

## Question 1

- Suppose the functions in Fig 1 are  $\mathbf{v}_a = x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}$ ,  $\mathbf{v}_b = \hat{\mathbf{z}}$ , and  $\mathbf{v}_c = z\hat{\mathbf{z}}$ . Calculate their divergence.

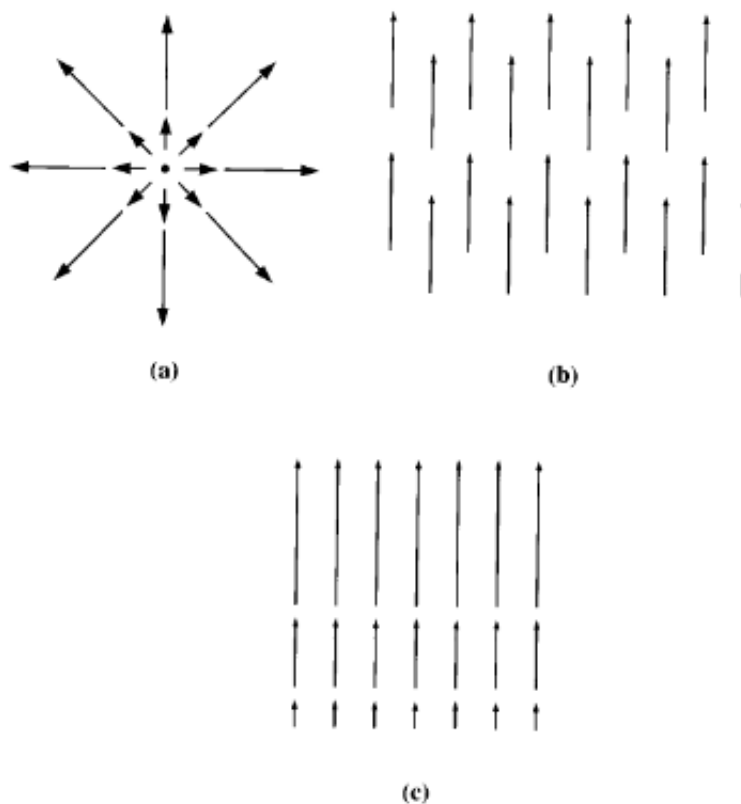


Figure 1: Figure for Question 1.1

- Suppose the function sketched in Fig 2a is  $\mathbf{v}_a = -y\hat{\mathbf{x}} + x\hat{\mathbf{y}}$ , and that in Fig 2b is  $\mathbf{v}_b = x\hat{\mathbf{y}}$ . Calculate their curls.

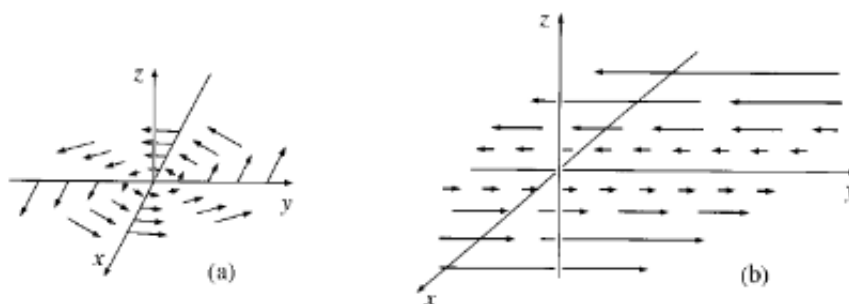


Figure 2: Figure for Question 1.2

- Construct a vector function that has zero divergence and zero curl everywhere. (A *constant* will do the job, of course, but make it something a little more interesting than that!)

## Question 2

The expression of the electric field for a point charge (Fig 3) is:

$$\mathbf{E}(\mathbf{r}) \equiv \frac{1}{4\pi\epsilon_0} \sum_{i=1}^n \frac{q_i}{r_i^2} \hat{\mathbf{r}}_i \quad (1)$$

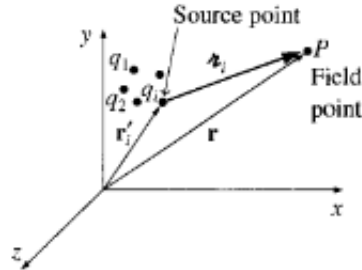


Figure 3: Figure for Eq 1

1. Find the electric field (magnitude and direction) a distance  $z$  above the midpoint between two equal charges,  $q$ , a distance  $d$  apart (Fig 4). Check that your result is consistent with what you'd expect when  $z \gg d$ .

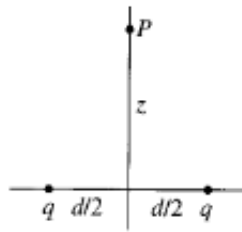


Figure 4: Figure for Question 2.1

2. Repeat part 1 in this question, only this time make the right-hand charge  $-q$  instead of  $+q$ .

## Question 3

The expression for the work required to bring together  $n$  charges is given by

$$W = \frac{1}{2} \sum_{i=1}^n q_i V(\mathbf{r}_i)$$

, where  $V(\mathbf{r}_i)$  is the potential.

1. Three charges are situated at the corners of a square (side  $a$ ), as shown in Fig 5. How much work does it take to bring in another charge,  $+q$ , from far away and place it in the fourth corners?

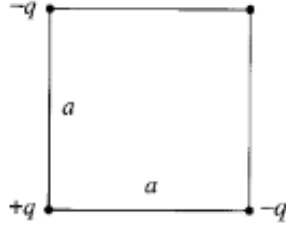


Figure 5: Figure for Question 3.1

2. How much work does it take to assemble the whole configuration of four charges?

## Question 4

The original form of the Gauss's law is

$$\oint_S \mathbf{E} \cdot d\mathbf{a} = \frac{1}{\epsilon_0} Q_{enc} \quad (2)$$

, where  $Q_{enc}$  is the total charge enclosed within the surface.

The electric displacement could be calculated by

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} \quad (3)$$

The revised form for the Gauss's law which considers the electric displacement is then

$$\oint \mathbf{D} \cdot d\mathbf{a} = Q_{fenc} \quad (4)$$

, where  $Q_{fenc}$  indicates the total free charge enclosed in the volume.

1. Show that the energy of an ideal dipole  $\mathbf{p}$  in an electric field  $\mathbf{E}$  is given by

$$U = -\mathbf{p} \cdot \mathbf{E}$$

2. A thick spherical shell (inner radius  $a$ , outer radius  $b$ ) is made of dielectric material with a "frozen-in" polarization

$$\mathbf{P}(\mathbf{r}) = \frac{k}{r} \hat{\mathbf{r}}$$

, where  $k$  is a constant and  $r$  is the distance from the center (Fig 6). (There is no *free* charge in the problem). Find the electric field in all three regions by two different methods:

- (a) Locate all the bound charge, and use Gauss's law (Eq 2) to calculate the field it produces.
- (b) Use Eq 4 to find  $\mathbf{D}$ , and then get  $\mathbf{E}$  from Eq 3. [Notice that the second method is much faster, and avoids any explicit reference to the bound charges.]

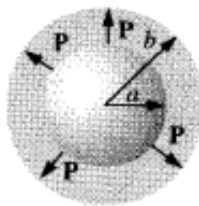


Figure 6: Figure for Question 4.2

## Question 5

Choose either question A or question B to answer. If both questions are answered, only question A would be graded.

### A

The Ampere's law with Maxwell's correction is given by:

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \quad (5)$$

Suppose  $\mathbf{J}(\mathbf{r})$  is constant in time but  $\rho(\mathbf{r}, t)$  is *not* — conditions that might prevail, for instance, during the charging of a capacitor.

1. Show that the charge density at any particular point is a linear function of time:

$$\rho(\mathbf{r}, t) = \rho(\mathbf{r}, 0) + \dot{\rho}(\mathbf{r}, 0)t,$$

, where  $\dot{\rho}(\mathbf{r}, 0)$  is the time derivative of  $\rho$  at  $t = 0$ . This is *not* an electrostatic or magnetostatic configuration; nevertheless — rather superisingly — both Coloumb's law and the Biot-Savart law hold, as you can confirm by showing that they satisfy Maxwell's equations. In particular:

2. Show that

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}') \times \hat{\mathbf{r}}}{r^2} d\tau'$$

obeys Ampere's law with Maxwell's displacement current term (Eq 5).

### B

The Maxwell's equations with magnetic charge is given by:

$$\nabla \cdot \mathbf{E} = \frac{1}{\epsilon_0} \rho_e, \quad (6a)$$

$$\nabla \cdot \mathbf{B} = \mu_0 \rho_m \quad (6b)$$

$$\nabla \times \mathbf{E} = -\mu_0 \mathbf{J}_m - \frac{\partial \mathbf{B}}{\partial t} \quad (6c)$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}_e + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \quad (6d)$$

1. Show that Maxwell's equations with magnetic charge (Eq 6) are invariant under the duality transformation

$$\left. \begin{aligned} \mathbf{E}' &= \mathbf{E} \cos \alpha + c \mathbf{B} \sin \alpha, \\ c \mathbf{B}' &= c \mathbf{B} \cos \alpha - \mathbf{E} \sin \alpha, \\ c q_e' &= c q_e \cos \alpha + q_m \sin \alpha, \\ q_m' &= q_m \cos \alpha - c q_e \sin \alpha, \end{aligned} \right\},$$

where  $c \equiv 1/\sqrt{\epsilon_0 \mu_0}$  and  $\alpha$  is an arbitrary rotation angle in " $\mathbf{E}/\mathbf{B}$ -space." Charge and current densities transform in the same way as  $q_e$  and  $q_m$ . [This means, in particular, that if you knows the fields produced by a configuration of electric charge, you can immediately (using  $\alpha = 90^\circ$ ) write down the fields produced by the corresponding arrangement of magnetic charge.]

2. Show that the force law

$$\mathbf{F} = q_e (\mathbf{E} + \mathbf{v} \times \mathbf{B}) + q_m (\mathbf{B} - \frac{1}{c^2} \mathbf{v} \times \mathbf{E})$$

is also invariant under the duality transformation.

## Question 6

Choose either question A or question B to answer. If both questions are answered, only question A would be graded.

The magnetic potential  $\mathbf{A}$  is introduced by

$$\mathbf{B} = \nabla \times \mathbf{A} \quad (T) \quad (7)$$

Also, we have

$$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t} \quad (V/m) \quad (8)$$

The nonhomogeneous wave equation for vector potential  $\mathbf{A}$  is:

$$\nabla^2 \mathbf{A} - \mu\epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} = -\mu \mathbf{J} \quad (9)$$

The nonhomogeneous wave equation for scalar potential  $V$  is :

$$\nabla^2 V - \mu\epsilon \frac{\partial^2 V}{\partial t^2} = -\frac{\rho}{\epsilon} \quad (10)$$

### A

The vector magnetic potential  $\mathbf{A}$  and scalar electric potential  $V$  are not unique in that it is possible to add to  $\mathbf{A}$  the gradient of a scalar  $\psi$ ,  $\nabla\psi$ , with no change in  $\mathbf{B}$  from Eq 7.

$$\mathbf{A}' = \mathbf{A} + \nabla\psi$$

In order not to change  $\mathbf{E}$  in using Eq 8,  $V$  must be modified to  $V'$ .

1. Find the relation between  $V'$  and  $V$ .
2. Discuss the condition that  $\psi$  must satisfy so that the new potential  $\mathbf{A}'$  and  $V'$  remain governed by the uncoupled wave equations Eq 9 and Eq 10

### B

Substitute Eq 7 and Eq 8 in Maxwell's equations to obtain wave equations for scalar potential  $V$  and vector potential  $\mathbf{A}$  for a linear, isotropic but inhomogeneous medium. Show that these wave equations reduce to Eq 9 and Eq 10 for simple media. (*Hint: Use the following gauge condition for potentials in an inhomogeneous medium:*

$$\nabla \cdot (\epsilon \mathbf{A}) + \mu\epsilon^2 \frac{\partial V}{\partial t} = 0$$

)

**Congratulations on finishing VE230!**



Figure 7: Figure for finishing VE230