On-Shell Methods for Tree-Level Amplitudes in (De)Constructed Gauge Theory

Su Yingze

Supervisor: ***Prof. Tanabashi Masaharu***

Theoretical Elementary Particle Physics Laboratory
Nagoya University

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Contents

Why We Study Scattering Amplitudes?

1. **Bridge between theory and experiment**
   - Core prediction targets for high-energy collider experiments such as the LHC
   - Any new theory (SUSY, GUTs, extra dimensions) must predict observable cross sections

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1. **Bridge between theory and experiment**
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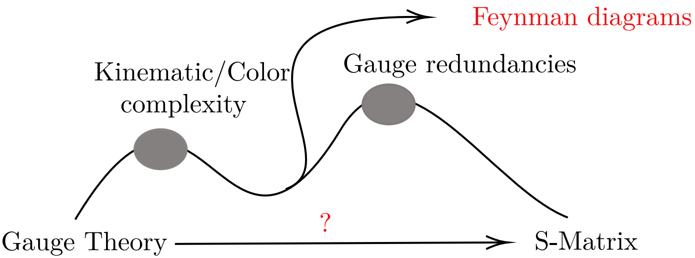
   - Any new theory (SUSY, GUTs, extra dimensions) must predict observable cross sections
2. **Reveal deep structures of quantum field theory**
   - Amplitudes exhibit hidden symmetries (e.g., dual conformal, Yangian) not visible in the Lagrangian
   - These symmetries suggest deeper theoretical frameworks, such as amplituhedra or AdS/CFT correspondence

Challenges we face befor



$n =$	4	5	6	7	8	9	10
	4	25	220	2485	34300	559405	10525900

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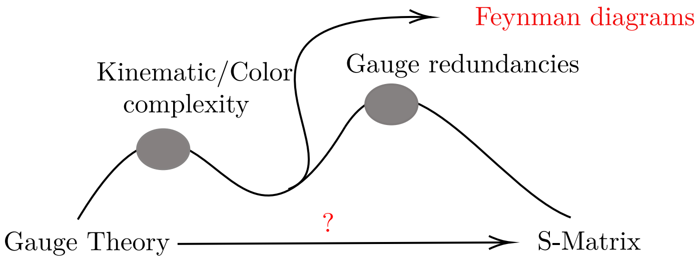

$n =$	4	5	6	7	8	9	10
	4	25	220	2485	34300	559405	10525900

The answer is On-shell method.

$$\text{Gauge Theory} \xrightarrow{\text{On-shell method}} \text{S-Matrix}$$

Contents

Conventional VS. Modern

Usually, when we compute the gluon amplitudes by using Feynman diagram, we will obtain something like

$$\begin{aligned}\mathcal{M}_s(p_1p_2 \to p_3p_4) &= -\frac{g_s^2}{s} f^{abe} f^{cde} \\ &\times \Big\{ -4\,\epsilon_1\cdot\epsilon_3^*\,\epsilon_2\cdot p_1\,\epsilon_4^*\cdot p_3 + 2\,\epsilon_1\cdot\epsilon_2\,\epsilon_3^*\cdot p_1\,\epsilon_4^*\cdot p_3 \\ &- 2\,\epsilon_1\cdot p_4\,\epsilon_2\cdot p_1\,\epsilon_3^*\cdot\epsilon_4^* + \epsilon_1\cdot\epsilon_2\,p_4\cdot p_1\,\epsilon_3^*\cdot\epsilon_4^* \\ &+ 4\,\epsilon_1\cdot\epsilon_4^*\,\epsilon_2\cdot p_1\,\epsilon_3^*\cdot p_4 - 2\,\epsilon_1\cdot\epsilon_2\,\epsilon_3^*\cdot p_4\,\epsilon_4^*\cdot p_1 \\ &- 2\,\epsilon_1\cdot p_2\,\epsilon_2\cdot p_3\,\epsilon_3^*\cdot\epsilon_4^* + \epsilon_1\cdot\epsilon_2\,\epsilon_3^*\cdot\epsilon_4^*\,p_2\cdot p_3 \\ &+ 4\,\epsilon_1\cdot p_2\,\epsilon_2\cdot\epsilon_3^*\,\epsilon_4^*\cdot p_3 - 2\,\epsilon_1\cdot\epsilon_2\,\epsilon_3^*\cdot p_2\,\epsilon_4^*\cdot p_3 \\ &+ 2\,\epsilon_1\cdot p_2\,\epsilon_2\cdot p_4\,\epsilon_3^*\cdot\epsilon_4^* - \epsilon_1\cdot\epsilon_2\,\epsilon_3^*\cdot\epsilon_4^*\,p_4\cdot p_2 \\ &+ 2\,\epsilon_1\cdot p_3\,\epsilon_2\cdot p_1\,\epsilon_3^*\cdot\epsilon_4^* - \epsilon_1\cdot\epsilon_2\,\epsilon_3^*\cdot\epsilon_4^*\,p_1\cdot p_3 \Big\}.\end{aligned}$$

If you consider 5point case, it will become worser:

★ There are nearly 10000 terms!

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Parke–Taylor Formula:

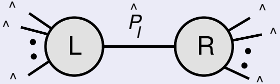
$$A_5(1^-, 2^-, 3^+, 4^+, 5^+) = \frac{\langle 12\rangle^4}{\langle 12\rangle\langle 23\rangle\langle 34\rangle\langle 45\rangle\langle 51\rangle}$$

★ It becomes quite simple and compact!

Fantasitic result from Cauchy Theorem

BCFW recursion relation

$$A_n = \sum_{\text{diagrams } I} \hat{A}_L(z_I)\frac{1}{P_I^2}\hat{A}_R(z_I) = \sum_{\text{diagrams } I}$$



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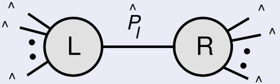
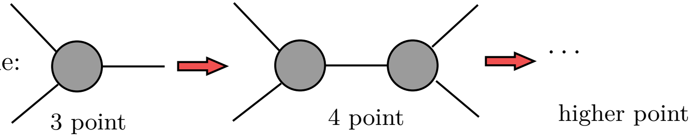


More understandable:



★ From lower point to higher point!!

Spinor-Helicity Formalism

Helicity

Helicity is defined as the projection of a particle's spin vector $\vec{S}$ onto the direction of its momentum $\vec{p}$:

$$h = \frac{\vec{S} \cdot \vec{p}}{|\vec{p}|}$$

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How can we catch the information of helicity?

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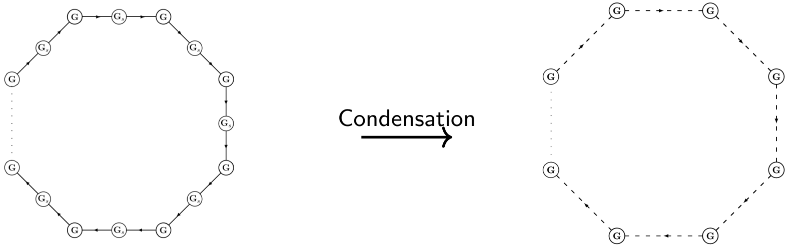
Massless Case:

- Momenta in spinor form:

$$p_\mu \sigma^\mu = p_{\alpha\dot{\alpha}} = p_\alpha \tilde{p}_{\dot{\alpha}} = |p\rangle[p|$$

Contents

Introduction of quiver gauge theory



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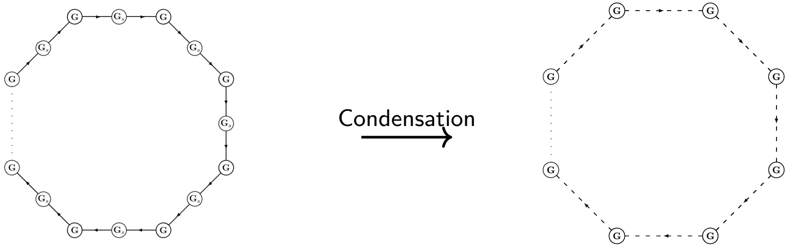


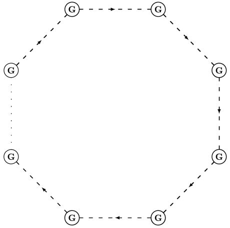
The Lagrangian can be written like

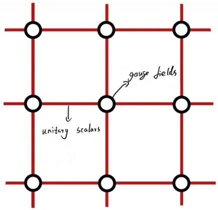
$$\mathcal{L} = -\sum_{i=1}^{N} \frac{1}{2}\mathrm{Tr}(F_i^2) + \sum_{i=1}^{N} \mathrm{Tr}\left[(D_\mu \Phi_i)^\dagger (D^\mu \Phi_i)\right],$$

here F_i refers to the ith gauge field strength. The scalar field Φ_i transforms under the bi-fundamental representation, and the covariant derivative equals to

$$D_\mu \Phi_i = \partial_\mu \Phi_i - i g_i A_{i\mu} \Phi_i + i g_{i+1} \Phi_i A_{i+1\mu}.$$

It has been proposed that this model actually discretized a five-dimension gauge theory with gauge group $SU(m)$, where only the fifth dimension are latticed. So it is an effective theory for 5d gauge theory.





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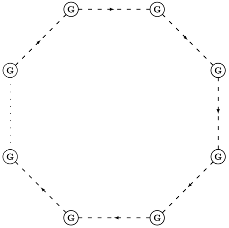

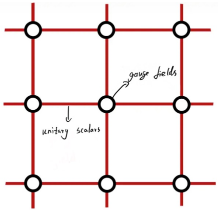

After higgsing the scalar field, we can obtain a spectrum

$$M_k^2 = 4g^2 f_s^2 \sin^2\left(\frac{\pi k}{N}\right)$$

This is precisely the Kaluza-Klein spectrum under S^2 compactification.

Scattering Amplitudes from BCFW

For simplicity, we start from the two-site gauge theory with gauge fields V_1, V_2 and scalar fields Φ, $\Phi^\dagger$.

$$\mathcal{L} = -\frac{1}{2}\mathrm{Tr}(F_1)^2 - \frac{1}{2}\mathrm{Tr}(F_2)^2 + \mathrm{Tr}[(D_\mu\Phi)^\dagger(D^\mu\Phi)],$$

We only foucus on the following amplitudes:

$$nV_1, \qquad nV_2, \qquad \Phi^\dagger nV_1\Phi, \qquad \Phi nV_2\Phi^\dagger, \qquad \Phi^\dagger\Phi\Phi^\dagger\Phi$$

here n can be any positive integer.

Basic building block – 3-point

From the previous section, we have known that there are only two kinds of 3 point amplitude

$$A[1,2,3^+] = \frac{[23][31]}{[12]}, \qquad A[1,2,3^-] = \frac{\langle 23\rangle\langle 31\rangle}{\langle 12\rangle}$$
$$A[3^+,4^+,5^-] = \frac{[34]^3}{[45][53]}, \qquad A[3^-,4^-,5^+] = \frac{\langle 34\rangle^3}{\langle 45\rangle\langle 53\rangle}$$

By using the 3 point building block, we can construct 4 point color-ordered amplitudes from BCFW recursion relation.

Gauge boson sector

- nV_1 or nV_2
  This part is completely the same as the pure gluon amplitude, so we can directly borrow the existing results.

$$\boxed{\text{Parke - Talyor Formula :} \quad A[\cdots, i^-, \cdots, j^-, \cdots] = \frac{\langle ij\rangle^4}{\langle 12\rangle\langle 23\rangle\cdots\langle n1\rangle}}$$

Notice that this formula only applies to MHV amplitudes, although the NMHV can be completely solved.

SQCD like sector

- $\Phi^\dagger V_1 V_1 \Phi$
  Here we compute the color-ordered amplitude $A[1,2,3^+,4^-]$. We choose $[2,3\rangle$ shift

$$|\hat{2}] = |2] - z|3], \qquad |\hat{2}\rangle = |2\rangle$$
$$|\hat{3}] = |3], \qquad |\hat{3}\rangle = |3\rangle + z|2\rangle$$

  The amplitudes can be computed

$$A[1,2,3^+,4^-] = (-1)\frac{\langle 14\rangle^2\langle 24\rangle^2}{\langle 12\rangle\langle 23\rangle\langle 34\rangle\langle 41\rangle}$$

- $\Phi^\dagger V_1 V_1 V_1 \Phi$

$$A[1,2,3^+,4^+,5^-] = \frac{\langle 15\rangle^2\langle 25\rangle^2}{\langle 12\rangle\langle 23\rangle\langle 34\rangle\langle 45\rangle\langle 51\rangle}$$

- $\Phi^\dagger(nV_1)\Phi$

$$A[1,2,\cdots,(n+2)^-] = (-1)^{n+1}\frac{\langle 1,n+2\rangle^2\langle 2,n+2\rangle^2}{\langle 12\rangle\langle 23\rangle\cdots\langle n+1,n+2\rangle\langle n+2,1\rangle}$$

$\star$ Bonus relation:

$$A[1,2,3^+,4^+]=0 \quad \Rightarrow \quad A[1,2,3^+,\cdots,n^+]=0$$

For the amplitude $\Phi(nV_2)\Phi^\dagger$, we can obtain nearly the same expression.

Pure 2-site amplitude

- $\Phi V_2 \Phi^\dagger V_1$

$$A[1,2,3_1^+,4_2^-] = \frac{\langle 14\rangle\langle 24\rangle}{\langle 13\rangle\langle 23\rangle}$$

- $\Phi V_2 \Phi^\dagger V_1 V_1$

$$A[1,2,3_1^+,4_1^+,5_2^-] = (-1)\frac{\langle 25\rangle^2\langle 15\rangle^2}{\langle 23\rangle\langle 34\rangle\langle 41\rangle\langle 25\rangle\langle 51\rangle}$$

- $\Phi V_2 V_2 \Phi^\dagger V_1 V_1$

$$A[1,2,3_1^+,4_1^+,5_2^+,6_2^-] = \frac{\langle 26\rangle^2\langle 16\rangle^2}{\langle 23\rangle\langle 34\rangle\langle 41\rangle\langle 25\rangle\langle 56\rangle\langle 61\rangle}$$

- Compact formula for general case

$$A = \frac{\langle 2a\rangle^2 \langle 1a\rangle^2}{\underbrace{\langle 2\star\rangle \cdots \langle \star 1\rangle}_{SU(N_1)} \underbrace{\langle 2*\rangle \cdots \langle * 1\rangle}_{SU(N_2)}}$$

Green: Particle with $-$ helicity

Blue: Particle belongs to the first gauge group

Red: Particle belongs to the second gauge group

$\star$: Order for gauge group 1

$*$: Order for gauge group 2

Contents

Summary

- Introduce the on-shell method, including BCFW recursion relation, spinor-helicity formalism, etc.
- Introduce a (de)constructed gauge theory model, which is an effective field theory for 5 dimension gauge theory.
- Much of the scattering amplitudes in this model can be recursively computed by BCFW, and some compact formulas are offered.

Thanks for your attention!

Brief explaination: We choose two momentum to be shifted oppositely

$$p_i \to \hat{p}_i(z) \equiv p_i - zk, \qquad p_j \to \hat{p}_j(z) \equiv p_j + zk$$

satisfying

$$k^2 = 0, \qquad p_i \cdot k = 0, \qquad p_j \cdot k = 0$$

We consider amplitude A_n in terms of shifted momentum $\hat{p}_i^{\mu}$ instead of original real momentum.

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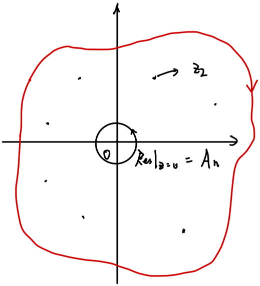
If we consider the meromorphic function $\frac{\hat{A}_n(z)}{z}$ in the complex plane. From Cauchy Theorem, we can ontain

$$A_n = -\sum_{z_I} \mathrm{Res}|_{z=z_I} \frac{\hat{A}_n(z)}{z} + B_n,$$

where B_n is the residue of the pole at $z = \infty$, called boundary term.



★ The most important point here is that

$$\boxed{\mathrm{Res}|_{z=0}\frac{\hat{A}_n(z)}{z} = \hat{A}_n(0) = A_n}$$

and

$$\mathrm{Res}|_{z=z_I} = \hat{A}_L(z_I)\frac{1}{P_I^2}\hat{A}_R(z_I)$$

Large z behavior

In the BCFW formula, the boundary term B_n affects a lot

$$A_n = -\sum_{z_I} \mathrm{Res}|_{z=z_I} \frac{\hat{A}_n(z)}{z} + B_n,$$

In most applications. one assumes or much better, proves $B_n = 0$. This is often justified by declaring a stronger condition

$$\hat{A}_n(z) \to 0 \quad \text{for} \quad z \to \infty$$

Here I show the large z behavior for gluon scattering

$[i \setminus j\rangle$	$+$	$-$
$+$	$1/z$	z^3
$-$	$1/z$	$1/z$

On-shell 3-point can be completely determined

For the complex momentum, we have

$$|1\rangle \propto |2\rangle \propto |3\rangle \qquad or \qquad |1] \propto |2] \propto |3]$$

$$\boxed{\begin{array}{ll} A_3^{h_1h_2h_3} = c\langle 12\rangle^{h_3-h_1-h_2}\langle 31\rangle^{h_2-h_1-h_3}\langle 23\rangle^{h_1-h_2-h_3} & h_1+h_2+h_3<0 \\ A_3^{h_1h_2h_3} = c'[12]^{h_1+h_2-h_3}[23]^{h_2+h_3-h_1}[31]^{h_3+h_1-h_2} & h_1+h_2+h_3>0 \end{array}}$$

$\star$ **All massless on-shell 3-point ampltides are completely determined by little group scaling!**

Example: 3-gluon amplitude

$$A_3(g_1^-, g_2^-, g_3^+) = g\frac{\langle 12\rangle^3}{\langle 23\rangle\langle 31\rangle}$$