

Space Geometry Nonlinear Total-Lagrange Timoshenko Beam Theory

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1 Rotation Vector and Matrix

First define two vector:

$$\vec{a} = (a_1, a_2, a_3) = a_1\vec{i} + a_2\vec{j} + a_3\vec{k} \quad (1)$$

$$\vec{b} = (b_1, b_2, b_3) = b_1\vec{i} + b_2\vec{j} + b_3\vec{k} \quad (2)$$

The cross product of the two vectors is:

$$\vec{a} \times \vec{b} = (a_2b_3 - a_3b_2)\vec{i} + (a_3b_1 - a_1b_3)\vec{j} + (a_1b_2 - a_2b_1)\vec{k} \quad (3)$$

In matrix form:

$$\vec{a} \times \vec{b} = \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix} = \tilde{a}\mathbf{b} \quad (4)$$

\tilde{a} is the skew-symmetric matrix of \mathbf{a} . And we can obtain from the vector multiply rule:

$$\vec{a} \times \vec{b} = -\vec{b} \times \vec{a} \quad (5)$$

$$\tilde{a}\mathbf{b} = -\tilde{b}\mathbf{a} = \tilde{b}^T\mathbf{a} \quad (6)$$

$$(\vec{a} \times \vec{b}) \times \vec{c} = \vec{a} \times (\vec{b} \times \vec{c}) - \vec{b} \times (\vec{a} \times \vec{c}) \quad (7)$$

$$\widetilde{\tilde{a}\mathbf{b}\mathbf{c}} = \tilde{a}\tilde{b}\mathbf{c} - \tilde{b}\tilde{a}\mathbf{c} \quad (8)$$

$$\widetilde{\tilde{a}\mathbf{b}} = \tilde{a}\tilde{b} - \tilde{b}\tilde{a} \quad (9)$$

The rotational spinor is:

$$\vec{w} = (w_1, w_2, w_3) \quad (10)$$

Normalize the spinor to get the rotation vector:

$$\vec{\theta} = \theta\vec{n} = \frac{\theta}{w}\vec{w} \quad (11)$$

And $w = |\vec{w}|, \theta = |\vec{\theta}|$. From Fig.1 we can see that:

$$\vec{x}_{\vec{\theta}} = \vec{x} + \vec{b} + \vec{c} = \vec{x} + (\vec{n} \times \vec{x}) \sin \theta + (\vec{n} \times (\vec{n} \times \vec{x}))(1 - \cos \theta) \quad (12)$$

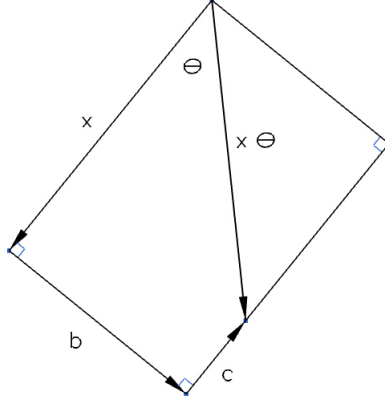


Figure 1: Rotation Vector

Because $\vec{n} \times \vec{x} = \frac{\tilde{w}}{w} \mathbf{x}, \mathbf{x}_\theta = \mathbf{R}^T \mathbf{x}$, we can get:

$$\begin{aligned} \mathbf{R}^T \mathbf{x} &= \mathbf{x} + \frac{\tilde{w}}{w} \mathbf{x} \sin \theta + \frac{\tilde{w}^2}{w^2} \mathbf{x} (1 - \cos \theta), \\ \mathbf{R}^T &= \mathbf{I} + \frac{\sin \theta}{w} \tilde{w} + \frac{1 - \cos \theta}{w^2} \tilde{w}^2 \end{aligned} \quad (13)$$

Here, \mathbf{R}^T is the rotation matrix and $\mathbf{R}^T \mathbf{R} = \mathbf{I}$. Differential this equation with X we can get:

$$\dot{\mathbf{R}}^T \mathbf{R} + \mathbf{R}^T \dot{\mathbf{R}} = 0 \quad (14)$$

That means $\dot{\mathbf{R}}^T \mathbf{R}$ is skew-symmetric matrix.

Then we consider the matrix exponential:

$$\text{Exp}(\gamma \tilde{w}) = \mathbf{I} + \gamma \tilde{w} + \frac{\gamma^2}{2!} \tilde{w}^2 + \frac{\gamma^3}{3!} \tilde{w}^3 + \dots \quad (15)$$

Using $\tilde{w}^n = -w^2 \tilde{w}^{n-2}$, we can get:

$$\text{Exp}(\gamma \tilde{w}) = \mathbf{I} + (\gamma - \frac{w^2}{3!} + \frac{w^4}{5!} + \dots) \tilde{w} + (\frac{\gamma^2}{2!} - \frac{\gamma^2 w^2}{4!} + \frac{\gamma^6 w^4}{6!} + \dots) \tilde{w}^2 \quad (16)$$

Given the series expansion on $\sin \theta$ and $1 - \cos \theta$:

$$\sin \theta = \theta - \frac{\theta^3}{6} + \frac{\theta^5}{120} + \dots \quad (17)$$

$$1 - \cos \theta = \frac{\theta^2}{2} - \frac{\theta^4}{24} + \frac{\theta^6}{720} + \dots \quad (18)$$

Substitute into the exponential equation:

$$\text{Exp}(\gamma \tilde{w}) = \mathbf{I} + \frac{\sin(\gamma w)}{w} \tilde{w} + \frac{1 - \cos(\gamma w)}{w^2} \tilde{w}^2 \quad (19)$$

This lead to the final equation:

$$\mathbf{R}^T = \text{Exp}\left(\frac{\theta}{w}\tilde{\mathbf{w}}\right) = \text{Exp}(\tilde{\boldsymbol{\theta}}) \quad (20)$$

In the same way, we can get:

$$\mathbf{R} = \text{Exp}(-\tilde{\boldsymbol{\theta}}) \quad (21)$$

And some more useful equations:

$$\delta \mathbf{R} = -\delta \tilde{\boldsymbol{\theta}} \mathbf{R} \quad (22)$$

$$\delta \tilde{\boldsymbol{\theta}} = -\delta \mathbf{R} \mathbf{R}^T \quad (23)$$

$$\frac{d\tilde{\boldsymbol{\theta}}}{dX_1} = -\frac{d\mathbf{R}}{dX_1} \mathbf{R}^T = \mathbf{R} \frac{d\mathbf{R}^T}{dX_1} = \tilde{\mathbf{k}} \quad (24)$$

Here \mathbf{k} is the curvature matrix.

$$\begin{aligned} \delta \tilde{\mathbf{k}} &= \delta \mathbf{R} \frac{d\mathbf{R}^T}{dX_1} + \mathbf{R} \frac{d\delta \mathbf{R}^T}{dX_1} \\ &= -\delta \tilde{\boldsymbol{\theta}} \mathbf{R} \frac{d\mathbf{R}^T}{dX_1} + \mathbf{R} \frac{d(\mathbf{R}^T \delta \tilde{\boldsymbol{\theta}})}{dX_1} \\ &= -\delta \tilde{\boldsymbol{\theta}} \tilde{\mathbf{k}} + \tilde{\mathbf{k}} \delta \tilde{\boldsymbol{\theta}} + \frac{d\delta \tilde{\boldsymbol{\theta}}}{dX_1} \end{aligned} \quad (25)$$

Multiply 2 sides by $\delta \boldsymbol{\theta}$,

$$\begin{aligned} \delta \tilde{\mathbf{k}} \delta \boldsymbol{\theta} &= -\delta \tilde{\boldsymbol{\theta}} \delta \mathbf{k} = -\delta \tilde{\boldsymbol{\theta}} \tilde{\mathbf{k}} \delta \boldsymbol{\theta} + \tilde{\mathbf{k}} \delta \tilde{\boldsymbol{\theta}} \delta \boldsymbol{\theta} + \frac{d\delta \tilde{\boldsymbol{\theta}}}{dX_1} \delta \boldsymbol{\theta} \\ &= -\delta \tilde{\boldsymbol{\theta}} \tilde{\mathbf{k}} \delta \boldsymbol{\theta} + 0 - \delta \tilde{\boldsymbol{\theta}} \frac{d\delta \boldsymbol{\theta}}{dX_1}, \\ \delta \mathbf{k} &= \tilde{\mathbf{k}} \delta \boldsymbol{\theta} + \frac{d\delta \boldsymbol{\theta}}{dX_1}. \end{aligned} \quad (26)$$

2 Basic Theory on Geometry Nonlinear FEM

2.1 Strain and Stress

Given the Green-Lagrange strain:

$$\mathbf{e} = \frac{1}{2}(\mathbf{G} + \mathbf{G}^T + \mathbf{G}^T \mathbf{G}) \quad (27)$$

\mathbf{G} is the displacement gradient tensor:

$$\mathbf{G} = [\mathbf{g}_1 \quad \mathbf{g}_2 \quad \mathbf{g}_3] \quad (28)$$

$$\mathbf{e} = \frac{1}{2}([\mathbf{g}_1 \quad \mathbf{g}_2 \quad \mathbf{g}_3] + \begin{bmatrix} \mathbf{g}_1^T \\ \mathbf{g}_2^T \\ \mathbf{g}_3^T \end{bmatrix} + \begin{bmatrix} \mathbf{g}_1^T \\ \mathbf{g}_2^T \\ \mathbf{g}_3^T \end{bmatrix} [\mathbf{g}_1 \quad \mathbf{g}_2 \quad \mathbf{g}_3]) \quad (29)$$

Define $\mathbf{h}_1, \mathbf{h}_2, \mathbf{h}_3$:

$$\mathbf{h}_1 = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}, \mathbf{h}_2 = \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}, \mathbf{h}_3 = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \quad (30)$$

Using $\mathbf{h}_1, \mathbf{h}_2, \mathbf{h}_3$ to Expand e:

$$\mathbf{e} = \frac{1}{2} \left(\begin{bmatrix} \mathbf{h}_1^T \mathbf{g}_1 & \mathbf{h}_1^T \mathbf{g}_2 & \mathbf{h}_1^T \mathbf{g}_3 \\ \mathbf{h}_2^T \mathbf{g}_1 & \mathbf{h}_2^T \mathbf{g}_2 & \mathbf{h}_2^T \mathbf{g}_3 \\ \mathbf{h}_3^T \mathbf{g}_1 & \mathbf{h}_3^T \mathbf{g}_2 & \mathbf{h}_3^T \mathbf{g}_3 \end{bmatrix} + \begin{bmatrix} \mathbf{h}_1^T \mathbf{g}_1 & \mathbf{h}_2^T \mathbf{g}_1 & \mathbf{h}_3^T \mathbf{g}_1 \\ \mathbf{h}_1^T \mathbf{g}_2 & \mathbf{h}_2^T \mathbf{g}_2 & \mathbf{h}_3^T \mathbf{g}_2 \\ \mathbf{h}_1^T \mathbf{g}_3 & \mathbf{h}_2^T \mathbf{g}_3 & \mathbf{h}_3^T \mathbf{g}_3 \end{bmatrix} + \begin{bmatrix} \mathbf{g}_1^T \mathbf{g}_1 & \mathbf{g}_1^T \mathbf{g}_2 & \mathbf{g}_1^T \mathbf{g}_3 \\ \mathbf{g}_2^T \mathbf{g}_1 & \mathbf{g}_2^T \mathbf{g}_2 & \mathbf{g}_2^T \mathbf{g}_3 \\ \mathbf{g}_3^T \mathbf{g}_1 & \mathbf{g}_3^T \mathbf{g}_2 & \mathbf{g}_3^T \mathbf{g}_3 \end{bmatrix} \right) \quad (31)$$

$$e_{ij} = \frac{1}{2} (\mathbf{h}_i^T \mathbf{g}_j + \mathbf{h}_j^T \mathbf{g}_i + \mathbf{g}_i^T \mathbf{g}_j) \quad (32)$$

Rearrange \mathbf{e} and \mathbf{G} in column form:

$$e_i = \bar{\mathbf{h}}_i^T \mathbf{g} + \frac{1}{2} \mathbf{g}^T \mathbf{H}_i^T \mathbf{g} \quad (33)$$

$$\mathbf{g} = \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} \quad (34)$$

And $\bar{\mathbf{h}}_i$ is the combination of \mathbf{h}_i , \mathbf{H}_i is symmetric matrix. Next, we define:

$$\mathbf{C}_i = \bar{\mathbf{h}}_i + \frac{1}{2} \mathbf{H}_i \mathbf{g} \quad (35)$$

Then:

$$e_i = \mathbf{C}_i^T \mathbf{g} \quad (36)$$

The conjugate second Piola-Kirchhoff(PK2) stress is:

$$s_i = E_{ij} e_j = E_{ij} \mathbf{C}_j^T \mathbf{g} \quad (37)$$

2.2 Energy Equation and Stiffness Matrix

The strain energy density is:

$$u = \frac{1}{2} s_i e_i = \frac{1}{2} \mathbf{g}^T \mathbf{C}_j E_{ij} \mathbf{C}_i^T \mathbf{g} = \frac{1}{2} \mathbf{g}^T \mathbf{S}^u \mathbf{g} \quad (38)$$

Here we get the core energy stiffness matrix:

$$\mathbf{S}^u = \mathbf{C}_j E_{ij} \mathbf{C}_i^T = E_{ij} \mathbf{C}_i \mathbf{C}_j^T \quad (39)$$

Variation of u, we can get:

$$\begin{aligned} \delta u &= \frac{1}{2} \delta \mathbf{g}^T \mathbf{S}^u \mathbf{g} + \frac{1}{2} \mathbf{g}^T \delta (\mathbf{S}^u \mathbf{g}) \\ &= \frac{1}{2} \delta \mathbf{g}^T \mathbf{S}^u \mathbf{g} + \frac{1}{2} \mathbf{g}^T (\delta \mathbf{S}^u \mathbf{g} + \mathbf{S}^u \delta \mathbf{g}) \\ &= \delta \mathbf{g}^T \mathbf{S}^u \mathbf{g} + \frac{1}{2} \mathbf{g}^T \delta \mathbf{S}^u \mathbf{g} = \delta \mathbf{g}^T \mathbf{S}^r \mathbf{g} = \delta \mathbf{g}^T \Phi \end{aligned} \quad (40)$$

This implicitly define the secant stiffness matrix \mathbf{S}^r . Φ is the core internal force vector.

Second variation of u, we can get:

$$\delta^2 u = \delta \mathbf{g}^T \delta \Phi + (\delta^2 \mathbf{g})^T \Phi \quad (41)$$

And because $\delta \Phi = \mathbf{S}^t \delta \mathbf{g}$, here \mathbf{S}^t is the core tangent stiffness matrix:

$$\delta^2 u = \delta \mathbf{g}^T \mathbf{S}^t \delta \mathbf{g} + (\delta^2 \mathbf{g})^T \Phi \quad (42)$$

Next, integration of energy density u to get the equation of energy U, also in second variation form:

$$\delta^2 U = \int_{V_0} \delta^2 u dV = \int_{V_0} (\delta \mathbf{g}^T \mathbf{S}^t \delta \mathbf{g} + (\delta^2 \mathbf{g})^T \Phi) dV \quad (43)$$

And the relation between displacement gradient \mathbf{g} and final DOF \mathbf{q} is:

$$\delta \mathbf{g} = \mathbf{T} \delta \mathbf{q} \quad (44)$$

Substitute into energy equation:

$$\delta^2 U = \int_{V_0} (\delta \mathbf{q}^T \mathbf{T}^T \mathbf{S}^t \mathbf{T} \delta \mathbf{q} + (\delta^2 \mathbf{g})^T \Phi) dV \quad (45)$$

From this equation, we can get the final tangent stiffness matrix. And the final internal force vector is the same way:

$$\delta U = \int_{V_0} \delta u dv = \int_{V_0} \delta \mathbf{g}^T \Phi dv = \int_{V_0} \delta \mathbf{q}^T \mathbf{T}^T \Phi dv = \delta \mathbf{q}^T \int_{V_0} \mathbf{T}^T \Phi dv \quad (46)$$

Next we derive the detailed expression of stiffness and force matrices:

$$\delta u = \delta \mathbf{g}^T \mathbf{S}^u \mathbf{g} + \frac{1}{2} \mathbf{g}^T \delta \mathbf{S}^u \mathbf{g} \quad (47)$$

We know that $\mathbf{S}^u = E_{ij} \mathbf{C}_i \mathbf{C}_j^T$, and we can get:

$$\begin{aligned} \frac{1}{2} \mathbf{g}^T \delta \mathbf{S}^u \mathbf{g} &= \frac{1}{2} \mathbf{g}^T \delta (E_{ij} \mathbf{C}_i \mathbf{C}_j^T) \mathbf{g} \\ &= \frac{E_{ij}}{2} \mathbf{g}^T (\delta \mathbf{C}_i \mathbf{C}_j^T + \mathbf{C}_i \delta \mathbf{C}_j^T) \mathbf{g} \\ &= \frac{E_{ij}}{2} \mathbf{g}^T \left(\frac{1}{2} \mathbf{H}_i \delta \mathbf{g} \mathbf{C}_j^T + \frac{1}{2} \mathbf{C}_i \delta \mathbf{g}^T \mathbf{H}_j^T \right) \mathbf{g} \\ &= \frac{E_{ij}}{4} \mathbf{g}^T \mathbf{H}_i \delta \mathbf{g} \mathbf{C}_j^T \mathbf{g} + \frac{E_{ij}}{4} \mathbf{g}^T \mathbf{C}_i \delta \mathbf{g}^T \mathbf{H}_j^T \mathbf{g} \\ &= \frac{1}{4} \mathbf{g}^T s_i \mathbf{H}_i \delta \mathbf{g} + \frac{1}{4} \delta \mathbf{g}^T s_j \mathbf{H}_j^T \mathbf{g} \\ &= \frac{1}{2} \delta \mathbf{g}^T s_i \mathbf{H}_i \mathbf{g} \end{aligned} \quad (48)$$

And from Eq.39 and Eq.40, we can get:

$$\mathbf{S}^r = E_{ij} \mathbf{C}_i \mathbf{C}_j^T + \frac{1}{2} s_i \mathbf{H}_i^T \quad (49)$$

And:

$$\begin{aligned}
\delta \Phi &= \delta \mathbf{S}^r \mathbf{g} + \mathbf{S}^r \delta \mathbf{g} \\
&= \delta (E_{ij} \mathbf{C}_i \mathbf{C}_j^T + \frac{1}{2} s_i \mathbf{H}_i^T) \mathbf{g} + \mathbf{S}^r \delta \mathbf{g} \\
&= (E_{ij} (\delta \mathbf{C}_i \mathbf{C}_j^T + \mathbf{C}_i \delta \mathbf{C}_j^T) + \frac{1}{2} \delta s_i \mathbf{H}_i^T) \mathbf{g} + \mathbf{S}^r \delta \mathbf{g} \\
&= (E_{ij} (\frac{1}{2} \mathbf{H}_i \delta \mathbf{g} \mathbf{C}_j^T + \frac{1}{2} \mathbf{C}_i \delta \mathbf{g}^T \mathbf{H}_j^T) + \frac{1}{2} \delta s_i \mathbf{H}_i^T) \mathbf{g} + \mathbf{S}^r \delta \mathbf{g} \\
&= \frac{E_{ij}}{2} \mathbf{H}_i \delta \mathbf{g} \mathbf{C}_j^T \mathbf{g} + \frac{E_{ij}}{2} \mathbf{C}_i \delta \mathbf{g}^T \mathbf{H}_j^T \mathbf{g} + \frac{E_{ij}}{2} \delta s_i \mathbf{H}_i^T \mathbf{g} + \mathbf{S}^r \delta \mathbf{g} \\
&= \frac{1}{2} s_i \mathbf{H}_i \delta \mathbf{g} + \frac{E_{ij}}{2} \mathbf{C}_i \mathbf{g}^T \mathbf{H}_j^T \delta \mathbf{g} + \frac{E_{ij}}{2} \mathbf{H}_i^T \mathbf{g} \delta (\mathbf{C}_j^T \mathbf{g}) + (E_{ij} \mathbf{C}_i \mathbf{C}_j^T + \frac{1}{2} s_i \mathbf{H}_i^T) \delta \mathbf{g}
\end{aligned} \tag{50}$$

And:

$$\begin{aligned}
\frac{E_{ij}}{2} \mathbf{H}_i^T \mathbf{g} \delta (\mathbf{C}_j^T \mathbf{g}) &= \frac{E_{ij}}{2} \mathbf{H}_i^T \mathbf{g} (\delta \mathbf{C}_j^T \mathbf{g} + \mathbf{C}_j^T \delta \mathbf{g}) \\
&= \frac{E_{ij}}{4} \mathbf{H}_i^T \mathbf{g} \delta \mathbf{g}^T \mathbf{H}_j^T \mathbf{g} + \frac{E_{ij}}{2} \mathbf{H}_i^T \mathbf{g} \mathbf{C}_j^T \delta \mathbf{g} \\
&= \frac{E_{ij}}{4} \mathbf{H}_i^T \mathbf{g} \mathbf{g}^T \mathbf{H}_j^T \delta \mathbf{g} + \frac{E_{ij}}{2} \mathbf{H}_i^T \mathbf{g} \mathbf{C}_j^T \delta \mathbf{g}
\end{aligned} \tag{51}$$

Substitute Eq.51 into Eq.50,we get:

$$\delta \Phi = (E_{ij} (\mathbf{C}_i + \frac{1}{2} \mathbf{H}_i \mathbf{g}) (\mathbf{C}_j^T + \frac{1}{2} \mathbf{g}^T \mathbf{H}_j^T) + s_i \mathbf{H}_i^T) \delta \mathbf{g} \tag{52}$$

Define $\mathbf{B}_i = \mathbf{C}_i + \frac{1}{2} \mathbf{H}_i \mathbf{g} = \mathbf{h}_i + \mathbf{H}_i \mathbf{g}$, then:

$$\mathbf{S}^t = E_{ij} \mathbf{B}_i \mathbf{B}_j^T + s_i \mathbf{H}_i^T \tag{53}$$

The core internal force matrix is:

$$\begin{aligned}
\Phi &= \mathbf{S}^r \mathbf{g} \\
&= (E_{ij} \mathbf{C}_i \mathbf{C}_j^T + \frac{1}{2} s_i \mathbf{H}_i^T) \mathbf{g} \\
&= E_{ij} \mathbf{C}_i \mathbf{C}_j^T \mathbf{g} + \frac{E_{ij}}{2} \mathbf{C}_j^T \mathbf{g} \mathbf{H}_i^T \mathbf{g} \\
&= E_{ij} \mathbf{C}_j^T \mathbf{g} (\mathbf{C}_i + \frac{1}{2} \mathbf{H}_i^T \mathbf{g}) = s_i \mathbf{B}_i
\end{aligned} \tag{54}$$

3 Space Total-Lagrange Beam Theory

3.1 Beam Deformation Theory

Assuming that the beam particle at (X_1, X_2, X_3) move to:

$$\mathbf{x}(\mathbf{X}) = \mathbf{x}_0(X_1) + \mathbf{R}^T(X_1) \boldsymbol{\zeta}(X_2, X_3) \tag{55}$$

Where $\boldsymbol{\zeta}^T = [0 \ X_2 \ X_3]$ is the cross-section coordinate. And the displacement is:

$$\mathbf{u} = \mathbf{x} - \mathbf{X} = \mathbf{x}_0 + \mathbf{R}^T \boldsymbol{\zeta} - (\mathbf{X}(X_1) + \boldsymbol{\zeta}) = \mathbf{u}_0 + (\mathbf{R}^T - \mathbf{I}) \boldsymbol{\zeta} \tag{56}$$

Here \mathbf{u}_0 is the centroidal displacement. Then, the displacement gradient is:

$$\mathbf{G} = \frac{\partial \mathbf{u}_0}{\partial \mathbf{X}} + \frac{\partial((\mathbf{R}^T - \mathbf{I})\boldsymbol{\zeta})}{\partial \mathbf{X}} \quad (57)$$

This can be rearrange as:

$$\mathbf{g}_i = \frac{\partial \mathbf{u}_0}{\partial X_i} + \frac{\partial((\mathbf{R}^T - \mathbf{I})\boldsymbol{\zeta})}{\partial X_i} \quad (58)$$

Then the formulation of \mathbf{g}_i is:

$$\mathbf{g}_1 = \frac{\partial \mathbf{u}_0}{\partial X_1} + \frac{\partial(\mathbf{R}^T \boldsymbol{\zeta})}{\partial X_1} = \frac{d\mathbf{u}_0}{dX_1} + \frac{d\mathbf{R}^T}{dX_1} \boldsymbol{\zeta} = \frac{d\mathbf{u}_0}{dX_1} + \mathbf{R}^T \mathbf{R} \frac{d\mathbf{R}^T}{dX_1} \boldsymbol{\zeta} \quad (59)$$

Using Eq.6 and Eq.25:

$$\begin{aligned} \mathbf{g}_1 &= \frac{d\mathbf{u}_0}{dX_1} + \mathbf{R}^T \tilde{\mathbf{k}} \boldsymbol{\zeta} = \frac{d\mathbf{u}_0}{dX_1} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k}, \\ \mathbf{g}_2 &= (\mathbf{R}^T - \mathbf{I}) \frac{\partial \boldsymbol{\zeta}}{\partial X_2} = (\mathbf{R}^T - \mathbf{I}) \mathbf{h}_2, \\ \mathbf{g}_3 &= (\mathbf{R}^T - \mathbf{I}) \frac{\partial \boldsymbol{\zeta}}{\partial X_3} = (\mathbf{R}^T - \mathbf{I}) \mathbf{h}_3. \end{aligned} \quad (60)$$

Obviously that $\mathbf{h}_2^T = [0 \ 1 \ 0]$ and $\mathbf{h}_3^T = [0 \ 0 \ 1]$. From Eq.32, we can know:

$$\begin{aligned} e_{11} &= \mathbf{h}_1^T \mathbf{g}_1 + \frac{1}{2} \mathbf{g}_1^T \mathbf{g}_1, \\ \gamma_{12} &= 2e_{12} = \mathbf{h}_2^T \mathbf{g}_1 + \mathbf{h}_1^T \mathbf{g}_2 + \mathbf{g}_1^T \mathbf{g}_2, \\ \gamma_{13} &= 2e_{13} = \mathbf{h}_3^T \mathbf{g}_1 + \mathbf{h}_1^T \mathbf{g}_3 + \mathbf{g}_1^T \mathbf{g}_3. \end{aligned} \quad (61)$$

Using Mathematica to simplify, we can know that e_{22} and e_{23} are 0. Substitute Eq.60 into Eq.61 and simplify:

$$\begin{aligned} e_{11} &= e_b + e_f, \\ \gamma_{12} &= \mathbf{h}_2^T \boldsymbol{\phi} + \mathbf{h}_2^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k}, \\ \gamma_{13} &= \mathbf{h}_3^T \boldsymbol{\phi} + \mathbf{h}_3^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k}, \\ e_b &= \left(\frac{d\mathbf{u}_0}{dX_1} \right)^T \left(\mathbf{h}_1 + \frac{d\mathbf{u}_0}{2dX_1} \right), \\ e_f &= \boldsymbol{\zeta}^T \mathbf{k}_e + \frac{1}{2} \mathbf{k}^T \tilde{\boldsymbol{\zeta}} \tilde{\boldsymbol{\zeta}}^T \mathbf{k} \simeq \boldsymbol{\zeta}^T \mathbf{k}_e, \\ \boldsymbol{\phi} &= \mathbf{R} \left(\mathbf{h}_1 + \frac{d\mathbf{u}_0}{dX_1} \right). \end{aligned} \quad (62)$$

The squared-curvature term of e_f is dropped.

3.2 Beam Core Equations

First we consider the core energy stiffness matrix \mathbf{S}^u , rewrite \mathbf{C}_i of Eq.35 in matrix form:

$$\begin{aligned} \mathbf{C}_1 &= \begin{bmatrix} \mathbf{h}_1 \\ 0 \\ 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} \mathbf{H} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{c}_1 \\ 0 \\ 0 \end{bmatrix}, \\ \mathbf{C}_2 &= \begin{bmatrix} \mathbf{h}_2 \\ \mathbf{h}_1 \\ 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 & \mathbf{H} & 0 \\ \mathbf{H} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{c}_2 \\ \mathbf{c}_1 \\ 0 \end{bmatrix}, \\ \mathbf{C}_3 &= \begin{bmatrix} \mathbf{h}_3 \\ 0 \\ \mathbf{h}_1 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 & 0 & \mathbf{H} \\ 0 & 0 & 0 \\ \mathbf{H} & 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{c}_3 \\ 0 \\ \mathbf{c}_1 \end{bmatrix}. \end{aligned} \quad (63)$$

Where $\mathbf{c}_i = \mathbf{h}_i + \frac{1}{2}\mathbf{H}\mathbf{g}_i$. \mathbf{H} is the identity matrix.

And we know that $\mathbf{S}^u = E_{11}\mathbf{C}_1\mathbf{C}_1^T + E_{22}\mathbf{C}_2\mathbf{C}_2^T + E_{33}\mathbf{C}_3\mathbf{C}_3^T$,
 $E_{11}=E, E_{22}=E_{33}=G$. Rewrite \mathbf{S}^u in matrix form:

$$\mathbf{S}^u = E \begin{bmatrix} \mathbf{c}_1\mathbf{c}_1^T & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + G \begin{bmatrix} \mathbf{c}_2\mathbf{c}_2^T & \mathbf{c}_2\mathbf{c}_1^T & 0 \\ \mathbf{c}_1\mathbf{c}_2^T & \mathbf{c}_1\mathbf{c}_1^T & 0 \\ 0 & 0 & 0 \end{bmatrix} + G \begin{bmatrix} \mathbf{c}_3\mathbf{c}_3^T & 0 & \mathbf{c}_3\mathbf{c}_1^T \\ 0 & 0 & 0 \\ \mathbf{c}_1\mathbf{c}_3^T & 0 & \mathbf{c}_1\mathbf{c}_1^T \end{bmatrix} \quad (64)$$

Define $\mathbf{S}_1^U = \mathbf{c}_1\mathbf{c}_1^T, \mathbf{S}_2^U = \mathbf{c}_2\mathbf{c}_2^T, \mathbf{S}_3^U = \mathbf{c}_3\mathbf{c}_3^T, \mathbf{S}_4^U = \mathbf{c}_2\mathbf{c}_1^T$ and $\mathbf{S}_5^U = \mathbf{c}_3\mathbf{c}_1^T$, and because of symmetry, we can get:

$$\mathbf{S}^u = \begin{bmatrix} E\mathbf{S}_1^U + G(\mathbf{S}_2^U + \mathbf{S}_3^U) & G\mathbf{S}_4^U & G\mathbf{S}_5^U \\ G\mathbf{S}_4^{UT} & G\mathbf{S}_1^U & 0 \\ G\mathbf{S}_5^{UT} & 0 & G\mathbf{S}_1^U \end{bmatrix} \quad (65)$$

And according to Eq.53, \mathbf{S}^t is in the same form:

$$\mathbf{S}^t = \begin{bmatrix} E\mathbf{S}_1 + G(\mathbf{S}_2 + \mathbf{S}_3) & G\mathbf{S}_4 & G\mathbf{S}_5 \\ G\mathbf{S}_4^T & G\mathbf{S}_1 & 0 \\ G\mathbf{S}_5^T & 0 & G\mathbf{S}_1 \end{bmatrix} + \begin{bmatrix} s_{11}\mathbf{H} & s_{12}\mathbf{H} & s_{13}\mathbf{H} \\ s_{12}\mathbf{H} & 0 & 0 \\ s_{13}\mathbf{H} & 0 & 0 \end{bmatrix} \quad (66)$$

3.3 Transformation to Section Result

First, we need the variation of \mathbf{g}_i based on Eq.60

$$\begin{aligned} \delta\mathbf{g}_1 &= \frac{d\delta\mathbf{u}_0}{dX_1} + \delta\mathbf{R}^T\tilde{\boldsymbol{\zeta}}^T\mathbf{k} + \mathbf{R}^T\tilde{\boldsymbol{\zeta}}^T\delta\mathbf{k} \\ &= \frac{d\delta\mathbf{u}_0}{dX_1} + \mathbf{R}^T\delta\tilde{\boldsymbol{\theta}}\tilde{\boldsymbol{\zeta}}^T\mathbf{k} + \mathbf{R}^T\tilde{\boldsymbol{\zeta}}^T\delta\mathbf{k} \\ &= \frac{d\delta\mathbf{u}_0}{dX_1} + \mathbf{R}^T(\tilde{\mathbf{k}}^T\tilde{\boldsymbol{\zeta}} - \tilde{\boldsymbol{\zeta}}\tilde{\mathbf{k}}^T)\delta\boldsymbol{\theta} + \mathbf{R}^T\tilde{\boldsymbol{\zeta}}^T(\tilde{\mathbf{k}}\delta\boldsymbol{\theta} + \frac{d\delta\boldsymbol{\theta}}{dX_1}) \\ &= \frac{d\delta\mathbf{u}_0}{dX_1} + \mathbf{R}^T\tilde{\boldsymbol{\zeta}}^T\frac{d\delta\boldsymbol{\theta}}{dX_1} + \mathbf{R}^T\tilde{\mathbf{k}}^T\tilde{\boldsymbol{\zeta}}\delta\boldsymbol{\theta}, \\ \delta\mathbf{g}_2 &= \delta\mathbf{R}^T\mathbf{h}_2 = \mathbf{R}^T\tilde{\mathbf{h}}_2^T\delta\boldsymbol{\theta}, \\ \delta\mathbf{g}_3 &= \delta\mathbf{R}^T\mathbf{h}_3 = \mathbf{R}^T\tilde{\mathbf{h}}_3^T\delta\boldsymbol{\theta}. \end{aligned} \quad (67)$$

This can be write in matrix form:

$$\delta \mathbf{g} = \begin{bmatrix} \delta \mathbf{g}_1 \\ \delta \mathbf{g}_2 \\ \delta \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{I} & \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T & \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} \\ 0 & 0 & \mathbf{R}^T \tilde{\mathbf{h}}_2^T \\ 0 & 0 & \mathbf{R}^T \tilde{\mathbf{h}}_3^T \end{bmatrix} \begin{bmatrix} \frac{d\delta \mathbf{u}_0}{dX_1} \\ \frac{d\delta \boldsymbol{\theta}}{dX_1} \\ \delta \boldsymbol{\theta} \end{bmatrix} = \mathbf{W} \delta \mathbf{w} \quad (68)$$

Here \mathbf{w} is the section gradients and will connect to the final DOF later. And the second variation is:

$$\begin{aligned} \delta^2 \mathbf{g}_1 &= (\delta^2 \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \delta^2 \mathbf{k}) + 2\delta \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \delta \mathbf{k} \\ &= (\delta^2 \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \delta^2 \mathbf{k}) + 2\mathbf{R}^T \delta \tilde{\boldsymbol{\theta}} \tilde{\boldsymbol{\zeta}}^T \delta \mathbf{k} \\ &= (\delta^2 \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \delta^2 \mathbf{k}) + 2\mathbf{R}^T \delta \tilde{\boldsymbol{\theta}} \tilde{\boldsymbol{\zeta}}^T \left(\frac{d\delta \boldsymbol{\theta}}{dX_1} + \tilde{\mathbf{k}} \delta \boldsymbol{\theta} \right), \\ \delta^2 \mathbf{g}_2 &= \delta^2 \mathbf{R}^T \mathbf{h}_2, \\ \delta^2 \mathbf{g}_3 &= \delta^2 \mathbf{R}^T \mathbf{h}_3. \end{aligned} \quad (69)$$

Then the section tangent stiffness matrix is:

$$\mathbf{S}^{st} = \int_A \mathbf{W}^T \mathbf{S}^t \mathbf{W} dA = \int_A \mathbf{W}^T (\mathbf{S}_M + \mathbf{S}_G) \mathbf{W} dA \quad (70)$$

The first part:

$$\mathbf{S}_M^s = \int_A \mathbf{W}^T (E \mathbf{B}_1 \mathbf{B}_1^T + G \mathbf{B}_2 \mathbf{B}_2^T + G \mathbf{B}_3 \mathbf{B}_3^T) \mathbf{W} dA = \mathbf{T}_1 + \mathbf{T}_2 + \mathbf{T}_3 \quad (71)$$

And

$$\begin{aligned} \mathbf{B}_1 &= \begin{bmatrix} \mathbf{h}_1 \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} \mathbf{H} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{b}_1 \\ 0 \\ 0 \end{bmatrix}, \\ \mathbf{B}_2 &= \begin{bmatrix} \mathbf{h}_2 \\ \mathbf{h}_1 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 & \mathbf{H} & 0 \\ \mathbf{H} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{b}_2 \\ \mathbf{b}_1 \\ 0 \end{bmatrix}, \\ \mathbf{B}_3 &= \begin{bmatrix} \mathbf{h}_3 \\ 0 \\ \mathbf{h}_1 \end{bmatrix} + \begin{bmatrix} 0 & 0 & \mathbf{H} \\ 0 & 0 & 0 \\ \mathbf{H} & 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \begin{bmatrix} \mathbf{b}_3 \\ 0 \\ \mathbf{b}_1 \end{bmatrix}. \end{aligned} \quad (72)$$

Using Mathematica to simplify \mathbf{T}_1 , and because \mathbf{T} is symmetric, we get:

$$\begin{aligned} \mathbf{T}_1(1, 1) &= \int_A E (\mathbf{g}_1 + \mathbf{h}_1) (\mathbf{g}_1 + \mathbf{h}_1)^T dA, \\ \text{and, } \mathbf{g}_1 + \mathbf{h}_1 &= \frac{d\mathbf{u}_0}{dX_1} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k} + \mathbf{h}_1 = \mathbf{R}^T \boldsymbol{\phi} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k}, \\ \mathbf{T}_1(1, 1) &= \int_A E (\mathbf{R}^T \boldsymbol{\phi} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k}) (\boldsymbol{\phi}^T \mathbf{R} + \mathbf{k}^T \tilde{\boldsymbol{\zeta}} \mathbf{R}) dA \\ &= \int_A E (\mathbf{R}^T \boldsymbol{\phi} \boldsymbol{\phi}^T \mathbf{R} + \mathbf{R}^T \boldsymbol{\phi} \mathbf{k}^T \tilde{\boldsymbol{\zeta}} \mathbf{R} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k} \boldsymbol{\phi}^T \mathbf{R} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k} \mathbf{k}^T \tilde{\boldsymbol{\zeta}} \mathbf{R}) dA, \\ \text{and, } \int_A \tilde{\boldsymbol{\zeta}} dA &= 0, \int_A \tilde{\boldsymbol{\zeta}} \tilde{\boldsymbol{\zeta}}^T dA = \mathbf{I}_s, \\ \mathbf{T}_1(1, 1) &= E \mathbf{R}^T (A \boldsymbol{\phi} \boldsymbol{\phi}^T + \tilde{\mathbf{k}}^T \mathbf{I}_s \tilde{\mathbf{k}}) \mathbf{R}. \end{aligned} \quad (73)$$

and

$$\begin{aligned}
T_1(1, 2) &= \int_A E(\mathbf{g}_1 + \mathbf{h}_1)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\zeta}^T dA \\
&= \int_A E(\mathbf{R}^T \phi \phi^T \tilde{\zeta}^T + \mathbf{R}^T \phi \mathbf{k}^T \tilde{\zeta} \tilde{\zeta}^T + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\zeta}^T + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k} \mathbf{k}^T \tilde{\zeta} \tilde{\zeta}^T) dA \\
&= \int_A E(\mathbf{R}^T \phi \mathbf{k}^T \tilde{\zeta} \tilde{\zeta}^T + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\zeta}^T) dA \\
&= E \mathbf{R}^T \phi \mathbf{k}^T \mathbf{I}_r + E \mathbf{R}^T \tilde{\mathbf{k}} \mathbf{I}_s \tilde{\phi}
\end{aligned} \tag{74}$$

and

$$\begin{aligned}
T_1(1, 3) &= \int_A E(\mathbf{g}_1 + \mathbf{h}_1)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\zeta} dA \\
&= \int_A E(\mathbf{R}^T \phi \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta} + \mathbf{R}^T \phi \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta} + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta} + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k} \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta}) dA \\
&= \int_A E(\mathbf{R}^T \phi \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta} + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta}) dA \\
&= \int_A E(\mathbf{R}^T \phi \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta}) dA + E \mathbf{R}^T \tilde{\mathbf{k}} \mathbf{I}_s \tilde{\mathbf{k}}_e
\end{aligned} \tag{75}$$

where $\mathbf{k}_e = \tilde{\phi} \mathbf{k}$, and

$$\begin{aligned}
T_1(2, 2) &= \int_A E \tilde{\zeta} \mathbf{R}(\mathbf{g}_1 + \mathbf{h}_1)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\zeta}^T dA \\
&= \int_A E(\tilde{\zeta} \phi \phi^T \tilde{\zeta}^T + \tilde{\zeta} \phi \mathbf{k}^T \tilde{\zeta} \tilde{\zeta}^T + \tilde{\zeta} \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\zeta}^T + \tilde{\zeta} \tilde{\zeta}^T \mathbf{k} \mathbf{k}^T \tilde{\zeta} \tilde{\zeta}^T) dA \\
&= E(\tilde{\phi}^T \mathbf{I}_s \tilde{\phi} + \mathbf{I}_r \mathbf{k} \mathbf{k}^T \mathbf{I}_r) \simeq E \tilde{\phi}^T \mathbf{I}_s \tilde{\phi}
\end{aligned} \tag{76}$$

and

$$\begin{aligned}
T_1(2, 3) &= \int_A E \tilde{\zeta} \mathbf{R}(\mathbf{g}_1 + \mathbf{h}_1)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\zeta} dA \\
&= \int_A E(\tilde{\zeta} \phi \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta} + \tilde{\zeta} \phi \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta} + \tilde{\zeta} \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta} + \tilde{\zeta} \tilde{\zeta}^T \mathbf{k} \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta}) dA \\
&= \int_A E(\tilde{\zeta} \tilde{\zeta}^T \mathbf{k} \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta}) dA + \tilde{\phi}^T \mathbf{I}_s \tilde{\mathbf{k}}_e \simeq \tilde{\phi}^T \mathbf{I}_s \tilde{\mathbf{k}}_e
\end{aligned} \tag{77}$$

and

$$\begin{aligned}
T_1(3, 3) &= \int_A E \tilde{\zeta}^T \tilde{\mathbf{k}} \mathbf{R}(\mathbf{g}_1 + \mathbf{h}_1)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\zeta} dA \\
&= \int_A E(\tilde{\zeta}^T \tilde{\mathbf{k}} \phi \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta} + \tilde{\zeta}^T \tilde{\mathbf{k}} \phi \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta} + \tilde{\zeta}^T \tilde{\mathbf{k}} \tilde{\zeta}^T \mathbf{k} \phi^T \tilde{\mathbf{k}}^T \tilde{\zeta} + \tilde{\zeta}^T \tilde{\mathbf{k}} \tilde{\zeta}^T \mathbf{k} \mathbf{k}^T \tilde{\zeta} \tilde{\mathbf{k}}^T \tilde{\zeta}) dA \\
&\simeq \tilde{\mathbf{k}}_e^T \mathbf{I}_s \tilde{\mathbf{k}}_e
\end{aligned} \tag{78}$$

Then using Mathematica to simplify $\mathbf{T}_2 + \mathbf{T}_3 = \mathbf{T}_\tau$, we get:

$$\begin{aligned}
\mathbf{T}_\tau(1, 1) &= \int_A G(\mathbf{g}_2 + \mathbf{h}_2)(\mathbf{g}_2 + \mathbf{h}_2)^T + G(\mathbf{g}_3 + \mathbf{h}_3)(\mathbf{g}_3 + \mathbf{h}_3)^T, \\
&\text{and } \mathbf{g}_2 + \mathbf{h}_2 = \mathbf{R}^T \mathbf{h}_2, \mathbf{g}_3 + \mathbf{h}_3 = \mathbf{R}^T \mathbf{h}_3, \\
\mathbf{T}_\tau(1, 1) &= \int_A G \mathbf{R}^T (\mathbf{h}_2 \mathbf{h}_2^T + \mathbf{h}_3 \mathbf{h}_3^T) \mathbf{R} dA = G \mathbf{A} \mathbf{R}^T \mathbf{I}_c \mathbf{R}, \\
&\text{where } \mathbf{I}_c = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}.
\end{aligned} \tag{79}$$

and

$$\begin{aligned}
\mathbf{T}_\tau(1, 2) &= \int_A G(\mathbf{g}_2 + \mathbf{h}_2)(\mathbf{g}_2 + \mathbf{h}_2)^T \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T + G(\mathbf{g}_3 + \mathbf{h}_3)(\mathbf{g}_3 + \mathbf{h}_3)^T \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \\
&= 0
\end{aligned} \tag{80}$$

and

$$\begin{aligned}
\mathbf{T}_\tau(1, 3) &= \int_A G((\mathbf{g}_2 + \mathbf{h}_2)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\mathbf{h}}_2^T + (\mathbf{g}_2 + \mathbf{h}_2)(\mathbf{g}_2 + \mathbf{h}_2)^T \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} \\
&\quad + (\mathbf{g}_3 + \mathbf{h}_3)(\mathbf{g}_1 + \mathbf{h}_1)^T \mathbf{R}^T \tilde{\mathbf{h}}_2^T + (\mathbf{g}_3 + \mathbf{h}_3)(\mathbf{g}_2 + \mathbf{h}_3)^T \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}}) dA \\
&= \int_A G(\mathbf{R}^T \mathbf{h}_2 (\mathbf{R}^T \boldsymbol{\phi} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k})^T \mathbf{R}^T \tilde{\mathbf{h}}_2^T + \mathbf{R}^T \mathbf{h}_3 (\mathbf{R}^T \boldsymbol{\phi} + \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T \mathbf{k})^T \mathbf{R}^T \tilde{\mathbf{h}}_3^T) dA \\
&= \int_A G(\mathbf{R}^T \mathbf{h}_2 \mathbf{h}_2^T \tilde{\boldsymbol{\phi}} + \mathbf{R}^T \mathbf{h}_3 \mathbf{h}_3^T \tilde{\boldsymbol{\phi}}) dA \\
&= G \mathbf{A} \mathbf{R}^T \mathbf{I}_c \tilde{\boldsymbol{\phi}}
\end{aligned} \tag{81}$$

and

$$\begin{aligned}
\mathbf{T}_\tau(2, 2) &= \int_A G \tilde{\boldsymbol{\zeta}} \mathbf{R} ((\mathbf{g}_2 + \mathbf{h}_2)(\mathbf{g}_2 + \mathbf{h}_2)^T + (\mathbf{g}_3 + \mathbf{h}_3)(\mathbf{g}_3 + \mathbf{h}_3)^T) \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T dA \\
&= \int_A G(\tilde{\boldsymbol{\zeta}} \mathbf{h}_2 \mathbf{h}_2^T \tilde{\boldsymbol{\zeta}}^T + \tilde{\boldsymbol{\zeta}} \mathbf{h}_3 \mathbf{h}_3^T \tilde{\boldsymbol{\zeta}}^T) dA \\
&= G \mathbf{I}_p, \\
&\text{where } \mathbf{I}_p = \tilde{\mathbf{h}}_2^T \mathbf{I}_s \tilde{\mathbf{h}}_2 + \tilde{\mathbf{h}}_3^T \mathbf{I}_s \tilde{\mathbf{h}}_3
\end{aligned} \tag{82}$$

and the equations of $\mathbf{T}_\tau(2, 3)$ and $\mathbf{T}_\tau(3, 3)$ are too long, so here give the results only.

$$\begin{aligned}
\mathbf{T}_\tau(2, 3) &= G \mathbf{I}_p \tilde{\mathbf{k}}, \\
\mathbf{T}_\tau(3, 3) &= G(\mathbf{A} \tilde{\boldsymbol{\phi}}^T \mathbf{I}_c \tilde{\boldsymbol{\phi}} + \tilde{\mathbf{k}}^T \mathbf{I}_p \tilde{\mathbf{k}}).
\end{aligned} \tag{83}$$

The final result in matrix form is:

$$\begin{aligned} \mathbf{S}_M^s = E & \begin{bmatrix} \mathbf{R}^T(A\phi\phi^T + \tilde{\mathbf{k}}^T \mathbf{I}_s \tilde{\mathbf{k}}) \mathbf{R} & \mathbf{R}^T \tilde{\mathbf{k}} \mathbf{I}_s \tilde{\phi} & \mathbf{R}^T \tilde{\mathbf{k}} \mathbf{I}_s \tilde{\mathbf{k}}_e \\ & \tilde{\phi}^T \mathbf{I}_s \tilde{\phi} & \tilde{\phi}^T \mathbf{I}_s \tilde{\mathbf{k}}_e \\ & & \tilde{\mathbf{k}}_e^T \mathbf{I}_s \tilde{\mathbf{k}}_e \end{bmatrix} \\ & + G \begin{bmatrix} \mathbf{A} \mathbf{R}^T \mathbf{I}_c \mathbf{R} & 0 & \mathbf{A} \mathbf{R}^T \mathbf{I}_c \tilde{\phi} \\ & \mathbf{I}_p & \mathbf{I}_p \tilde{\mathbf{k}} \\ & & \mathbf{A} \tilde{\phi}^T \mathbf{I}_c \tilde{\phi} + \tilde{\mathbf{k}}^T \mathbf{I}_p \tilde{\mathbf{k}} \end{bmatrix} \end{aligned} \quad (84)$$

Next consider the high order term $(\delta^2 \mathbf{g})^T \Phi$ in Eq.45, this can be split in two parts:

$$(\delta^2 \mathbf{g})^T \Phi = \begin{bmatrix} \delta^2 \mathbf{g}_1^T & \delta^2 \mathbf{g}_2^T & \delta^2 \mathbf{g}_3^T \end{bmatrix} \begin{bmatrix} s_{11} \mathbf{b}_1 + s_{12} \mathbf{b}_2 + s_{13} \mathbf{b}_3 \\ s_{12} \mathbf{b}_1 \\ s_{13} \mathbf{b}_1 \end{bmatrix} \quad (85)$$

Before giving the detailed expressions, we define some section quantities:

$$\begin{aligned} P &= A \sigma_b, \\ \mathbf{Q} &= G \int_A (\mathbf{h}_2 \gamma_{12} + \mathbf{h}_3 \gamma_{13}) dA, \\ \mathbf{M}_\sigma &= E \int_A \zeta e_{11} dA = E \int_A \zeta \zeta^T \mathbf{k}_e dA = E \mathbf{I}_s \mathbf{k}_e, \\ \mathbf{M}_\tau &= G \int_A (\tilde{\zeta} \mathbf{h}_2 \gamma_{12} + \tilde{\zeta} \mathbf{h}_3 \gamma_{13}) dA \\ &= G \int_A (\tilde{\mathbf{h}}_2^T \zeta \zeta^T \tilde{\mathbf{h}}_2 \mathbf{k} + \tilde{\mathbf{h}}_3^T \zeta \zeta^T \tilde{\mathbf{h}}_3 \mathbf{k}) dA \\ &= G \mathbf{I}_p \mathbf{k}. \\ \mathbf{I}_p &= \tilde{\mathbf{h}}_2^T \mathbf{I}_s \tilde{\mathbf{h}}_2 + \tilde{\mathbf{h}}_3^T \mathbf{I}_s \tilde{\mathbf{h}}_3. \end{aligned} \quad (86)$$

Here $P, \mathbf{Q}, \mathbf{M}_\sigma, \mathbf{M}_\tau$ are axial force, transverse shear force, bending moments and torsional moments. Then we consider the first part of $(\delta^2 \mathbf{g})^T \Phi$:

$$\begin{aligned} L_1 &= \int_A \delta^2 \mathbf{g}_1^T s_{11} \mathbf{b}_{11} dA \\ &= \int_A s_{11} 2 \left(\frac{d\delta \boldsymbol{\theta}^T}{dX_1} + \delta \boldsymbol{\theta}^T \tilde{\mathbf{k}}^T \right) \tilde{\zeta}^T \delta \tilde{\boldsymbol{\theta}}^T \mathbf{R} (\mathbf{R}^T \phi + \mathbf{R}^T \tilde{\zeta}^T \mathbf{k}) dA \\ &= \int_A s_{11} 2 \left(\frac{d\delta \boldsymbol{\theta}^T}{dX_1} + \delta \boldsymbol{\theta}^T \tilde{\mathbf{k}}^T \right) \tilde{\zeta} \delta \tilde{\boldsymbol{\theta}} \phi dA \\ &= \int_A s_{11} 2 \left(\frac{d\delta \boldsymbol{\theta}^T}{dX_1} \tilde{\zeta} \delta \tilde{\phi} \delta \boldsymbol{\theta} + \delta \boldsymbol{\theta}^T \tilde{\mathbf{k}}^T \tilde{\zeta} \tilde{\phi} \delta \boldsymbol{\theta} \right) dA \\ &= 2 \left(\frac{d\delta \boldsymbol{\theta}^T}{dX_1} \widetilde{\mathbf{M}_\sigma} \delta \tilde{\phi} \delta \boldsymbol{\theta} + \delta \boldsymbol{\theta}^T \tilde{\mathbf{k}}^T \widetilde{\mathbf{M}_\sigma} \tilde{\phi} \delta \boldsymbol{\theta} \right) \end{aligned} \quad (87)$$

And because of symmetric, this can be written in matrix form:

$$\mathbf{L}_1 = \begin{bmatrix} \frac{d\delta \mathbf{u}_0^T}{dX_1} & \frac{d\delta \boldsymbol{\theta}^T}{dX_1} & \delta \boldsymbol{\theta}^T \end{bmatrix} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & \widetilde{\mathbf{M}_\sigma} \tilde{\phi} \\ 0 & (\widetilde{\mathbf{M}_\sigma} \tilde{\phi})^T & 2 \tilde{\mathbf{k}}^T \widetilde{\mathbf{M}_\sigma} \tilde{\phi} \end{bmatrix} \begin{bmatrix} \frac{d\delta \mathbf{u}_0}{dX_1} \\ \frac{d\delta \boldsymbol{\theta}}{dX_1} \\ \delta \boldsymbol{\theta} \end{bmatrix} \quad (88)$$

The second part:

$$\mathbf{L}_2 = \int_A s_{12}(\delta^2 \mathbf{g}_1^T \mathbf{b}_2 + \delta^2 \mathbf{g}_2^T \mathbf{b}_1) + s_{13}(\delta^2 \mathbf{g}_1^T \mathbf{b}_3 + \delta^2 \mathbf{g}_3^T \mathbf{b}_1) dA \quad (89)$$

This part has 4 sub-member, we call them \mathbf{L}_{21} - \mathbf{L}_{24} ,

$$\begin{aligned} \mathbf{L}_{21} &= \int_A s_{12} \delta^2 \mathbf{g}_1^T \mathbf{b}_2 dA \\ &= \int_A s_{12} 2 \left(\frac{d\delta \boldsymbol{\theta}}{dX_1} + \delta \boldsymbol{\theta}^T \tilde{\mathbf{k}}^T \right) \tilde{\boldsymbol{\zeta}} \delta \tilde{\boldsymbol{\theta}}^T \mathbf{h}_2 dA \\ &= \int_A s_{12} 2 \left(\frac{d\delta \boldsymbol{\theta}}{dX_1} + \delta \boldsymbol{\theta}^T \tilde{\mathbf{k}}^T \right) \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_2 \delta \boldsymbol{\theta} dA \end{aligned} \quad (90)$$

This can also be written in matrix form, and the central matrix is:

$$\begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & \int_A s_{12} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_2 dA \\ 0 & (\int_A s_{12} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_2 dA)^T & 2 \int_A s_{12} \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_2 dA \end{bmatrix} \quad (91)$$

And \mathbf{L}_{23} is the same form:

$$\begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & \int_A s_{13} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_3 dA \\ 0 & (\int_A s_{13} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_3 dA)^T & 2 \int_A s_{13} \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_3 dA \end{bmatrix} \quad (92)$$

and the \mathbf{L}_{22} and \mathbf{L}_{24} have high order terms, so we drop them here. Finally consider the section geometric stiffness matrix:

$$\mathbf{S}_{GP} = \int_A \mathbf{W}^T s_i \mathbf{H}_i^T \mathbf{W} dA \quad (93)$$

Using Mathematica to carry out and split the result in σ and τ part:

$$\begin{aligned} \mathbf{S}_{GP\sigma}(1,1) &= \int_A s_{11} \mathbf{I} dA = P \mathbf{I} \\ \mathbf{S}_{GP\sigma}(1,2) &= \int_A s_{11} \mathbf{R}^T \tilde{\boldsymbol{\zeta}}^T dA = \mathbf{R}^T \widetilde{\mathbf{M}}_{\sigma}^T \\ \mathbf{S}_{GP\sigma}(1,3) &= \int_A s_{11} \mathbf{R}^T \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} dA = \mathbf{R}^T \tilde{\mathbf{k}}^T \widetilde{\mathbf{M}}_{\sigma}^T \\ \mathbf{S}_{GP\sigma}(2,2) &= \mathbf{S}_{GP\sigma}(2,3) = \mathbf{S}_{GP\sigma}(3,3) = 0 \end{aligned} \quad (94)$$

And

$$\begin{aligned} \mathbf{S}_{GP\tau}(1,1) &= \mathbf{S}_{GP\tau}(1,2) = \mathbf{S}_{GP\tau}(2,2) = 0 \\ \mathbf{S}_{GP\tau}(1,3) &= \mathbf{R}^T \int_A s_{12} dA \tilde{\mathbf{h}}_2^T + \mathbf{R}^T \int_A s_{13} dA \tilde{\mathbf{h}}_3^T, \\ \mathbf{S}_{GP\tau}(2,3) &= \int_A s_{12} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_2^T dA + \int_A s_{13} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{h}}_3^T dA, \\ \mathbf{S}_{GP\tau}(3,3) &= \int_A s_{12} \tilde{\mathbf{h}}_2 \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} dA + \int_A s_{13} \tilde{\mathbf{h}}_3 \tilde{\mathbf{k}}^T \tilde{\boldsymbol{\zeta}} dA \\ &\quad + \int_A s_{12} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{k}} \tilde{\mathbf{h}}_2^T dA + \int_A s_{13} \tilde{\boldsymbol{\zeta}} \tilde{\mathbf{k}} \tilde{\mathbf{h}}_3^T dA \end{aligned} \quad (95)$$

Then combine \mathbf{L}_1 and $\mathbf{S}_{GP\sigma}, \mathbf{L}_2$ and $\mathbf{S}_{GP\tau}$, we get:

$$\mathbf{S}_{GC\sigma} = \mathbf{L}_1 + \mathbf{S}_{GP\sigma} = \begin{bmatrix} \mathbf{PI} & \mathbf{R}^T \widetilde{\mathbf{M}}_{\sigma}^T & \mathbf{R}^T \tilde{\mathbf{k}}^T \widetilde{\mathbf{M}}_{\sigma}^T \\ 0 & \widetilde{\mathbf{M}}_{\sigma} \tilde{\phi} & \\ \text{symm} & 2\tilde{\mathbf{k}}^T \widetilde{\mathbf{M}}_{\sigma} \tilde{\phi} & \end{bmatrix} \quad (96)$$

$$\mathbf{S}_{GC\tau} = \mathbf{L}_2 + \mathbf{S}_{GP\tau} = \begin{bmatrix} 0 & 0 & \mathbf{R}^T \tilde{\mathbf{Q}}^T \\ 0 & 0 & 0 \\ \text{symm} & 0 & \end{bmatrix} \quad (97)$$

And the final section tangent stiffness matrix is:

$$\mathbf{S}_C^t = \mathbf{S}_M^s + \mathbf{S}_{GC\sigma} + \mathbf{S}_{GC\tau} \quad (98)$$

3.4 Transformation to Final DOFs

First we define the linear shape function matrix \mathbf{N} :

$$\mathbf{N} = \frac{1}{2} \begin{bmatrix} 1-\epsilon & 0 & 0 & 1+\epsilon & 0 & 0 \\ 0 & 1-\epsilon & 0 & 0 & 1+\epsilon & 0 \\ 0 & 0 & 1-\epsilon & 0 & 0 & 1+\epsilon \end{bmatrix} \quad (99)$$

where $\epsilon = \frac{2X_1}{L} - 1$ is the local coordinate from -1 to 1, and we can apply the one point Gauss Integration Rule on it.

Then the section result is connected to the final DOFs:

$$\mathbf{w} = \begin{bmatrix} \frac{d\mathbf{N}}{dX_1} & 0 \\ 0 & \frac{d\mathbf{N}}{dX_1} \\ 0 & \mathbf{N} \end{bmatrix} \begin{bmatrix} \mathbf{d}_n \\ \mathbf{a}_n \end{bmatrix} = \mathbf{D}\mathbf{v} \quad (100)$$

and $\delta\mathbf{w} = \mathbf{D}\delta\mathbf{v}$. The final internal force and tangent stiffness matrix are:

$$\begin{aligned} \mathbf{f} &= \int_0^L \mathbf{D}^T \mathbf{f}_z dX = \frac{L}{2} \int_{-1}^1 \mathbf{D}^T \mathbf{f}_z d\epsilon, \\ \mathbf{K}^t &= \int_0^L \mathbf{D}^T \mathbf{K}_z \mathbf{D} dX = \frac{L}{2} \int_{-1}^1 \mathbf{D}^T \mathbf{K}_z \mathbf{D} d\epsilon \end{aligned} \quad (101)$$

We can then use the one point Gauss Integration Rule to numerically get the result.