## Lecture 25

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## 1 Ampère's Law

**Example 1.1.** An infinitely long metallic strip of negligible thickness and width 2a carries a total current of I. The strip lies in the xy-plane, is centered about the z-axis, and is infinitely long in the x-direction. The current is uniformly distributed over the width of the strip and flows in the +x-direction. Determine the magnetic field due to this current strip at a point P(0,0,z).

Using Biot-Savart law,

$$d\vec{B} = \frac{\mu_0 \vec{J} \times (\vec{R} - \vec{R}') dS}{4\pi |\vec{R} - \vec{R}'|^3}$$

Plugging this into an integral,

$$\vec{B} = \int_{-\infty}^{\infty} \int_{-a}^{a} \frac{\mu(\frac{I}{2a})\hat{a}_{x} \times (-x\hat{a}_{x} - y\hat{a}_{y} + z\hat{a}_{z})}{(x^{2} + y^{2} + z^{2})^{1.5}} dy dx$$

$$= \frac{\mu_{0}I}{8\pi a} \int_{-\infty}^{\infty} \int_{-a}^{a} \frac{-z\hat{a}_{y} - y\hat{a}_{z}}{(x^{2} + y^{2} + z^{2})^{1.5}} dy dx$$

$$= \frac{\mu_{0}I}{8\pi a} \int_{-\infty}^{\infty} \int_{-a}^{a} \frac{-z\hat{a}_{y}}{(x^{2} + y^{2} + z^{2})^{1.5}} dy dx$$

$$= -\frac{\mu_{0}I}{2\pi a} \tan^{-1} \left(\frac{a}{z}\right) \hat{a}_{y}$$

Note that the  $\hat{a}_z$  term is ignored. By the right hand rule, we know the resulting field is parallel to  $\hat{a}_y$ . Moreover, the  $\hat{a}_z$  term is odd, so the integral yields a 0.

The forms of **Ampère's Law** are its differential form

$$\vec{\nabla} \times \vec{H} = \vec{J}$$

and its integral form

$$\oint_C \vec{H} \cdot d\vec{l} = \iint_S \vec{J} \cdot d\vec{S} = I$$

This means at any point in space, the magnetic field has a non zero curl if and only if a current density  $\vec{J}$  is present.

## 1.1 Integral Form

In the integral form, we will use an **Ampèrian Loop**. We choose a loop such that

- $\vec{H}$  is either tangential or normal to the loop, so  $\vec{H} \cdot d\vec{l}$  is Hdl or 0
- $\vec{H}$  has a constant value when  $\vec{H}$  is tangential, so  $\int H dl = H \int dl = H L$

**Example 1.2.** Find the magnetic fields within each region of a coaxial cable. First observe that  $\vec{H} = H_{\phi}\hat{a}_{\phi}$ . Now for 0 < r < a, the inner cabe,

$$\oint_C \vec{H} \cdot d\vec{l} = H_\phi(2\pi r) = I_{\rm enc} = J\pi r^2 = \left(\frac{I_0}{\pi}a^2\right)\pi r^2 \Rightarrow \vec{H} = \frac{I_0a^2r}{2\pi}$$