

# Lecture 11

niceguy

March 23, 2023

## 1 Ideal Gas

$$S(E, N, V) = k \ln \Omega(E, N, V)$$

and

$$\frac{1}{T} = \left( \frac{\partial S}{\partial E} \right)_{V, N}$$

We want to show this for the ideal gas.

$$U = \sum_i \frac{\vec{p}_i^2}{2m} \Rightarrow \sum_i \vec{p}_i^2 = 2m\vec{U}$$

Now for a sphere in  $n$  dimensions,

$$\sum_i x_i^2 = R^2$$

Both are similar, i.e.  $2mU \propto R^2$ , in the sense that they are the sum of squares. Using this analogy, we have a *radius* for  $\vec{p}_i$ . Now, the area of  $S^{3N-1}$  of radius  $R$  is

$$A = \frac{2\pi^{\frac{3N}{2}}}{\left(\frac{3N}{2} - 1\right)!} R^{3N-1}$$

Now, it seems that

$$\Omega(N, U) \propto \frac{2\pi^{\frac{3N}{2}}}{\left(\frac{3N}{2} - 1\right)!} (2mU)^{\frac{3N-1}{2}}$$

Note that  $2mU$  shares the same units as  $\vec{p}^2$ , which is  $\frac{\pi\hbar}{L}$ . To nondimensionalise this,

$$\Omega(N, U) \propto \frac{2\pi^{\frac{3N}{2}}}{\left(\frac{3N}{2} - 1\right)!} \left( \frac{\sqrt{2mU}}{\pi\hbar/L} \right)^{3N-1}$$

Secondly, since all  $\vec{p}_i$  are positive, we need a factor of  $\frac{1}{2^{3N}}$ .

$$\Omega(N, U) \propto \frac{1}{2^{3N}} \frac{2\pi^{\frac{3N}{2}}}{\left(\frac{3N}{2} - 1\right)!} \left( \frac{\sqrt{2mU}}{\pi\hbar/L} \right)^{3N-1}$$

There is also the indistinguishability of the particles. We only have access to the different  $\vec{p}_i$ , but not the corresponding particle. There are  $N!$  permutations which we have to divide by. Then

$$\Omega(N, U) \propto \frac{1}{2^{3N} N!} \frac{2\pi^{\frac{3N}{2}}}{\left(\frac{3N}{2} - 1\right)!} \left( \frac{\sqrt{2mU}}{\pi\hbar/L} \right)^{3N-1}$$

Using more approximations,

$$\begin{aligned} \Omega(N, U) &\approx \left( \frac{e}{N} \right)^N \frac{1}{8^N} \frac{2(\pi^{1.5})^N}{\left(\frac{3N}{2e}\right)^{1.5N}} ((2mU)^{1.5})^N \left( \frac{V}{(\pi\hbar)^3} \right)^N \\ &\approx \left[ \frac{e^{1.5}}{N} \frac{1}{8} \frac{\pi^{1.5} 2^{1.5}}{3^{1.5}} \frac{V}{(\pi\hbar)^3} \frac{2^{1.5} m^{1.5}}{N^{1.5}} U^{1.5} \right]^N \times 2 \\ &= 2 \left[ e^{1.5} \frac{V}{N} \frac{(4\pi mU)^{1.5}}{(3N(2\pi\hbar)^2)^{1.5}} \right]^N \\ \ln \Omega(N, U) &\approx kN \left[ \ln \left( \frac{(4\pi mU)^{1.5}}{(3N(2\pi\hbar)^2)^{1.5}} \frac{V}{N} \right) + \frac{5}{2} \right] \end{aligned}$$

Differentiating,

$$\frac{1}{T} = \left( \frac{\partial S}{\partial U} \right)_{N, V} = kN \frac{3}{2} \frac{1}{U} \Rightarrow U = \frac{3}{2} kNT$$

We conclude that

- Entropy is extensive (if  $V, N, U$  are all doubled, so is  $S$ )
- Entropy is huge, because of the factor of  $N$
- $U$  increases with  $S$
- $C_V > 0$