

2.3.2 High-Dimensional Vectors

For an n -dimensional space, it has a basis of n basis vectors. Every vector,

$$\vec{a} = a_0 \hat{x}_0 + a_1 \hat{x}_1 + \cdots + a_{n-1} \hat{x}_{n-1}, \quad (2.16)$$

or in bra-ket notation, it is,

$$\begin{aligned} |0\rangle &= a_0 |x_0\rangle + a_1 |x_1\rangle + \cdots + a_{n-1} |x_{n-1}\rangle, \\ &= a_0 |0\rangle + a_1 |1\rangle + \cdots + a_{n-1} |n-1\rangle, \\ &= \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_{n-1} \end{pmatrix}, \end{aligned} \quad (2.17)$$

where in line 2, we changed the names of the basis states to emphasize that whatever is put inside the ket is just name. As long as it is not confusing, it does not matter how we name it. When we write the vector in a column form, we have assumed that the basis vectors are orthonormal, which will be discussed soon in Sect. 2.3.4.

Each vector has a corresponding vector in the bar-space (dual correspondence). This is similar to the fact that every object has an image in the mirror. The bra of $|b\rangle$ is written as $\langle b|$. And to construct the *bra* version of $|b\rangle$ in matrix form, we need to **perform conjugate transpose**. That is to swap the rows and columns and then apply complex conjugate to each element. Therefore, the column vector has a row vector in its *bra* version. For example, vector $|b\rangle$, which is expressed as

$$|b\rangle = b_0 |x_0\rangle + b_1 |x_1\rangle + \cdots + b_{n-1} |x_{n-1}\rangle = \begin{pmatrix} b_0 \\ b_1 \\ \vdots \\ b_{n-1} \end{pmatrix} \quad (2.18)$$

has its *bra* version expressed as

$$\begin{aligned} \langle b| &= b_0^* \langle x_0| + b_1^* \langle x_1| + \cdots + b_{n-1}^* \langle x_{n-1}|, \\ &= (b_0^* b_1^* \cdots b_{n-1}^*) \end{aligned} \quad (2.19)$$

The *bra-ket* notation is very useful in linear algebra. For example, the inner product of two vectors, $|b\rangle$ and $|a\rangle$, is just the multiplication between the *bra* of $|b\rangle$ and the ket of $|a\rangle$,

$$\langle b|a\rangle = a_0 b_0^* + a_1 b_1^* + \cdots + a_{n-1} b_{n-1}^*, \quad (2.20)$$

which is the same as how we wrote it in Eq.(2.3).

Example 2.3 For $|a\rangle = \begin{pmatrix} 3i+2 \\ 0 \\ 5 \\ 4-2i \end{pmatrix}$ and $|b\rangle = \begin{pmatrix} 2 \\ i \\ 0 \\ 2i \end{pmatrix}$ find $\langle a|b\rangle$.

$$\begin{aligned} \langle a|b\rangle &= (-3i+2 \ 0 \ 5 \ 4-2i) \begin{pmatrix} 2 \\ i \\ 0 \\ 2i \end{pmatrix}, \\ &= (-6i+4) + 0 + (8i-4) = 2i, \end{aligned} \quad (\text{ex.2.3})$$

2.3.3 Measurement of a quantum state

Measurement is not a part of linear algebra. However, I would like to inject this topic so that we can understand the following sections better. The measurement of a quantum state results in the **collapse of the state** to one of its basis states. That means that the measurement outcome is one of the basis states and the original quantum state no longer exists. The process is completely random except that the probability it will collapse to a certain basis vector is the square of the magnitude of the corresponding coefficient (Fig.2.3). For example, for $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$, the probability it will collapse to $|0\rangle$ is

$$Prob(|0\rangle) = \alpha\alpha^* = |\alpha|^2 \quad (2.21)$$

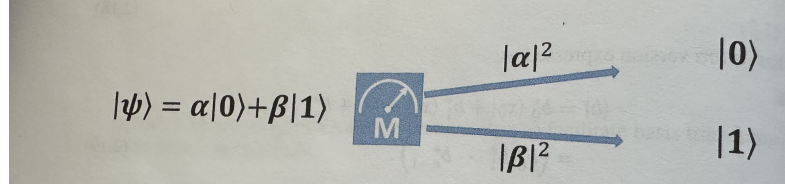


Fig.2.3 Upon measurement, a state will collapse to one of the basis states with a probability equal to the square of the magnitude of the corresponding coefficient

and the probability it will collapse to $|1\rangle$ is

$$Prob(|1\rangle) = \beta\beta^* = |\beta|^2 \quad (2.22)$$