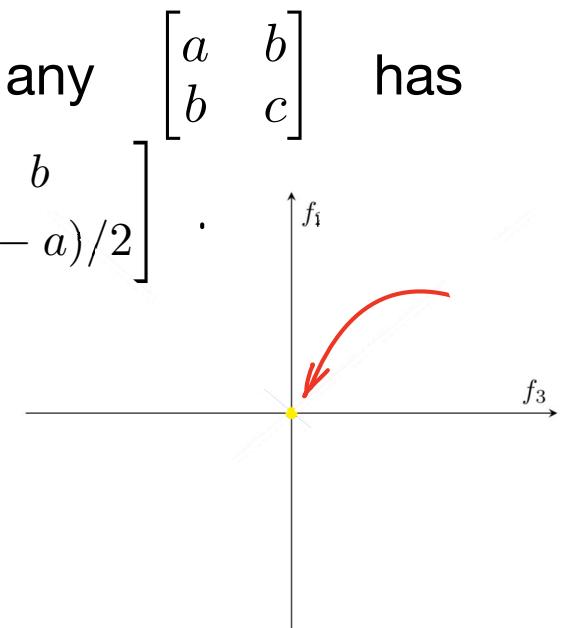
It represents all symmetric 2×2 matrices spectrally, since any $\begin{bmatrix} a & b \\ b & c \end{bmatrix}$ has the same eigenvalues and eigenvectors as $\begin{bmatrix} (a-c)/2 & b \\ b & (c-a)/2 \end{bmatrix}$.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_1 \\ f_1 & -f_3 - \omega \end{vmatrix} = \omega^2 - f_1^2 - f_3^2$$

has a double root if and only if $f_1 = f_3 = 0$.

The parameter $f_1 f_3$ -plane thus has an **isolated singular point** $(0, 0)$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

The Hermitian case is simple, as the singularity is isolated, yet has profound physical implications already known to Arnold. Let us consider the real Hamiltonian

$$H(f_1, f_3) = \begin{bmatrix} f_3 & f_1 \\ f_1 & -f_3 \end{bmatrix}$$

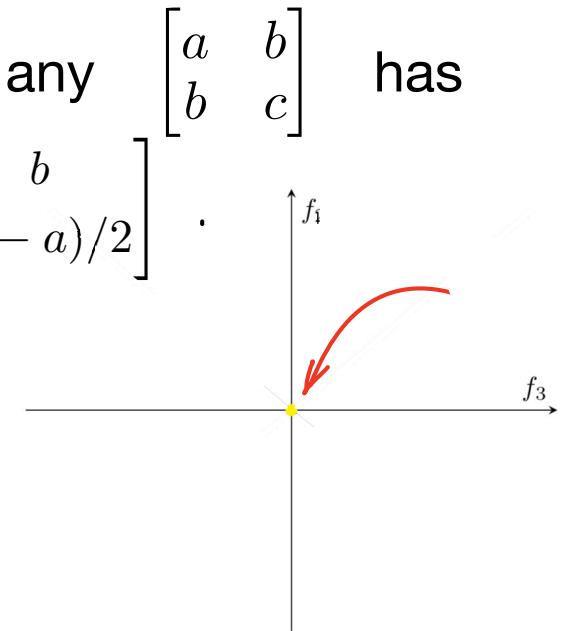
It represents all symmetric 2×2 matrices spectrally, since any $\begin{bmatrix} a & b \\ b & c \end{bmatrix}$ has the same eigenvalues and eigenvectors as $\begin{bmatrix} (a-c)/2 & b \\ b & (c-a)/2 \end{bmatrix}$.

Its characteristic polynomial

$$\begin{vmatrix} f_3 - \omega & f_1 \\ f_1 & -f_3 - \omega \end{vmatrix} = \omega^2 - f_1^2 - f_3^2$$

has a double root if and only if $f_1 = f_3 = 0$.

The parameter $f_1 f_3$ -plane thus has an **isolated singular point** $(0, 0)$ and is a particularly simple **stratified space**.



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained **non-Abelian** “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1\left(SO(3)/(O(1) \times O(1))\right) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained non-Abelian “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1\left(SO(3)/(O(1) \times O(1))\right) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

However, their explanation for the appearance of the $O(1)^n$ -action in constructing the moduli spaces was rather ad hoc.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained non-Abelian “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1\left(SO(3)/(O(1) \times O(1))\right) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

However, their explanation for the appearance of the $O(1)^n$ -action in constructing the moduli spaces was rather ad hoc.

One of our key steps is a more conceptual understanding of the above moduli spaces in the case of $n=2$ through **bundle theory**.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained non-Abelian “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1\left(SO(3)/(O(1) \times O(1))\right) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

However, their explanation for the appearance of the $O(1)^n$ -action in constructing the moduli spaces was rather ad hoc.

One of our key steps is a more conceptual understanding of the above moduli spaces in the case of $n=2$ through **bundle theory**. To see how they rotate, let us compute the unit eigenvectors explicitly.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained non-Abelian “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1\left(SO(3)/(O(1) \times O(1))\right) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

However, their explanation for the appearance of the $O(1)^n$ -action in constructing the moduli spaces was rather ad hoc.

One of our key steps is a more conceptual understanding of the above moduli spaces in the case of $n=2$ through **bundle theory**. To see how they rotate, let us compute the unit eigenvectors explicitly.

$$\begin{vmatrix} f_3 - \omega & f_1 \\ f_1 & -f_3 - \omega \end{vmatrix} = \omega^2 - f_1^2 - f_3^2$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

How does the eigenframe rotate over this stratified parameter plane?

As our starting point, previous work of Wu et al. [**Science**, 2019] classified the eigenframe rotation by

$$\pi_1(SO(2)/O(1)) \cong \pi_1(S^1) \cong \mathbb{Z}$$

Moreover, they obtained non-Abelian “topological charge” for n -band Hermitian systems when $n > 2$, such as

$$\pi_1\left(SO(3)/(O(1) \times O(1))\right) \cong \mathbb{H} = \{\pm 1, \pm i, \pm j, \pm k\}$$

However, their explanation for the appearance of the $O(1)^n$ -action in constructing the moduli spaces was rather ad hoc.

One of our key steps is a more conceptual understanding of the above moduli spaces in the case of $n=2$ through **bundle theory**. To see how they rotate, let us compute the unit eigenvectors explicitly.

$$\begin{vmatrix} f_3 - \omega & f_1 \\ f_1 & -f_3 - \omega \end{vmatrix} = \omega^2 - f_1^2 - f_3^2 = 0 \implies \omega_{\pm} = \pm \sqrt{f_1^2 + f_3^2}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

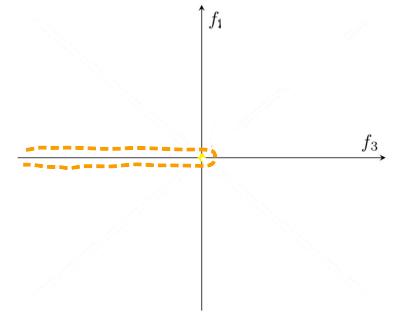
To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0}$$
$$\begin{bmatrix} \left(f_3 - \sqrt{}\right)\left(-f_3 - \sqrt{}\right) & f_1\left(-f_3 - \sqrt{}\right) \\ f_1 & -f_3 - \sqrt{} \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$
$$\begin{bmatrix} \left(f_3 - \sqrt{\quad}\right)\left(-f_3 - \sqrt{\quad}\right) & f_1\left(-f_3 - \sqrt{\quad}\right) \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$

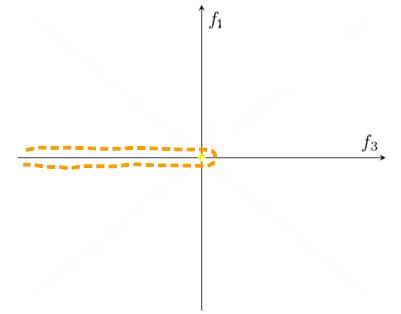


Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$

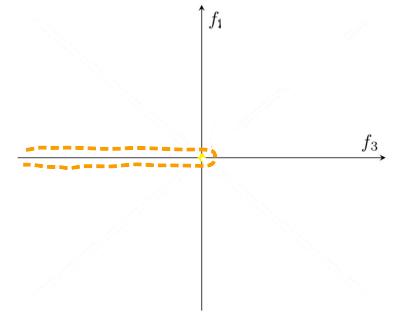
$$\begin{bmatrix} (f_3 - \sqrt{\quad})(-f_3 - \sqrt{\quad}) & f_1(-f_3 - \sqrt{\quad}) \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \rightarrow \begin{bmatrix} f_1^2 & -f_1f_3 - f_1\sqrt{\quad} \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$



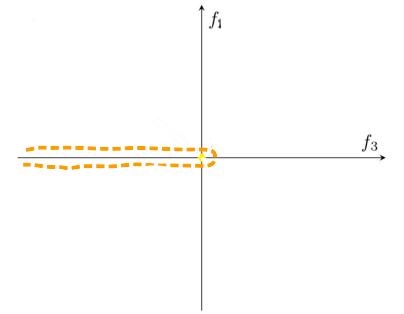
$$\begin{bmatrix} (f_3 - \sqrt{\quad})(-f_3 - \sqrt{\quad}) & f_1(-f_3 - \sqrt{\quad}) \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \rightarrow \begin{bmatrix} f_1^2 & -f_1 f_3 - f_1 \sqrt{\quad} \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$

$$\rightarrow \begin{bmatrix} 0 & 0 \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$



$$\begin{bmatrix} (f_3 - \sqrt{\quad})(-f_3 - \sqrt{\quad}) & f_1(-f_3 - \sqrt{\quad}) \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \rightarrow \begin{bmatrix} f_1^2 & -f_1 f_3 - f_1 \sqrt{\quad} \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$

$$\rightarrow \begin{bmatrix} 0 & 0 \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \Rightarrow v_+ = \begin{bmatrix} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{bmatrix}$$

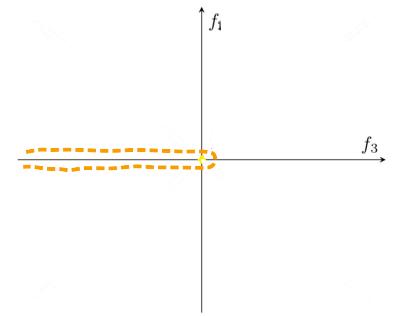
Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$

$$\begin{bmatrix} (f_3 - \sqrt{})(-f_3 - \sqrt{}) & f_1(-f_3 - \sqrt{}) \\ f_1 & -f_3 - \sqrt{} \end{bmatrix} \rightarrow \begin{bmatrix} f_1^2 & -f_1 f_3 - f_1 \sqrt{} \\ f_1 & -f_3 - \sqrt{} \end{bmatrix}$$

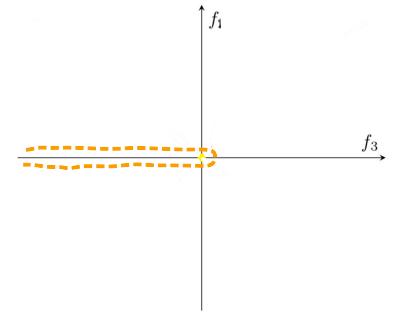
$$\rightarrow \begin{bmatrix} 0 & 0 \\ f_1 & -f_3 - \sqrt{} \end{bmatrix} \Rightarrow v_+ = \begin{bmatrix} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{bmatrix} \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} -\pi < \theta < \pi}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$



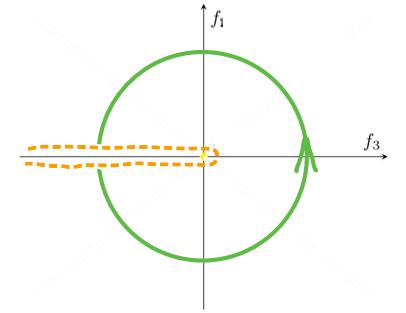
$$\begin{bmatrix} (f_3 - \sqrt{\quad})(-f_3 - \sqrt{\quad}) & f_1(-f_3 - \sqrt{\quad}) \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \rightarrow \begin{bmatrix} f_1^2 & -f_1 f_3 - f_1 \sqrt{\quad} \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$

$$\rightarrow \begin{bmatrix} 0 & 0 \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \Rightarrow v_+ = \begin{bmatrix} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{bmatrix} \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} \quad -\pi < \theta < \pi} \begin{bmatrix} \cos \theta + 1 \\ \sin \theta \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{bmatrix} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{bmatrix} \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0}$$



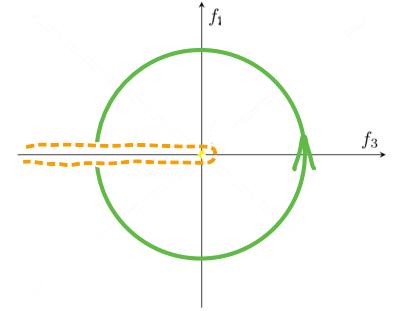
$$\begin{bmatrix} (f_3 - \sqrt{\quad})(-f_3 - \sqrt{\quad}) & f_1(-f_3 - \sqrt{\quad}) \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \rightarrow \begin{bmatrix} f_1^2 & -f_1 f_3 - f_1 \sqrt{\quad} \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix}$$

$$\rightarrow \begin{bmatrix} 0 & 0 \\ f_1 & -f_3 - \sqrt{\quad} \end{bmatrix} \Rightarrow v_+ = \begin{bmatrix} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{bmatrix} \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} \quad -\pi < \theta < \pi} \begin{bmatrix} \cos \theta + 1 \\ \sin \theta \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{aligned}
 & \left[\begin{array}{cc} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{array} \right] \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0} \\
 & \left[\begin{array}{cc} \left(f_3 - \sqrt{}\right) \left(-f_3 - \sqrt{}\right) & f_1 \left(-f_3 - \sqrt{}\right) \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \rightarrow \left[\begin{array}{cc} f_1^2 & -f_1 f_3 - f_1 \sqrt{} \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \\
 & \rightarrow \left[\begin{array}{cc} 0 & 0 \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \Rightarrow v_+ = \left[\begin{array}{c} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{array} \right] \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} \quad -\pi < \theta < \pi} \left[\begin{array}{c} \cos \theta + 1 \\ \sin \theta \end{array} \right]
 \end{aligned}$$

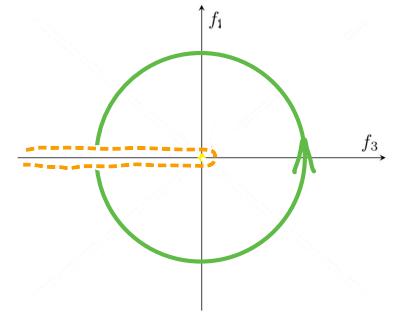


Observe that when $\theta \rightarrow (-\pi)_+$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_-$,

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{aligned}
 & \left[\begin{array}{cc} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{array} \right] \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0} \\
 & \left[\begin{array}{cc} (f_3 - \sqrt{})(-f_3 - \sqrt{}) & f_1(-f_3 - \sqrt{}) \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \rightarrow \left[\begin{array}{cc} f_1^2 & -f_1 f_3 - f_1 \sqrt{} \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \\
 & \rightarrow \left[\begin{array}{cc} 0 & 0 \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \Rightarrow v_+ = \left[\begin{array}{c} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{array} \right] \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} \quad -\pi < \theta < \pi} \left[\begin{array}{c} \cos \theta + 1 \\ \sin \theta \end{array} \right]
 \end{aligned}$$

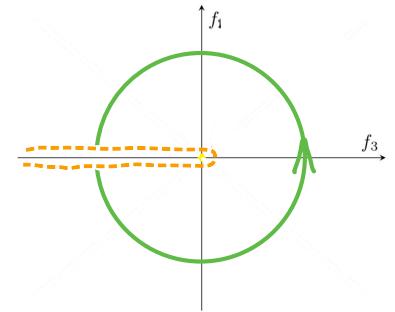


Observe that when $\theta \rightarrow (-\pi)_+$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_-$, whereas when $\theta \rightarrow \pi_-$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_+$.

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{aligned}
 & \left[\begin{array}{cc} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{array} \right] \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0} \\
 & \left[\begin{array}{cc} (f_3 - \sqrt{\quad})(-f_3 - \sqrt{\quad}) & f_1(-f_3 - \sqrt{\quad}) \\ f_1 & -f_3 - \sqrt{\quad} \end{array} \right] \rightarrow \left[\begin{array}{cc} f_1^2 & -f_1 f_3 - f_1 \sqrt{\quad} \\ f_1 & -f_3 - \sqrt{\quad} \end{array} \right] \\
 & \rightarrow \left[\begin{array}{cc} 0 & 0 \\ f_1 & -f_3 - \sqrt{\quad} \end{array} \right] \Rightarrow v_+ = \left[\begin{array}{c} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{array} \right] \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} \atop -\pi < \theta < \pi} \left[\begin{array}{c} \cos \theta + 1 \\ \sin \theta \end{array} \right]
 \end{aligned}$$



Observe that when $\theta \rightarrow (-\pi)_+$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_-$, whereas when $\theta \rightarrow \pi_-$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_+$.

We compute that

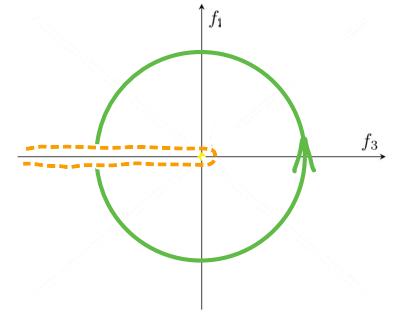
$$\lim_{\theta \rightarrow (-\pi)_+} \frac{v_+}{|v_+|} = \begin{bmatrix} 0 \\ -1 \end{bmatrix}$$

$$\lim_{\theta \rightarrow \pi_-} \frac{v_+}{|v_+|} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

$$\begin{aligned}
 & \left[\begin{array}{cc} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{array} \right] \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0} \\
 & \left[\begin{array}{cc} (f_3 - \sqrt{})(-f_3 - \sqrt{}) & f_1(-f_3 - \sqrt{}) \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \rightarrow \left[\begin{array}{cc} f_1^2 & -f_1 f_3 - f_1 \sqrt{} \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \\
 & \rightarrow \left[\begin{array}{cc} 0 & 0 \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \Rightarrow v_+ = \left[\begin{array}{c} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{array} \right] \xrightarrow{\begin{cases} f_3 = \cos \theta \\ f_1 = \sin \theta \end{cases} \quad -\pi < \theta < \pi} \left[\begin{array}{c} \cos \theta + 1 \\ \sin \theta \end{array} \right]
 \end{aligned}$$



Observe that when $\theta \rightarrow (-\pi)_+$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_-$, whereas when $\theta \rightarrow \pi_-$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_+$.

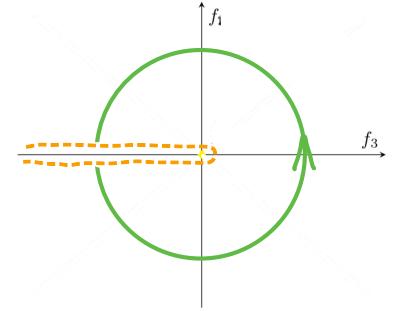
We compute that

$$\left. \begin{array}{l} \lim_{\theta \rightarrow (-\pi)_+} \frac{v_+}{|v_+|} = \begin{bmatrix} 0 \\ -1 \end{bmatrix} \\ \lim_{\theta \rightarrow \pi_-} \frac{v_+}{|v_+|} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \end{array} \right\} \text{Half Möbius band!}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

To solve for eigenvectors v_+ corresponding to ω_+ , perform Gaussian elimination through elementary row operations:

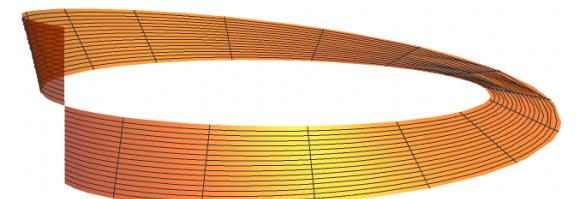
$$\begin{aligned}
 & \left[\begin{array}{cc} f_3 - \sqrt{f_1^2 + f_3^2} & f_1 \\ f_1 & -f_3 - \sqrt{f_1^2 + f_3^2} \end{array} \right] \xrightarrow{f_3 - \sqrt{f_1^2 + f_3^2} \neq 0 \iff f_1 = 0, f_3 \leq 0} \\
 & \left[\begin{array}{cc} (f_3 - \sqrt{})(-f_3 - \sqrt{}) & f_1(-f_3 - \sqrt{}) \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \rightarrow \left[\begin{array}{cc} f_1^2 & -f_1 f_3 - f_1 \sqrt{} \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \\
 & \rightarrow \left[\begin{array}{cc} 0 & 0 \\ f_1 & -f_3 - \sqrt{} \end{array} \right] \Rightarrow v_+ = \left[\begin{array}{c} f_3 + \sqrt{f_1^2 + f_3^2} \\ f_1 \end{array} \right] \xrightarrow{\begin{array}{l} \left\{ \begin{array}{l} f_3 = \cos \theta \\ f_1 = \sin \theta \end{array} \right. \\ -\pi < \theta < \pi \end{array}} \left[\begin{array}{c} \cos \theta + 1 \\ \sin \theta \end{array} \right]
 \end{aligned}$$



Observe that when $\theta \rightarrow (-\pi)_+$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_-$, whereas when $\theta \rightarrow \pi_-$, we have $\cos \theta + 1 \rightarrow 0_+$ and $\sin \theta \rightarrow 0_+$.

We compute that

$$\left. \begin{array}{l} \lim_{\theta \rightarrow (-\pi)_+} \frac{v_+}{|v_+|} = \begin{bmatrix} 0 \\ -1 \end{bmatrix} \\ \lim_{\theta \rightarrow \pi_-} \frac{v_+}{|v_+|} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \end{array} \right\} \text{Half Möbius band!}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$\begin{aligned} S^0 &\hookrightarrow S^1 \rightarrow S^1 & \mathbb{R} \\ S^1 &\hookrightarrow S^3 \rightarrow S^2 & \text{if the Hamiltonian is over } \mathbb{C} \\ S^3 &\hookrightarrow S^7 \rightarrow S^4 & \mathbb{H} \end{aligned}$$

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

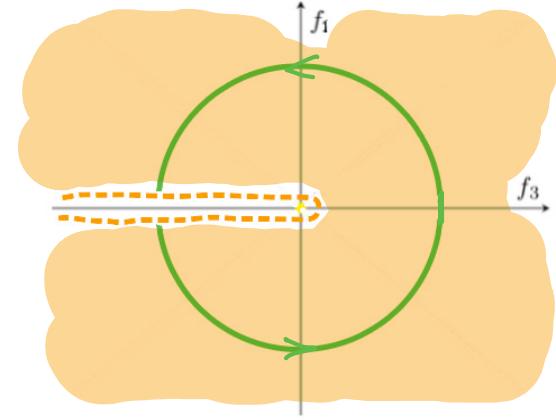
Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$S^0 \hookrightarrow S^1 \rightarrow S^1$$

$$S^1 \hookrightarrow S^3 \rightarrow S^2 \quad \text{if the Hamiltonian is over } \mathbb{R}$$

$$S^3 \hookrightarrow S^7 \rightarrow S^4 \quad \mathbb{C}$$

$$S^7 \hookrightarrow S^{15} \rightarrow S^8 \quad \mathbb{H}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

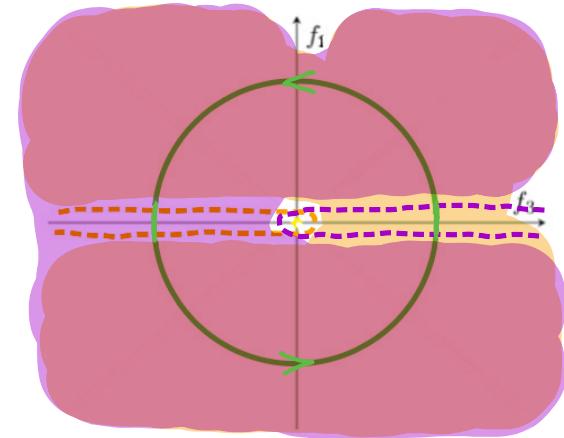
Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$S^0 \hookrightarrow S^1 \rightarrow S^1$$

$$S^1 \hookrightarrow S^3 \rightarrow S^2 \quad \text{if the Hamiltonian is over } \mathbb{R}$$

$$S^3 \hookrightarrow S^7 \rightarrow S^4 \quad \mathbb{C}$$

$$S^7 \hookrightarrow S^{15} \rightarrow S^8 \quad \mathbb{H}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

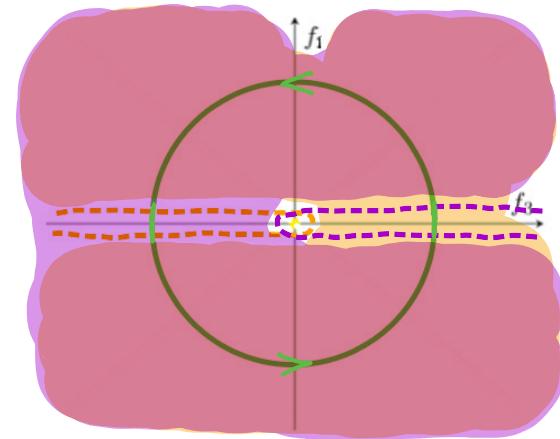
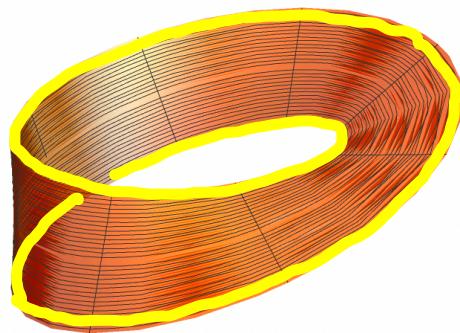
Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$S^0 \hookrightarrow S^1 \rightarrow S^1$$

$$S^1 \hookrightarrow S^3 \rightarrow S^2 \quad \text{if the Hamiltonian is over } \mathbb{R}$$

$$S^3 \hookrightarrow S^7 \rightarrow S^4 \quad \text{if the Hamiltonian is over } \mathbb{C}$$

$$S^7 \hookrightarrow S^{15} \rightarrow S^8 \quad \text{if the Hamiltonian is over } \mathbb{H}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

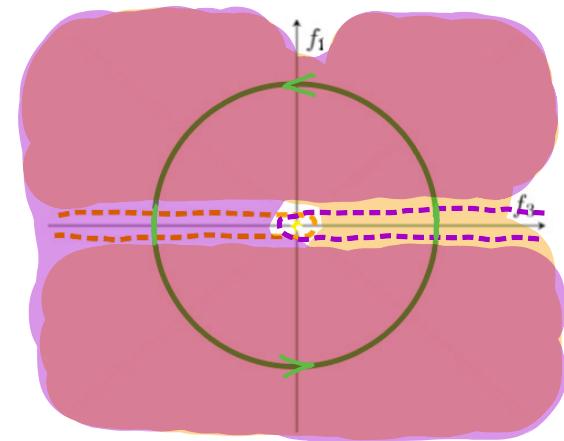
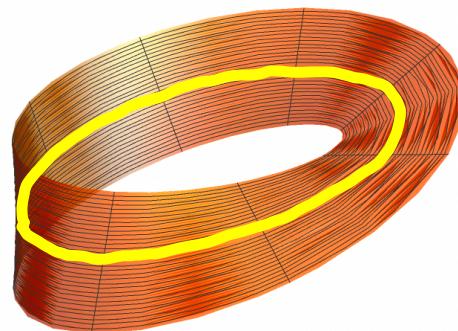
Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$S^0 \hookrightarrow S^1 \rightarrow \textcolor{red}{S}^1$$

$$S^1 \hookrightarrow S^3 \rightarrow S^2 \quad \text{if the Hamiltonian is over } \mathbb{R}$$

$$S^3 \hookrightarrow S^7 \rightarrow S^4 \quad \mathbb{C}$$

$$S^7 \hookrightarrow S^{15} \rightarrow S^8 \quad \mathbb{H}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

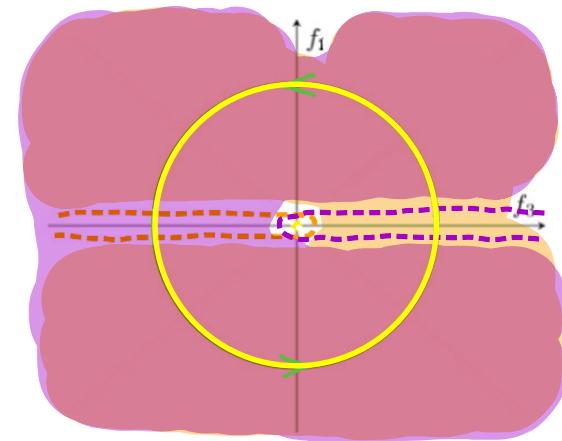
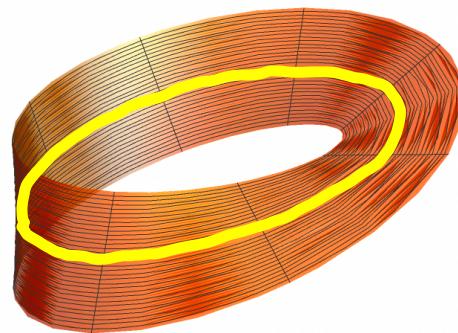
Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$S^0 \hookrightarrow S^1 \rightarrow \textcolor{red}{S}^1$$

$$S^1 \hookrightarrow S^3 \rightarrow S^2 \quad \text{if the Hamiltonian is over } \mathbb{R}$$

$$S^3 \hookrightarrow S^7 \rightarrow S^4 \quad \mathbb{C}$$

$$\mathbb{H}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

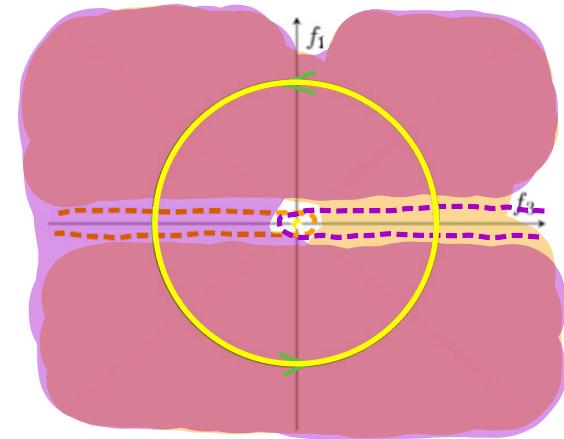
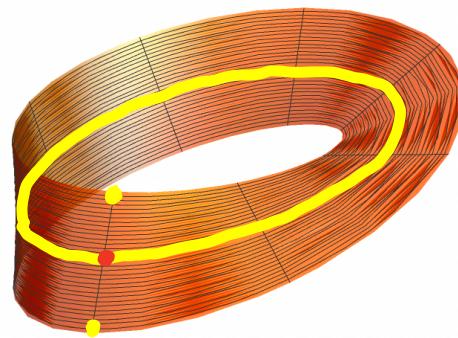
Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$S^0 \hookrightarrow S^1 \rightarrow S^1$$

$$S^1 \hookrightarrow S^3 \rightarrow S^2 \quad \text{if the Hamiltonian is over } \mathbb{R}$$

$$S^3 \hookrightarrow S^7 \rightarrow S^4 \quad \mathbb{C}$$

$$S^7 \hookrightarrow S^{15} \rightarrow S^8 \quad \mathbb{H}$$



Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$\begin{aligned} S^0 &\hookrightarrow S^1 \rightarrow S^1 & \mathbb{R} \\ S^1 &\hookrightarrow S^3 \rightarrow S^2 & \text{if the Hamiltonian is over } \mathbb{C} \\ S^3 &\hookrightarrow S^7 \rightarrow S^4 & \mathbb{H} \end{aligned}$$

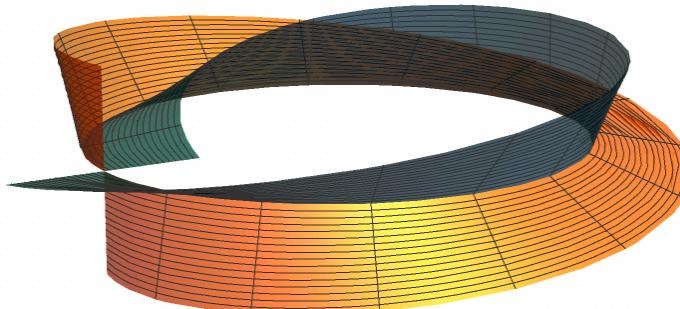
Corollary. The universal eigenbundle for real Hermitian 2-band systems is given by a pair of orthogonally intersecting half Möbius bands

Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$\begin{aligned} S^0 \hookrightarrow S^1 &\rightarrow S^1 & \mathbb{R} \\ S^1 \hookrightarrow S^3 &\rightarrow S^2 \quad \text{if the Hamiltonian is over} & \mathbb{C} \\ S^3 \hookrightarrow S^7 &\rightarrow S^4 & \mathbb{H} \end{aligned}$$

Corollary. The universal eigenbundle for real Hermitian 2-band systems is given by a pair of orthogonally intersecting half Möbius bands

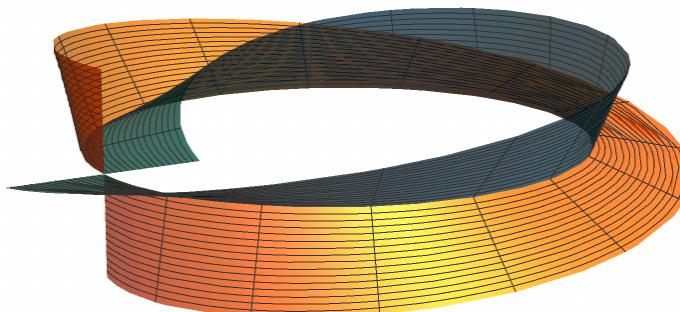


Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$\begin{aligned} S^0 &\hookrightarrow S^1 \rightarrow \textcolor{red}{S^1} & \mathbb{R} \\ S^1 &\hookrightarrow S^3 \rightarrow S^2 & \text{if the Hamiltonian is over } \mathbb{C} \\ S^3 &\hookrightarrow S^7 \rightarrow S^4 & \mathbb{H} \end{aligned}$$

Corollary. The universal eigenbundle for real Hermitian 2-band systems is given by a pair of orthogonally intersecting half Möbius bands over the **unit circle** in the punctured parameter plane.

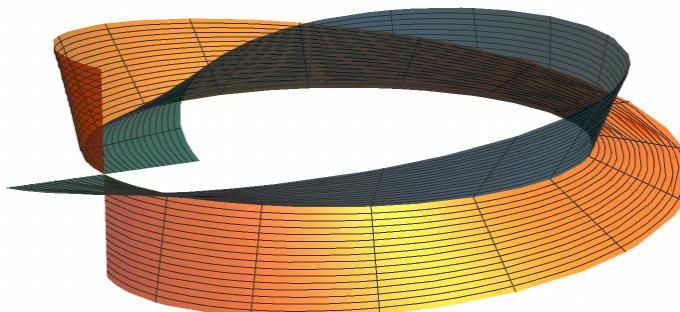


Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$\begin{aligned} S^0 \hookrightarrow S^1 &\rightarrow \textcolor{red}{S^1} & \mathbb{R} \\ S^1 \hookrightarrow S^3 &\rightarrow S^2 \quad \text{if the Hamiltonian is over} & \mathbb{C} \\ S^3 \hookrightarrow S^7 &\rightarrow S^4 & \mathbb{H} \end{aligned}$$

Corollary. The universal eigenbundle for real Hermitian 2-band systems is given by a pair of orthogonally intersecting half Möbius bands over the **unit circle** in the punctured parameter plane. In particular, eigenframe rotations along a generic loop in the moduli space are classified by $\pi_1(S^1) \cong \mathbb{Z}$.

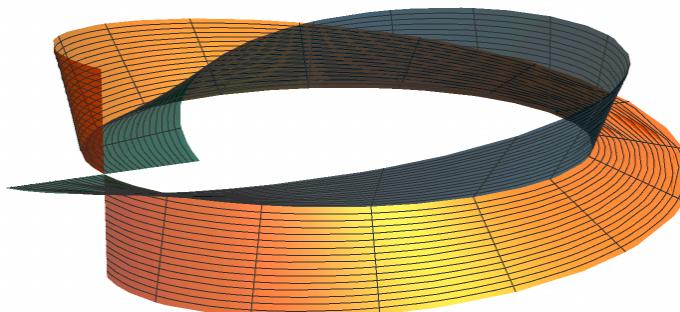


Eigenframe rotation as vector bundles: Revisiting the Hermitian case

Lemma. The universal rank-1 eigenbundle for Hermitian 2-band systems is given by the Hopf bundle

$$\begin{aligned} S^0 \hookrightarrow S^1 &\rightarrow \textcolor{red}{S^1} & \mathbb{R} \\ S^1 \hookrightarrow S^3 &\rightarrow S^2 \quad \text{if the Hamiltonian is over} & \mathbb{C} \\ S^3 \hookrightarrow S^7 &\rightarrow S^4 & \mathbb{H} \end{aligned}$$

Corollary. The universal eigenbundle for real Hermitian 2-band systems is given by a pair of orthogonally intersecting half Möbius bands over the unit circle in the **punctured** parameter plane. In particular, eigenframe rotations along a **generic** loop in the moduli space are classified by $\pi_1(S^1) \cong \mathbb{Z}$.



Eigenframe rotation as Higgs bundles: The non-Hermitian case



Eigenframe rotation as Higgs bundles: The non-Hermitian case

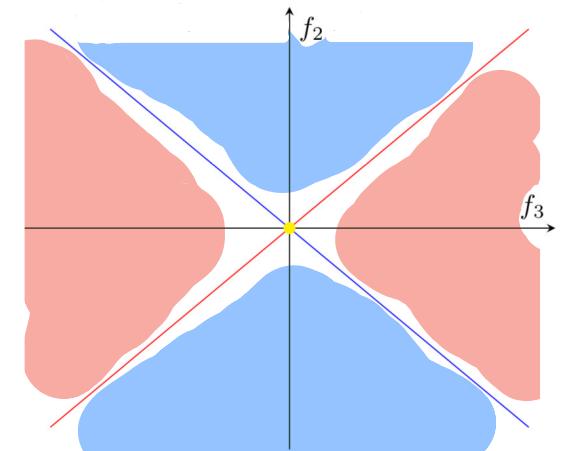
Recall that non-Hermitian 2-band systems

$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

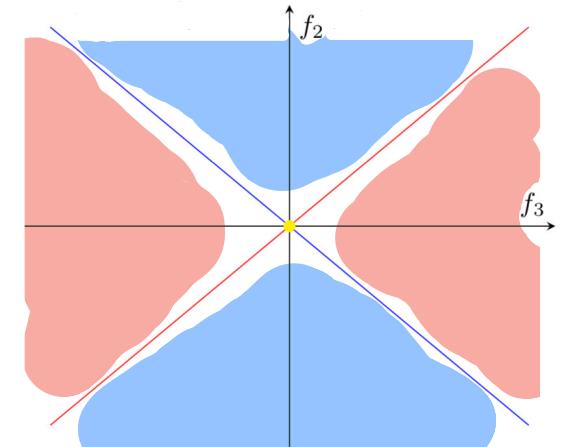


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.

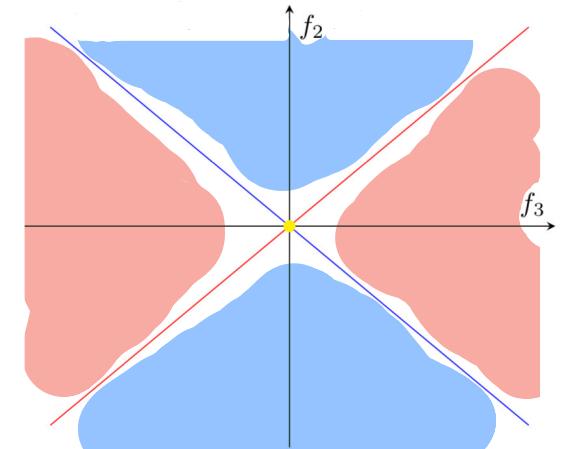


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.
1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.

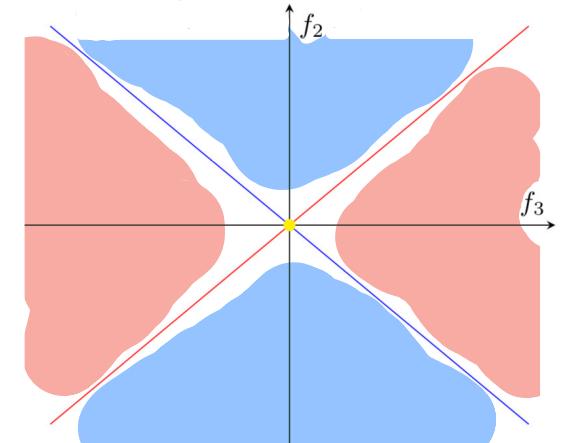


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.
1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.
2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are **real**. When $|f_2| > |f_3|$, the eigenvectors are **not real**.

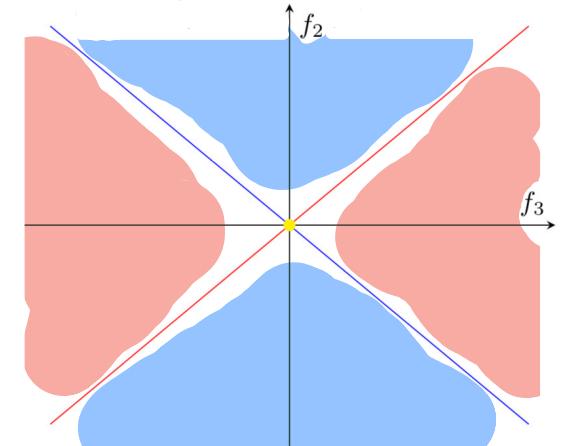


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.
1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.
2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are **real**. When $|f_2| > |f_3|$, the eigenvectors are **not real**.



A *Higgs bundle* $(E, \phi) \rightarrow C$ is essentially a family of matrices

Peter Higgs (bosons)

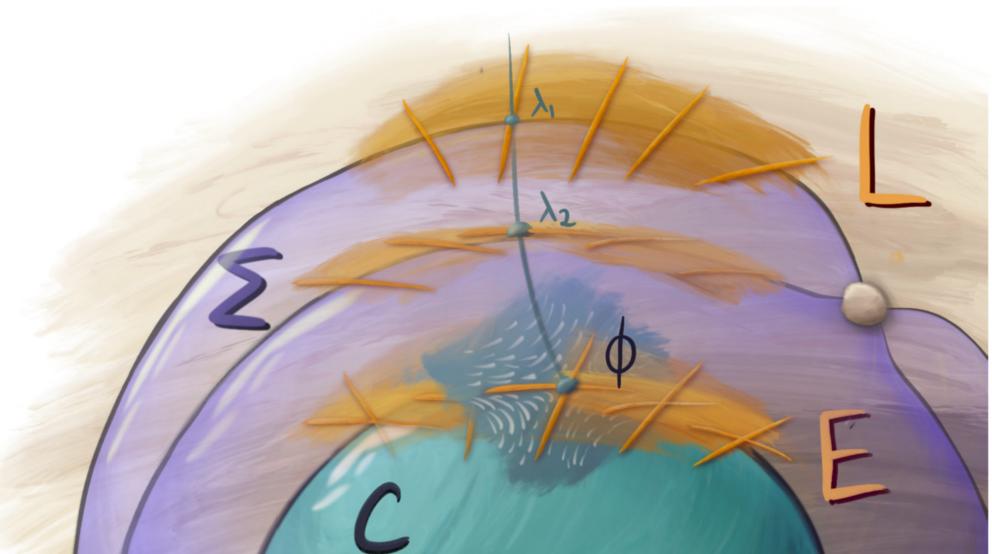
Nigel Hitchin 1987

Carlos Simpson

C compact Riemann surface (or more generally Kähler manifold)

E holomorphic vector bundle

ϕ Higgs field: a holomorphic 1-form taking values in the bundle of endomorphisms of E such that $\phi \wedge \phi = 0$

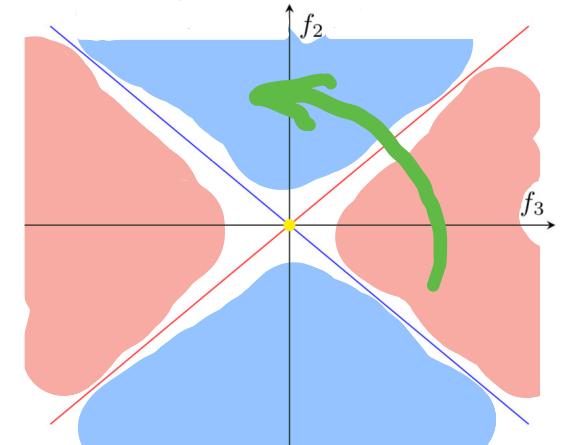


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

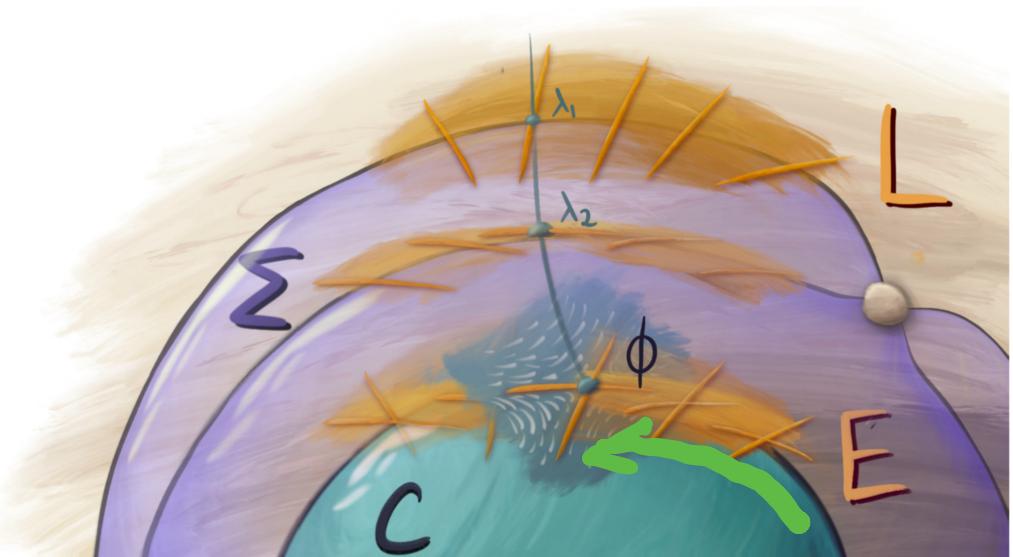
$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.
1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.
2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are **real**. When $|f_2| > |f_3|$, the eigenvectors are **not real**.



A *Higgs bundle* $(E, \phi) \rightarrow C$ is essentially a family of matrices, and if you try to diagonalize one you get a *spectral cover*.

$$\phi_x \in \text{End}(E_x), x \in C$$

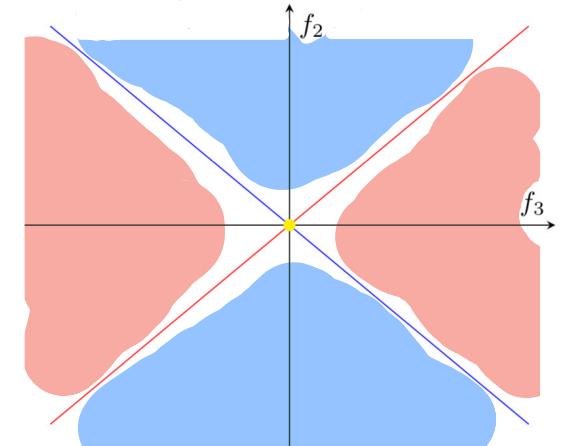


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

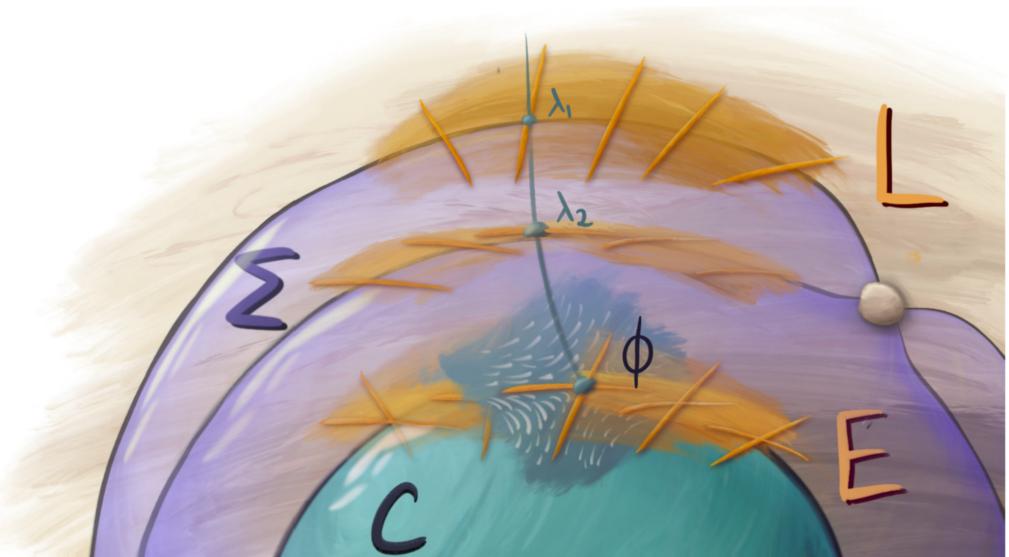
$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.
1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.
2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are **real**. When $|f_2| > |f_3|$, the eigenvectors are **not real**.



A *Higgs bundle* $(E, \phi) \rightarrow C$ is essentially a family of matrices, and if you try to diagonalize one you get a *spectral cover*.

Portrait from Kienzle and Rayan,
Hyperbolic band theory through Higgs bundles, **Adv. Math.**, 2022.

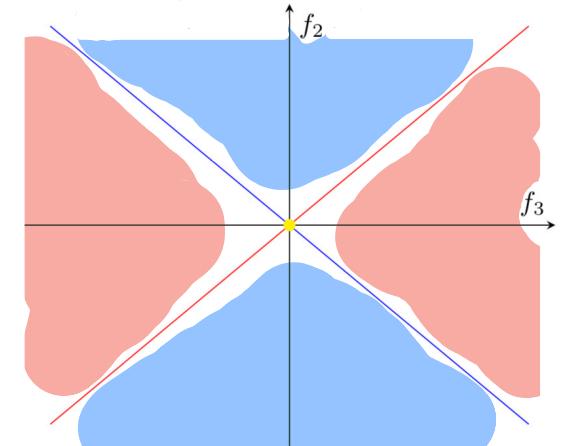


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Recall that non-Hermitian 2-band systems have a stratified parameter plane:

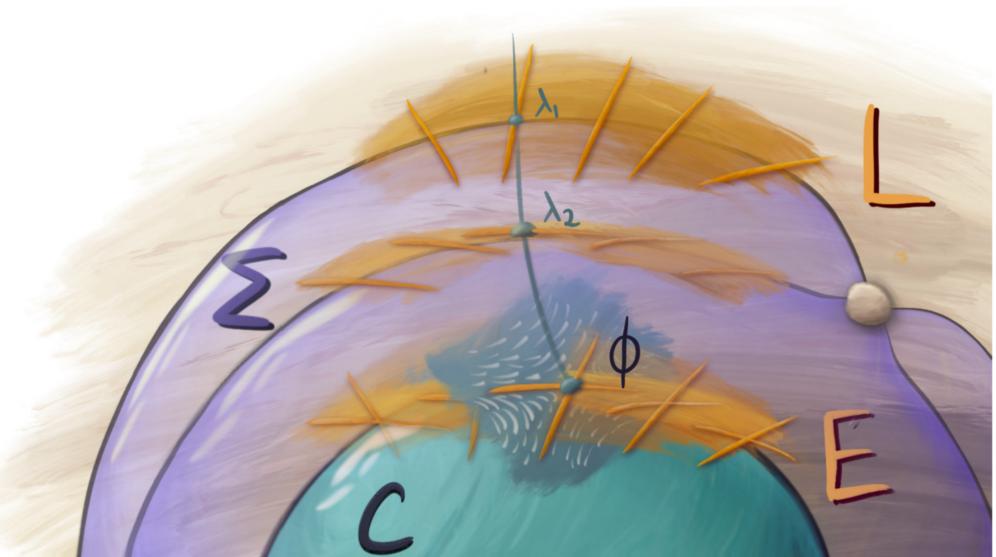
$$H(f_2, f_3) = \begin{bmatrix} f_3 & f_2 \\ -f_2 & -f_3 \end{bmatrix}$$

0. Over $\{(0, 0)\}$, H has a **double** eigenvalue, whose eigenspace is 2-dimensional.
1. Over $\{f_2 = \pm f_3\} - \{(0, 0)\}$, again H has a double eigenvalue, but its eigenspace is of **dimension 1**.
2. Over $\{f_2 \neq \pm f_3\}$, H has 2 distinct eigenvalues. When $|f_2| < |f_3|$, the eigenvectors are **real**. When $|f_2| > |f_3|$, the eigenvectors are **not real**.



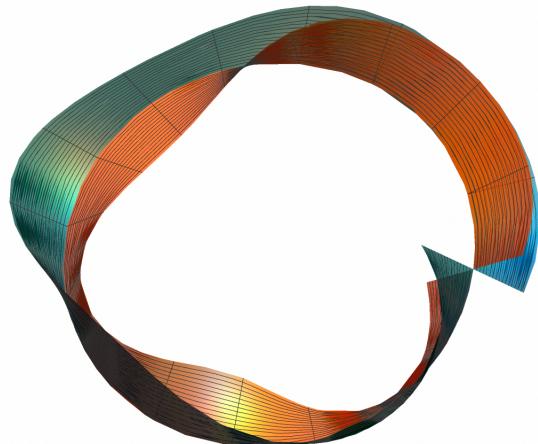
A *Higgs bundle* $(E, \phi) \rightarrow C$ is essentially a family of matrices, and if you try to diagonalize one you get a *spectral cover*.

Portrait from Kienzle and Rayan,
Hyperbolic band theory through Higgs bundles, *Adv. Math.*, 2022.



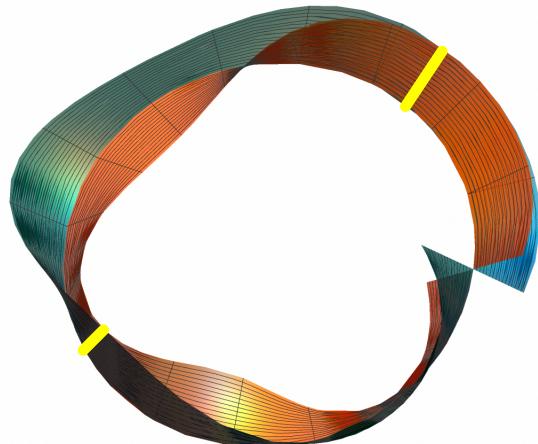
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands



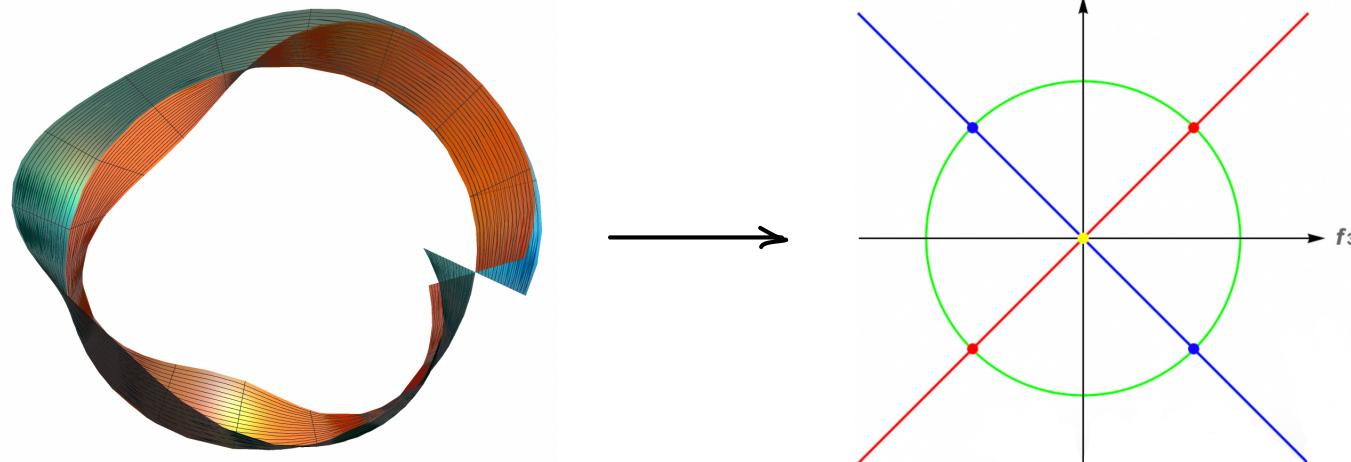
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of **kissing** half Möbius bands



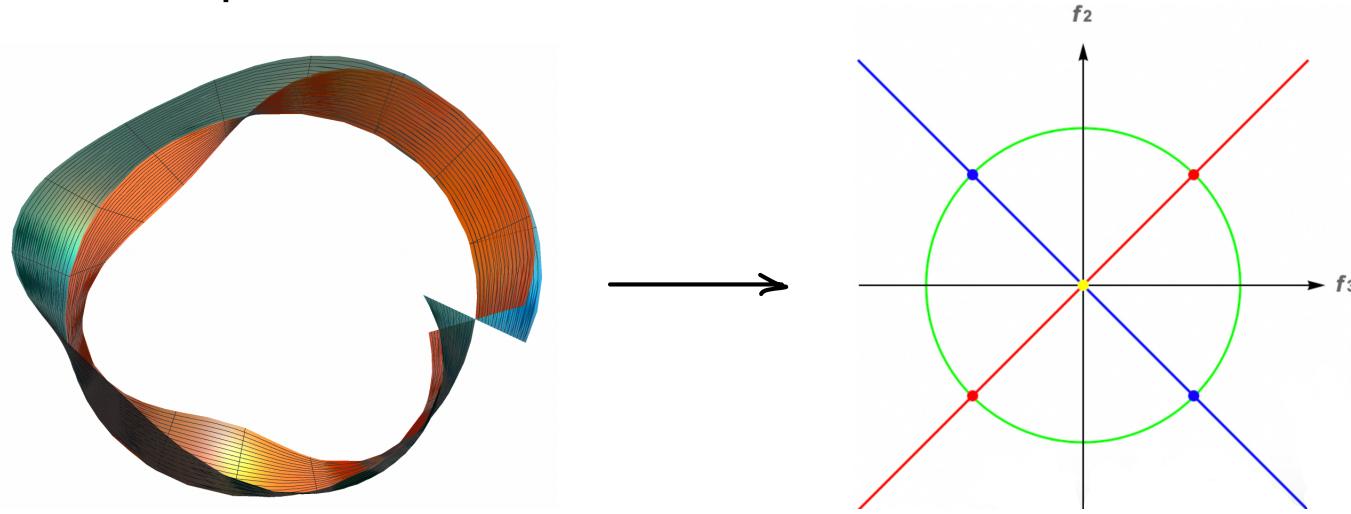
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the **stratified** unit circle in the punctured parameter plane



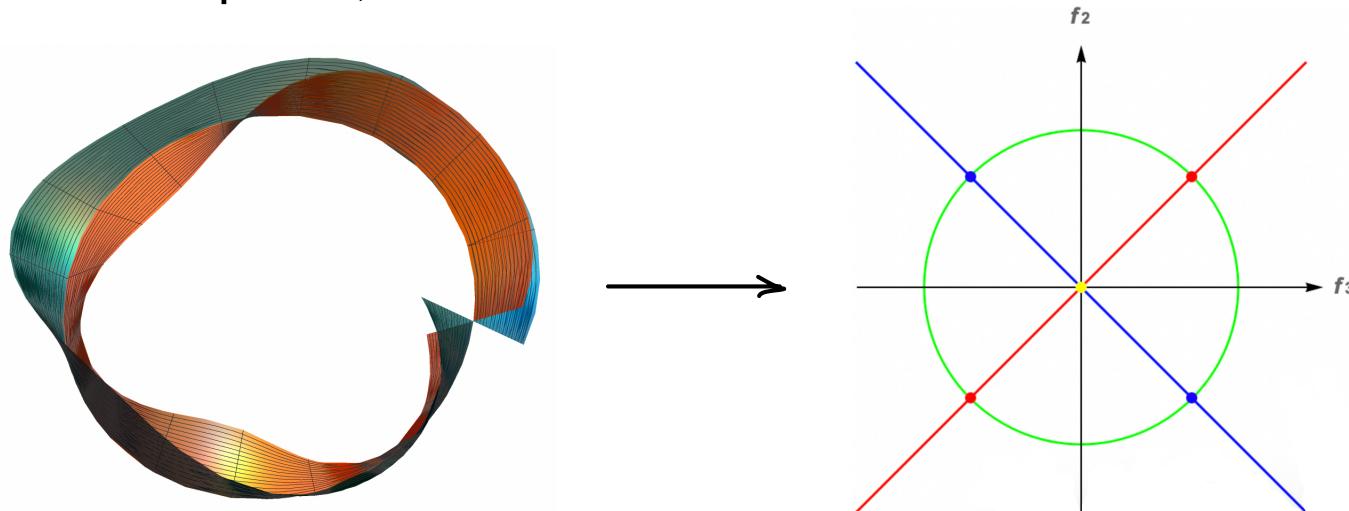
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the **stratified** unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.



Eigenframe rotation as Higgs bundles: The non-Hermitian case

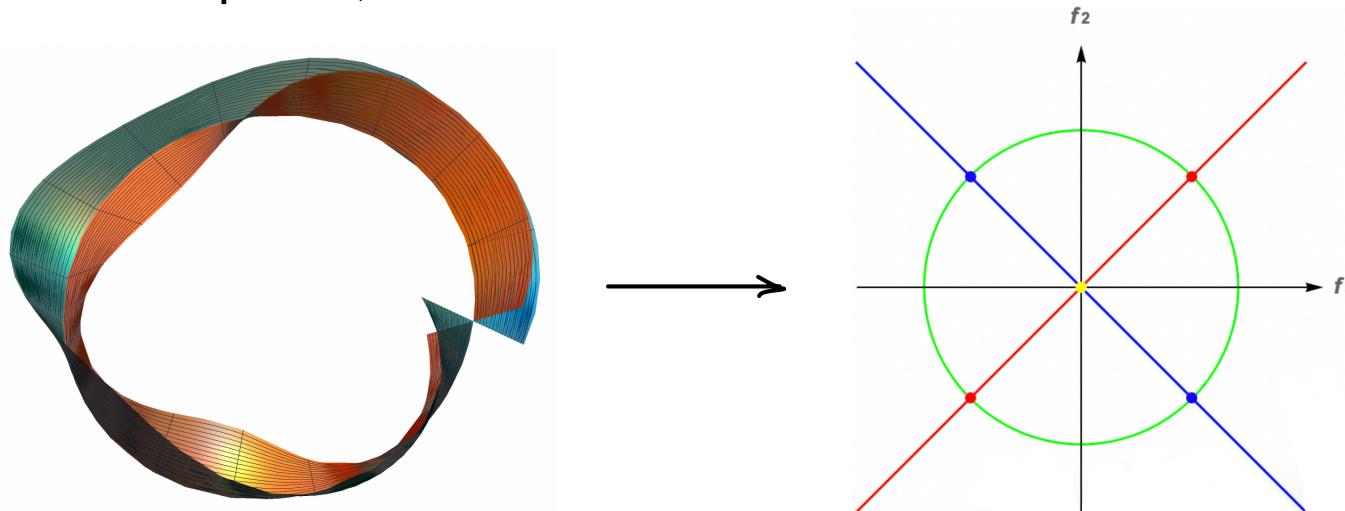
Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the stratified unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.



Here is a video showing the eigenframe rotation: <https://yifeizhu.github.io/swallowtail/rotate.mp4>

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the stratified unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.

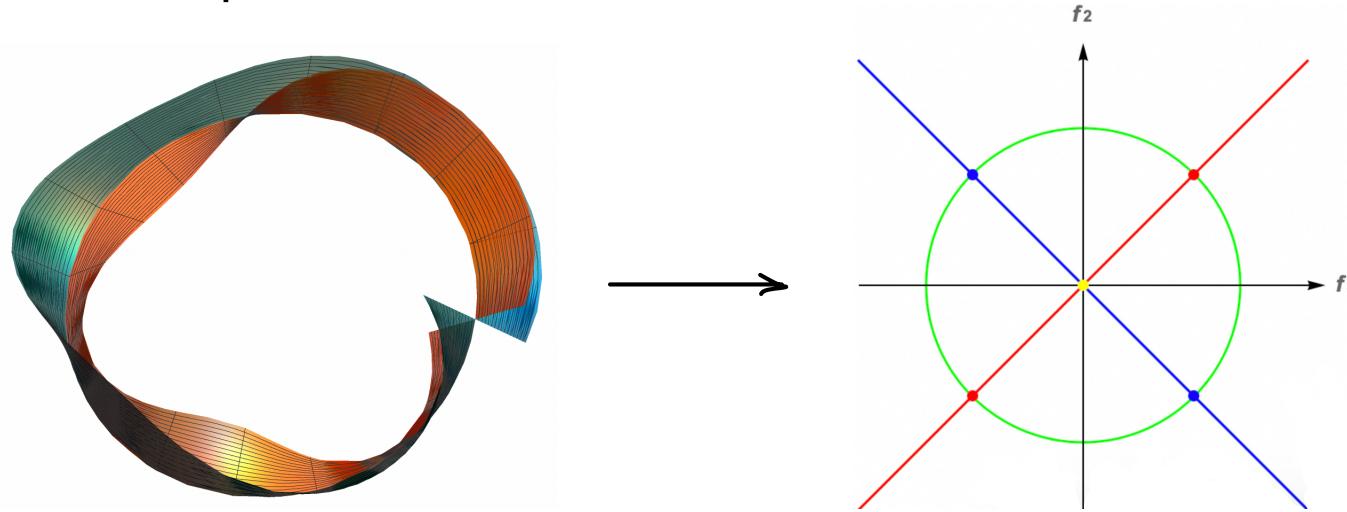


Here is a video showing the eigenframe rotation: <https://yifeizhu.github.io/swallowtail/rotate.mp4>

Note. In the non-Hermitian case, since the eigenvectors are in \mathbb{C}^2

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the stratified unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.

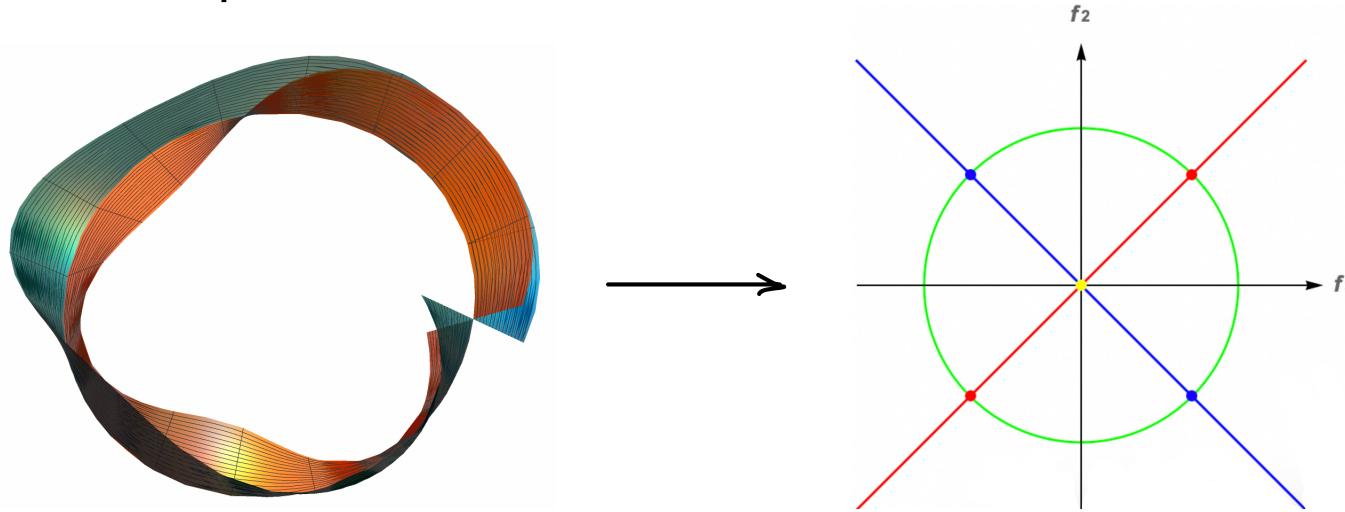


Here is a video showing the eigenframe rotation: <https://yifeizhu.github.io/swallowtail/rotate.mp4>

Note. In the non-Hermitian case, since the eigenvectors are in \mathbb{C}^2 , we have adopted (a variant of) the *Hermitian angle* to properly characterize the eigenframe rotation and degeneration

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the stratified unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.



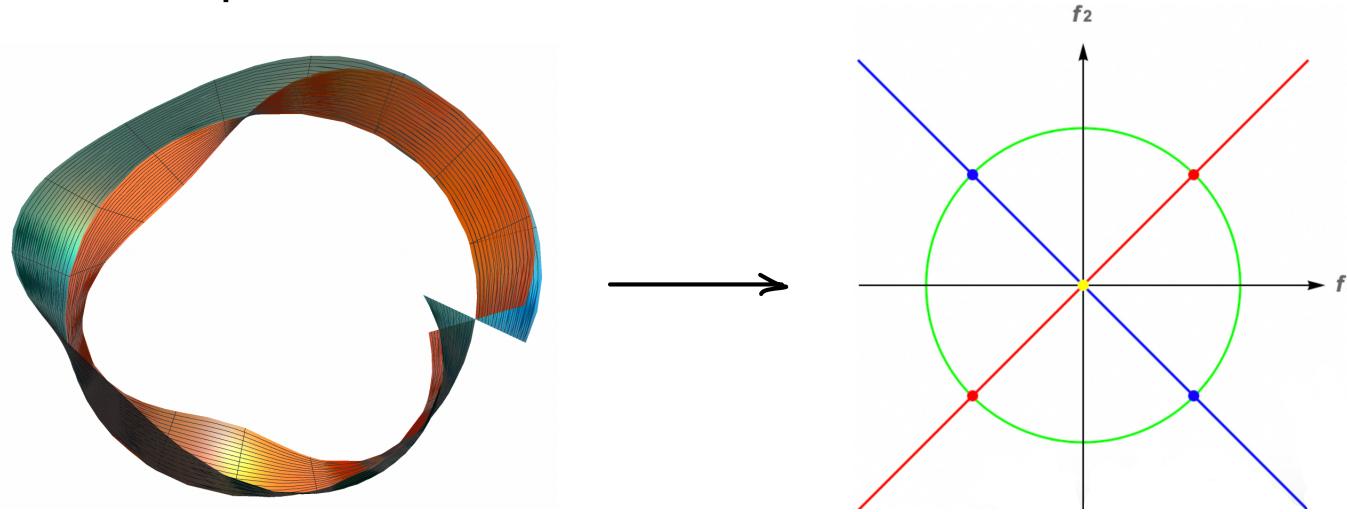
Here is a video showing the eigenframe rotation: <https://yifeizhu.github.io/swallowtail/rotate.mp4>

Note. In the non-Hermitian case, since the eigenvectors are in \mathbb{C}^2 , we have adopted (a variant of) the *Hermitian angle* to properly characterize the eigenframe rotation and degeneration:

$$\frac{\langle v_+, v_- \rangle_{\mathbb{C}}}{|v_+| |v_-|} = \rho e^{i\psi}$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the stratified unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.



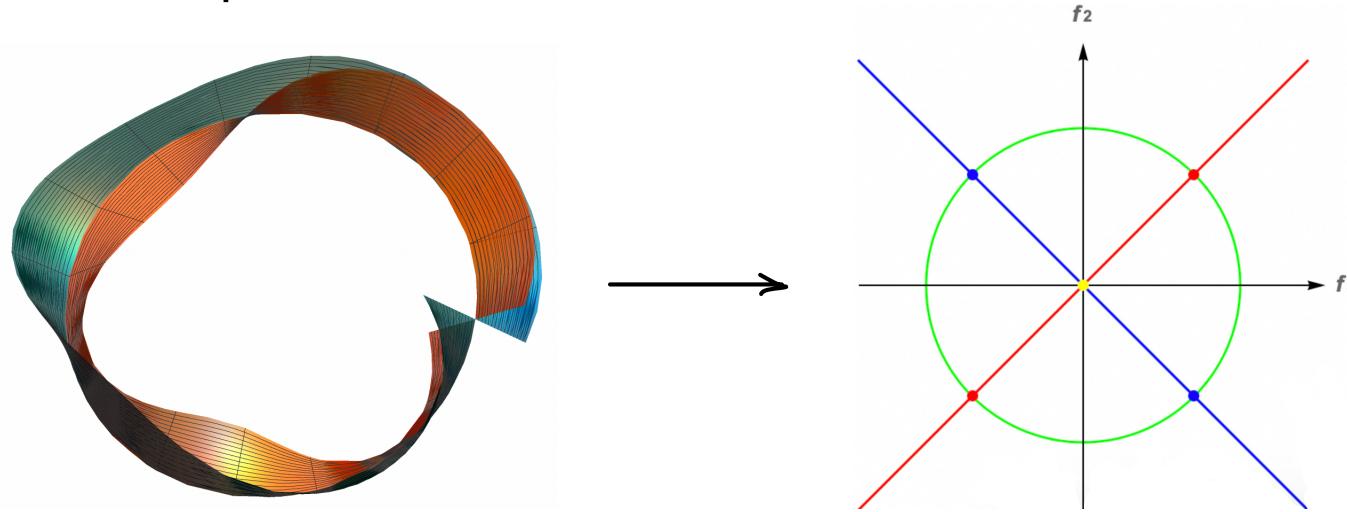
Here is a video showing the eigenframe rotation: <https://yifeizhu.github.io/swallowtail/rotate.mp4>

Note. In the non-Hermitian case, since the eigenvectors are in \mathbb{C}^2 , we have adopted (a variant of) the *Hermitian angle* to properly characterize the eigenframe rotation and degeneration:

$$\frac{\langle v_+, v_- \rangle_{\mathbb{C}}}{|v_+| |v_-|} = \rho e^{i\psi}, \quad \cos(v_+, v_-)_{\text{Herm}} := \rho$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Proposition. The universal eigenbundle for non-Hermitian 2-band systems is given by a pair of kissing half Möbius bands over the stratified unit circle in the punctured parameter plane, whose 0-dimensional stratum consists of 4 points.



Here is a video showing the eigenframe rotation: <https://yifeizhu.github.io/swallowtail/rotate.mp4>

Note. In the non-Hermitian case, since the eigenvectors are in \mathbb{C}^2 , we have adopted (a variant of) the *Hermitian angle* to properly characterize the eigenframe rotation and degeneration:

$$\frac{\langle v_+, v_- \rangle_{\mathbb{C}}}{|v_+| |v_-|} = \rho e^{i\psi}, \quad \cos(v_+, v_-)_{\text{Herm}} := \rho$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the **stratified** moduli space.

*Gajer, The intersection Dold–Thom theorem,
Topology, 1996. (Ph.D. student of Blaine Lawson, 1993)*

Goresky and MacPherson, 1974.

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the stratified moduli space.

- 0'th **intersection homology group** recovers the Hermitian 2-band charge of \mathbb{Z} .

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the stratified moduli space.

- 0'th intersection homology group recovers the Hermitian 2-band charge of \mathbb{Z} .
- Need compatibility with our earlier ad hoc classification:

$$\pi_1(S^1 \vee S^1 \vee S^1) \cong \mathbb{Z} * \mathbb{Z} * \mathbb{Z}$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the stratified moduli space.

- 0'th intersection homology group recovers the Hermitian 2-band charge of \mathbb{Z} .
- Need compatibility with our earlier ad hoc classification:

$$\pi_1(S^1 \vee S^1 \vee S^1) \cong \mathbb{Z} * \mathbb{Z} * \mathbb{Z}$$

Question. How does the eigenframe rotation in the **non-Hermitian case** relate to that in the **Hermitian case**?

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the stratified moduli space.

- 0'th intersection homology group recovers the Hermitian 2-band charge of \mathbb{Z} .
- Need compatibility with our earlier ad hoc classification:

$$\pi_1(S^1 \vee S^1 \vee S^1) \cong \mathbb{Z} * \mathbb{Z} * \mathbb{Z}$$

Question. How does the eigenframe rotation in the non-Hermitian case relate to that in the Hermitian case?

Conjecture. It does so through a *deformation* (or homotopy) of Riemannian metrics

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the stratified moduli space.

- 0'th intersection homology group recovers the Hermitian 2-band charge of \mathbb{Z} .
- Need compatibility with our earlier ad hoc classification:

$$\pi_1(S^1 \vee S^1 \vee S^1) \cong \mathbb{Z} * \mathbb{Z} * \mathbb{Z}$$

Question. How does the eigenframe rotation in the non-Hermitian case relate to that in the Hermitian case?

Conjecture. It does so through a *deformation* (or homotopy) of Riemannian metrics, i.e., a 1-parameter continuous family $\{\eta_t\}_{0 \leq t \leq 1}$ of metrics with

$$\eta_t = \begin{bmatrix} e^{i\pi t} & 0 \\ 0 & 1 \end{bmatrix}$$

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How to compute the topological charge?

In progress: Need to compute the *intersection fundamental group* of the stratified moduli space.

- 0'th intersection homology group recovers the Hermitian 2-band charge of \mathbb{Z} .
- Need compatibility with our earlier ad hoc classification:

$$\pi_1(S^1 \vee S^1 \vee S^1) \cong \mathbb{Z} * \mathbb{Z} * \mathbb{Z}$$

Question. How does the eigenframe rotation in the non-Hermitian case relate to that in the Hermitian case?

Conjecture. It does so through a *deformation* (or homotopy) of Riemannian metrics, i.e., a 1-parameter continuous family $\{\eta_t\}_{0 \leq t \leq 1}$ of metrics with

$$\eta_t = \begin{bmatrix} e^{i\pi t} & 0 \\ 0 & 1 \end{bmatrix}$$

Here is a video of the eigenbundle deformation: <https://yifeizhu.github.io/swallowtail/deform.mp4>

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian **3-band** systems?

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered **cancellation of charges**

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered **cancellation of charges** as well as **reduction to the 2-band case**.

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

$$\begin{bmatrix} 1 & f_1 & f_2 \\ -f_1 & -1 & f_3 \\ -f_2 & f_3 & -1 \end{bmatrix}$$

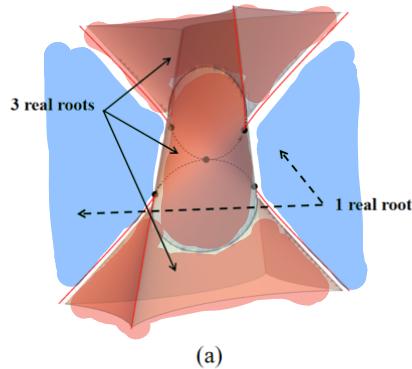
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

$$\begin{bmatrix} 1 & f_1 & f_2 \\ -f_1 & -1 & f_3 \\ -f_2 & f_3 & -1 \end{bmatrix}$$



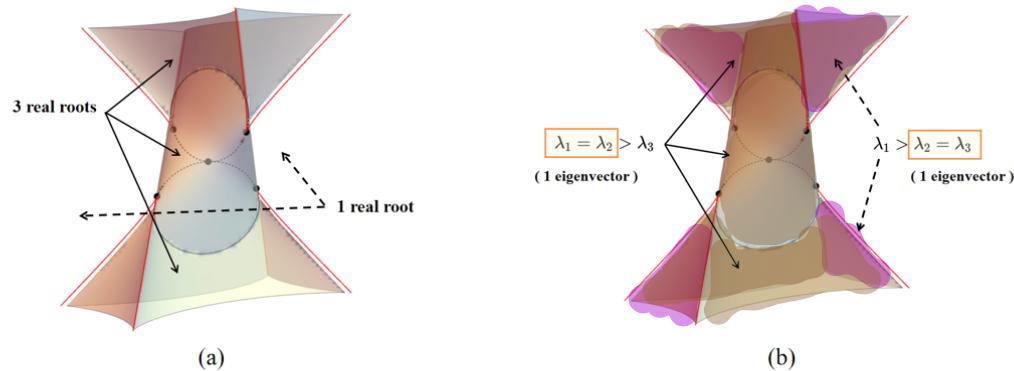
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

$$\begin{bmatrix} 1 & f_1 & f_2 \\ -f_1 & -1 & f_3 \\ -f_2 & f_3 & -1 \end{bmatrix}$$



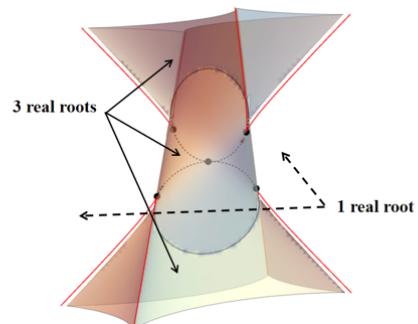
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

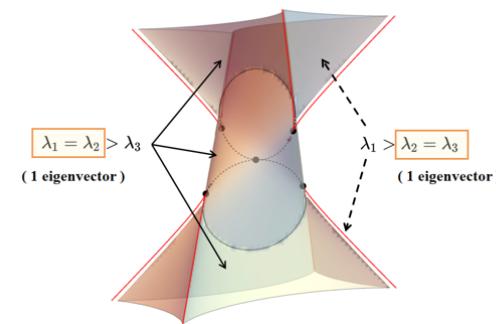
In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

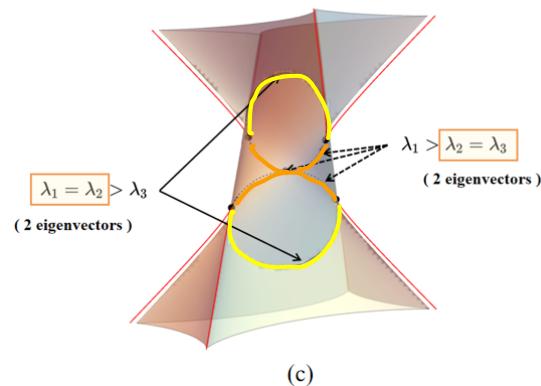
$$\begin{bmatrix} 1 & f_1 & f_2 \\ -f_1 & -1 & f_3 \\ -f_2 & f_3 & -1 \end{bmatrix}$$



(a)



(b)



(c)

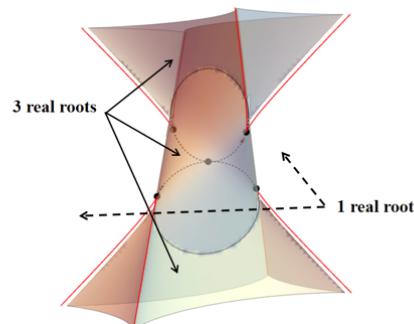
Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

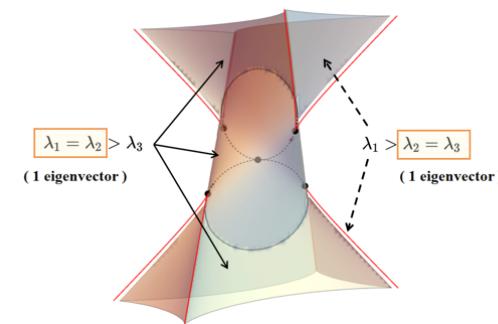
In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

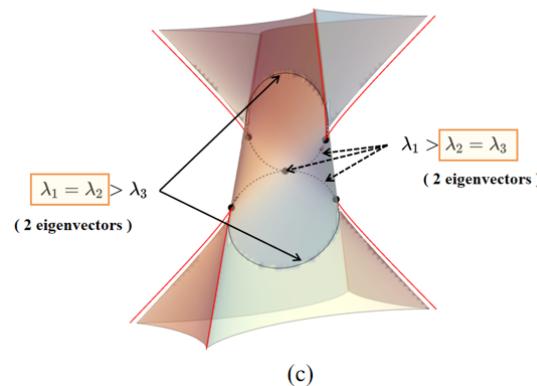
$$\begin{bmatrix} 1 & f_1 & f_2 \\ -f_1 & -1 & f_3 \\ -f_2 & f_3 & -1 \end{bmatrix}$$



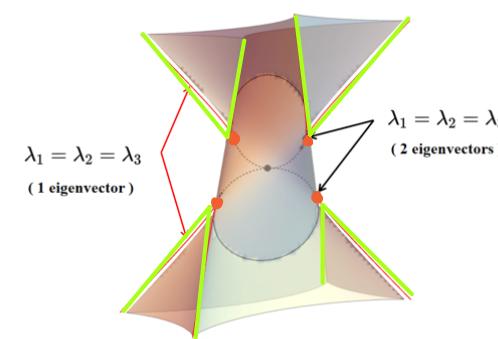
(a)



(b)



(c)



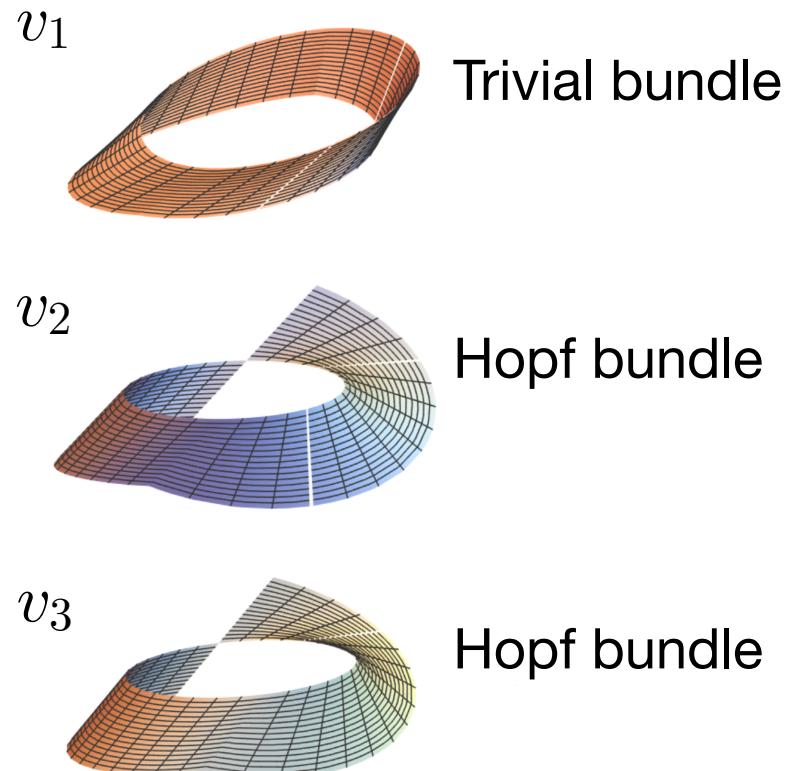
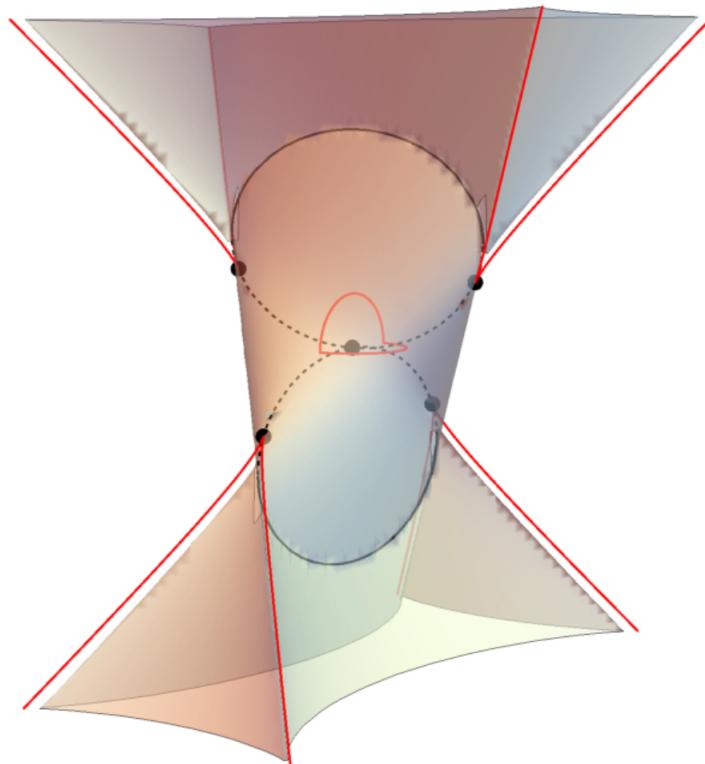
(d)

Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

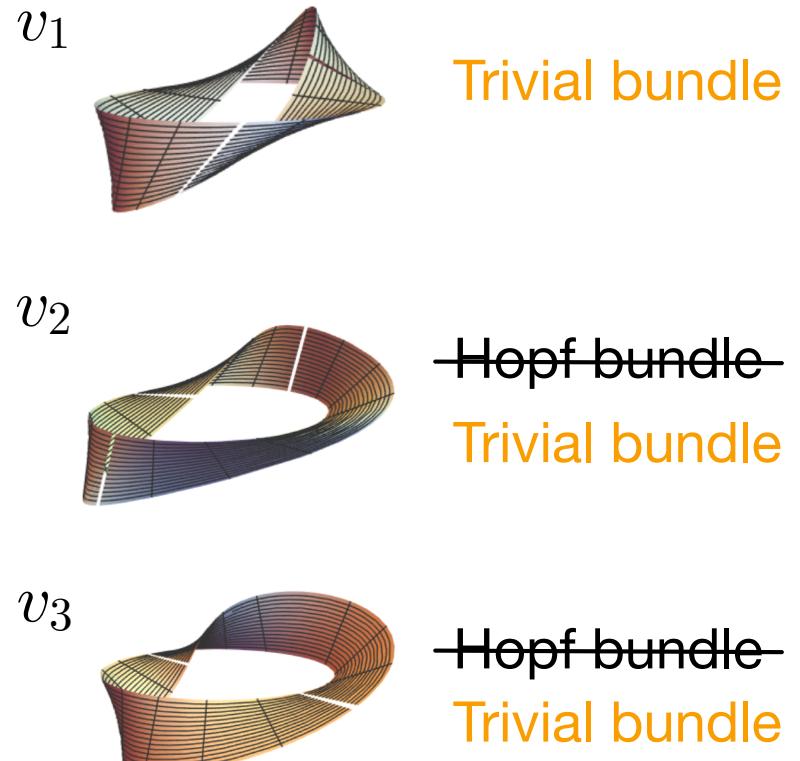
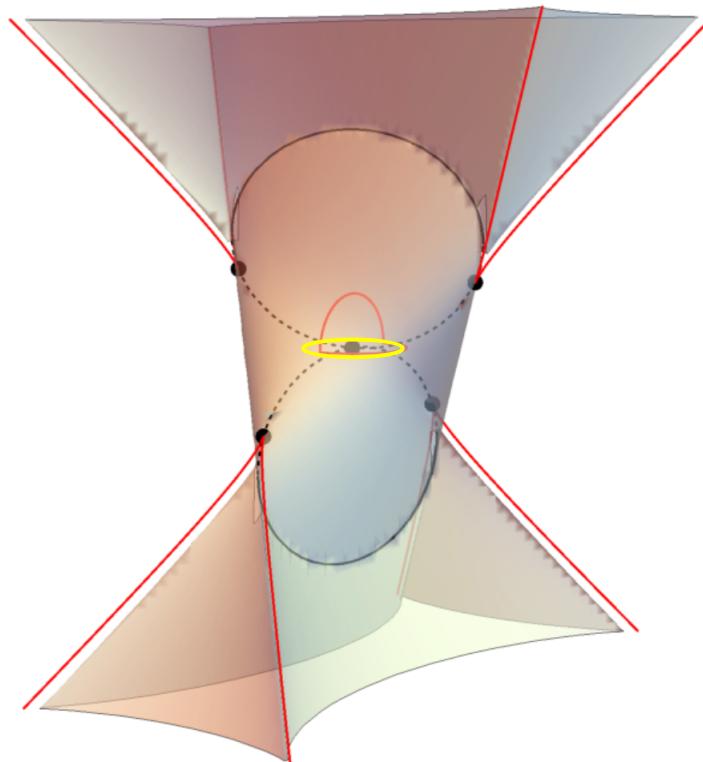


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered **cancellation of charges** as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

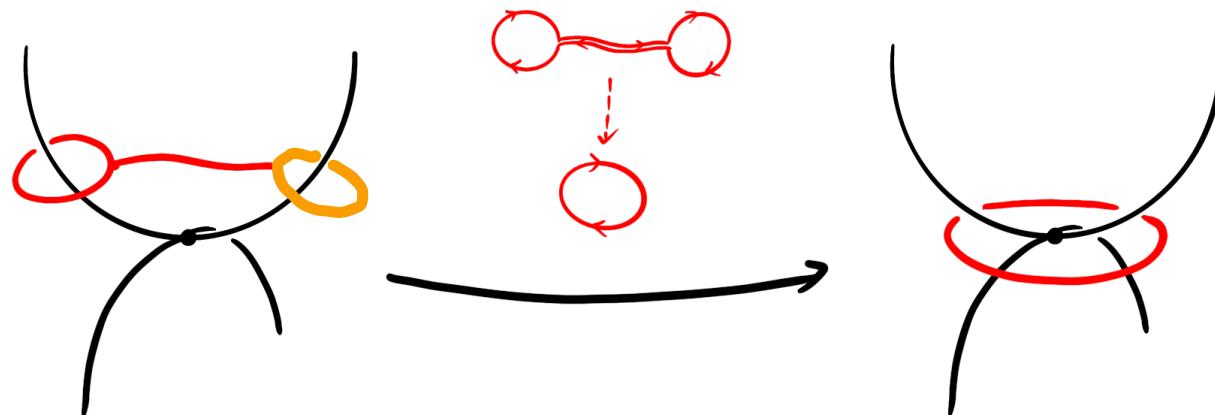


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered **cancellation of charges** as well as reduction to the 2-band case.

Example (Swallowtail quadruple sw4).

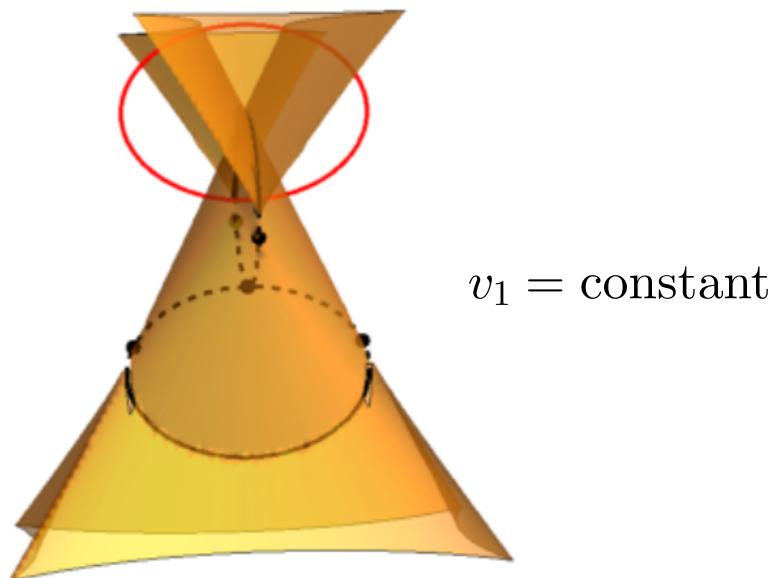


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as **reduction to the 2-band case**.

Example (Swallowtail quadruple sw4).

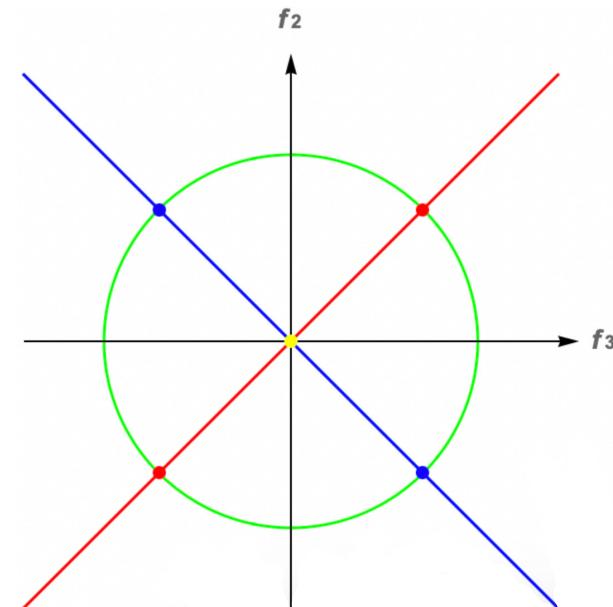
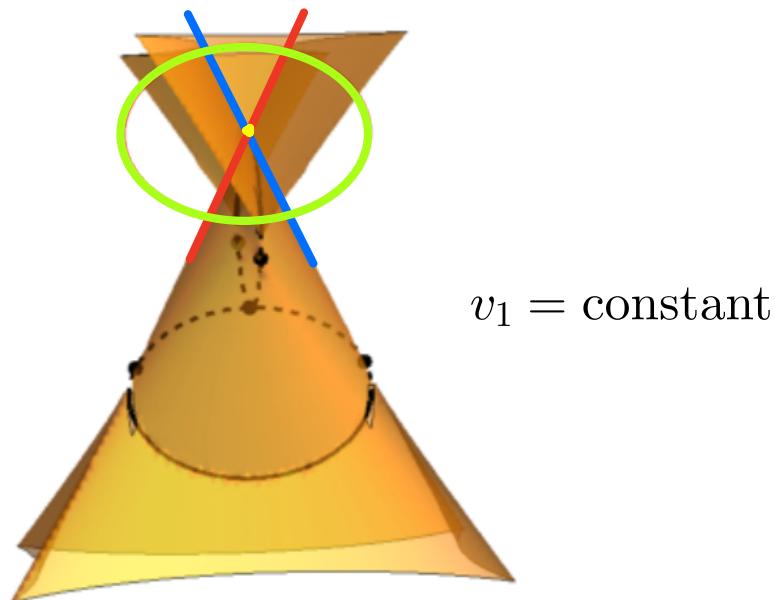


Eigenframe rotation as Higgs bundles: The non-Hermitian case

Question. How does eigenframe rotate in non-Hermitian 3-band systems?

In progress: We have investigated *slices* of the 3D moduli spaces containing swallowtails, and discovered cancellation of charges as well as **reduction to the 2-band case**.

Example (Swallowtail quadruple sw4).



Bulk-edge correspondence

We have been **experimentally** investigating the *bulk–edge correspondence* for hypersurface singularities stemmed from our theoretical analysis with the swallowtail moduli spaces

Bulk-edge correspondence

We have been **experimentally** investigating the **bulk–edge correspondence** for hypersurface singularities stemmed from our theoretical analysis with the swallowtail moduli spaces, i.e., the topology of **bulk states** (moduli space) determines the numerology of **edge states** (parametrized system).

Bulk-edge correspondence

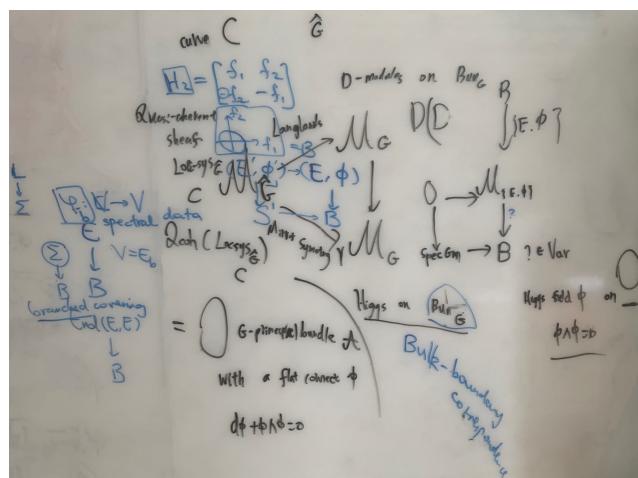
We have been **experimentally** investigating the **bulk–edge correspondence** for hypersurface singularities stemmed from our theoretical analysis with the swallowtail moduli spaces, i.e., the topology of **bulk states** (moduli space) determines the numerology of **edge states** (parametrized system).

There has not been a rigorous mathematical explanation for such a correspondence in general, but it is reminiscent of the **Langlands duality**.

Bulk-edge correspondence

We have been **experimentally** investigating the ***bulk-edge correspondence*** for hypersurface singularities stemmed from our theoretical analysis with the swallowtail moduli spaces, i.e., the topology of ***bulk states*** (moduli space) determines the numerology of ***edge states*** (parametrized system).

There has not been a rigorous mathematical explanation for such a correspondence in general, but it is reminiscent of the [Langlands duality](#).

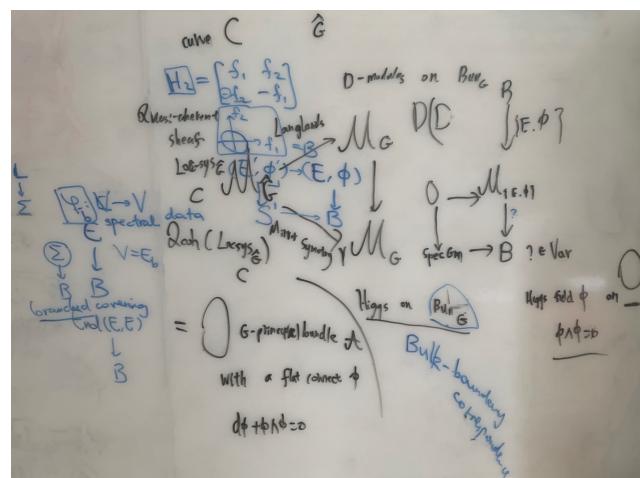


Bulk-edge correspondence

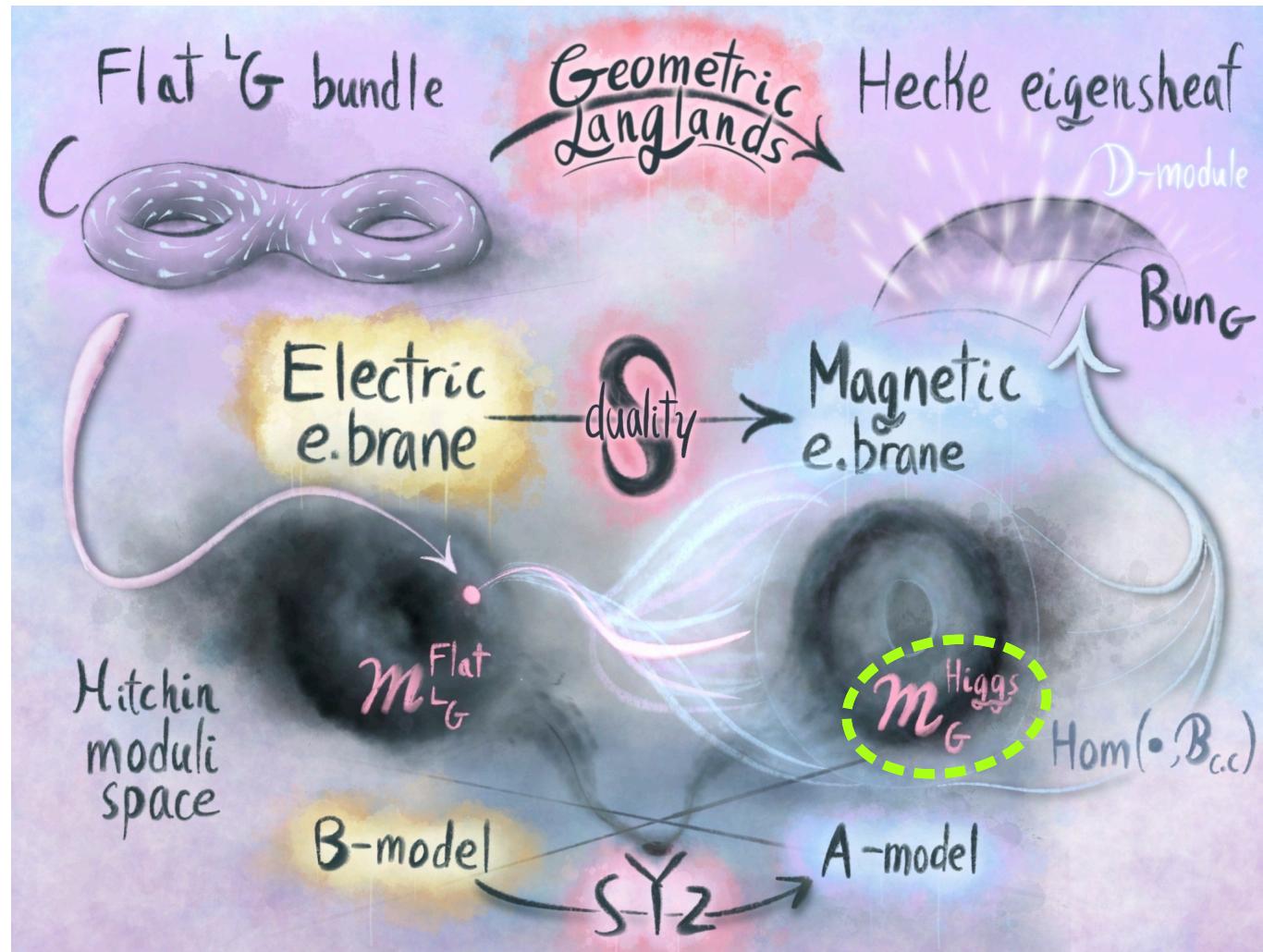
We have been **experimentally** investigating the ***bulk-edge correspondence*** for hypersurface singularities stemmed from our theoretical analysis with the swallowtail moduli spaces, i.e., the topology of ***bulk states*** (moduli space) determines the numerology of ***edge states*** (parametrized system).

There has not been a rigorous mathematical explanation for such a correspondence in general, but it is reminiscent of the [Langlands duality](#).

Indeed, Higgs bundles sit on one side of the geometric Langlands duality!
We've at least found some testing ground.



Thank you.



Portrait by Elliot Kienzle

Illustration credits

- p. 14, Weidong Luo
- p. 21, Zhou Fang
- p. 26, Zhou Fang
- pp. 58–59, Boris Khesin and Sergei Tabachnikov. Vladimir Igorevich Arnold, 12 June 1937 — 3 June 2010. Biogr. Mem. Fell. R. Soc. 64, 7–26, 2018.
- p. 61, I.M. Gelfand, M.M. Kapranov, and A.V. Zelevinsky, Discriminants, resultants and multidimensional determinants. Birkhäuser, 1994.
- p. 114, Elliot Kienzle and Steven Rayan. Hyperbolic band theory through Higgs bundles. Adv. Math., 409:Paper No. 108664, 53, 2022.
- pp. 141–148, Chenlu Huang
- p. 153, Xuecai Ma
- p. 155, Elliot Kienzle

Plots are made possible thanks to Mathematica and Notability.