

# Weak Solution: Convex Integration

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# 1 Introduction

First let's give an example to help define what is the weak solution.

**Example 1.1.**

$$\Delta u = f$$

If  $u$  don't have to be continuous take a test function:  $\forall \phi \in C_c^\infty(\mathbb{R}^d)$ ,  $\int_{\mathbb{C}} \phi \Delta u = \int_{\mathbb{C}} \phi f$ . If  $u$  is holomorphic  $u : \mathbb{C} \rightarrow \mathbb{C}$

$$\frac{d}{dt} u(z + t_\alpha) = \frac{\partial u}{\partial z} \alpha + \frac{\partial u}{\partial \bar{z}} \bar{\alpha} = \frac{\partial u}{\partial x} \operatorname{Re} \alpha + \frac{\partial u}{\partial y} \operatorname{Im} \alpha$$

$$\text{i.e. } du = \frac{\partial u}{\partial z} dz + \frac{\partial u}{\partial \bar{z}} d\bar{z}.$$

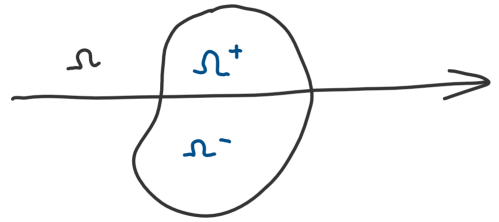
**Theorem 1.1.** If  $\frac{\partial u}{\partial \bar{z}} = 0$  or  $\Delta u = 0$  in the weak sense (against all test function). Then  $u \in C^\infty(\mathbb{R})$  and satisfies the equation in the classical sense.

Useful:  $f_k(z)$  holomorphic  $u(z) = \sum_{k=0}^{\infty} f_k(z)$ , the series is absolutely convergent.

$$\begin{aligned} \text{Fubini} &= - \int_{\mathbb{C}} \frac{\partial \phi}{\partial \bar{z}} u(z) \\ &= - \sum_{k=0}^{\infty} \int_{\mathbb{C}} \frac{\partial \phi}{\partial \bar{z}} f_k(z) \\ &= \sum_{k=0}^{\infty} \int_{\mathbb{C}} \phi \frac{\partial f_k}{\partial \bar{z}} = 0 \end{aligned}$$

(???)

**Theorem 1.2.** (Swartz reflection principle.) If  $f$  is holomorphic on  $\Omega \cap \{y > 0\}$  and  $\Omega \cap \{y < 0\}$ . If  $f$  is continu-



ous on  $\Omega$  on  $\Omega$  including  $\Omega \cap \{y = 0\}$ . Then  $f$  is holomorphic on  $\Omega$ .

In  $D' f = \lim_{\delta \rightarrow 0} f(H(y - \varepsilon) + H(\varepsilon - y))$ , here  $H$  is heaviside function.

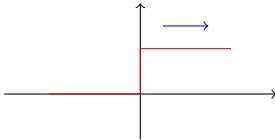
$$\frac{\partial f}{\partial \bar{z}} = \lim_{\varepsilon \rightarrow 0} \frac{\partial f}{\partial \bar{z}} + f \left( \frac{\partial y}{\partial \bar{z}} \delta(y - \varepsilon) - \frac{\partial y}{\partial \bar{z}} \delta(\varepsilon - y) \right)$$

Since  $f$  is continuous  $\lim_{\varepsilon \rightarrow 0} f(\delta(y - \varepsilon) - \delta(\varepsilon - y)) = 0$

$$\square u = 0 \quad \text{where } \square := -\partial_t^2 + \partial_1^2 + \dots + \partial_n^2 \quad (\text{Wave})$$

$$\forall \phi \in C_c^\infty(\mathbb{R}^{d-1}) \quad \int_{\mathbb{R}} \square \phi u = 0$$

(Wave) has a solution on  $\mathbb{R}^{1+1}$  given by  $u(t, x) = u(t - x)$ , a traveling wave:



**Example 1.2.**

$$u(t, x) = H(t - x) - H(t + x)$$

is the unique solution to (Wave) on  $\mathbb{R}^{1+d}$

The green lines are smooth approximation. After some time, it is still good enough to approximate the real world solution.

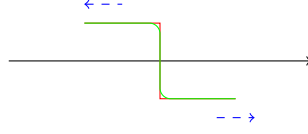


Figure 1:  $t = 0$

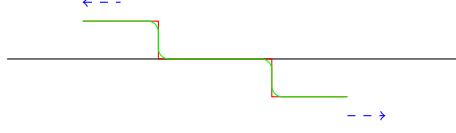


Figure 2:  $t = 1$

## 2 Fluid Mechanics

### 2.1 Incompressible Euler Equation

Define 2 velocity field:  $v : \mathbb{R} \times \mathbb{R}^d \rightarrow \mathbb{R}$   $p : \mathbb{R} \times \mathbb{R}^d \rightarrow \mathbb{R}$

$$\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p = 0$$

$$\nabla_j v^j = 0 \quad \text{divergence free}$$

This system obvious make sense for  $v \in L^2_{loc}$ . Let's recall the derivation of Euler equation.  $\forall \Omega$  with  $C^1 \partial\Omega$

$$\int_{\partial\Omega} v \cdot \vec{n} d\sigma = 0 \quad \forall t$$

meaning water coming in is exactly the same as water going out.

$$\frac{d}{dt} \left[ \begin{array}{c} \text{total momentum} \\ m \cdot v \end{array} \right] = [\text{Force on } \Omega] + \left[ \begin{array}{c} \text{Flux of} \\ \text{momentum} \end{array} \right] \Rightarrow \frac{d}{dt} \int_{\Omega} v^\ell dx = \int_{\partial\Omega} p \vec{n}^\ell dx \quad \forall t$$

These integral gives also the weak form of equation, let's say if  $p$  is good enough. If  $v, p \in C^1$ , use  $\int_{\partial\Omega} f \vec{n}_j d\sigma = - \int_{\Omega} \nabla_j f dx$

$$\frac{d}{dt} \int_{\Omega} v^\ell = - \left( \int_{\Omega} \nabla^\ell p + \nabla_j (v^j v^\ell) dx \right)$$

$$\int_{\Omega} (\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p) dx = 0 \quad \forall \Omega, \quad \forall t$$

Here comes a natural question: Are weak solution to the Euler equation physical meaningful?

Some physical properties are required. Take  $\Omega = \mathbb{R}^d$  and  $v \in L^2_{t,x}(I \times \mathbb{R}^d)$

$$\frac{d}{dt} \int_{\mathbb{R}^d} K_\ell v^\ell(t, x) dx = 0$$

If  $(1 + |x|)v \in L^1_{t,x}(I \times \mathbb{R}^d)$ , then linear angular momentum conserved?

Here  $\forall K^\ell$  s.t.  $\nabla_j K_\ell + \nabla_\ell K_j = 0$  on  $\mathbb{R}^d$

**Example 2.1.**  $K = e_{(i)}$  the basis vector,  $\int_{\mathbb{R}^d} K_\ell v^\ell dx = \int_{\mathbb{R}^d} v^{(i)} dx$

**Example 2.2.** Rotation  $K(a, b) = x^a e_b - x^b e_a$ ,  $1 \leq a < b \leq d$ .

$$\int_{\mathbb{R}^d} K_\ell (\partial_t v^\ell + \nabla : (v^j v^\ell) + \nabla^\ell p) dx = 0$$

$$\frac{d}{dt} \int_{\mathbb{R}^d} K_\ell v^\ell - \int_{\mathbb{R}^d} \nabla : K_\ell (v^j v^\ell) - \int_{\mathbb{R}^d} \nabla^\ell K_\ell p dx = 0$$

where  $\operatorname{div} K = \nabla^\ell K_\ell = \delta^{j\ell} \nabla_j K_\ell = \frac{1}{2} \delta^{j\ell} (\nabla_j K_\ell + \nabla_\ell K_j) = 0$  by assumption.

$$\partial v^\ell + \nabla : (v^j v^\ell) + \nabla^\ell p) dx = 0$$

Test against a space cut-off function  $K^\ell(B) := q(t) \varphi(\frac{|x|}{B})(x^b e_m - x^a e_b)$ . Here  $K$  is rotationally symmetric, so it is divergence-free.

$$- \int_{\mathbb{R}^+} \eta'(t) \left[ \int_{\mathbb{R}^d} K_\ell^{(\beta)} v^\ell dx \right] dt - \int_{\mathbb{R}} \eta \int_{\mathbb{R}^d} \nabla_j K_\ell^{(\beta)} v^j v^\ell - \int_{\mathbb{R}^+} \eta(t) \int_{\mathbb{R}^d} \nabla^\ell K_\ell^{(\beta)} p dx dt = 0$$

Here, due to divergence-free, like what we did previously,  $\int_{\mathbb{R}^+} \eta(t) \int_{\mathbb{R}^d} \nabla^\ell K_\ell^{(\beta)} p dx dt = 0$ . i.e.

$$- \int_{\mathbb{R}^+} \eta'(t) \left[ \int_{\mathbb{R}^d} K_\ell^\beta v^\ell dx \right] dt - \frac{1}{2} \int_{\mathbb{R}^+} \int_{\mathbb{R}^d} \left( \nabla_j K_\ell^{(\beta)} + \nabla_\ell K_j^{(\beta)} \right) v^j v^\ell dx dt - \frac{1}{2} \int_{\mathbb{R}} \eta(t) \int_{\mathbb{R}^d} \nabla_j \varphi\left(\frac{|x|}{B}\right) K_\ell v^j v^\ell dx dt = 0$$

The 1st term is dominated by  $|x| \cdot v \in L_{t,x}^1$  by assumption.  $\frac{1}{|x|} \cdot v$  dominated the derivative and integrant.

## 2.2 Conservation of Energy

If  $(1 + |x|)v \in L_{t,x}^1(I \times \mathbb{R}^d)$ ,  $v \in L_{t,x}^2(I \times \mathbb{R}^d)$ , then  $\forall K^\ell$ ,  $\nabla_j K_l + \nabla_l K_j = 0$ , then we have the conservation of angular momentum:

$$\frac{d}{dt} \int_{\mathbb{R}^d} K_\ell v^\ell(t, x) dx = 0,$$

where  $K \in \operatorname{span}\{x^a e_b = x^b e_a : e_i, 1 \leq i \leq d, 1 \leq a < b \leq d\}$ . Here decay assumption is needed but not the regularity assumption. If  $f \in \mathcal{D}'(\mathbb{R})$ ,  $\frac{df}{dt} = 0 \Rightarrow f = c$  limit of constant.

$$\delta_j^\ell = \nabla_j w^{j\ell} \quad w^{j\ell} = -w^{lj} \quad \text{antisymmetric}$$

Approximate by  $\nabla(\phi(\frac{|x|}{B}) w^{j\ell})$

$$w^{j\ell} = x^2 (\delta_2^j \delta_1^\ell - \delta_1^j \delta_2^\ell)$$

$$\begin{aligned} & \nabla_j (x^2 (\delta_2^j \delta_1^\ell - \delta_1^j \delta_2^\ell)) \\ &= \delta_j^2 (\delta_2^j \delta_1^\ell - \delta_1^j \delta_2^\ell) = \delta_1^\ell \end{aligned}$$

If  $w^{il} = -w^{lj}$

$$\nabla_j \nabla_j w^{j\ell} = -\nabla_\ell \nabla_j w^{lj} = -\nabla_j \nabla_\ell w^{lj} = -\nabla_\ell \nabla_j w^{j\ell}$$

Conservation of energy means that  $\frac{d}{dt} \int_{\mathbb{R}^d} \frac{|v|^2}{2}(t, x) dx = 0$ . Note that energy is nonlinear.

$$\partial_t \left( \frac{|v|^2}{2} \right) + \nabla_j \left( \left( \frac{|v|^2}{2} + p \right) v^j \right) = 0$$

If  $v \in C^1 \cap L_{t,x}^2 \cap L_{t,x}^3(I \times \mathbb{R}^d)$  both local and global conservation of energy hold. Note that here  $B$  could be  $\infty$ . Multiply the local energy by  $\eta(t) \varphi(\frac{|x|}{B})$

$$\int \eta \frac{d}{dt} \int \varphi\left(\frac{|x|}{B}\right) \frac{|v|^2}{2}(t, x) dx dt \quad - \int \eta(t) \int \nabla_j \left[ \varphi\left(\frac{|x|}{B}\right) \right] \left( \frac{|v|^2}{2} + p \right) v^j dx dt$$

(1)
(2)

$$(1) = - \int \underbrace{\eta'(t)} \int \varphi\left(\frac{|x|}{B}\right) \frac{|v|^2}{2}(t, x) dx dt \quad \text{Integral by parts} \quad \text{(Local)}$$

$$\text{dominated by } \frac{|x|^2}{2} |\eta'| \in L_{t,x}^1$$

term (2) converge to 0 pointwisely when  $B \rightarrow \infty$  and dominated by  $|\eta t| \left( \frac{|v|^3}{2} + |p||v| \right)$ .  
Let's recall Euler equation.

$$\begin{cases} \partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p &= 0, \\ \nabla_j v^j &= 0 \end{cases} \quad \text{(Euler)}$$

Take divergence over (Euler),  $\Rightarrow$

$$\nabla_j \nabla_\ell (v^j v^\ell) + \nabla_\ell \nabla^\ell p = 0$$

i.e.

$$\begin{aligned} \Delta p &= -\nabla_\ell \nabla_j (v^j v^\ell) \\ p &= \underbrace{(-\Delta)^{-1} \nabla_\ell \nabla_j (v^j v^\ell)}_{\text{zero order operator}} \underbrace{(v^j v^\ell)}_{\in L_{t,x}^{3/2}} \end{aligned}$$

Thus naturally  $p \in L_x^{3/2}$  a.e.  $t \in \mathbb{R}^+$

$$\|p\|_{L_x^{3/2}(L_t^{3/2})} = \|p\|_{L_{t,x}^{3/2}} < \infty$$

$$v_\ell (\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p) = 0$$

$$\nabla_j v^j = 0$$

Thus

$$\Rightarrow \partial_t \left( \frac{|v|^2}{2} \right) + v_\ell v^j \nabla_j v^\ell + v_\ell \nabla^\ell p = 0$$

$$\Rightarrow \partial_t \left( \frac{|v|^2}{2} \right) + v^j \nabla_j \left( \frac{|v|^2}{2} + v_j v^j p \right) = 0$$

$$\nabla_j v^j = 0$$

$$\partial_t \left( \frac{|v|^2}{2} + \nabla_j \left( \left( \frac{|v|^2}{2} + p \right) v^j \right) \right) = 0$$

Using  $\nabla_j v^j = 0$  and product rule, conservation of energy is proved for sufficient regular solutions. But how sufficient do we need?

In turbulence situation (Navier-Stokes equations) with  $\nu \ll 1$

$$v_\ell (\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p) = \nu v_\ell \Delta v^\ell$$

$$\frac{d}{dt} \int \frac{|v|^2}{2}(t, x) dx = -\nu \int |\nabla v|^2 dx = \nu \int v_\ell \nabla_i \nabla^i v^\ell$$

Taking a formal limit,  $\exists$  incompressible Euler flows with

$$\frac{d}{dt} \int \frac{|v|^2}{2}(t, x) dx < -\varepsilon < 0$$

**Theorem 2.1.** *Onsager's Conjecture*

(+) If  $\alpha > 1/3$  and  $(v(t, x + \Delta x) - v(t, x)) \leq c|\Delta x|$  where  $x \in \mathbb{T}^3 (v \in L_t^\infty C_x^\alpha)$ , then the energy conserved.

(-) (K41) If  $\alpha \leq 1/3 \exists$  incompressible Euler flows with  $v \in L_t^\infty L_x^\alpha$  s.t.  $\int_{\mathbb{T}^d} \frac{|v|^2}{2}(t, x) dx$  is not constant.

Now we follow [2] and discuss the (+) part first.

$$\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p = 0$$

In order to get into Onsager's explanation of how this might be possible, we expand the velocity  $v$  in Fourier series,

$$v(x, t) = \sum_{k \in \mathbb{Z}^3} a_k(t) e^{ik \cdot x}.$$

Obviously  $a_{-k} = \overline{a_k}$ , because  $v$  is real-valued. Moreover the divergence-free constraint translates into the identity  $k \cdot a_k = 0$ . We then rewrite the remaining equations of (2.2) as an infinite-dimensional system of ODEs for the  $a_k$ :

$$\frac{da_k}{dt} = i \sum_\ell a_{k-\ell} \cdot \ell \left[ -a_\ell + \frac{(a_\ell \cdot k) k}{|k|^2} \right] - \nu |k|^2 a_k \quad (1)$$

The total kinetic energy is (up to a constant factors)  $\sum_k |a_k|^2$ .

(Don't understand) Energy starts at low wave numbers and moves to higher wave numbers in finite number.

$\sum_{\frac{\lambda}{2} \leq |k| \leq 2\lambda} |a_k|^2 \sim \lambda^{-2/3}$  matches (K41), corresponding to exactly 1/3 regularity for solutions.

Low frequency energy will goes to all frequency and when it goes to infinity, it will disappear.

(K 41)  $E \lim_{v \rightarrow 0} \left\langle v \int |\nabla v|^2 dx \right\rangle$  and  $v$  determine all statistic properties of turbulent flows.

$$\langle |v(x + \Delta x) - v(x)|^p \rangle^{1/p} \sim \varepsilon^p |\Delta x|^{1/3}$$

Try to find  $|\Delta x| \leq L \sim \varepsilon^a v^b$ .

Now (+) is solved by [4] and [1] with the goal  $L_t^3 B_{3,C(N)}^{1/3}, L_t^3 B_{3,\infty}^{1/3+\varepsilon}$ .

(-) is solved ( $d \geq 3$ ) with  $\alpha = \frac{1}{3}$ , using convex integration by Phillip Isett [7].

Convex integration originated from the Nash–Kuiper Paradox(50's) for  $C^1$  isometric embedding. Connection to Euler equation discovered by Camillo De Lellis and László Székelyhidi (08,12). First result towards Onsager conjecture is in [8]. And  $\alpha < \frac{1}{5}$  by [10]. The non-uniqueness example was first given by [12] and then Shnirelman give a different proof in [13].

## 2.3 Another way of proving (+)

(+) (Eyink, Constantin, E, Titi 94')  $L^3(B_{3,\infty}^\alpha)$

$$\|v\|_{C^\alpha} = \sup_{h \neq 0} \frac{\|v(x+h) - v(x)\|_{L^\infty}}{|h|^\alpha}$$

$$\|v\|_{B_{3,\infty}^\alpha} = \sup_{h \neq 0} \frac{\|v(x+h) - v(x)\|_{L^3}}{|h|^\alpha}$$

**Lemma 2.2.** *Commutator Estimate*

$$R_\varepsilon^{j\ell} = \eta_\varepsilon * (v^j v^\ell) - (v_\varepsilon^j v_\varepsilon^\ell)$$

$$\|R_\varepsilon\|_{L^{3/2}} \lesssim \varepsilon^{2\alpha} \|v\|_{B_{3,\infty}^\alpha}^2$$

Let's think  $R_\varepsilon^{j\ell}$  as an expectation with the idea:

$$R = \mathbb{E}[v^2] - (\mathbb{E}[v])^2 = \mathbb{E}[(v - \mathbb{E}(v))^2],$$

which is quadratic.

$$R_\varepsilon^{j\ell} = \int v^i(x-h) v^\ell(x-h) \eta_\varepsilon(h) dh - \int v^j(x-h_1) \eta_\varepsilon(h_1) dh_1 \int v^\ell(x-h_2) \eta_\varepsilon(h_2) dh_2$$

$$\text{Using } \int \eta_\varepsilon(h) dh = 1$$

$$= \int (v^j(x-h) - v_\varepsilon^j(x)) (v^\ell(x-h) - v_\varepsilon^\ell(x)) \eta_\varepsilon(h) dh$$

By Lemmas in [1], we decompose above equation into  $\sum_{i=1}^4 R_{\varepsilon i}^{j\ell}$ , where

$$R_{\varepsilon 1} = \int (v^j(x-h) - v_\varepsilon^j(x-h)) (v^\ell(x-h) - v_\varepsilon^\ell(x-h)) \eta_\varepsilon(h) dh$$

$$R_{\varepsilon 2} = \int (v_\varepsilon^j(x-h) - v_\varepsilon^j(x)) (v^\ell(x-h) - v_\varepsilon^\ell(x-h)) \eta_\varepsilon(h) dh$$

$$R_{\varepsilon 3} = \int (v^j(x-h) - v_\varepsilon^j(x-h)) (v_\varepsilon^\ell(x-h) - v_\varepsilon^\ell(x)) \eta_\varepsilon(h) dh$$

$$R_{\varepsilon 4} = \int (v_\varepsilon^j(x-h) - v_\varepsilon^j(x)) (v_\varepsilon^\ell(x-h) - v_\varepsilon^\ell(x)) \eta_\varepsilon(h) dh$$

For example,

$$R_{\varepsilon 2} = \int_{\mathbb{R}^d} \int_0^1 \frac{d}{d\sigma} v_\varepsilon^j(x - \sigma h) d\sigma (v^\ell(x-h) - v_\varepsilon^\ell(x)) \eta_\varepsilon(h) dh$$

$$= \int_{\mathbb{R}^d} \int_0^1 d\sigma \nabla_i v_\varepsilon^j(x - \sigma h) h^i (v^\ell(x-h) - v_\varepsilon^\ell(x)) \eta_\varepsilon(h) dh$$

$$\|R_{\varepsilon 2}^j\| \leq_{\mathbb{R}^d} \int_0^1 \|\nabla v_\varepsilon\|_{L^3} |h| \|v(\cdot - h) - v(\cdot)\|_{L^3} |\eta_\varepsilon(h)| dh$$

Modify the equation with modifier  $\eta_\varepsilon$ :

$$\begin{aligned} \eta_\varepsilon * (\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p) &= 0 \\ \partial_t v_\varepsilon^\ell + \nabla_j (v_\varepsilon^j v_\varepsilon^\ell) + \nabla^\ell p_\varepsilon &= -\nabla_j R_\varepsilon^{j\ell} \end{aligned}$$

(Thus we need smoothness in time)  $\times v_\varepsilon$  then integral by parts:

$$\partial_t \left( \frac{|v_\varepsilon|^2}{2} \right) + v_{\varepsilon l} \nabla_j (v_\varepsilon^j v_\varepsilon^\ell) + v_{\varepsilon l} \nabla^\ell p_\varepsilon = -v_{\varepsilon l} \nabla_j R_\varepsilon^{j\ell} = \int_{\mathbb{R}^d} \nabla_j \left| \frac{v_\varepsilon^2}{2} v_\varepsilon^j \right| \rightarrow 0$$

with assumption.

$$\frac{d}{dt} \int_{\mathbb{R}^d} \frac{|v_\varepsilon|^2}{2}(t, x) dx + \int_{\mathbb{R}^d} v_\varepsilon^j \nabla_j v_\varepsilon^\ell v_{\varepsilon l} + \nabla^\ell v_{\varepsilon l} p_\varepsilon = \int_{\mathbb{R}^d} \nabla v_{\varepsilon l} R_\varepsilon^{j\ell}$$

$\nabla^\ell v_{\varepsilon l} p_\varepsilon = 0$  for divergence-free.

LHS converges to  $\frac{d}{dt} \int \frac{|v|^2}{2}(t, x) dx$  in  $\mathcal{D}'(\mathbb{R})$  since  $v_\varepsilon \rightarrow v$  in  $L_{t,x}^2$ .

$$\begin{aligned} \left\| \frac{d}{dt} \int \frac{|v|^2}{2}(t, x) dx \right\|_{L_t^1} &\leq \limsup_{\varepsilon \rightarrow 0} \int \int |\nabla_j v(t, x) R_\varepsilon^{j\ell}| dx dt \\ &\leq \limsup_{\varepsilon \rightarrow 0} \int \|\nabla v_\varepsilon(t, \cdot)\|_{L_x^3} \|R_\varepsilon\|_{L^{3/2}} dt \\ &\leq \limsup_{\varepsilon \rightarrow 0} \int \varepsilon^{-1+\alpha} \|v(t)\|_{B_{3,\infty}^\alpha} \|v(t, \cdot)\|_{B_{3,\infty}^\alpha}^2 dt \\ &< \limsup_{\varepsilon \rightarrow 0} \varepsilon^{-1+3\alpha} \int \|v(t, \cdot)\|_{B_{3,\infty}^\alpha}^3 dt \rightarrow 0 \quad \text{with } \alpha > \frac{1}{3} \end{aligned}$$

If  $\alpha = \frac{1}{3}$  and  $v$  bounded in  $L_t^1(I)$  for some finite time period.

$$\frac{d}{dt} \int \frac{|v|^2}{2}(t, x) dx = \lim_{\varepsilon \rightarrow 0} \frac{d}{dt} \int \frac{|v_\varepsilon|^2}{2}(t, x) dx$$

$v\phi \in C_c^\infty(I)$

$$\frac{d}{dt} < \int \frac{|v|^2}{2}(t, x) dx, \phi > \leq \|\phi\|_{L^\infty(I)}$$

LHS is of finite measure.  $e(t) = \int \frac{|v|^2}{2}(t, x)$  is of bounded variation. IN fact  $\frac{d}{dt} e(t)$  is finite.

(???) If  $v \in L^r B_{3,\infty}^{1/3}$ , consider  $\left\| \frac{d}{dt} e(t) \right\|_{L_t^{r'/t}}$  using duality.  $u \in L_t^\infty B_{3,\infty}^{1/3}$  uniformly  $\left\| \frac{d}{dt} e(t) \right\|_{L^{i_nfty} t} \leq C$  and also  $\frac{d}{dt} e(t) \leq -\varepsilon < 0$  is stable under perturbation. If not, the dissipation  $\int_{\mathbb{R}^d} \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell} dx$  can be really big.

*Remark.* The singular support of a generalized function  $u$  is the complement of the largest open set on which  $u$  is smooth. Roughly speaking, it is the closed set where the distribution does not correspond to a smooth function.

## 2.4 Local energy conservation

$$\partial_t v_\varepsilon^\ell + \nabla_l (v_\varepsilon^\ell v_\varepsilon^\ell) + \nabla^\ell p_\varepsilon = -\nabla_j R_\varepsilon^{j\ell}$$

where  $R_\varepsilon^{j\ell} = \eta_\varepsilon * (v^j v^\ell) - v_\varepsilon^j v_\varepsilon^\ell$

$$\begin{aligned} \|R_\varepsilon(t, \cdot)\|_{L_t^{3/2}} &\leq \varepsilon^{2\alpha} \|v(t)\|_{B_{3,\infty}^\alpha}^2 \\ \frac{1}{2} \int \frac{|v_\varepsilon|^2}{2}(t, x) dx &= \lim_{\varepsilon \rightarrow 0} \int \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell} dx \end{aligned}$$

Here to clarify the space:

$$B_{3,c(N)}^{1/3} = (\overline{C^\infty})^{B_{3,\infty}^{1/3}} = B_{3,\infty}^{1/3} \cap \left\{ \lim_{h \rightarrow 0} \frac{|v(x+h) - v(x)|}{|h|^{1/3}} = 0 \right\}$$

The "Holder Continuity" is the reason for smooth approximation. Define

$$c^{1/3} = (\overline{C^\infty})^{C^{1/3}}$$

Note that, here  $c^{1/3}$  is not dense in  $C^{1/3}$ . Let  $\varphi(x)$  be a smooth cut off function, then,  $|x|^{1/3} \in C^{1/3} \setminus c^{1/3}$ , but  $\varphi(x)|x|^{1/3} \notin C^{1/3} \setminus c^{1/3}$

**Lemma 2.3.**  $\|\nabla v_\varepsilon\|_{L^3} = o(\varepsilon^{-1+\alpha})$  if  $v \in B_{3,c(N)}^\alpha$

*Proof.* Claim:  $\varepsilon^{1-\alpha}\nabla(\eta_\varepsilon * \cdot) : B_{3,\infty}^\alpha \rightarrow L^3$  is uniformly bounded.

$$\|\nabla v_\varepsilon\|_{L_x^3} \lesssim \varepsilon^{-1+\alpha} \|v\|_{B_{3,\infty}^\alpha}$$

Let  $\delta > 0$  be given, choose  $\tilde{v} \in C^\infty$  s.t.  $\|v - \tilde{v}\|_{B_{3,\infty}^\alpha} < \frac{\delta}{2C_2}$ .

$$\begin{aligned} \|\varepsilon^{1-\alpha}\nabla\eta_\varepsilon * v\|_{L_x^3} &\leq \|\varepsilon^{1-\alpha}\nabla\eta_\varepsilon * (v - \tilde{v})\|_{L_x^3} + \|\varepsilon^{1-\alpha}\nabla\eta_\varepsilon * \tilde{v}\|_{L_x^3} \\ &\leq \frac{\delta}{2} + \varepsilon^{1-\alpha}\|\nabla\eta_\varepsilon * \tilde{v}\|_{L_x^3} \\ &\leq \frac{\delta}{2} + \varepsilon^{1-\alpha}\tilde{C} \quad \text{for } \varepsilon^{1-\alpha} < \frac{\delta}{2\tilde{C}} \text{ and } \|\varepsilon^{1-\alpha}\nabla\eta_\varepsilon * v\|_{L_x^3} < \delta \end{aligned}$$

$$\begin{aligned} \int_I \frac{d}{dt} \int \frac{|v_\varepsilon|^2}{2}(t, x) dx dt &\leq \limsup_{\varepsilon \rightarrow 0} \int_I \int |\nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell}| dx dt \\ &\leq \limsup_{\varepsilon \rightarrow 0} \int_I \|\nabla v_\varepsilon(t)\|_{L_x^3} \varepsilon^{2\alpha} \|v\|_{B_{3,\infty}^\alpha}^2 dt \end{aligned}$$

For a.e.  $t$ ,  $v \in B_{3,C(N)}^{1/3}$ , the integrand is bounded by  $o(\varepsilon^{-1+2/3})\varepsilon^{3/2} = o(1)$ . Thus above integral is dominated by:

$$\int_I \varepsilon^{-1 \times 1/3 + 2/3} \|v(t)\|_{B_{3,\infty}^{1/3}}^3 dt \leq \int_I \|v(t)\|_{B_{3,\infty}^{1/3}}^3 dt$$

By assumption and DCT, bounded.  $\square$

**Theorem 2.4.** (Isett 18') An energy dissipating solution whose singularities have 0 Lebesgue measure in  $\mathbb{R}^4$  cannot be of class  $L_t^r B_{\zeta,\infty}^{1/3}$  if  $r > 3$ .

Compared with Meneveau-Sreenivasan [11],

$$< |v(x + \Delta x) - v(x)|^r > = |\Delta x|^{\xi_r}$$

singular support in  $L_t^3 B_{3,C(N)}^{1/3}$ . (K41) implies  $\xi_r \sim \frac{r}{3}$  (only correct when  $r = 3$ ).

**Lemma 2.5.** (Local energy conservation Duchon-Robert[3] formula  $D[v, p] = \partial_t(\frac{|v|^2}{2}) + \nabla(\frac{|v|^2}{2} + p)v^j = \lim_{\varepsilon \rightarrow 0} \nabla_j v_{\varepsilon \rightarrow 0} \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell}$  dissipation distribution  $v \in L_{t,x}^3$ . If  $D[v, p] = 0$  and  $v \in L_{t,x}^2 \cap L_{t,x}^3$ , then  $\int \frac{|v|^2}{2}(t, x) dx$  is constant and  $D[v, p]$  if  $v, p \in C^1$ .

If  $v \in L_t^r B_{r,\infty}^{1/3}$  with  $r > 3$  is energy dissipating, since  $v \in L_{t,x}^2 \cap L_{t,x}^3$ ,  $D[v, p] \neq 0$ . Since  $\frac{d}{dt} \int \frac{|v|^2}{2} = \int D[v, p] dx$  But we claim that

$$\|D[v, p]\|_{L_{t,x}^{r/3}} < \infty$$

using Duchon-Robert formula:  $\|\nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell}\|_{L_{t,x}^{r/3}}$  is bounded uniformly in  $\varepsilon > 0$ .

$$\begin{aligned} \|\nabla_j v_\varepsilon\|_{L_x^r} &\lesssim \varepsilon^{-1+1/3} \|v(t)\|_{B_{r,\infty}^{1/3}} \\ \|R_\varepsilon^{j\ell}\| &\lesssim \varepsilon^{2/3} \|v(t)\|_{B_{r,\infty}^{1/3}}^2 \\ \Rightarrow \quad \forall \phi \in C^i nftyty_c(I \times \mathbb{R}^d) \quad < D < [v, p], p > \lesssim C \|p\|_{L_{t,x}^s} \end{aligned}$$

Then  $D[v, p]$  is in the dual of  $L_{t,x}^s$  which is  $L_{t,x}^{r/3}$  provided  $r > 3$ . Let  $\frac{1}{s} + \frac{3}{r} = 1$ .  $\text{supp } D[v, p]$  has positive Lebesgue measure, but  $\text{supp } D[u, p] \subset \text{sing}(\text{supp } U)$  also has positive measure.

There is an open problem to find a function  $f(r)$  s.t. the condition  $\frac{\zeta_r}{r} < \frac{1}{3} - f(r)$  works.

*Proof.* (Proof of Duchon-Robert formula) Considering Euler equation(Euler)

$$\eta_{\varepsilon\delta} * u := J_\zeta *_t \eta_\varepsilon *_x u$$

Let's test against  $w_\varepsilon \delta = \eta_{\varepsilon\delta} * (\phi \eta_{\varepsilon\delta} * v)$ .

$$0 = - \int_{I \times \mathbb{R}^d} v^\ell \partial_t \eta_{\varepsilon\delta} + (\phi \eta * v_\ell) + v^j v^\ell \nabla_j \eta_{\varepsilon\delta} * (\phi \eta_{\varepsilon\delta} * v) + p \nabla^\ell \eta_{\varepsilon\delta} * (\phi \eta_{\varepsilon\delta} * v_\ell) dx dt,$$



where  $\phi \in C_c^\infty(I \times \mathbb{R}^d)$ . Use the definition of self adjointness solution  $\eta_{\varepsilon\delta}*$  and divergence-free properties of  $\eta_{\varepsilon\delta} * v_\ell$ . Then Euler equation becomes

$$0 = - \int \partial_t \phi \frac{|\eta_{\varepsilon\delta} * v|^2}{2} + v^j v^\ell \eta_{\varepsilon\delta} \nabla_j \phi \eta_{\varepsilon\delta} * v_\ell + p \eta_{\varepsilon\delta} * (\nabla^\ell \phi \eta_{\varepsilon\delta} * v_\ell) dx dt$$

Let  $\varepsilon \rightarrow 0$  using uniform boundedness of  $\eta_\varepsilon*$  and  $\nabla_j \eta_\varepsilon*$ . As  $\delta \rightarrow 0$ , thanks to  $\nabla_j \eta_\varepsilon*$ , the time derivative naturally goes away. Then

$$0 = - \int \partial_t \phi \frac{v_\varepsilon^2}{2} + \nabla_j \phi \left( \frac{|v_\varepsilon|^2}{2} v^j + \eta_\varepsilon * p v_\varepsilon^j \right) dx dt \quad (2)$$

$$+ \int \phi \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell} + Z_\varepsilon, \quad (3)$$

where  $Z_\varepsilon = \int \nabla_j \phi R_\varepsilon^{j\ell} v_{\varepsilon l}$ . Take both time and space derivative of  $\phi$ . Using  $v \in L_{t,x}^2 \cap L_{t,x}^2$  and  $p = (-\Delta)^{-1} \nabla_j \nabla_\ell (v^j v^\ell) \in L_{t,x}^{3/2}$   
(2)  $\Rightarrow \langle D[u, p], \phi \rangle$  as  $\varepsilon \rightarrow 0$ .  
(3)  $\Rightarrow \lim_{\varepsilon \rightarrow 0} \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell} + Z_\varepsilon$

$$Z_\varepsilon = \int \nabla_j \phi (\eta_\varepsilon * (v^j v^\ell) - v_\varepsilon^j v_\varepsilon^\ell) v_{\varepsilon l} dx dt =: B_\varepsilon[v, v]$$

Here we define the commutator  $B_\varepsilon[\cdot, \cdot]$ .

$$\|B_\varepsilon[u, w]\|_{L_{t,x}^{3/2}} \leq C \|u\|_{L_{t,x}^3} \|w\|_{L_{t,x}^3}$$

which is independent of  $t$ .

If  $u$  or  $w \in C_c^\infty$   $\|B_\varepsilon[u, w]\|_{L_{t,x}^{3/2}} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . By approximation  $\|B_\varepsilon[v, v]\|_{L_{t,x}^{3/2}} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ .

By Holder inequality,

$$|Z_\varepsilon| \leq \|\nabla \phi\|_{L^\infty} \|B_\varepsilon[v, v]\|_{L_{t,x}^{3/2}} \|v_\varepsilon\|_{L^2} \rightarrow 0 \text{ as } \varepsilon \rightarrow 0$$

□

*Remark.* Improvement:

- Calderon-Zygmund Theorem.
- More regularity in time.

*Proof.* ( $\text{supp } D[u, p] \subset \text{sing}(\text{supp } v)$ ) In fact  $\text{supp } D[u, p] \subset \text{sing}(\text{supp } L_\varepsilon^3 B_{3,C(N)}^{1/3} v)$ .

What's good for not using Littlewood-Paley definition of Besov space? the solution above can be defined locally.

$\phi \in C_c^\infty(I \times B_q)$  and  $B'_q \subset B_q$  a smaller ball with same center  $q$ . Let  $q \notin \text{sing}(\text{supp } B_{3,C(N)}^{1/3} v)$

$$\langle \phi, D[v, p] \rangle = \lim_{\varepsilon \rightarrow 0} \int_I \int_{B'_q} \phi(t, x) \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell} \lesssim \|\phi\|_{L^\infty} \int_I \|\nabla v_\varepsilon(t)\|_{L(B'_q)} \|R_\varepsilon\|_{L_x^{2/3}}$$

dominated by  $\|\phi\|_{L^\infty} \int_I \|v(t, \cdot)\|_{B_{3,\infty}^{1/3}(B_\varepsilon)}^3$ . For a.e.  $t$ , we have  $\|\nabla v_\varepsilon\|_{L^3} \cdot \|R_\varepsilon\|_{L^{3/2}} = o(1)$  and  $D[u, p] \rightarrow 0$  by dominate convergence theorem. □

*Remark.* 1. Heat flow approach can also be applied to this problem. The proof is quite different[9] and is on a compact Riemann manifold(no convolution can be used).

2. Compressible Euler Case. The problem lies when dealing with commutator estimation.

### 3 Holder Continuity

The following lecture are basic on [9].

Note that  $B_{3,C(\mathbb{N})}^{1/3} \subsetneq B_{3,\infty}^{1/3}$  and we can find a function in  $B_{3,\infty}^{1/3} \setminus B_{3,C(N)}^{1/3}$ .

$\phi(x)\chi_{\{x'>0\}} \in B_{p,\infty}^{1/p} \quad \forall 1 < p < \infty$  Let's consider

$$\partial_t u + \partial_x \left( \frac{u^2}{2} \right) = 0 \quad L_t^\infty B_{3,\infty}^{1/3}$$

Energy dissipation at time  $t = 0$

$$\frac{d}{dt} e(t) = \int_{\mathbb{T}^d} \nabla_j v_{\varepsilon l} R_\varepsilon^{j\ell}(0, x) dx$$

Eyink proved that there exists a divergence-free vector field in the space  $C^{1/3} B_{3,C(N)}^{1/3}$ , s.t.  $\frac{d}{dt}|_{t=0} e(t) < 0$ . We have a useful counter example:

$$v(x) = \sum_q 2^{2q\alpha} \sin(2^{2q}x) \in \dot{B}_{3,C(N)}^{1/3} \setminus \dot{B}_{3,\infty}^{1/3}$$

Now we consider this problem on compact Riemannian manifold for the conclusion  $L_t^3 B_{3,C(N)}^{1/3}$ . Consider (Euler), instead of  $\eta_\varepsilon * v^\ell$ , we consider

- Estimates(Commutator)
- Keeping divergence-free property

Define the opetator  $-\Delta_H = d_\delta + \delta_d$ , which looks like a 1-form. In Hodge heat flow equation,

$$\partial_s v^\ell = \Delta_H v^\ell = \nabla_j \nabla^j v^\ell - \text{Ric}_\ell^k v^k$$

Since we know what the solution exactly is,

$$\eta_\varepsilon * v^\ell \rightarrow e^{s\Delta_H} v^\ell$$

The square root of heat time  $s^{1/2} \sim \varepsilon$  and solution at time  $s$   $S_{[s]}v = e^{s\Delta_H}v$ .

To estimate

$$\eta_\varepsilon * \nabla_j (v^j v^\ell) - \nabla_j (\eta_\varepsilon v^j \eta_\varepsilon * v^\ell),$$

we would need the commutator

$$w^\ell(s) = S_{[s]} \nabla_j (v^j v^\ell) - \nabla_j (S_{[s]} v^j S_{[s]} v^\ell)$$

and Riemannian manifold  $M$  will be always assumed to be smooth.

$$s \in (0, 1] \quad \int_{I \times M} \eta(t) [S_{[s]} \nabla_j (v^j v^\ell) - \nabla_j (S_{[s]} v^j S_{[s]} v^\ell) S_{[s]} v_\ell] d^{1+d} \text{vol},$$

here volume is in time  $\times$  space.

Let's calculate

$$\begin{aligned} (\partial_s - \Delta_H) w^\ell &= N^\ell(t, s) \\ w^\ell(s) &= \int_0^s e^{(s-s')\Delta_H} N^\ell(t, s') ds', \end{aligned}$$

by d'Alembert's formula.

$$\begin{aligned} w^\ell(t, s) &= S_{[s]} \nabla_j (v^j v^\ell) - \nabla_j (S_{[s]} v^j S_{[s]} v^\ell) \\ (\partial_s - \Delta_H) w^\ell &= (\partial_s - \Delta_H) \nabla_j (S_{[s]} v^j S_{[s]} v^\ell) \\ &= (\partial_s - \nabla_i \nabla^i) \nabla_j (S_{[s]} v^j S_{[s]} v^\ell) + \text{curvature terms} \\ &= -2 \nabla_j (\nabla_i S_{[s]} v^j \nabla^i S_{[s]} v^\ell) + \text{low order terms} \end{aligned}$$

$$\begin{aligned}
\text{Commutator} &= -2 \int_{I \times M} \eta(t) \int_0^s e^{(s-s')\Delta} \nabla_j (\nabla_i S_{[s]} v^j \nabla^i S_{[s]} v^\ell) ds' S_{[s]} v_\ell d^{1+d} \text{vol} \\
\text{Integral by parts} &= 2 \int_{I \times M} \eta(t) \int_0^s \nabla_i S_{[s]} v^j \nabla^i S_{[s]} \underbrace{S_{[2s-s']} v_\ell ds'}_{\text{very low frequency}} d^{1+d} \text{vol}
\end{aligned}$$

Assume that  $v \in L_t^3 B_{3,\infty}^\alpha$ , claim that  $\|\nabla S_{[s]} v\|_{L^3} \lesssim s^{-\frac{1+\alpha}{2}} \|v\|_{L_t^3 B_{3,\infty}^\alpha}$  with  $\alpha > 1/3$ . First we can try  $v \in L_t^3 C^\alpha$  or  $v \in L_t^3 W^{\alpha,3}$ .

$$\begin{aligned}
|\text{Commutator}| &\lesssim \|\eta\|_{L^\infty} \int_0^s (2s-s')^{-\frac{1+\alpha}{2}} (s')^{-2\frac{1+\alpha}{2}} ds' \\
&\lesssim s^{-1/2+3\alpha/2} \int_0^1 (2-\sigma)^{-\frac{1+\alpha}{2}} \sigma^{-1+\alpha} d\sigma \rightarrow 0 \text{ as } \varepsilon \rightarrow 0 \text{ with } \alpha < 1/3
\end{aligned}$$

No derivatives that we can take over the heat flow. In order to prove the claim:  $\|\nabla S_{[s]} v\|_{L^3} \lesssim s^{-\frac{1+\alpha}{2}} \|v\|_{W^{\alpha,3}}$  with  $\alpha > 1/3$ .

*Proof.*

$$\begin{aligned}
\|\nabla S_{[s]} v\|_{L^3} &\leq s^{-\frac{1}{2}} \|v\|_{L^3} \\
\|\nabla S_{[s]} v\|_{L^3} &\leq \|v\|_{W^{1,3}}
\end{aligned}$$

Since  $u(s) = e^{s\Delta} u$ ,  $\|u\|_{L^\infty} \lesssim \|u\|_{L^p}$   $s \in (0, 1]$  On compact manifold we have that

$$\|\nabla u\|_{L^p} \lesssim \|\nabla u\|_{L^r} + \|u\|_{L^p}$$

Here let  $p = 2z$  where  $z \geq 2$  is an integer. □

$$\begin{aligned}
\partial_s |u|^2 - \Delta |u| + |\nabla u|^2 &= 0 \text{ (or } -2 \text{Ric}_{jk} u^j u^k) \\
&\Rightarrow \int_M |u|^2 d\text{vol} \searrow \text{ and } \int_M u^{2z} d\text{vol} \searrow \\
\frac{1}{z} \partial_s \int |u|^{2z} d\text{vol} &= \int_M (\Delta |u|^2 - 2|\nabla u|^2) |u|^{2(z-1)} d\text{vol} \\
&= - \int_M \nabla^j |u|^2 \nabla_j |u|^{2(z-2)} - 2|\nabla u|^2 |u|^{2(z-1)} \leq 0
\end{aligned} \tag{4}$$

For curvature terms, they can be bounded by  $\|\text{Ric}\|_{L^\infty} \int_M |u|^{2z} d\text{vol}$  remains bounded. So  $\int_M u^{2z} d\text{vol} \leq \int_M u^{2z}(t=0) d\text{vol}$ .

$$\partial_s |\nabla u|^2 - \Delta |\nabla u|^2 + |\nabla \nabla u|^2 (\text{BAD}) = 0 \text{ (or } \text{Riem}(\nabla u \nabla u) + \nabla \text{Riem} u \nabla u) \tag{5}$$

Multiply by  $|\nabla u|^{2(z-1)}$ ,

$$\frac{1}{z} \partial_j \int_M |\nabla u|^{2z} d\text{vol} \searrow$$

or

$$\begin{aligned}
&\int_M |\nabla u|^{2(z-1)} \nabla_{\text{Riem}} u \nabla u d\text{vol} \quad \text{by integration by parts} \\
&\leq 2 \underbrace{\int_M |\nabla \nabla u|^2 |\nabla u|^{2(z-1)}}_{\text{BAD}} + C \underbrace{z \int_M \|R\|_{L^\infty} \|\nabla u\|^{2z} + \|\text{Riem}\|^{z+1} \|u\|^{2z}}_{\text{GOOD}}
\end{aligned}$$

which can cancel bad terms from (5).

$$\partial_s (s |\nabla u|^2) - \Delta (s |\nabla u|^2) + 2s |\nabla \nabla u|^2 - |\nabla u|^2 (\text{BAD})$$

But here the bad term can cancel with (4). Let  $\Phi(s) = s |\nabla u|^2 + \frac{1}{2} |u|^2$ , then we have

$$\partial_z \Phi_s - \Delta \Phi_s = 2s |\nabla \nabla u|^2$$

Thus

$$\begin{aligned}
\frac{1}{z} \partial_s \int_M |\Phi_s|^z (\searrow) &\leq \frac{1}{z} \int_M |\Phi_s|^z d\text{vol} \leq \frac{1}{z} \int_M |\Phi_0|^z = \frac{1}{z} \int_M |u|^{2z} \\
\frac{1}{z} \left( \int_M |\nabla u|^{2z} \right)^{1/2z} &\leq s^{-1/2} \left( \int_M |u|^{2z} \right)^{1/2z}
\end{aligned}$$

Hodge Laplacian commute with derivatived and divergenced.  $(\partial_s - \Delta_H) \delta w_\ell = \delta(\partial_s - \Delta_H) w_\ell = 0$  with 0 initial condition.

### 3.1 Isentropic Compressible Euler

$$\begin{aligned} \text{Mass} \quad & \partial_t \rho + \nabla_j (\rho v^j) = 0 \\ \text{Momentum} \quad & \partial_t (\rho v^\ell) + \nabla_j (\rho v^j v^\ell) + \nabla^\ell (P\rho) = 0 \end{aligned} \quad (6)$$

We made some assumption  $p(\rho) = \rho^2 \gamma$  and  $\rho \in C^2$  and away from  $\rho \equiv 0$ . Here exists a problem: modifier doesn't commute with nonlinearity.

$$\partial_t \left( \frac{1}{2} \rho |v| + p(\rho) \right) + \nabla_j \left( \left( \frac{1}{2} \rho |v| + p(\rho) + p(\rho) v^j \right) p(\rho) \right) = \rho \int_1^\rho \frac{p(r)}{r^2} dr$$

If conservation holds in  $B_{3,\infty}^\alpha$  in both  $(t, x)$ , we need to estimate the commutator:  $\eta_\varepsilon * (p(\rho)) - p\eta * \rho$ .

*Remark.* • Heat flow also works.

- [5] gives another method

$$\begin{aligned} \text{Commutator} &= \int p(\rho(x-h)) \eta_\varepsilon(h) dh - p \left( \int \rho(x-h) \eta_\varepsilon(h) dh \right) \\ &= \int p(\rho(x-h)) \eta_\varepsilon(h) dh - p \left( \int (\rho(x-h) \eta_\varepsilon(h)) dh \right) = \int (1-\sigma) \int p''((1-\eta)\rho_\varepsilon(x) + \sigma\rho(x-h)) (\rho(x-h) - \rho_\varepsilon(x)) \end{aligned}$$

Therefore bounded.

$$\overline{f(X)} - f(\overline{X}) = \mathbb{E}[f(X)] - f(\mathbb{E}[X]) \quad (7)$$

Since

$$\begin{aligned} \phi(1) &= \phi(0) + \int_0^1 \frac{d}{d\sigma} \phi(\sigma) d\sigma \\ &= \phi(0) + \frac{d}{d\sigma} \Big|_{\sigma=0} \phi(\sigma) + \int_0^1 (1-\sigma) \frac{d^2}{d\sigma^2} \phi(\sigma) d\sigma \end{aligned}$$

RHS of (7) becomes

$$\begin{aligned} & \int +0^1 \frac{d}{d\sigma} \mathbb{E}[f(1-\sigma)\overline{X} + \sigma(X)] \\ &= \frac{d}{d\sigma} \Big|_{\sigma=0} \mathbb{E}[f(1-\sigma)\overline{X} + \sigma(X)] + \int_0^1 (1-\sigma) \frac{d^2}{d\sigma^2} \mathbb{E}[f(1-\sigma)\overline{X} + \sigma(X)] d\sigma \\ &= \mathbb{E}[\nabla_i f(\overline{X})(X^i - \overline{X}^i)] + \int_0^1 (1-\sigma) \mathbb{E}[\nabla_a \nabla_b f((1-\sigma)\overline{X} + \sigma X)(X^a - \overline{X}^a)(X^b - \overline{X}^b)] d\sigma \end{aligned}$$

Here last term is a quadratic form.

## 4 Convex Integration

Here "convex" refers to convex linear combination.

**Theorem 4.1.** (Old conclusion)  $\forall E(t) \in C^\infty, E(t) \geq C, \forall \alpha < \frac{1}{10}, \exists v \in C_{t,x}^\alpha(I \times \mathbb{T}^3)$  s.t.  $\int \frac{|v|^2}{2}(t, x) dx = E(t)$  conserved.

**Theorem 4.2.** (Isset [6])  $\forall \alpha < 1/5, \exists v \in C_{t,x}^\alpha(I \times \mathbb{T}^3, p \in C_{t,x}^{2,\alpha})$ . A non trivial solution with compact support in time. (0 is not the only solution stays 0 implies non-uniqueness).

- Question: How to construct continuous solution?
- Idea: Euler-Reynolds flows

For  $R$  a symmetric tensor  $(v, p, R)$  that solves

$$\begin{aligned} \partial_t v^l + \nabla_j (v^j v^\ell) + \nabla^\ell p &= \nabla_j R^{j\ell} \\ \nabla_j v^j &= 0 \end{aligned}$$

Here  $R^{j\ell} = R^{\ell j}$ . If  $R = 0$ , we have the Euler equation(Euler). If  $R \neq 0$ , we have a smooth approximation of Euler Equation.

Construction  $(v_q, p_q, R_q)$  the Euler-Reynolds flow:

$$R_q \rightarrow 0 \text{ as } q \rightarrow \infty$$

and

$$(v_q, p_q)$$

converge uniformly.

$$-\int \partial_t \phi_\ell v_q^l + \nabla_j \phi_\ell + v_q^j v_q^\ell + \nabla^\ell \phi \cdot p_q = \nabla_j \phi R_q^{j\ell} \rightarrow 0$$

The idea would make sense if every continuous solution  $(v, p)$  is a uniform limit of  $(v_\varepsilon, p_\varepsilon, R_\varepsilon)$  as  $\varepsilon \rightarrow 0$ . I.e. every Euler flow is a limit of Euler-Reynolds flow. Proof will also be given by modification. To check weather it is a E-R flow, we need to check weather it conserves energy.

**Lemma 4.3.** (Main Lemma) Given  $(v, p, R) \ni$  new

$$(\mathring{v}, \mathring{p}, \mathring{R})$$

with

$$\|\mathring{R}\|_0 << \|R\|_0,$$

where

$$\mathring{v} = v + V \quad \mathring{p} = p + P. \quad (8)$$

We expect that  $\|V\|_0 \leq \|R\|_0^{1/2} \leftarrow [\frac{m}{s}]$  and  $\|P\|_0 \leq \|R\|_0 \leftarrow [\frac{m}{s}]$  by dimension analysis.

Apply 4.3 over and over again generate  $(v_q, p_q, R_q)$  with  $\|R_q\|_0 \rightarrow 0$  rapidly. Set  $v = v(0) + \sum_q v_q$  and  $p = p(0) + \sum_q p_q$ . Plug in new 8

$$\partial_t \mathring{v}^{\ell} + \nabla_j (\mathring{v}^i \mathring{v}^{*l}) + \nabla^\ell \mathring{p} = \nabla_j R^{j\ell} + \partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla_j (v^j v^\ell) + \nabla^\ell p + \nabla_j (v^j v^\ell)$$

i.e.

$$\text{RHS} = \nabla v^\ell + \nabla_j (v^j V^l) + \nabla_j (V^j V^l + P \delta^{jl} + R^{jl}) + \nabla_j (v^l V^j)$$

where  $R^{jl}$  is the old error. We want  $\nabla_j R^{*jl} = \text{RHS}$  with  $R^*$  small. <sup>1</sup>

The rest is not small. Let them equals to  $\nabla_j Q^{jl}$  with  $Q^{jl} = Q^{lj}$ . Either  $\nabla_j (v^{jl} + p \delta^{jl} + R^{jl})$  or  $\nabla_j (v^j V^l)$  is not small.

## 4.1 Non-stationary Phase Lemma

When you can find a small solution  $Q^{j\ell} = Q^{\ell j}$  to the symmetric divergence equation

$$\nabla_j Q^{j\ell} = u^\ell$$

Necessary for  $u^\ell$  to be high-frequency  $n \geq 1$  and  $\int_{\mathbb{T}^3} u^\ell dx = 0$ .

What does  $v^\ell$  look like? [2] Beltrami flows.

$$v^\ell = \sum_{|k|=R} a_k B_k^\ell e^{i\lambda k \cdot x}, \quad k \in \mathbb{R}^d$$

$\nabla \times v = R\lambda v$  if  $(ik) \times B_k = |k|B_k$  and  $a_k B_k$  are constant.

$$v^l = \sum_{|k|=R} a_k(R, E, v, \lambda t) B_k^\ell e^{i\lambda k x}$$

Construction of frequency:

$$v^\ell = \sum_I v_I^l, \quad v_I^\ell = \bar{v}_I^l, \quad \nabla_j v_I^j = 0$$

$$v_I^l = \nabla_j w_I^{jl}, \quad w_I^{jl} = -w_I^{lj} \Leftrightarrow$$

---

<sup>1</sup>Here the position of  $*$  is decided by whether it is a equation or just math script. The different position of  $*$  doesn't have different meaning

$w^{jl}$  is a curl. For high frequency wave

$$\begin{aligned} v_I^\ell + \delta v_I^\ell & \quad \text{with } \lambda \gg 1 \\ v_I \cdot \nabla \xi_I &= 0 \quad \text{for divergence free} \end{aligned}$$

Here  $\xi_I$  is always tangent to level phase function.

$$w_I^{jl} = \frac{e^{i\lambda\xi_I}}{i\lambda} w_I^{j\ell} = \frac{e^{i\lambda\xi}}{i\lambda} (\nabla^j \xi_I v_I^\ell - v_I^j \nabla^l \xi_I)$$

then it is anti-symmetric. Thus

$$\nabla_j w_I^{j\ell} = e^{i\lambda\xi_I} (\nabla_j \xi_I w_I^{il} + \frac{\nabla_j w_I^{j\ell}}{i\lambda}) = e^{i\lambda\xi_I} (v_I^\ell + \delta v_I^\ell),$$

where  $\delta v_I = \mathcal{O}(\frac{1}{\lambda})$ , thus small. The transport term:

$$\nabla_j R_T^{jl} = \partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla_j (V^j v^\ell)$$

want to solve

$$\nabla_j Q^{jl} = u^\ell = e^{i\lambda\xi_I} v_I^\ell$$

with small  $Q^{j\ell} = Q^{\ell j}$  on  $\mathbb{T}^3$ . The equations above are under-determined. So the solution is not unique and we are try to find a small one.

**Lemma 4.4.** *Non-stationary Phase Lemma* If  $\left\| |\nabla \xi_I|^{-1} \right\|_0 \leq A$  (no zero frequency) and  $\int_{\mathbb{T}^3} u^\ell(x) dx = 0$ . We note that if  $|\nabla \xi_I| \rightarrow 0$ , then we just rotating a constant but there is no waves. Then  $\exists Q^{j\ell} = Q^{\ell j}$  of size  $\|Q\| \leq \mathcal{O}(\lambda^{-1}) \|u\|_0 + \mathcal{O}_{(u, \nabla q)}(\lambda^{-2})$ . Since  $\int \partial_t V^l dx = 0 \Rightarrow \frac{d}{dt} \int v^\ell dx = 0$  conserve the momentum, consistent with solution of Euler.

*Proof.* d=1.

$$\frac{dQ}{dx} = e^{i\lambda\xi(x)} u(x) \quad \text{on } \mathbb{T} = \mathbb{R}/\mathbb{Z}$$

$\Rightarrow Q(x) = \int_0^x e^{i\lambda\xi(s)} u(s) ds$  is periodic because  $\int_{\mathbb{T}} u(x) dx = 0$ . By integration by parts,

$$= \frac{1}{i\lambda\xi(s)} e^{i\lambda\xi(s)} u(s) \Big|_{s=0}^x - \frac{1}{i\lambda} \int_0^x e^{i\lambda\xi(s)} \frac{d}{ds} \left( \frac{u(s)}{\xi(s)} \right) ds = \mathcal{O}\left(\frac{1}{\lambda} + \frac{1}{\lambda} \mathcal{O}(\lambda^{-2})\right)$$

Let's look at the transport term

$$\begin{aligned} \nabla R_T^{j\ell} &= \partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla_j (v^j v^\ell) = \partial_t v^\ell + v^j \nabla_j v^\ell + v^j \nabla_j v^\ell \\ &= (i\lambda) \sum_I e^{i\lambda\xi_I} (\partial_t \xi_I) v_q^\ell + (v^j \nabla_j) \xi_I v^\ell + \sum_I e^{i\lambda\xi_q} \end{aligned}$$

□

Apply the Non-stationary Phase Lemma 4.4

$$\|R_T\|_0 \lesssim \lambda^{-1} \|RHS\|_0 \lesssim \lambda^{-1} (\lambda \|(\partial_t + v^j \nabla_j) \xi_I\|_0 \|V_I\|_0) + \mathcal{O}(\lambda^{-2})$$

Impose the phase function:

$$\begin{aligned} (\partial_t + V \cdot \nabla) \xi_I &= 0 \\ D_t \xi_I &= 0 \end{aligned} \tag{9}$$

Taylor hypothesis: High frequency flows are actually carried by low frequency flows. So it make sense with our assumption. If  $\left\| |\nabla \xi_I|^{-1} \right\|_0 \leq A$  fails to control the amplitude, we include time cut-offs in  $v_I$ . Look at  $\nabla_j (v^j v^\ell + p \delta^{jl} + R^{jl})$ , want it to be div of something small. Want  $R_{small}^{j\ell} = \sum_I v_I^j \bar{v}_I^\ell + P_0 \delta + R^{j\ell}$

$$\sum_I v_I^j \bar{v}_I^\ell + P_0 \delta^{j\ell} + \sum_{j \neq \bar{I}} v_I^j v_j^\ell + P_{Ij} \delta^{j\ell} = \sum_I v_I^j \bar{v}_I^\ell + P_0 \delta^{j\ell} + R^{j\ell} + \sum_I \delta v_I^j v^\ell \left( = 0 + \mathcal{O}\left(\frac{1}{\lambda}\right) \right)$$

Let's pick a  $t$  to make this = 0 pointwisely.

$$\sum_I v_I^j \bar{v}_j^\ell = -P_0 \delta^{j\ell} - R^{j\ell} = e(t) \delta^{j\ell} - E^{j\ell}$$

Choose  $e(t) \geq 10^4 \|R\|_0$  on  $\text{supp } R$ , here  $e(t)$  is the lifting function.

$$\Rightarrow \|V_I\|_0 \lesssim e^{1/2} \lesssim \|R\|_0^{1/2}$$

We need to choose time depending on amplitude. Note that  $\text{supp}_t e \subset B(\text{supp } R, \varepsilon)$ ,

$$\Rightarrow \text{supp}_t(\text{New Error}) \subset \text{supp}_t e$$

For (8), we rewrite our (Euler) as:

$$\begin{aligned} \partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p &= \nabla_j R^{jl} \\ \nabla_j v^j &= 0 \end{aligned} \quad (10)$$

where we decompose  $v^\ell = \sum_I v_I^l$  and  $\nabla_j V_I^j = 0$  by  $v^l = \nabla_j w_I^{jl}$  since  $w_I^{jl}$  is anti-symmetric tensor.  $V_I^l = e^{i\lambda \xi_I} (v_I^\ell + \delta v_I^\ell)$ , where  $\delta v_I^\ell \sim \mathcal{O}(\lambda^{-1})$ . As for phase  $\xi_I \in \mathbb{C}$ , we have  $\left\| |\nabla \xi_I|^{-1} \right\|_0 \leq A$ ,  $V_I \cdot \nabla \xi_I = 0$ . Let's impose  $(\partial_t + v \cdot \nabla) \xi_I = 0 \Rightarrow$  We need time cut offs inside  $v_I$  to maintain the non-stationary phase. Let  $\|R\|_0 \leq e_R$

$$R^* = \begin{array}{cccc} R_T^{j\ell} + & R_L^{j\ell} + & R_S^{j\ell} + & R_H^{j\ell} \\ \text{Transport} & \text{Low Frequency} & \text{Stress} & \text{High Frequency} \end{array}$$

$$\begin{aligned} \nabla_j R_T^{j\ell} &= \partial_t v^\ell + \nabla_j (v^j V^\ell) \\ \nabla_j R_L^{j\ell} &= \nabla_j (v^j V^L) \\ \nabla_j R_S^{j\ell} &= \nabla_j (\sum_I v_I^j \bar{v}_I^\ell + p \delta^{j\ell} + R^{j\ell}) \\ \nabla_j R_H^{j\ell} &= \sum_{j \neq I} \nabla_j \end{aligned}$$

$$R_S^{j\ell} = \sum_I V_I^j \bar{v}_I^\ell + \rho \delta^{j\ell} + R^{j\ell} + \mathcal{O}(\delta v_I)$$

(Phase cancel thus high frequency disappear)

$$\sum_I V_I^j V_I^{-\ell} = -P_0 \delta^{j\ell} - R^{j\ell}$$

Let's set  $p_0 = -e(t)$ ,  $e(t) \geq K e_R(\text{error})$ .  $\sum_I v_I^j \bar{v}_I^{-\ell} = e(t) \delta^{j\ell} - R^{j\ell}$ .  $e(t) \geq K e_\ell$  on  $\text{supp } R$ .

What about the high frequency?

$$\begin{aligned} \nabla_j R_H^{j\ell} &= \sum_{j \neq I} \nabla_j (V_I^j \bar{V}_I^\ell) + \nabla^\ell P_{IJ} \\ &= \frac{1}{2} \sum_{j \neq I} V_I^j \nabla_j V_j^\ell + V_j^j \nabla_j V_I^\ell + \nabla^\ell P_{IJ} \\ &= (i\lambda) \sum_{j \neq J} e^{i\lambda(\xi_I + \xi_J)} \nabla_j \xi_J V_j^\ell + \text{L.O.T} \\ \|R_H\|_0 &\lesssim \lambda^{-1} \|\text{RHS}\|_0 + \mathcal{O}(\lambda^{-2} \|\text{RHS}\|) \\ &= \lambda^{-1} \lambda \|V_I\|_0 \|\nabla \xi_I\| \|V_I\|_0 \lesssim e_R^{1/2} 1 e_R^{1/2} \end{aligned} \quad (11)$$

Better way "write  $V_I$  as a steady state solution".

Idea: If  $\nabla \cdot V_I = \lambda V_I$  and  $\nabla \cdot V_J = \lambda V_J$  then  $V_I + V_J$  is steady state for E-R flow with appropriate pressure  $P_{IJ}$  s.t.

$$V_I^j \nabla_j V_J^i + V_J^j \nabla_j V_I^i + \nabla^i P_{IJ} = 0$$

To make  $\nabla \times V_I \sim \lambda V_I$

$$\begin{aligned} \nabla \times (e^{i\lambda \xi_I} V_I) &= \lambda e^{i\lambda \xi_I} ((i \nabla \xi_I) \times V_I) \\ (i \nabla \xi_I) \times V_I &= |\nabla \xi_I| V_I \sim \lambda e^{i\lambda \xi_I} |\nabla \xi_I| \times V_I \end{aligned}$$

Here  $|\nabla \xi_I|$  is eigenvalue. If  $\| |\nabla \xi_I| - 1 \|_0 < 1$ , (by sharp time cut-off)

$$\begin{aligned} (i\nabla \xi_I) \times V_I &= -(\nabla \xi) \times b_I + i(\nabla \xi_I) \times a_I \\ |\nabla \xi_I| \cdot V_I &= |\nabla \xi_I| a_I + i|\nabla \xi_I| b_I \end{aligned}$$

choose  $b_I \in \langle \nabla \xi_I \rangle^+$  and set  $a_I = \frac{-(\nabla \xi_I) \times b_I}{|\nabla \xi_I|}$ , thus

$$\begin{aligned} (\nabla \xi_I) \times a_I &= -\nabla \xi_I \times \frac{(\nabla \xi_I \times b_I)}{|\nabla \xi_I|} = \frac{-(-|\nabla \xi|^2 b_I)}{|\nabla \xi_I|} = |\nabla \xi_I| b_I \\ e^{i\lambda \xi_I} V_I^l &= e^{i\lambda \xi_I} (a_I^l + i b_I^l + I) = \cos \lambda \xi_I a_I^l - \sin \lambda \xi_I b_I^l \\ a_I &= \frac{-\nabla \xi_I \times b_I}{|\xi_I|} \end{aligned}$$

By (11)

$$\frac{1}{2} \sum_{J+I} (V_I)_j (\nabla^j V_J^j) + (V_J)_j (\nabla^j V_I^\ell - \nabla^\ell v_I^j) + \nabla^\ell \frac{V_i V_J}{2} + \nabla^\ell P_{IJ},$$

$\nabla^\ell \frac{V_i V_J}{2} + \nabla^\ell P_{IJ} = 0$  by our choice of  $P_{IJ}$ .

$$P_{IJ} = -\frac{1}{2} V_I V_J = -\frac{1}{2} \sum_{J \neq \bar{I}} (V_I \times (\nabla \times V_J) + V_J \times (\nabla \times V_I)) = 0$$

If  $V_J, V_I$  are eigen-function of  $\nabla \times$ .

$$\begin{aligned} P_{IJ} &= -\frac{1}{2} V_I V_J = \lambda e^{\lambda(\xi_I \xi_J)} (V_I \times ((i\nabla \xi_J) \times V_J) + V_J \times (i\nabla \xi_I) V_I) + \text{L.O.T} \\ &= \lambda e^{\lambda(\xi_I \xi_J)} (V_I \times (|\nabla \xi_J| - 1) V_J) - V_J \times (|\nabla \xi_I| - 1) V_I + \text{L.O.T} \end{aligned}$$

for  $V_I \times V_J + V_J \times V_I = 0$ . Since non-stationary phase  $\| |\nabla \xi_I + \nabla \xi_J|^{-3} \| \leq A$ .

$$\begin{aligned} \|R_H\|_0 &\lesssim \lambda \cdot \lambda \max_I \|V_I\|_0^2 \cdot \max_I \| |\nabla \xi_I| - 1 \|_0 \\ &\lesssim e_R \max_I \| |\nabla \xi_I| - 1 \|_0 + \text{L.O.T} \end{aligned}$$

Can be made small with a sharp time cut off in each  $V_I$ :  $\tau \sim b|\nabla V|^{-1}$ . Then we need to pay the price for time cut-off.

$$\begin{aligned} \nabla_j R_T^{j\ell} &= \partial_t v^\ell + \nabla_j (v^j V^\ell) = \partial_t v^\ell + v^j \nabla_j V^\ell \\ &= \sum_I \exp(i\lambda \xi_I) (\partial_t + v^j \nabla_j) V_I^\ell \\ \|R_T\|_0 &\lesssim \lambda^{-1} \text{RHS} + \mathcal{O}(\lambda^{-2}) \lesssim \lambda^{-1} \tau^{-1} \end{aligned}$$

Choose  $\lambda \gg \tau^{-1} \Rightarrow \|R_T\|_0$  is arbitrarily small.

*Remark.* • Convex linear combination problem  $\sum v^j \bar{v}_I^\ell = e(t) \delta^{jl} - R^{jl}$

- Integration means using  $\nabla R^{j\ell}$  to approximate  $R^{j\ell}$ .

**Lemma 4.5.** (Main Lemma)  $\exists K, \exists C > 1$  s.t.  $\forall \varepsilon > 0 \forall (v, p, R)$  uniformly  $C^3$  solution to Euler-Reynolds with  $\text{supp } R \subset I \times \mathbb{T}^3$ , and  $\|R\|_0 \leq e_R$ . Let  $e(t) : \mathbb{R} \rightarrow \mathbb{R}^*$ , s.t.

$$\frac{d}{dt} e^{\frac{1}{2}}(t) \in C^\infty$$

and  $e(t) \geq K e_R$  on a neighbourhood of  $\text{supp } R$ .  $\exists (v^*, p^*, R^*)$ ,  $v^* = v + V$ ,  $p^* = p + P$ .  $\text{supp}(V, P, R) \subset \text{supp } I \times \mathbb{T}^3$ , with  $\|V\|_0 \lesssim e_R^{1/2}$ ,  $\|P\|_0 \lesssim e_R$  and  $\|R^*\|_0 < \varepsilon \forall \varepsilon > 0$  by choosing arbitrarily big  $\lambda$ .  $\left\| \int_{\mathbb{T}^3} |v^*|^2 - \int_{\mathbb{T}^3} |v|^2 + e(t) dx \right\|_0 \leq \varepsilon$  (i.e. nice low frequency).



*Proof.*

$$\int_{\mathbb{T}^3} \left| \begin{smallmatrix} * \\ \mathbf{v} \end{smallmatrix} \right|^2 - \left| \begin{smallmatrix} * \\ \mathbf{V} \end{smallmatrix} \right|^2 = \int_{\mathbb{T}^3} |v + V|^2 - |v|^2 dx = \int_{\mathbb{T}^3} 2vV + V^2 dx$$

The cross term is the correction of high frequency and is small because  $v, V$  are almost orthogonal.

Claim:  $\int_{\mathbb{T}^3} |V|^2 dx + \mathcal{O}(\lambda^{-1})$ . Let  $V^\ell = \nabla_j W^{j\ell}$ , where  $W^{j\ell} = \frac{1}{\lambda} e^{i\lambda \xi_I} (w_I^{j\ell})$ .

$$\begin{aligned} \int vV &= - \int \nabla_j v_\ell W_I^{j\ell} dx = \mathcal{O}(\lambda^{-1}) \\ \int |V|^2 dx &= \sum_I \int V_I \bar{V}_I + \sum_{J \neq I} \int V_I \cdot V_J \end{aligned}$$

$\sum_{J \neq I} \int V_I \cdot V_J$  is small because for  $\sum_{J \neq I} \int e^{i\lambda_I + \lambda_J} V_I V_J dx$ , since  $\left\| |\nabla \xi_I + \nabla \xi_J|^{-1} \right\|_0 \leq A$ . If  $|\nabla \xi|^{-1} \leq A$  failed.

$$\begin{aligned} e^{i\lambda \xi} &= \frac{\nabla^a \xi_I \nabla_a \xi_J}{(i\lambda) |\nabla \xi|^2} \quad \xi = \xi_I + \xi_J \\ &= \sum_{J \neq I} \int e^{i\lambda_I + \lambda_J} \frac{\nabla^a \xi}{|\nabla \xi|^2} V_I V_J dx \\ &= \mathcal{O}(\lambda^{-1}) \quad \text{by IBP} \end{aligned}$$

$$\begin{aligned} \sum_I \int V^I \bar{V}_I, \quad V_I &= e^{i\lambda \xi_I} (V_I^I + \delta V_I^\ell) \\ \int \delta_{jl} \sum_I V_I^j \bar{V}_I^\ell dx &\quad \text{since } \sum_I V_I^j \bar{V}_I^\ell = -P_0 \lambda^{j\ell} - R^{j\ell} \\ &= \delta_{jl} \int (-P_0 \delta^{j\ell} - R^{j\ell}) \\ &= \delta_{jl} \int (e(t) \frac{\delta^{j\ell}}{d} - \hat{R}^{j\ell}) dx \quad \text{by choose } p_0 = -\frac{e(t)}{d} - \frac{\delta_{ij} R^{j\ell}}{d} \end{aligned}$$

□

Since  $V_I^J \nabla_j V_J^\ell = (V_I)_j (\nabla^j V_J^I - \nabla_J^j) = -V_I \times (\nabla \times V_J)$ . Define  $(\nabla \times V_J)^c = \varepsilon^{cdf} \nabla_d V_{Jf}$ . we notice that here:

$$\varepsilon^{abs} = -\varepsilon^{bac} = -\varepsilon^{cbs}$$

etc. And  $\varepsilon^{123} = 1$ ,

$$\varepsilon^{\ell jc} \varepsilon^{cdf} = \delta_d^\ell \delta_f^j - \delta_f^\ell \delta_d^j$$

We find it anti-symmetric in  $\ell j$  and  $df$ . Proof:

$$\varepsilon^{12c} \varepsilon^{cdf} = \begin{cases} 0 & \text{if } df \notin \{(12) \text{ or } (21)\} \\ 1 & \text{if } df = (12) \\ -1 & \text{if } df = (21) \end{cases} = \delta_d^\ell \delta_f^j - \delta_f^\ell \delta_d^j$$

$$\begin{aligned} V_I \times \nabla \times V_J &= \varepsilon^{\ell jc} (V_I)_j (\nabla \times V_J)_c \\ &= \varepsilon^{\ell jc} (V_I)_j \varepsilon^{cdf} \nabla^d V_J^f \\ &= (\delta_d^\ell \delta_f^j - \delta_f^\ell \delta_d^j) (V_I)_j \nabla^d V_J^\ell \\ &= (V_I)_j \nabla^\ell V_J^j - (V_I)_j \nabla^j V_J^\ell \\ &= (V_I)_j (\nabla^\ell V_J^j - \nabla^j V_J^\ell) \end{aligned}$$

**Lemma 4.6.** (*Main Lemma*)  $\exists k \geq 1 \exists C \geq 1$  s.t.  $\forall \varepsilon > 0, \forall (v, p, R)$  uniform  $C^1$  Euler-Reynolds flow with  $\text{supp } R \subset I \times \mathbb{T}^3, \|R\|_0 \leq e_R$ . For any  $e(t) : \mathbb{R}^* \rightarrow \mathbb{R}^+$

- $e^{1/2}(t)$  is  $C_c^\infty$
- $e(t) \geq K e_R$  on a neighbourhood of  $I$ .

$\exists (v^*, p^*, R^*)$ , where  $v^* = v + V, p^* = p + P, \|R\|_0 < \varepsilon \text{ supp}(v, p, R) \subset \text{supp } I \times \mathbb{T}^d$

*Proof.*

$$\left\| \int \left| \begin{smallmatrix} * \\ \mathbf{v} \end{smallmatrix} \right| dx - \int (|v|^2 + e(t)) dx \right\|_0 \leq \varepsilon$$

$$\int |v|^2 dx = \int \sum_I V_I \bar{V}_I + \sum_{J \neq \bar{I}} V_J \cdot V_I$$

Since  $V_I = e^{i\lambda\xi_I}(V_I + \delta V_I) = e^{i\lambda\xi_I} \tilde{V}_I$ ,

$$\sum_{J \neq \bar{I}} \int e^{i\lambda(\xi_I + \xi_J)} \tilde{V}_I \cdot \tilde{V}_J dx \quad \text{not exactly } \tilde{V}_I \cdot \tilde{V}_J$$

$$\int \nabla_a e^{i\lambda(\xi_a + \xi_J)} \frac{\nabla^a \xi_I + \nabla^a \xi_J}{|\nabla^a \xi_I + \nabla^a \xi_J|^2} \tilde{V}_I \tilde{V}_J dx$$

They are oscillations in orthogonal direction, thus integral by parts gives cancellation. Set  $(v, p, R) = (0, 0, 0)$  Let  $e^{1/2}(t)$  be a function smooth and supported in a simple connected area, and apply lemma 4.6 with  $\varepsilon_{(1)} = e_R(1) = \frac{1}{10k} \|e^{1/2}\|_0^2$ .

$$\dots \leq \varepsilon_{(2)} \leq \varepsilon_{(1)} \leq \varepsilon_{(0)}.$$

By choosing the lifting function to have larger  $L^\infty$  norm and smaller support. Then  $\|R_{(k)}\| \leq e_{R,(k)} \leq \varepsilon_{(k)} \searrow 0$  rapidly.  $\|v_{(k)}\|_0 \leq C e_{R,(k)}^{1/2} \searrow 0$  rapidly.  $\|P_{(k)}\| \leq C e_{R,(k)} \searrow 0$  rapidly. Thus  $\sum_k v_{(k)}^\ell$  converges to some solution  $v^\ell$  and  $\sum_k P_{(k)}$  converge to some  $P$ . And since  $\|R_{(k)}\|_0 \searrow 0$   $(v, p)$  is a weak solution to Euler

$$\begin{aligned} \partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p &= \nabla_j R^{j\ell} \\ \nabla_j v^j &= 0 \end{aligned} \tag{12}$$

$$\int |v_k|^2 - |v_{(k-1)}|^2(t, x) dx \geq \int e_k(t) dx - e(0)$$

Choose  $e_{(k)} < \frac{1}{2} \int e_k(0) dx$

$$\int |v_k|^2 - |v_{(k-1)}|^2(0, x) dx \geq \frac{1}{2} \int e_k(0) dx \geq 0,$$

by assumption. Thus the continuous solution we found was nontrivial and compactly supported.  $\square$

We have a To-Do list: For non-stationary phase

$$\nabla_j Q^{j\ell} = u_j^\ell = e^{i\lambda\xi} u_j^\ell \quad \left\| |\nabla\xi|^{-1} \right\|_0 \leq A, \quad \int u^\ell dx = 0$$

$$\exists Q^{j\ell} = Q^{\ell j}, \quad \|Q\|_0 \leq \lambda^{-1} \|u\|_0 + \mathcal{O}_{a, \nabla\xi}(\lambda^{-2})$$

1. How to solve  $\sum_I V_I^j \bar{V}_I^\ell = e(t) \delta^{j\ell} / d - R^{oj\ell}$
2. How to ensure  $\left\| |\nabla\xi|^{-1} \right\|_0, \left\| (\nabla\xi_I \cdot \nabla\xi_J)^{-1} \right\|_0 \leq A$  with  $J \neq \bar{I}$  that interact and  $\left\| |\nabla\xi_I| - 1 \right\|_0 \lesssim b$  small.

Suppose  $u^\ell = e^{i\lambda\xi} u^\ell \left\| |\nabla\xi|^{-1} \right\|_0 \leq A \int u^\ell dx = 0$ .

(Pair symmetric) From a solution of the form

$$\mathcal{Q}^{j\ell} = \frac{e^{i\lambda\xi}}{i\lambda} q_{(1)}^{j\ell} + \tilde{Q}_{(1)}^{j\ell}$$

Take  $\nabla$ ,

$$\nabla_j \mathcal{Q}^{j\ell} = e^{j\lambda\xi} \nabla_j \xi q_{(1)}^{j\ell} + \frac{e^{i\lambda\xi} \nabla_j q_{(1)}^{j\ell}}{i\lambda} + \nabla_j \tilde{Q}_{(1)}^{j\ell}.$$

Choose  $q_{(1)}^{j\ell}$  s.t.

$$\nabla_j \xi q_{(1)}^{j\ell} = u^\ell$$

Choose correction  $\tilde{Q}_{(1)}^{j\ell}$  s.t.

$$\nabla_j \tilde{Q}_{(1)}^{j\ell} = e^{i\lambda\xi} \left( \frac{-\nabla q_{(1)}^{j\ell}}{i\lambda} \right) = \nabla_j \left( \frac{e^{i\lambda\xi} q_{(1)}^{j\ell}}{i\lambda} \right) - e^{i\lambda\xi} u^\ell - u^\ell, \text{ with integral 0}$$

How to solve  $\nabla_j \xi q_{(1)}^{j\ell} = u^\ell$ . If we did not need  $q$  be symmetric, i.e.  $q_{(1)}^{j\ell} = q_{(1)}^{\ell j}$  could use  $q_{(1)}^{j\ell} = \frac{\nabla^j \xi u^\ell}{|\nabla \xi|^2}$ .

Decomposing  $u^\ell = u_\perp^\ell + u_\parallel^\ell = u_\perp^\ell + \frac{(u \cdot \nabla \xi) \cdot \nabla^\ell \xi}{|\nabla \xi|^2}$ .

Choose  $q_{(1)}^{j\ell} = \frac{1}{|\nabla \xi|^2} (\nabla^j \xi u_\perp^\ell + \nabla^\ell \xi u_\perp^j) + \frac{(u \cdot \nabla \xi)}{|\nabla \xi|^2} \delta^{j\ell}$ . Check that  $\nabla_j \xi q_{(1)}^{j\ell} = u_\perp^\ell + 0 + u_\parallel^\ell = u^\ell$ . This equation is under-determinate.

$$q_{(1)}^{j\ell} = q_a^{(jl)} (\nabla \xi) u^a$$

$$q_a^{j\ell}(\alpha p) = \alpha^{-1} q^{j\ell}(p) \quad \text{Homogeneous}$$

As for solving  $\nabla_j \tilde{Q}_{(1)}^{j\ell} = e^{i\lambda\xi} \frac{\nabla_j q_{(1)}^{j\ell}}{i\lambda}$ ,  $\|\tilde{Q}\|_0 \lesssim \frac{1}{\lambda}$ .  $\nabla_j \tilde{Q}_{(1)}^{j\ell} = e^{i\lambda\xi} u_{(2)}^\ell$ . By induction,  $\tilde{Q}^{j\ell} = \frac{e^{i\lambda\xi}}{i\lambda} q_{(2)}^{j\ell} + \tilde{Q}_{(2)}^{j\ell}$ ,  $q_{(2)}^{j\ell} = q_{(a)}^{j\ell} (\nabla \xi) u_{(2)}^a$

$$\nabla_j \tilde{Q}_{(1)}^{j\ell} = -e^{i\lambda\xi} \frac{\nabla_j q_{(2)}^{j\ell}}{(i\lambda)}$$

gives  $\mathcal{O}(\lambda^{-2})$  at the cost of one more derivative.

If we didn't require that  $\tilde{Q}^{j\ell} = \tilde{Q}^{\ell j}$ ,  $\nabla_j \tilde{Q}_{(2)}^{j\ell} = u_{(2)}^\ell$ ,  $\tilde{Q}_{(2)}^{j\ell} = \mathbb{P}_{I_0} u_{(2)}^{\ell 2}$ . Instead we decompose  $u_{(2)}^\ell = \mathcal{H}u^\ell + \nabla^\ell (\text{divergence free} + \nabla^\ell \nabla^{-1} \nabla_a u^a \text{gradient part, so } \nabla_\ell \mathcal{H}u^\ell = 0$ .

Set

$$\tilde{Q}_{(2)}^{j\ell} = \Delta^{-1} (\nabla^j \mathcal{H}u^\ell + \nabla^\ell \mathcal{H}u^j) + \Delta^{-1} \nabla_a u^a \delta^{j\ell} := R_a^{j\ell} [u_{(2)}^a]$$

$$\|\tilde{Q}_{(2)}\|_0 \lesssim \|u_{(2)}\|_0.$$

Since  $\Delta^{-1} \nabla$  is an order  $-1$  operator, it's bounded on  $C^0$ . Check  $\nabla_j \tilde{Q}_{(2)}^{j\ell} = \mathbb{P}_{I_0} \mathcal{H}u_{(2)}^\ell + 0 + \nabla^\ell \Delta^{-1} \nabla_a u_{(2)}^a = \mathbb{P}_{I_0} u_{(2)}^\ell = u_{(2)}^\ell$ .

To summarize  $Q^{j\ell} = \sum_{k=1}^2 \frac{e^{i\lambda\xi}}{i\lambda} q_{(k)}^{j\ell} + \tilde{Q}_{(2)}^{j\ell}$ ,  $q_{(k)}^{j\ell} = q_\alpha^{j\ell} (\nabla \xi) u^a$ ,  $q_\alpha^{j\ell}(\alpha p) = \alpha^{-1} q_\alpha^{j\ell}(p)$ .  $\nabla_j \xi q_{(k)}^{j\ell} = u_{(k)}^\ell$ ,  $u_{(1)}^\ell = -u_{(1)}^\ell$ ,

$$u_{(k+1)}^\ell = \frac{-\nabla_j q_{(k)}^{j\ell}}{(i\lambda)}$$

*Proof.* Error term

$$\begin{aligned} & \nabla_j (V^j V^\ell + p \delta^{j\ell} + R^{j\ell}) \\ &= \nabla_j (\sum_I V_I^j V_I^\ell + p \delta^{j\ell} + R^{j\ell}) \\ &= \nabla_j (\sum_I \tilde{V}_I^j \tilde{V}_I^\ell + p \delta^{j\ell} + R^{j\ell}) \\ &= \nabla_j (\sum_I V_I^j V_I^\ell + p \delta^{j\ell} + R^{j\ell} + \mathcal{O}(\delta V_I V_I) + \dots) \\ &= \nabla_j (\underbrace{(\sum_I V_I^j V_I^\ell + p \delta^{j\ell} + R^{j\ell})}_{\text{small because low frequency cancels.}} + \underbrace{(R^{j\ell} - R_\varepsilon^{j\ell})}_{\text{small with small } \varepsilon} + \mathcal{O}(\delta V_I \bar{V}_I)) \end{aligned}$$

□

$$\sum W_I \bar{V}_I^\ell + p_0 \delta^{j\ell} + R_\varepsilon^{j\ell} = 0$$

$$\partial_t V^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p = \nabla_j R^{j\ell}$$

$$\partial_t v^\ell + \nabla_j (v^j v^\ell) + \nabla^\ell p = \nabla_j E^{j\ell} + \text{linear term}(v) + \nabla_j (v v^\ell) + \nabla^\ell p$$

---

<sup>2</sup> $I_0$  means integral 0.

Since  $V_I^\ell = e^{i\lambda\xi_I}(V_I^\ell + \delta V_I^\ell)$ ,

$$\nabla_j(\sum W_I \bar{V}_I^\ell + p_0 \delta^{j\ell} + R_\varepsilon^{j\ell}) = \sum_I V_I^j \bar{V}_I^\ell + p_0 \delta^{j\ell} + R_\varepsilon^{j\ell} - (R^\ell - R_\varepsilon^{j\ell}) + \mathcal{O}(\delta V_I)$$

$$\sum_I V_I^j \bar{V}_I^\ell e(t) \frac{\delta^{j\ell}}{d} - \overset{o}{R}_\varepsilon^{j\ell} \quad (\text{Must be positive definite}) \quad (13)$$

Because  $p_0 = \frac{e(t)}{d} - \frac{\text{tr } R_\varepsilon}{d}$

$$\int \left| \overset{o}{V} \right|^2 - |v|^2 dx = \int e(t) dx + \mathcal{O}(\lambda^{-1})$$

$e(t) \geq K e_R$  on  $\text{supp } R_\varepsilon$ . We make a choice of  $K$  later. Then we localize and renormalize it by choosing,

$$\eta_{K_0}(t) = \eta\left(\frac{t - K_0\tau}{\tau}\right),$$

$\tau$  is the time cut off and  $\sum_{K_0} \eta_{K_0}^2(t) = 1$ . (Start with  $\tilde{\eta}(t) \geq 1$  on  $[-\frac{2}{3}, \frac{2}{3}]$  and  $\tilde{\eta}(t) \geq 0$  and  $\in C_c^\infty((-\frac{3}{4}, \frac{3}{4}))$ ).

$\eta := \frac{\tilde{\eta}(t)}{(\sum_{K \in \mathbb{Z}} \tilde{\eta}(t - K_0)^2)^{1/2}}$ . Thus  $\eta^2 := \frac{\tilde{\eta}(t)^2}{(\sum_{K \in \mathbb{Z}} \tilde{\eta}(t - K_0)^2)} = 1$  gives us a partition of unity. Let's localize (13) to have

$$\text{RHS} = \sum_{K_0} \eta_{K_0}^2(t) (t(t) \frac{\delta^{j\ell}}{d} - R_\varepsilon^{j\ell})$$

Let's write  $V_I^\ell = \eta_{K_0} V_I^{\circ\ell}$  and  $I \in \mathbb{I}(K_0)$  if  $V_I$  is positive at time  $K_0\tau$ . so

$$\text{LHS} = \sum_{K_0} \eta_{K_0}^2 \sum_{I \in \mathbb{I}(K_0)} \overset{o}{V}^j \bar{\overset{o}{V}}_I^\ell.$$

So it's sufficient to solve

$$\sum_{I \in \mathbb{I}(K_0)} \overset{o}{V}^j \bar{\overset{o}{V}}_I^\ell = e(t) \frac{\delta^{j\ell}}{d} - \overset{o}{R}_\varepsilon^{j\ell}$$

on every time slice.  $\eta_{K_0}$  only lives for  $t \sim \tau$ .

Cut off the space.

I need  $|\nabla \hat{\xi}_I| = 1$  at initial time.

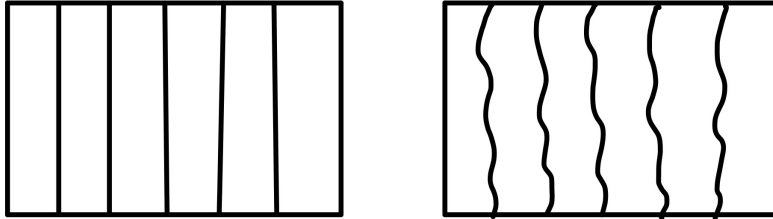


Figure 3: Initial and later phase function

At  $t = 0$  Assume  $\begin{cases} \nabla \hat{\xi}_I \in \{e_1, e_2, e_3\} \\ \hat{\xi}_I(t(I), x) = e_1 \cdot x \end{cases}$  is smooth enough.  $e^{i\lambda\xi(t,x)}$   $\lambda \in 2\pi\mathbb{Z}$  and  $e^{i\lambda(\xi_I + \xi_J)}$  with  $\left\| |\nabla \xi_I + \nabla x i_j|^{-1} \right\|_{C^0} \leq$

A. Let  $\Psi_I$  be spacial cut off. Write  $v_I^\ell = \eta_{k_0} \hat{v}_I^\ell = \eta_{k_0} \Psi_I \hat{v}_I^\ell$ , where  $\eta_{k_0} \hat{v}_I^\ell$  is the active part. Also  $\sum_{I \in \mathcal{I}(k_0)} = 1$ .

---

<sup>3</sup>trace-free

The space cut off is depending on time  $\sum_{I \in \mathcal{I}(k_0)} \dot{v}_I^j \bar{v}_I^\ell = e(\xi) \frac{\delta^{j\ell}}{d} - \dot{R}_\varepsilon^{j\ell}$  on time slice.

$$\sum_{I \in \mathcal{I}(k_0)} \Psi_I^2 \dot{v}_I^j \bar{v}_I^\ell = \sum_{I \in \mathcal{I}(k_0)} \Psi_I^2 (e(t) \frac{\delta^{j\ell}}{d} - \dot{R}_\varepsilon^{j\ell})$$

$$\Psi_I(t(I), x) = \Psi(k_0 \tau, x)$$

$$\sum_{K \in (\mathbb{Z}/2\mathbb{Z})^3} \Psi(k_0 t, x) = 1$$

Partition of unit is also transported

$$(\partial_t + v \cdot \nabla) \Psi_I(t, x) = 0$$

$$\sum_{K \in (\mathbb{Z}/2\mathbb{Z})^3} \Psi^2 = 1$$

by uniqueness.

$$\text{Apply the cut off } \Psi_I^2 \left( \sum_{I \in \mathbb{I}(K)} \dot{v}_I^j \bar{v}_I^\ell \right) = \Psi_I^2 \left( e(t) \frac{\delta^{j\ell}}{d} - \dot{R}_\varepsilon^{j\ell} \right)$$

$$v_I^\ell = \eta_{t_0}(t) \Psi_I(t, x) \dot{v}_I^\ell$$

where  $I = (K, f) = (\text{time, direction of oscillation})$ ,  $K \in \mathbb{Z} \times (\mathbb{Z}/2\mathbb{Z})^3$  location index. Reason that we can use partition of unity: Homogeneous in amplitude.

$$v_I^\ell = \eta_{K_0} \Psi_I e^{1/2}(t) \dot{v}_I^\ell$$

LHS  $\sum_I e(t) \dot{v}_I^j \bar{v}_I^\ell = e(t) \frac{\delta^{j\ell}}{d} - \dot{R}_\varepsilon^{j\ell}$ . On  $\text{supp } R_\varepsilon$ ,  $e(t) > K e_R > 0$

$$\sum_I \dot{v}_I^j \bar{v}_I^\ell = \frac{\delta^{j\ell}}{d} - \frac{\dot{R}_\varepsilon^{j\ell}}{e(t)} = \frac{\delta^{j\ell}}{d} + \varepsilon^{j\ell}$$

where  $\frac{\delta^{j\ell}}{d}$  is the dominated term.

$$\sum_I \dot{v}_I^j \bar{v}_I^\ell = \frac{\delta^{j\ell}}{d} + \varepsilon^{j\ell} \quad \varepsilon = \mathcal{O}\left(\frac{1}{K}\right)$$

$$\sum_I \dot{v}_I^j \bar{v}_I^\ell = \frac{\delta^{j\ell}}{d} - \varepsilon^{j\ell} \tag{14}$$

What we know about  $v_I$ ?

1.  $v_I^\ell \cdot \xi_I = 0$
2.  $(i \nabla \xi_I) \times V_I = |\nabla \xi_I| v_I$

$$\dot{b}_I \in \langle \nabla \xi_I \rangle^\perp$$

$$\dot{a}_I = -\frac{\nabla \xi_I}{|\xi_I|} \times b_I \Leftarrow \text{rotation in plane } \langle \nabla \xi \rangle^\perp$$

Thus (14) =  $\sum_I (\dot{a}_I + i \dot{b}_I^j) (\dot{a}_I^\ell - i \dot{b}_I^\ell) = 2 \sum_I (a_I^j \dot{a}_I^\ell + \dot{b}_I^j b_I^\ell) = \sum (\delta^{j\ell} - \frac{\nabla^j \xi_I \nabla^\ell \xi_I}{|\nabla \xi_I|^2} \cdot \left| \dot{b}_I \right|^2)$ , is true because the imaginary part cancels.  $(\frac{\dot{a}_I}{|\dot{b}_I|}, \frac{\dot{b}_I}{|\dot{b}_I|}, \frac{\nabla^i \xi_I}{|\nabla \xi_I|})$  are orthogonal frame. Thus the identity

$$\delta^{j\ell} = \frac{\dot{a}_I^j \dot{a}_I^\ell}{\left| \dot{b}_I \right|^2} + \frac{\dot{b}_I^j \dot{b}_I^\ell}{\left| \dot{b}_I \right|^2} + \frac{\nabla^j \xi_I \nabla^\ell \xi_I}{|\nabla \xi_I|^2} \tag{15}$$

by renormalization. Choose  $\mathring{b}_I^\ell = \gamma_I \mathbb{P}_I^\perp(\nabla \xi_{\sigma_I}) = \gamma_I(\nabla^\ell \xi_I - \frac{(\nabla \xi_{\sigma_I} \nabla \xi_I) \nabla^\ell \xi_I}{|\nabla \xi_I|^2})$ , since  $\nabla \xi_{\sigma_I}$  is not parallel to  $\nabla \xi_{\sigma_I}$ .

$$\begin{aligned}
|\mathring{b}_I|^2 &= \gamma_I^2 \left( |\nabla \xi_{\sigma_I}|^2 - 2 \frac{(\nabla \xi_{\sigma_I} \cdot \nabla \xi_I)^2}{|\nabla \xi_I|^2} + \frac{(\nabla \xi_{\sigma_I} \cdot \nabla \xi_I)^2}{|\nabla \xi_I|^2} \right) \\
&= \gamma_I^2 \left( \frac{|\nabla \xi_{\sigma_I}|^2 |\nabla \xi_I|^2 - (\nabla \xi_{\sigma_I} \cdot \nabla \xi_I)^2}{|\nabla \xi_I|^2} \right) \\
&= \gamma_I^2 \left( \frac{|\nabla \xi_{\sigma_I} \wedge \nabla \xi_{\sigma_I}|^2}{|\nabla \xi_I|^2} \right) \\
\sum_I \gamma_I^2 \frac{|\nabla \xi_{\sigma_I} \wedge \nabla \xi_I|^2}{|\nabla \xi_I|^2} \left( \delta^{j\ell} - \frac{\nabla^j \xi_I \nabla^\ell \xi_I}{|\nabla \xi_I|^2} \right) &= \frac{\delta^{j\ell}}{d} + \varepsilon^{j\ell}
\end{aligned} \tag{16}$$

where  $(\nabla \xi_I^\perp)^{j\ell}$ .

$$\sum \gamma_I^2 \frac{|\nabla \xi_{\sigma_I} \wedge \nabla \xi_{\sigma_I}|^2}{|\nabla \xi_I|^2} (\nabla \xi_I^\perp)^{j\ell} \nabla_J \xi_J \nabla_\ell \xi_J = \frac{|\nabla \xi_J|^2}{d} + \varepsilon^{j\ell} \nabla_J \xi_J \nabla_\ell \xi_J$$

Plug (16), (15) back in

$$\sum_I \gamma_I^2 A(\nabla \xi)_J^I = \frac{|\nabla \xi_J|^2}{d} + \varepsilon^{j\ell} \nabla_J \xi_J \nabla_\ell \xi_J$$

We need 12 waves in 6 direction and their conjugate. Letting  $I = (k, f) \in \mathbb{Z} \times (\mathbb{Z}/\mathbb{Z}\tau)^3 \times \mathbb{F}$ ,  $\mathbb{F}$  are the faces of regular dodecahedron.

$$\sum_{I \in K \times \mathbb{F}} \gamma_I^2 A(\nabla \xi_J^I) = \frac{|\nabla \xi_J|^2}{2d} + \varepsilon^{j\ell} \frac{\nabla_J \xi_J \nabla_\ell \xi_J}{2}$$

$\exists c > 0$ , s.t.  $\|A_J^T - A(\nabla \hat{x}_J^I)\| \leq C$ ,  $\left\| \xi_J - \frac{|\nabla \xi_J|^2}{2d} \right\| \leq C$ . The equation  $\sum_{I \in \mathbb{F}} A_J^I \gamma_I^2 = y_I$  has a solution  $\gamma_I = \gamma_f(A_J^I, y_I)$  depending smoothly with uniform bounds on  $\partial_A \gamma_f$  and  $\gamma_y \gamma_f$  and higher derivative.

$$\sum_I A(\nabla \xi)_J^I \gamma_I^2 = \frac{|\nabla \xi_J|^2}{2d} + \varepsilon^{j\ell} \nabla_J \xi_J \nabla_\ell \xi_J$$

where  $\varepsilon^{j\ell} \sim \mathcal{O}(\frac{1}{k})$ .

*Proof. (inverse function theorem)* Rewrite as  $F(\gamma, A, y) = 0$ ,  $F_J(\gamma, A, y) = \sum_I A_J^I \gamma_I^2 - y_J$ . It is sufficient to check that  $\frac{\partial F_J}{\partial \gamma_\gamma}(\gamma)$  is invertible.

$$\gamma_I, A_J^T, y_J) = (\hat{\gamma}_f, A(\nabla \hat{\xi})_J^I, \frac{|\nabla \xi_J|^2}{2d})$$

If  $h_I$  is in the null space at this point

$$\frac{\partial f_J}{\partial \gamma} H_I = 2A(\nabla \hat{\xi})_J^I \hat{\gamma}_I h_I = 0 \quad \Rightarrow \quad \hat{\gamma}_I h_I = 0 \quad \forall I$$

since  $A(\nabla \hat{\xi})^I$  is invertible.  $\Rightarrow h_I = 0 \quad \forall I$  since  $\hat{\gamma}_I \neq 0 \quad \forall I$ . □

$$\begin{aligned}
V_I &= e^{i\lambda \xi_I} (v_I^\ell + \delta v_I^\ell) \\
I_M(V_I) &= \eta_{K_0}(t) \Psi_I e^{1/2}(t) (\gamma_I P_I^\perp(\delta \xi_{\sigma_I})) \\
R_\ell v_I &= -\frac{\nabla \xi_I}{|\nabla \xi_I|} \times I_m V_I \\
\|V_I\|_0 &\leq \|\eta\| \|\Psi\| \left\| e^{1/2} \right\| \|\gamma_I\| \|\nabla \xi_{\sigma_I}\|
\end{aligned}$$

$$\sum_{K_0} \eta_{K_0}^2(t) = 1 \quad e(t) \geq K e_R$$

$$\sum_I \gamma_I^2 \frac{|\nabla \xi_I \wedge \nabla \xi_{\sigma_I}|^2}{|\nabla \xi_I|^2} \left( \delta^{j\ell} - \frac{\nabla^j \xi_I \nabla^\ell \xi_I}{|\nabla \xi_I|^2} \right) = \frac{\delta^{j\ell}}{d} + \varepsilon^{j\ell}$$

where  $\varepsilon^{j\ell} = \frac{-\mathring{R}_\varepsilon^{j\ell}}{e(t)}$ .

## Introduce Frequency Energy levels

Let  $\Theta \geq 2$   $e_v > e_R > 0$ ,  $L \geq 1$ . Say  $(v, p, R) \leq (\Theta, e_v, e_R)$  if

$$\|\nabla_{\vec{a}} v\|_0 \leq \Theta^{|\vec{a}|} e_v^{1/2} \quad 1 \leq |\vec{a}| \leq L$$

$$\|\nabla_{\vec{a}} R\|_0 \leq \Theta^{|\vec{a}|} e_R$$

we always begin with Renyolds flow from previous stage.  $\|v\|_0$  close to/converge to real solutions.

$$\|\nabla \check{V}\|_0 = \|\nabla V\|_0 + \|\nabla V\|_0$$

The correction  $\|V\|_0 \lesssim e_R^{1/2}$ ,  $\|\nabla v\|_0 \lesssim \lambda e^{1/2}$  from the previous stage.  $\|\nabla v\| \rightarrow \Theta e_v^{1/2}$ . Taking derivative cost a factor of  $\Theta$ .

**Lemma 4.7.** (Practice main lemma)  $\exists \hat{c} \geq 1$  such that the following holds. If  $(v, p, R) \leq (\Theta, e_v, e_R)$  if  $N \geq (\frac{e_v}{e_R}) \exists (\check{v}, \check{p}, \check{R}) \leq (\check{\Theta}, \check{e}_v, \check{e}_R)$ . Here we want  $\check{e}_R$  as small as possible and  $\check{v} = v + V$   $\check{p} + P$

$$\|v\|_0 \leq \hat{c} e_R^{1/2} \quad \|p\|_0 \leq c e_R$$

$$\|\nabla v\| \leq c(N\Theta) e_R^{1/2} \quad \|\nabla P\|_0 \leq c(N\Theta) e_R$$

How big is  $\check{e}_R$ ?

1. High-Low interaction term

$$\nabla_j R_L^{j\ell} = \nabla_j (v^j v^\ell) = v^j \nabla_j v^\ell = \Theta_I e^{i\lambda \xi_I} \tilde{v}_I^j \nabla_j v^\ell$$

Non-stationary phase

$$\|R_L\|_0 \lesssim \frac{1}{\lambda} \|v\|_0 \|\nabla v\|_0 \leq (N\Theta)^{-1} e_R^{1/2} (\Theta e_v^{1/2}) = \frac{e_v^{1/2} e_R^{1/2}}{N}$$

We hope that  $\check{e}_R = \frac{e_v^{1/2} e_R^{1/2}}{N}$

2. transport  $\nabla_j R_T^{j\ell} = \partial_t v^\ell + \nabla_j (v^j v^\ell) = (\partial_t + v^j \nabla_j) (v^j v^\ell) = (\partial_t + v^j \nabla_j) v^\ell$

$$\|R_T\|_0 \leq \lambda^{-1} \|v\|_0 \|\nabla v\|_0 \leq (N\Theta)^{-1} \cdot 1 \cdot \Theta e_R^{1/2} = N^{-1} e^{1/2}$$

where  $R_T^{j\ell} = \sum_I \frac{e^{i\lambda \xi_I}}{i\lambda} q_a^{j\ell} (\nabla \xi_I) D_t V_1$ . In the result we missing  $e_v^{1/2}$  which make this even worse to  $\frac{1}{3+\sqrt{8}}$ .

To optimizing, we keep the transport term.

$$\|R_T\|_0 \lesssim \lambda^{-1} \|(\partial_t + v^j \nabla_j) V_I\|_0 = (N\Theta)^{-1} (\|D_t\|) e_R^{1/2}$$

$$(\partial_t + v^i \nabla_i) \xi_I = 0$$

$$(\partial_t + v^i \nabla_i) \nabla_a \xi_I = -\nabla_a v^i \nabla_i \xi \sim \Theta e_v^{1/2} \cdot 1$$

The result  $\|R_T\|_0 = \frac{e_v^{1/2} e_R^{1/2}}{N}$  is ideal

$$\sum_I v_I^j v_I^\ell = e(t) \frac{\delta^{j\ell}}{d} - \tilde{R}_\varepsilon^{j\ell}$$

And more for frequency level:

$$\|\nabla_{\vec{a}}(\partial_t + v\nabla)R\|_0 \leq \Theta^{|\vec{a}|+1} e_v^{1/2} e_R \quad 0 \leq |\vec{a}| \leq L-1$$

better for space derivative.

## Actual Frequency Levels

Note that  $e_R, e_v \sim R, p \sim \frac{m^2}{s^2}$ , and  $\Theta \sim \nabla \sim m^{-1} \partial_t \sim s^{-1} v \sim \frac{m}{s}$  and  $(\Theta e_v^{1/2}) \sim (\partial_t + v \cdot \nabla) \sim s^{-1}$ . No assumed bound for  $\|v\|_0$ . The bounds are invariant under

$$\begin{cases} v^\ell & c^\ell + v^\ell(t, x - tc) \\ p(t, x) & \Rightarrow p(t, x - tc) \\ R^{j\ell}(t, x) & R^{j\ell}(t, x - tc) \end{cases}$$

Thus well-posed.

$$\begin{aligned} \nabla_j R_L^{j\ell} &= v^j \nabla_j v^\ell = \sum_I e^{i\lambda \xi_I} \tilde{v}_I^j \nabla_j v^\ell = \sum_I e^{i\lambda \xi_I} u_I^\ell \\ R^{j\ell} &= \sum_I \frac{e^{i\lambda \xi_I}}{i\lambda} q_a^{j\ell}(\nabla \xi_I) u_I^a + L.O.T \\ \overline{D}_t R_T^{j\ell} &= \frac{e^{i\lambda \xi_I}}{i\lambda} D_t(q_a^{j\ell}(\nabla \xi_I) u_I^a) \end{aligned}$$

seems not give a  $\lambda$ .

*Remark.*

$$\begin{aligned} \overline{D}_t R_L^{j\ell} &= \frac{e^{i\lambda \xi_I}}{i\lambda} D_t(q_a^{j\ell}(\nabla \xi_I) \tilde{v}_I^b \nabla_b v^a) \\ D_t \nabla_b v^a &= (\partial_t + v^j \nabla_j) \nabla_b v^a \end{aligned}$$

If we use Euler-Reynolds equation (12)

$$(\partial_t + v^j \nabla_j) \nabla_b v^\ell = -\nabla_b v^j \nabla_j v^\ell - \nabla_b \nabla^\ell P - \nabla_b \nabla_j R^{j\ell}$$

where  $|D_t| \leq \Theta e_v^{1/2}$ . The bound that I want is  $(\Theta e^{1/2})(\Theta e^{1/2})$ .

$$\tilde{D}_t = \partial_t + v \cdot \nabla + V \cdot \nabla$$

where  $V$  is small and there hidden a  $\lambda$  in the last  $\nabla$  from  $e^{i\lambda \xi_I}$ .

## Frequency energy levels

We say  $(v, p, R)$  have energy levels below  $(\Theta, e_0, e_R)$  to order  $L$ .  $(v, p, R) \leq (\Theta, e_v, e_R)$  is  $\|\nabla_{\vec{a}} v\|_0 \leq \Theta^{|\vec{a}|} e^{1/2}$ ,  $1 \leq |\vec{a}| \leq L$ .  $\|\nabla_{\vec{a}} p\|_0 \leq \Theta^{|\vec{a}|} e_v$  and  $\|\nabla_{\vec{a}} R\|_0 \leq \Theta^{|\vec{a}|} e_R$ .  $\|\nabla_{\vec{a}}(\partial_t + v\nabla)R\|_0 \leq \Theta^{|\vec{a}|} (\Theta e^{1/2}) e_R$ .  $0 \leq |\vec{a}| \leq L$  and on the RHS we need  $0 \leq |\vec{a}| \leq L-1$

**Lemma 4.8.** *If  $L \geq 2$   $\eta > 0$  and  $M > 0$   $K > 0$  is as determined previously  $\exists \hat{c} = \hat{c}(\eta, L, M)$  s.t, the following holds: Let  $(v, p, R) \leq (\Theta, e_v, e_R)$  to order  $L$ .  $\text{supp } R \subset I \cdot \mathbb{T}^3$ .*

*Let  $e(t) : \mathbb{R} \rightarrow \mathbb{R}_+$  satisfying  $e(t) \geq K e_R \forall t \in N(I; (\Theta, e^{1/2})^{-1})$  and for lifting function*

$$\left\| \frac{d^r}{dt^r} e^{1/2} \right\|_0 \leq M (\Theta e_0^{1/2})^r e_R^{1/2}$$

*Let  $N \geq 0$   $N \geq (\frac{e_v}{e_R})$  be given  $N > \Theta^\eta$  where  $\eta$  is small.*

*Then  $\exists (\check{v}, \check{p}, \check{R}) \leq (\check{\Theta}, \check{e}_v, \check{e}_R) = (\check{c} N \Theta, e_R, \frac{e_v^{1/2} e_R^{1/2}}{N})$*



$$\check{v} = v + V \check{p} + P$$

$$\|v\|_0 \leq \hat{c} e_R^{1/2} \quad \|p\|_0 \leq \check{c} e_R$$

$$\|\nabla v\|_0 \leq \check{c}(N\Theta) e_R^{1/2} \quad \|\nabla P\|_0 \leq \check{c}(N\Theta) e_R$$

$$\text{supp}(V, p, R) \subset N(\text{supp } e, (\Xi e^{1/2})^{-1})$$

$$\left\| \int |v|^2 dx - \int e(t) dx \right\| \leq \frac{e_R}{N}$$

$$\text{small error. Letting } V^\ell = \nabla_j w^{j\ell} \quad \|w\|_0 \leq \frac{e_R^{1/2}}{N\Xi}$$

*Proof. (of Onsager's conjecture)* Set  $(v, p, R)_{(0)} = (0, 0, 0)$  initial (12) flow. Let  $\Xi_{(0)} > 1$  (very big)  $e_v \geq e_R(0)$

be given. So we want new error  $e_{R(K+1)} = \frac{e_{R(k)}^{1+\delta}}{z}$   $z$  to be chosen. Force  $\frac{e_{R(k)}^{1/2} e_{R(k)}^{1/2}}{N_{(k)}} : \frac{e_R^{1-\delta}}{z} \Rightarrow N_{(k)} = z \cdot \left(\frac{e_v}{e_R}\right)_K^{1/2} e^{-\delta}$  shrinking exponentially.  $\square$

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