

Smoothed Particle Hydrodynamics

An Overview

Yogeshwaran R

FSI Lab, Department of Applied Mechanics
Indian Institute of Technology, Madras

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Introduction

- Formulated by Lucy (1977) and Gingold and Monaghan (1977).
- Originally invented for astrophysical applications.
- Later, extended and applied to wide range of applications such as :
 - Dynamic Fluid flows
 - Explosions
 - Hydrodynamics with material strength etc.

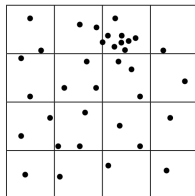
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Foundations of SPH : Calculating Density

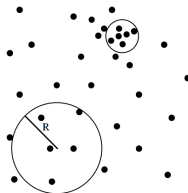
SPH :

"A meshfree lagrangian particle based method for solving equation of hydrodynamics"

How to compute density from an arbitrary distribution of point mass particles ?

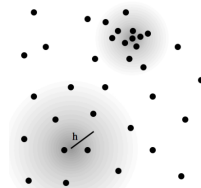


(a) $\rho = m/V$



(b) based on local sampling

$$\rho(r) = \frac{\sum_{b=1}^{N_{neigh}} m_b}{\frac{4}{3} \pi R^3}$$



(c) SPH density estimator

$$\rho(r) = \sum_{b=1}^{N_{neigh}} m_b W(r-r_b, h)$$

FIGURE : Computing a continuous density field from a collection of point mass particles.

The SPH density estimate

The density is computed using a weighted summation over nearby particles, which is given by,

$$\rho(r) = \sum_{b=1}^{N_{neigh}} m_b W(r - r_b, h)$$

Accuracy of density estimate \rightarrow depends on choice of weight function

A good density kernel should have following properties :

- $W \rightarrow$ '+ve, decrease monotonically and have smooth derivatives.
- Symmetric w.r. to $|r-r'|$ (i.e.,) $W(r-r', h) \equiv |W(|r' - r|, h)|$
- Compact support. $[O(N^2) \rightarrow O(N_{neigh} N) \Rightarrow \text{Computationally efficient}]$

Basics :

$$A(r) = \int A(r') \underbrace{\delta(r - r')}_{\text{dirac delta fn.}} dr'$$

Two major approximations in SPH :

1) Kernel approximation

2) Particle approximation

1) Kernel approximation

$$A(r) = \int A(r') W(r - r', h) dr' + O(h^2)$$

2) Particle approximation

$$\langle A(r) \rangle = \int \frac{A(r')}{\rho(r')} W(r - r', h) \rho(r') dr'$$

$$\langle A(r) \rangle \approx \sum_{b=1}^{N_{\text{neigh}}} m_b \frac{A_b}{\rho_b} W(r - r_b, h)$$

Error

i) Error in Kernel approximation :

$$\langle A(r) \rangle = \int A(r') W(r - r', h) dV'$$

Expanding $\rho(r')$ in Taylor's series about r ,

$$\langle A(r) \rangle = A(r) \underbrace{\int W(r - r', h) dV'}_{\approx 1 \text{ (Normalization cond.)}} + \nabla A(r) \underbrace{\int (r' - r) W(r - r', h) dV'}_{= 0 \text{ (symmetric)}} + O(h^2)$$

ii) Error in particle approximation :

$$\langle A_a \rangle = \sum_b m_b \frac{A_b}{\rho_b} W_{ab}$$

Expanding A_b in Taylor's series,

$$\langle A_a \rangle = A_a \underbrace{\sum_b \frac{m_b}{\rho_b} W_{ab}}_{\approx 1} + \nabla A_a \underbrace{\sum_b \frac{m_b}{\rho_b} (r_b - r_a) W_{ab}}_{= 0 \text{ (anti-symm.)}} + O(h^2)$$

From Density to Hydrodynamics

$$L = \sum_b m_b \left[\frac{1}{2} v_b^2 - u_b(\rho_b, s_b) \right] \rightarrow \text{Lagrangian}$$

+

$$du = \frac{p}{\rho^2} d\rho \rightarrow 1^{st} \text{ Law of thermodynamics}$$

+

$$\rho(r) = \sum_{b=1}^{N_{neigh}} m_b W(r - r_b, h) \rightarrow \text{Density Sum}$$

+

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \mathbf{v}} \right) - \frac{\partial L}{\partial \mathbf{r}} = 0 \rightarrow \text{Euler-Lagrange equations}$$

=

$$\frac{d\mathbf{v}}{dt} = - \sum_j m_j \left(\frac{p_i}{\rho_i^2} + \frac{p_j}{\rho_j^2} \right) \nabla_i W_{ij}(h) \Rightarrow \text{Eqn. of motion}$$

Generalised first derivative operators

Two basic types :

$$(i) \quad \langle \nabla A \rangle \approx \sum_b \frac{m_b}{\rho_b} \frac{\phi_b}{\phi_a} (A_b - A_a) \nabla_a W_{ab}$$

$$(ii) \quad \langle \nabla A \rangle \approx \sum_b \frac{m_b}{\rho_b} \left(\frac{\phi_b}{\phi_a} A_a + \frac{\phi_a}{\phi_b} A_b \right) \nabla_a W_{ab}$$

where ϕ is any arbitrary, differentiable scalar quantity defined on the particles.

Various '*alternative formulations*' of SPH have been proposed based on the choice of ϕ

For example.,

Taking $\phi = \rho$ in (i)

$$\nabla A \approx \frac{1}{\rho_a} \sum_b m_b (A_b - A_a) \nabla_a W_{ab}$$

→ **exact derivative**

Taking $\phi = \rho$ in (ii)

$$\nabla A \approx \rho_a \sum_b m_b \left(\frac{A_a}{\rho_a^2} + \frac{A_b}{\rho_b^2} \right) \nabla_a W_{ab}$$

→ **exact conservation**

Which is preferable : exact derivative or exact conservation ?

Let us consider momentum eqn.,

$$\frac{d\mathbf{v}}{dt} = -\frac{\nabla p}{\rho}$$

Based on the *exact derivative approach* :

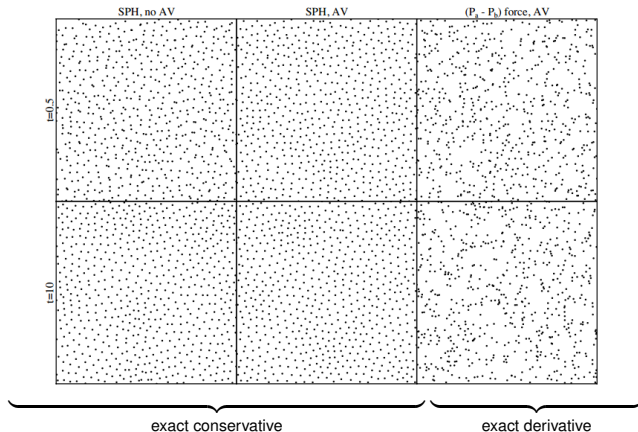
$$\frac{d\mathbf{v}_a}{dt} = \sum_b m_b \frac{p_a - p_b}{\rho_a \rho_b} \nabla_a W_{ab}$$

Based on the *exact conservation approach* :

$$\frac{d\mathbf{v}_a}{dt} = -\sum_b m_b \left(\frac{p_a}{\rho_a^2} + \frac{p_b}{\rho_b^2} \right) \nabla_a W_{ab}$$

For illustration, let us consider a distribution of particles in a closed box.

What happens if the pressure is constant ?



The exact conservative approach contains "*intrinsic remeshing*" which cares about bad arrangement of particles.

Insensitive to particle arrangement leads to *interpolation error*.

Tensile instability

Consider a mom. eqn. with pressure gradient is of the form,

$$\frac{d\mathbf{v}}{dt} = - \sum_b m_b \left(\frac{p_a - p_0}{\rho_a^2} + \frac{p_b - p_0}{\rho_b^2} \right) \underbrace{\nabla_a W_{ab}}_{\text{'-' ve (for '+' ve definite kernels)}}$$

When $p_0 > p$, the pairwise force \rightarrow '-ve \Rightarrow particles clump together unphysically
 due to attractive force

This is known as "*Tensile instability*".

Why one cannot simply use '*more neighbors*' ?

- Consider a 2D setup with particles are initially placed on a close packed lattice.
- If we consider **M4 cubic spline kernel with a large ratio of smoothing length to particle**, what will happen ?

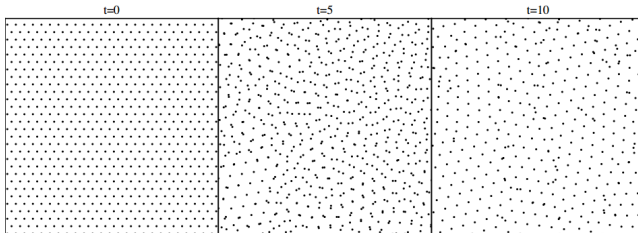
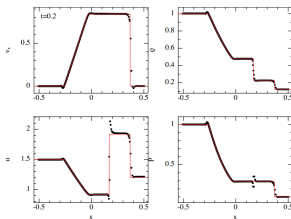
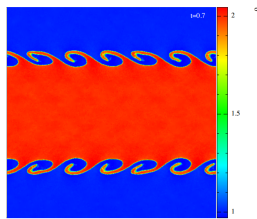


FIGURE : Pairing instability in action

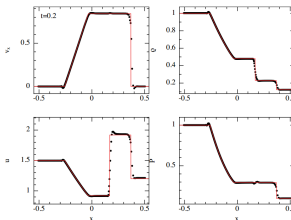


(a) 1D Sod Shock tube problem

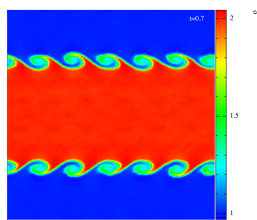


(b) 2D Kh instability

FIGURE : With artificial viscosity



(a) 1D Sod Shock tube problem



(b) 2D Kh instability

FIGURE : With artificial viscosity + artificial conductivity

Thank you !!!