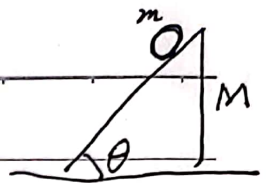


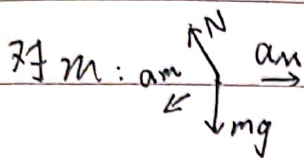
小球 (m)，斜木块 (M) 及地面三者均光滑 求 a_m, a_M (对地)

对 M 选地面作为参考系，对 m 选斜面做参考系 (非惯性系)



对 M :

$$\begin{array}{c} \uparrow N \\ \downarrow mg \end{array} \quad N \sin \theta = M a_M \quad (1)$$



(设向右为正)

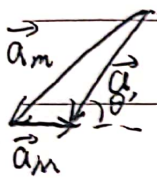
水平方向: $-N \sin \theta - m a_M = -m a_m \cos \theta \quad (2)$

惯性力

竖直方向 (设向上为正): $N \cos \theta - mg = -m a_m \sin \theta \quad (3)$

联立 (1)(2)(3) 得:

$$\begin{cases} a_M = \frac{m \sin \theta \cos \theta}{M + m \sin^2 \theta} g \\ a_m = \frac{(M + m) \sin \theta}{M + m \sin^2 \theta} g \quad (\star \text{ 此为对斜面的速度}) \end{cases}$$



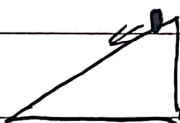
$\vec{a} = \vec{a}_m + \vec{a}_M$

$a_x = a_m \cos \theta - a_M = \frac{M \sin \theta \cos \theta}{M + m \sin^2 \theta} g$

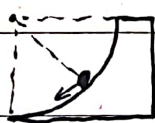
$a_y = a_m \sin \theta = \frac{(M + m) \sin^2 \theta}{M + m \sin^2 \theta} g$

$\tan \theta' = \frac{a_y}{a_x} = \frac{(M + m) \sin \theta}{M \cos \theta} = \frac{M + m}{M} \tan \theta > \tan \theta$

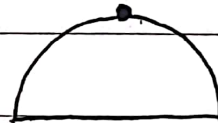
$\therefore \theta'$ 比 θ 陡.



最简单

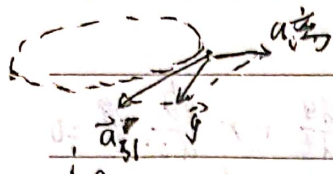


角度会变



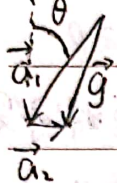
$\vec{v} = \vec{v}_{mB} + \vec{v}_B$

例：地球自转对地球作为惯性系的影响



$$a_{\text{离心}} = r\omega^2$$

$$a_{31} \gg a_{\text{离心}}, g \approx a_{31}$$



$$g = a_1 - a_2 \sin \theta \quad \text{推导:}$$

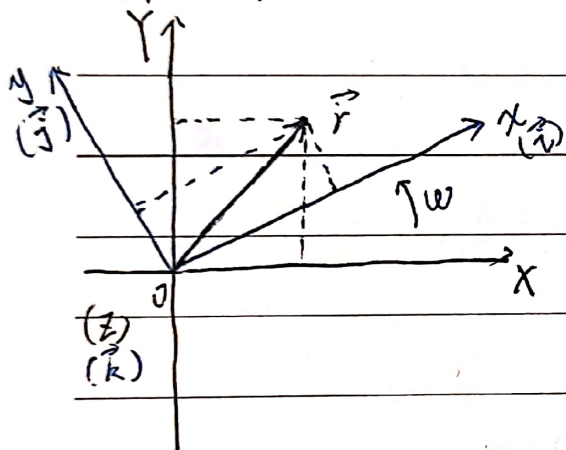
$$g = [a_1^2 + a_2^2 - 2a_1 a_2 \sin \theta]^{\frac{1}{2}}$$

$$\approx \sqrt{a_1^2 - 2a_1 a_2 \sin \theta}$$

$$= a_1 \sqrt{1 - 2 \frac{a_2}{a_1} \sin \theta}$$

$$\approx a_1 \left(1 - \frac{a_2}{a_1} \sin \theta\right) \quad \downarrow \text{Taylor展开}$$

科里奥利力在具体问题中的应用



$$\vec{r} = x\vec{i} + y\vec{j}$$

$$\vec{v} = \left(\frac{dx}{dt}\vec{i} + \frac{dy}{dt}\vec{j}\right) + \left(x\frac{d\theta}{dt}\vec{j} + y\frac{-d\theta}{dt}\vec{i}\right)$$

$$\text{对 } XOY = \vec{v}_r + \vec{\omega} \times \vec{r}$$

对xoy
(相对速度)

xoy对XOY
(牵连速度)

$$\begin{aligned} \vec{a} &= \frac{d^2x}{dt^2}\vec{i} + \frac{d^2y}{dt^2}\vec{j} + \frac{dx}{dt}\frac{d\theta}{dt}\vec{j} - \frac{dy}{dt}\frac{d\theta}{dt}\vec{i} \\ &+ \frac{dx}{dt}\frac{d\theta}{dt}\vec{j} - \frac{dy}{dt}\frac{d\theta}{dt}\vec{i} + x\frac{d^2\theta}{dt^2}\vec{j} - y\frac{d^2\theta}{dt^2}\vec{i} \\ &- x\left(\frac{d\theta}{dt}\right)^2\vec{i} - y\left(\frac{d\theta}{dt}\right)^2\vec{j} \end{aligned}$$

\vec{i}	\vec{j}	\vec{k}
0	0	ω
x	y	0

$$= \vec{a}_r + 2\vec{\omega} \times \vec{v}_r + \underbrace{\vec{\beta} \times \vec{r}}_{\vec{a}_\tau} - \underbrace{\omega^2 \vec{r}}_{\vec{a}_n} \quad (\beta \text{ 是角加速度})$$

$a_{\text{对xoy}}$

\vec{a}_τ
xoy对XOY

科里奥利加速度

$$F = \vec{F} + \vec{F}_i = m\vec{a}_r$$

$$\vec{F}_i = -m[\underbrace{\vec{\beta} \times \vec{r}}_{m\omega^2 r \text{ (离心力)}} - \underbrace{\omega^2 \vec{r}}_{m\omega^2 r \text{ (离心力)}} + 2\vec{\omega} \times \vec{v}_r]$$

$m\omega^2 r$
(离心力)

$+ 2m\vec{v}_r \times \vec{\omega}$