

# 1 Poisson's Equation and Laplace's Equation

## 1.1 Poisson's Equation:

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

- In Cartesian System:

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

- In Cylindrical System:

$$\nabla^2 V = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 V}{\partial \varphi^2} + \frac{\partial^2 V}{\partial z^2}$$

- In Spherical System:

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \varphi^2}$$

## 1.2 Laplace's Equation:

For a simple medium where there is no free charge:

$$\nabla^2 V = 0$$

For problem involving conductors:

- Use Laplace's Equation to obtain electric potential  $V$ .
- Use  $\mathbf{E} = -\nabla V$  to work out  $E$ .
- Use  $\rho_s = \epsilon E$  to get charge density on the conductor surface.

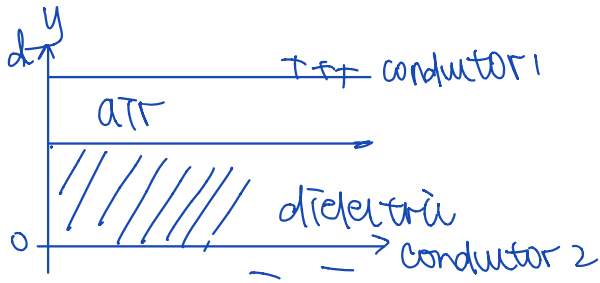
## 1.3 Uniqueness of Electrostatic Solutions

A solution of Poisson's Equation or Laplace's Equation that satisfies the given boundary conditions is a unique solution.

## 1.4 Examples

1. The upper and lower conducting plates of a large parallel-plate capacitor are separated by a distance  $d$  and maintained at potentials  $V_0$  and 0, respectively. A dielectric slab of dielectric constant 6.0 and uniform thickness of  $0.8d$  is placed over the lower plate. Assuming negligible fringing effect, determine
  - a) the potential and electric field distribution in the dielectric slab,
  - b) the potential and electric field distribution in the air space between the dielectric slab and the upper plate,

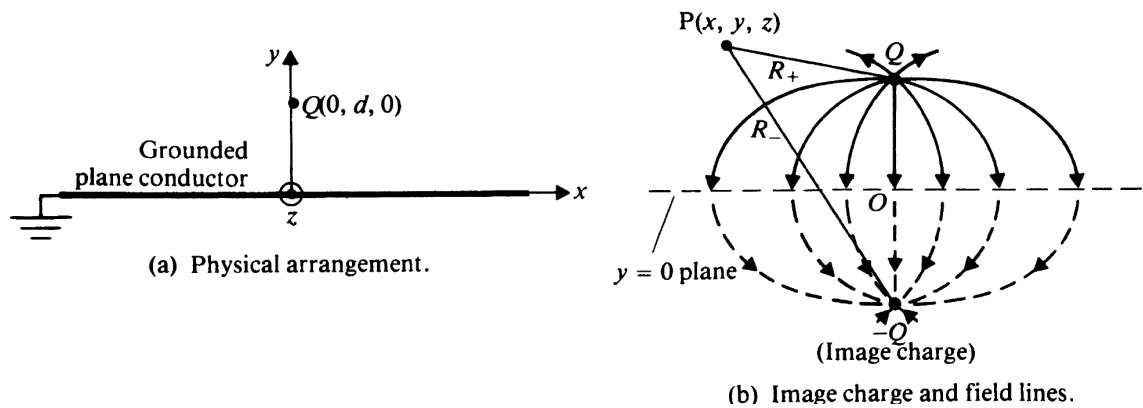
- c) the surface charge densities on the upper and lower plates.
- d) Compare the results in part (b) with those without the dielectric slab.



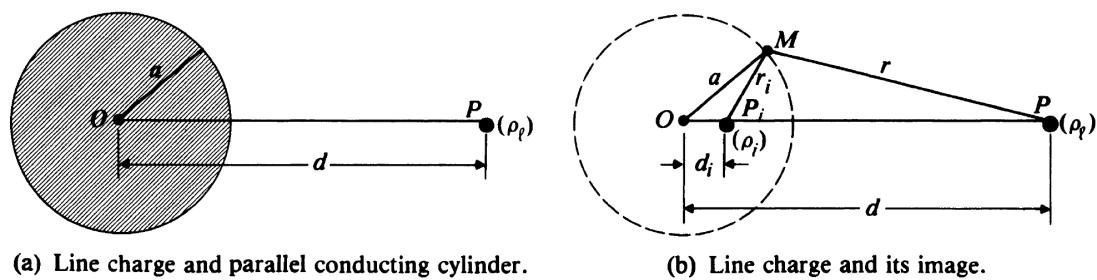
**EXAMPLE 4-2** Determine the  $\mathbf{E}$  field both inside and outside a spherical cloud of electrons with a uniform volume charge density  $\rho = -\rho_0$  (where  $\rho_0$  is a positive quantity) for  $0 \leq R \leq b$  and  $\rho = 0$  for  $R > b$  by solving Poisson's and Laplace's equations for  $V$ .

# 2 Method of Images

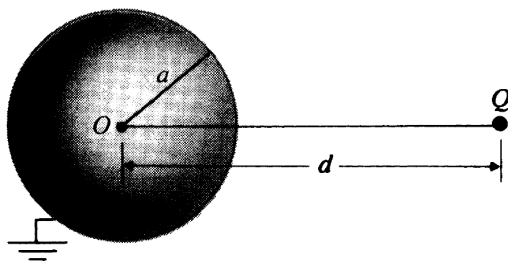
## 2.1 Point Charge and Conducting Planes



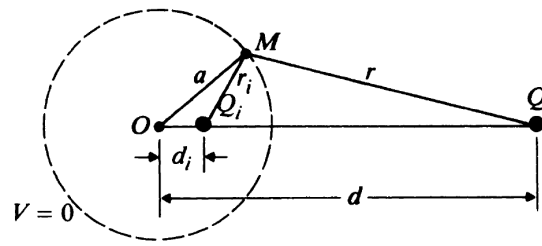
## 2.2 Line Charge and Parallel Conducting Cylinder



## 2.3 Point Charge and a Conducting Sphere

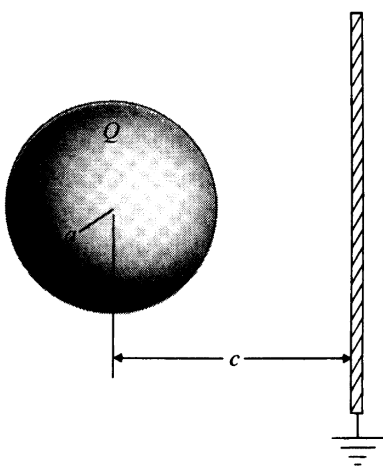


(a) Point charge and grounded conducting sphere.

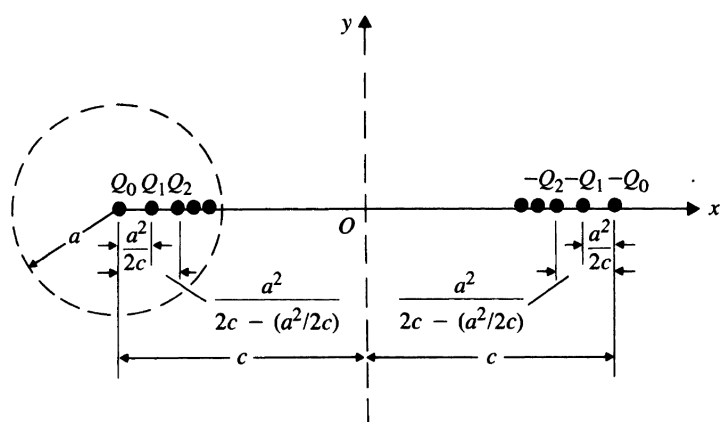


(b) Point charge and its image.

## 2.4 Charge Sphere and Grounded Plane



(a) Physical arrangement.

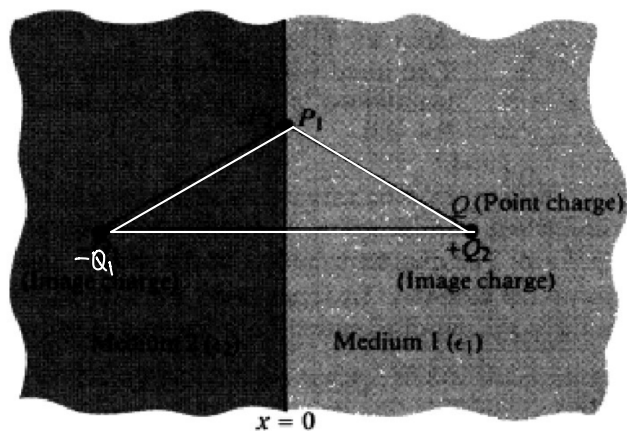


(b) Two groups of image point charges.

# 2.5 Examples

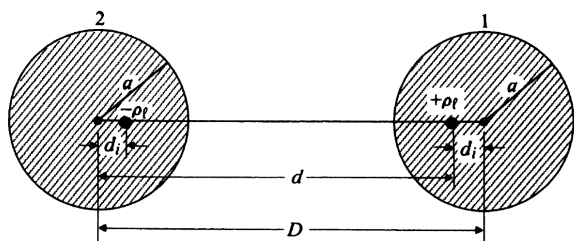
**P.4-17** Two dielectric media with dielectric constants  $\epsilon_1$  and  $\epsilon_2$  are separated by a plane boundary at  $x = 0$ , as shown in Fig.4-23. A point charge  $Q$  exists in medium 1 at distance  $d$  from the boundary.

c) Determine  $Q_1$  and  $Q_2$ .



**FIGURE 4-23**  
Image charges in dielectric media (Problem P.4-17).

- Determine the capacitance per unit length between two long, parallel, circular conducting wires of radius  $a$ . The axes of the wires are separated by a distance  $D$



# 1 Boundary Value problem in Cartesian Coordinates

## 1.1 Boundary Condition Problem

- In order to find specific voltage on conductor systems without isolated free charge.
- **General idea:** Use boundary condition to find coefficients for general solution form from Laplace equation.
- **Types of boundary condition:** (1) Dirichlet:  $V$  is specific; (2) Neumann:  $\frac{dV}{dn}$  is specified on boundaries (3) Mixed:  $V$  is specific on some boundaries;  $\frac{dV}{dn}$  is specified on some boundaries.
- **Solution Form:** Separation of variables, which means  $V(x, y, z) = X(x)Y(y)Z(z)$ . When the potential or normal derivative is specified, and it coincide with coordinate surfaces of an orthogonal, curvilinear coordinate system.

## 1.2 Boundary condition value in Cartesian Coordinate

(1) Laplace's Equation for  $V$  in Cartesian coordinates is

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

(2) To use the Separation of variables and take it into Laplace's Equation.

$$\frac{1}{X(x)} \frac{d^2 X(x)}{dx^2} + \frac{1}{Y(y)} \frac{d^2 Y(y)}{dy^2} + \frac{1}{Z(z)} \frac{d^2 Z(z)}{dz^2} = 0$$

(3) In order to satisfied all  $x, y, z$  values, these three parts should be constant. Then we can get

$$\frac{1}{X(x)} \frac{d^2 X(x)}{dx^2} = -k_x^2, \frac{1}{Y(y)} \frac{d^2 Y(y)}{dy^2} = -k_y^2, \frac{1}{Z(z)} \frac{d^2 Z(z)}{dz^2} = -k_z^2$$
$$k_x^2 + k_y^2 + k_z^2 = 0$$

(4) List the boundary conditions we got.

(5) The general solution formats for above differential equation  $\frac{d^2 X(x)}{dx^2} + k_x^2 X(x) = 0$  are:

$k_x^2$	$k_x$	$X(x)$	Exponential forms <sup>†</sup> of $X(x)$
0	0	$A_0 x + B_0$	
+	$k$	$A_1 \sin kx + B_1 \cos kx$	$C_1 e^{jkx} + D_1 e^{-jkx}$
-	$jk$	$A_2 \sinh kx + B_2 \cosh kx$	$C_2 e^{kx} + D_2 e^{-kx}$

We need to choose the proper form of solution given boundary condition.

If  $V$  is independent of  $x$ , We can see  $X(x)=0$ ;

If  $V$  goes to infinity or 0 as  $x$  goes to infinity, we choose  $k_x^2$  is negative.

(6) Find the coefficients through boundary condition.

## Quiz 2

Two infinite grounded metal plates lie parallel to the  $xz$  plane, one at  $y=0$ , the other at  $y=a$ (Fig.2). The left end, at  $x=0$ , is closed off with an infinite strip insulated from the two plates and maintained at a specific potential  $V_0(y)$ . Find the potential inside this "slot."

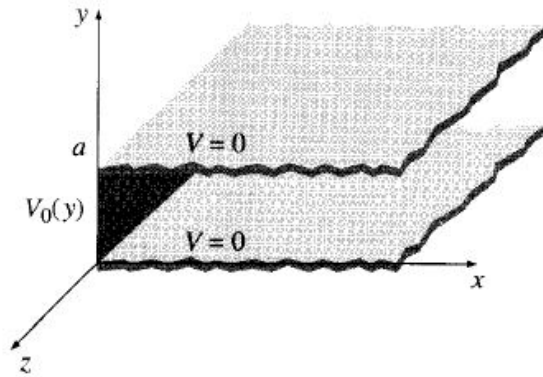


Figure 1: Question 2

### Hint:

1. After using separation of variables you can get

$$\frac{1}{X} \frac{d^2 X}{dx^2} = C_1 \quad \text{and} \quad \frac{1}{Y} \frac{d^2 Y}{dy^2} = C_2, \quad \text{with} \quad C_1 + C_2 = 0$$

This implies

$$\frac{d^2 X}{dx^2} = k^2 X, \quad \frac{d^2 Y}{dy^2} = -k^2 Y$$

2. Mathematical equations you may use:

$$\sum_{n=1}^{\infty} C_n \int_0^a \sin(n\pi y/a) \sin(n'\pi y/a) dy = \int_0^a V_0(y) \sin(n'\pi y/a) dy$$

With the left side:

$$\int_0^a \sin(n\pi y/a) \sin(n'\pi y/a) dy = \begin{cases} 0, & \text{if } n' \neq n \\ \frac{a}{2}, & \text{if } n' = n \end{cases}$$

And then:

$$C_n = \frac{2}{a} \int_0^a V_0(y) \sin(n\pi y/a) dy$$

## 2 Boundary-value Problems in Cylindrical Coordinates

(1) Laplace Equation:

$$\frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

(2) General solution: Assuming  $V$  is independent of  $Z$ .

$$\frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 V}{\partial \phi^2} = 0$$

(3) Separation of variables:  $V(r, \phi) = R(r)\Phi(\phi)$

(4) Equations for  $\Phi(\phi)$

$$\frac{d^2 \Phi(\phi)}{d\phi^2} + k^2 \Phi(\phi) = 0$$

Since the solution should be periodic among  $\phi$ , we can get  $k=n$  and we should use  $k^2 > 0$  be like

$$\Phi(\phi) = A_\phi \sin n\phi + B_\phi \cos n\phi$$

(5) Equations for  $R(r)$ : After using separation of variables, we get

$$\frac{r}{R(r)} \frac{d}{dr} \left[ r \frac{dR(r)}{dr} \right] = k^2$$

Which is a second order differential Equation

$$r^2 \frac{d^2 R(r)}{dr^2} + r \frac{dR(r)}{dr} - n^2 R(r) = 0$$

And the general solution is

$$R(r) = A_r r^n + B_r r^{-n}$$

If we study the area including  $r=0$ ,  $B_r=0$ , Otherwise,  $V$  goes to infinity at  $r=0$

If we study the area including  $r = \infty$ ,  $A_r=0$

(6) Equations for  $V_n(r, \phi)$ ,

$$V_n(r, \phi) = r^n (A_n \sin n\phi + B_n \cos n\phi) + r^{-n} (A'_n \sin n\phi + B'_n \cos n\phi), \quad n \neq 0$$

(7) Special case: if  $V$  is independent of  $\phi$ ,  $k=0$ . Then we get

$$\frac{d}{dr} \left[ r \frac{dR(r)}{dr} \right] = 0$$

$$V(r) = C_1 \ln r + C_2$$



### 3 Boundary-value Problem in Spherical Coordinates

- (1) Since we only consider the situation that  $V$  is independent of  $\phi$ , the Laplace Equation in Spherical coordinates is simplified to

$$\frac{1}{R^2} \frac{\partial}{\partial R} \left( R^2 \frac{\partial V}{\partial R} \right) + \frac{1}{R^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial V}{\partial \theta} \right) = 0$$

- (2) By using separation of variables, we assign  $V(R, \theta) = \Gamma(R)\Theta(\theta)$ . Then it looks like

$$\frac{1}{\Gamma(R)} \frac{d}{dR} \left[ R^2 \frac{d\Gamma(R)}{dR} \right] + \frac{1}{\Theta(\theta) \sin \theta} \frac{d}{d\theta} \left[ \sin \theta \frac{d\Theta(\theta)}{d\theta} \right] = 0$$

- (3) General solutions for  $\Gamma(R)$ . Firstly, we assume the part for  $\Gamma(R)$  equals to  $k^2$ :

$$\frac{1}{\Gamma(R)} \frac{d}{dR} \left[ R^2 \frac{d\Gamma(R)}{dR} \right] = k^2$$

It is actually second differential Equation:

$$R^2 \frac{d^2 \Gamma(R)}{dR^2} + 2R \frac{d\Gamma(R)}{dR} - k^2 \Gamma(R) = 0$$

The solution form is

$$\Gamma_n(R) = A_n R^n + B_n R^{-(n+1)}, \text{ where } k = n(n+1), n > 0$$

- (4) General Solutions for  $\theta$ . Similarly, we can get

$$\frac{1}{\Theta(\theta) \sin \theta} \frac{d}{d\theta} \left[ \sin \theta \frac{d\Theta(\theta)}{d\theta} \right] = -k^2$$

Since we already know  $n(n+1)=k^2$ , we can get the second differential equation:

$$\frac{d}{d\theta} \left[ \sin \theta \frac{d\Theta(\theta)}{d\theta} \right] + n(n+1) \Theta(\theta) \sin \theta = 0$$

. It is called Legendre's equation and for  $\theta \in [0, \pi]$ , the solution has special forms called Legendre's polynomials:

$$\Theta_n(\theta) = P_n(\cos \theta)$$

There are some solutions forms for usual  $n$ .

$n$	$P_n(\cos \theta)$
0	1
1	$\cos \theta$
2	$\frac{1}{2} (3 \cos^2 \theta - 1)$
3	$\frac{1}{2} (5 \cos^3 \theta - 3 \cos \theta)$

- (5) By Combining them together,

$$V_n(R, \theta) = [A_n R^n + B_n R^{-(n+1)}] P_n(\cos \theta)$$

## 4 Steady Electric Currents

Types of electric currents caused by the motion of free charges:

1. **conduction currents:** drift motion of conduction electrons and/or holes in conductors/semi-conductors.
2. electrolytic currents: migration of positive and negative ions.
3. convection currents: motion of electrons and/or ions in a vacuum.

## 5 Current Density and Ohm's Law

$$I = \int_S \mathbf{J} \cdot d\mathbf{s} \quad (A)$$

where  $\mathbf{J}$  is the volume current density or current density, defined by

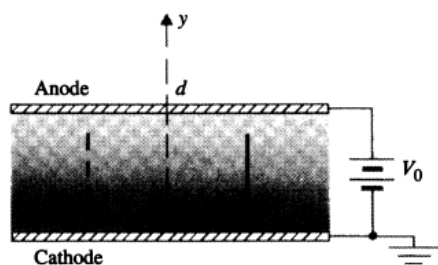
$$\mathbf{J} = Nqu \quad (A/m^2)$$

where  $N$  is the number of charge carriers per unit volume, each of charges  $q$  moves with a velocity  $u$ .

Since  $Nq$  is the free charge per unit volume, by  $\rho = Nq$ , we have:

$$\mathbf{J} = \rho u \quad (A/m^2)$$

**Ex.5-1** In vacuum-tube diodes, electrons are emitted from a hot cathode at zero potential and collected by an anode maintained at a potential  $V_0$ , resulting in a convection current flow. Assuming that the cathode and the anode are parallel conducting plates and that the electrons leave the cathode with a zero initial velocity (space-charge limited condition), find the relation between the current density  $J$  and  $V_0$ .



**FIGURE 5-2**  
Space-charge-limited vacuum diode (Example 5-1).

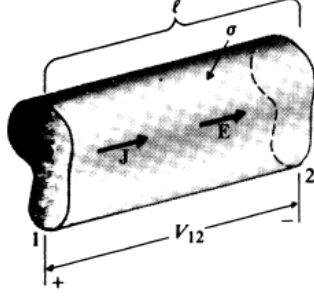
For conduction currents,

$$J = \sigma E \quad (A/m^2)$$

where  $\sigma = \rho_e \mu_e$  is conductivity, a macroscopic constitutive parameter of the medium.  $\rho_e = -Ne$  is the charge density of the drifting electrons and is negative.  $u = -\mu_e E$  ( $m/s$ ) where  $\mu_e$  is the electron mobility measured in ( $m^2/V \cdot s$ ).

Materials where  $J = \sigma E$  ( $A/m^2$ ) holds are called ohmic media. The form can be referred as the point form of Ohm's law.

**Derivation of voltage-current relationship of a piece of homogeneous material by the point form of Ohm's law.**



**FIGURE 5-3**  
Homogeneous conductor with a constant cross section.

Thus, the resistance is defined as

$$R = \frac{l}{\sigma S} \quad (\Omega)$$

where  $l$  is the length of the homogeneous conductor,  $S$  is the area of the uniform cross section. The conductance  $G$  (reciprocal of resistance), is defined by

$$G = \frac{1}{R} = \sigma \frac{S}{l} \quad (S)$$

1. Resistance in series:

$$R_{sr} = R_1 + R_2$$

2. Resistance in parallel:

$$\frac{1}{R_{||}} = \frac{1}{R_1} + \frac{1}{R_2}$$

, or

$$G_{||} = G_1 + G_2$$

## 6 Electromotive Force and Kirchhoff's Voltage Law

A steady current cannot be maintained in the same direction in a closed circuit by an electrostatic field, which is:

$$\oint_C \frac{1}{\sigma} \mathbf{J} \cdot d\mathbf{l} = 0$$

Kirchhoff's voltage law: around a closed path in an electric circuit, the algebraic sum of the emf's (voltage rises) is equal to the algebraic sum of the voltage drops across the resistance, which is:

$$\sum_j V_j = \sum_k R_k I_k \quad (V)$$

## 7 Equation of Continuity and Kirchhoff's Current Law

Equation of continuity:

$$\nabla \cdot \mathbf{J} = -\frac{\partial \rho}{\partial t} \quad (A/m^3)$$

,

where  $\rho$  is the volume charge density.

For steady currents, as  $\partial \rho / \partial t = 0$ ,  $\nabla \cdot \mathbf{J} = 0$ . By integral, we have Kirchhoff's current law, stating that the algebraic sum of all the currents flowing out of a junction in an electric circuit is zero:

$$\sum_j I_j = 0$$

For a simple medium conductor, the volume charge density  $\rho$  can be expressed as:

$$\rho = \rho_0 e^{-(\rho/\epsilon)t} \quad (C/m^3)$$

where  $\rho_0$  is the initial charge density at  $t = 0$ . The equation implies that the charge density at a given location will decrease with time exponentially.

Relaxation time: an initial charge density  $\rho_0$  will decay to  $1/e$  or 36.8% of its original value:

$$\tau = \frac{\epsilon}{\sigma} \quad (s)$$

## 8 Power Dissipation and Joule's Law

For a given volume  $V$  that the total electric power converted to heat is:

$$P = \int_V \mathbf{E} \cdot \mathbf{J} dv$$

$$P = \int_L E d\ell \int_S J ds = VI = I^2 R$$

## 9 Boundary Conditions

### 9.1 Governing Equations for Steady Current Density

- Differential form:

$$\begin{aligned}\nabla \cdot \mathbf{J} &= 0 \\ \nabla \times \left( \frac{\mathbf{J}}{\sigma} \right) &= 0\end{aligned}$$

- Integral form:

$$\begin{aligned}\oint_S \mathbf{J} \cdot d\mathbf{s} &= 0 \\ \oint_C \frac{1}{\sigma} \mathbf{J} \cdot d\boldsymbol{\ell} &= 0\end{aligned}$$

### 9.2 Boundary Conditions:

- Normal Component:

$$J_{1n} = J_{2n}$$

- Tangential Component:

$$\frac{J_{1t}}{J_{2t}} = \frac{\sigma_1}{\sigma_2}$$

Combining with boundary conditions of electric field:

$$\begin{aligned}J_{1n} = J_{2n} &\rightarrow \sigma_1 E_{1n} = \sigma_2 E_{2n} \\ D_{1n} - D_{2n} = \rho_s &\rightarrow \epsilon_1 E_{1n} = \epsilon_2 E_{2n}\end{aligned}$$

Surface charge density on the interface:

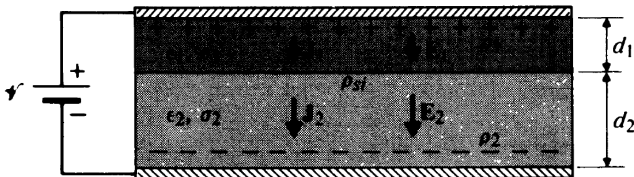
$$\rho_s = \left( \epsilon_1 \frac{\sigma_2}{\sigma_1} - \epsilon_2 \right) E_{2n} = \left( \epsilon_1 - \epsilon_2 \frac{\sigma_1}{\sigma_2} \right) E_{1n}$$

If medium 2 is a much better conductor than medium 1:

$$\rho_s = \epsilon_1 E_{1n} = D_{1n}$$

**EXAMPLE 5 – 4** An emf  $\mathcal{V}$  is applied across a parallel-plate capacitor of area  $S$ . The space between the conducting plates is filled with two different lossy dielectrics of thicknesses  $d_1$  and  $d_2$ , permittivities  $\epsilon_1$  and  $\epsilon_2$ , and conductivities  $\sigma_1$  and  $\sigma_2$ , respectively. Determine

- the current density between the plates,
- the electric field intensities in both dielectrics, and
- the surface charge densities on the plates and at the interface.



## 10 Resistance Calculation

$$C = \frac{Q}{V} = \frac{\oint_S \mathbf{D} \cdot d\mathbf{s}}{-\int_L \mathbf{E} \cdot d\boldsymbol{\ell}} = \frac{\oint_S \epsilon \mathbf{E} \cdot d\mathbf{s}}{-\int_L \mathbf{E} \cdot d\boldsymbol{\ell}}$$

$$R = \frac{V}{I} = \frac{-\int_L \mathbf{E} \cdot d\boldsymbol{\ell}}{\oint_S \mathbf{J} \cdot d\mathbf{s}} = \frac{-\int_L \mathbf{E} \cdot d\boldsymbol{\ell}}{\oint_S \sigma \mathbf{E} \cdot d\mathbf{s}}$$

If  $\sigma$  and  $\epsilon$  have the same space dependence or the medium is homogeneous:

$$RC = \frac{C}{G} = \frac{\epsilon}{\sigma}$$

Method of resistance calculation:

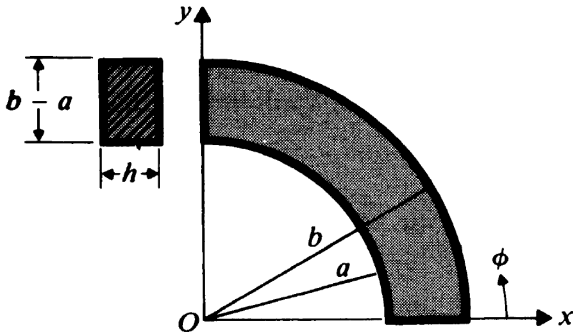
1. Choose a proper coordinate for the system.
2. Assume a voltage different  $V_0$  between the two conductors.
3. Find the electric field  $\mathbf{E}$  within the conductors (for homogenous medium, solve Laplace's Equation  $\Delta^2 V = 0$  and get  $E = -\Delta V$ ).

4. Find the total current

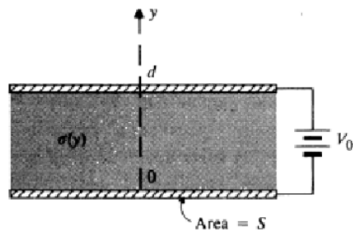
$$I = \int_S \mathbf{J} \cdot d\mathbf{s} = \int_S \sigma \mathbf{E} \cdot d\mathbf{s}$$

5. The resistance is  $R = \frac{V_0}{I}$ .

EXAMPLE 5-6 A conducting material of uniform thickness  $h$  and conductivity  $\sigma$  has the shape of a quarter of a flat circular washer, with inner radius  $a$  and outer radius  $b$ , as shown in Fig. 5 – 8. Determine the resistance between the end faces.



- P.5-10 The space between two parallel conducting plates each having an area  $S$  is filled with an inhomogeneous ohmic medium whose conductivity varies linearly from  $\sigma_1$  at one plate ( $y = 0$ ) to  $\sigma_2$  at the other plate ( $y = d$ ). A d-c voltage  $V_0$  is applied across the plates as in Fig.5-11. Determine
- the total resistance between the plates,
  - the surface charge densities on the plates,
  - the volume charge density and the total amount of charge between the plates.



**FIGURE 5-11**  
Inhomogeneous ohmic medium with conductivity  $\sigma(y)$  (Problem P.5-10).