



北京大学



Spin field effect transistor

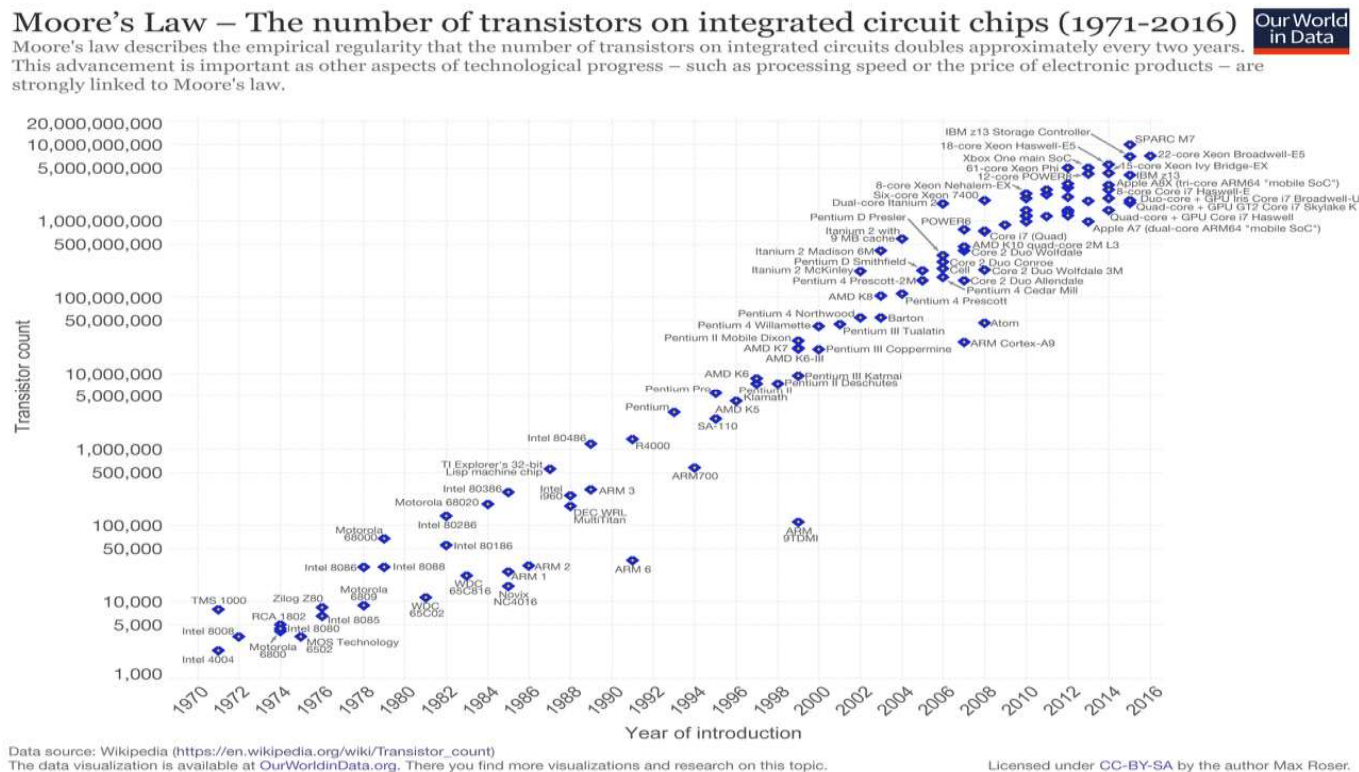
2018.12.21

主讲人：张仕雄

组员：杨洁、朱鹏飞、赵嘉佶、梁栋、丁石磊

Background

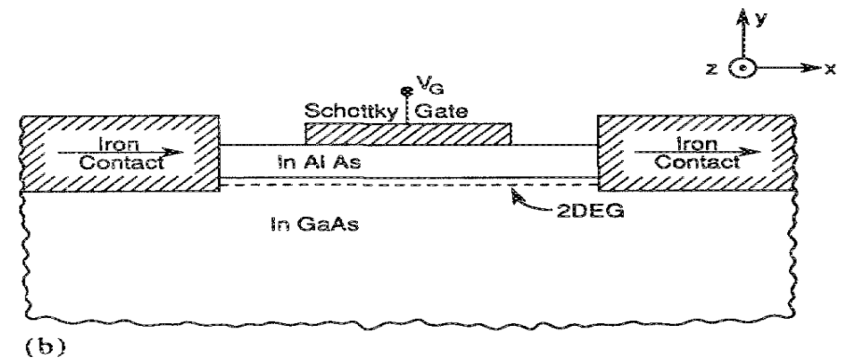
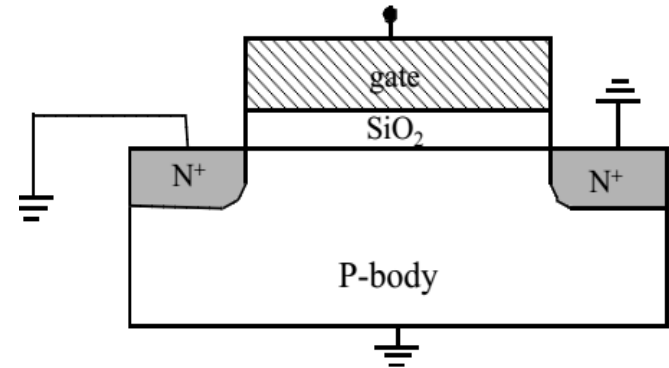
- Moore's law : The number of transistors in a dense integrated circuit doubles about every two years



https://en.wikipedia.org/wiki/Moore%27s_lawv

MOSFET

- Challenge for MOSFET
 - Heat dissipation problems
 - Quantum tunneling
- One of solutions :spin field effect transistor



Rashba SOC in 2DEG

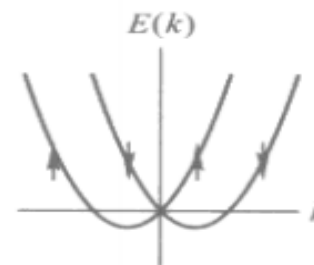
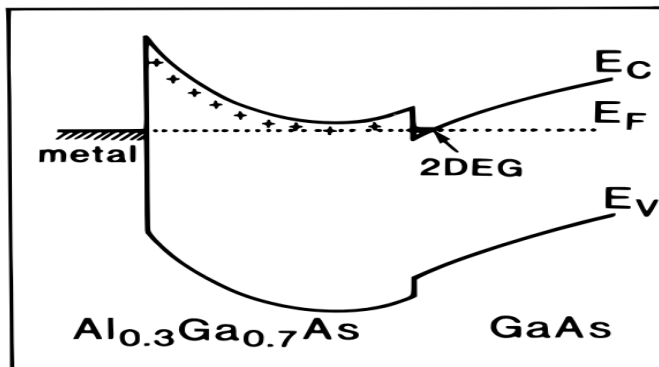
- Rashba spin orbit coupling(structure inversion asymmetry):

$$H_R = \eta(\sigma_z k_x - \sigma_x k_z)$$

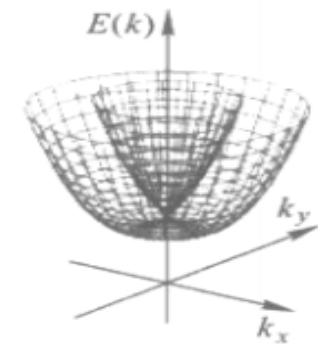
- Spin split energy band:

$$E(z \text{ pol.}) = \hbar^2 k_{x1}^2 / 2m^* - \eta k_{x1},$$

$$E(-z \text{ pol.}) = \hbar^2 k_{x2}^2 / 2m^* + \eta k_{x2}$$



(a)



(b)

Electro-optic modulator

- Different spin polarized electrons have different wave vector

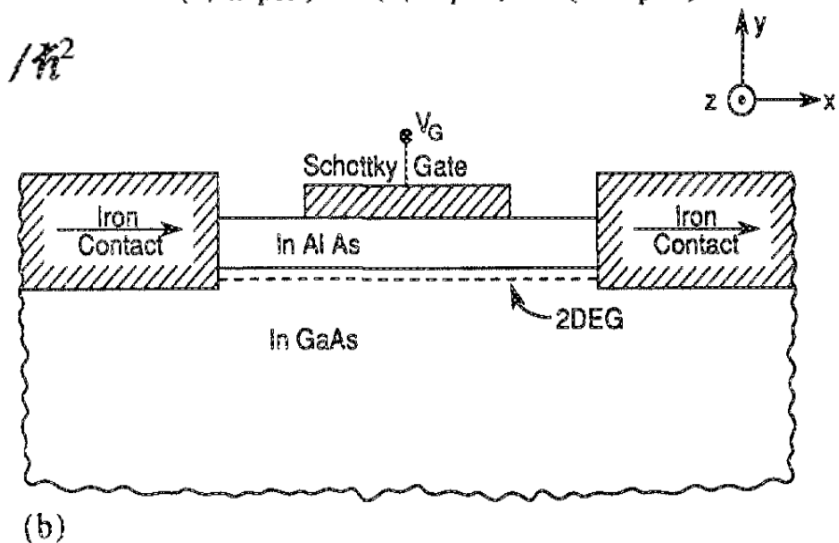
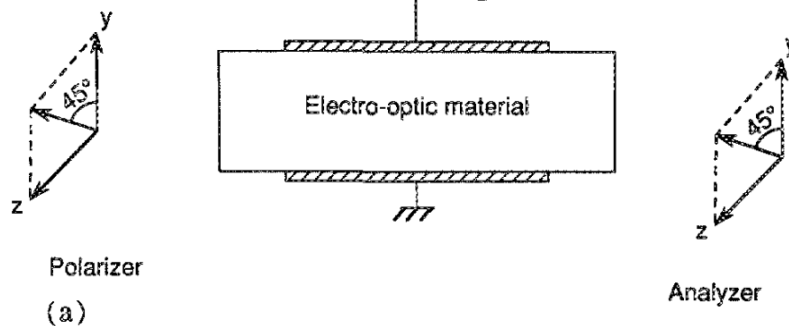
$$k_{x1} - k_{x2} = 2m^*\eta/\hbar^2$$

- A differential phase shift

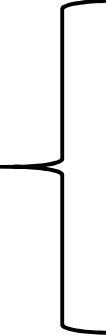
$$\Delta\theta = (k_{x1} - k_{x2})L = 2m^*\eta L / \hbar^2$$

$$\begin{pmatrix} 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ 1 \end{pmatrix} .$$

(+ x pol.) (+ z pol.) (- z pol.)



Why choose spin FET?

- Property 
 - high speed
 - low power consumption
 - high level integration

III-V semiconductor

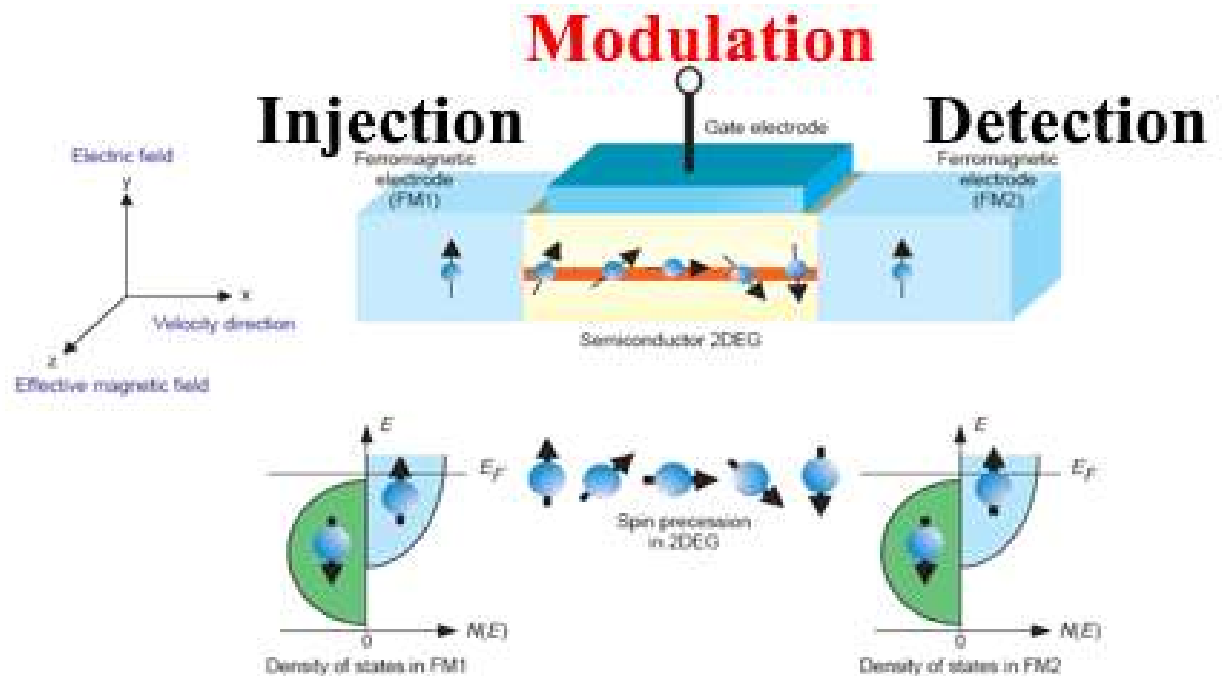
- Curie temperature above RT
- Long spin relaxation time
- Strong spin orbit coupling

Graphene

- Long spin diffusion length
- Weak spin orbit coupling

How to work?

- Spin injection
- Spin modulation
- Spin detection

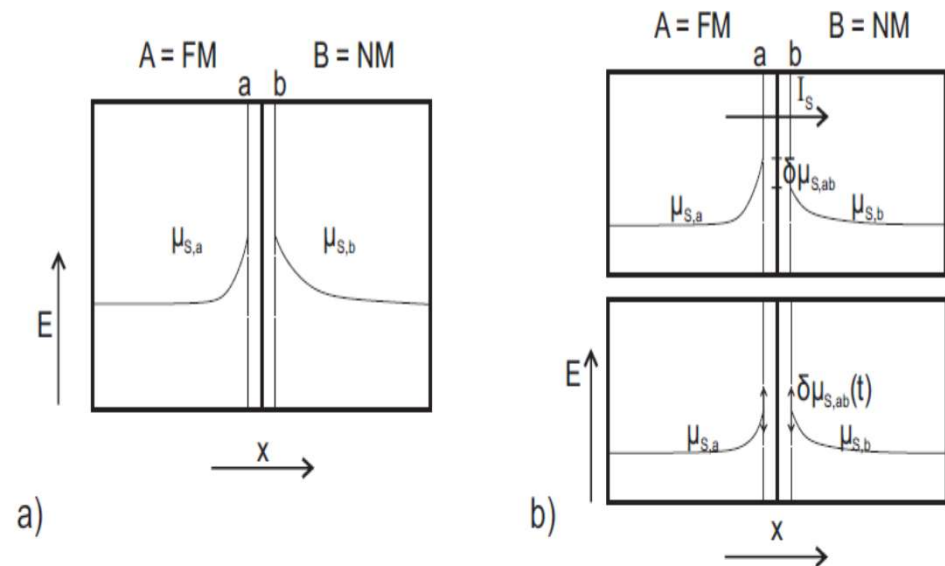
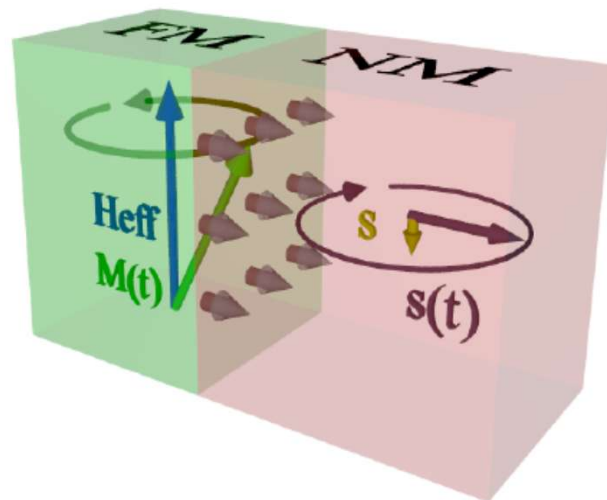


Spin injection

- Spin dynamic pumping
- Electrical injection
 - Schottky barrier
 - Oxide tunneling barrier
 - Spin Esaki diode
 - Hot electron spin injection

Spin dynamic pumping

- Processing magnetization in FM layer pumps spin current into NM layer.

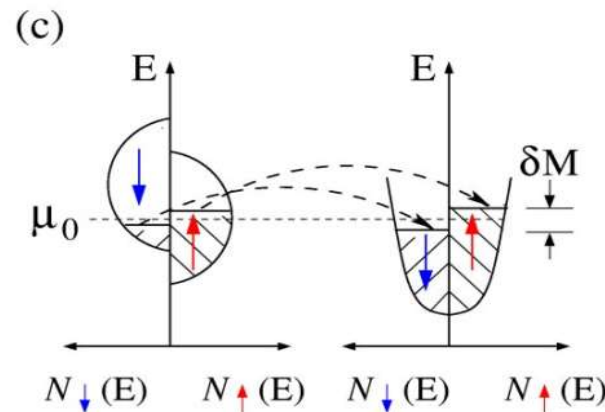
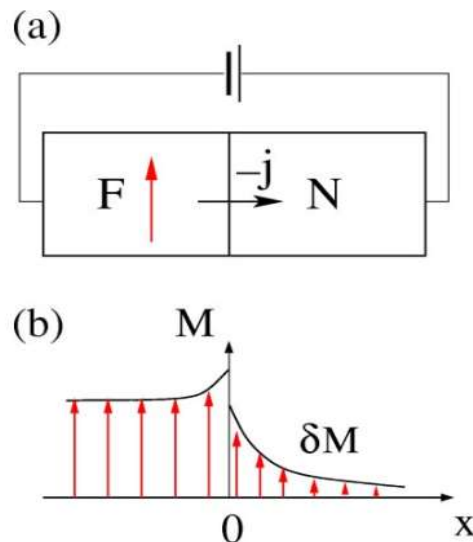


$$\frac{d\mathbf{M}}{dt} = -|\gamma|\mathbf{M} \times \mathbf{H}_{\text{eff}} + \frac{\alpha_G}{M_s}\mathbf{M} \times \frac{d\mathbf{M}}{dt}$$

$$\mathbf{I}_{s, \text{net}}^{\text{pump}} = \frac{\hbar}{4\pi} \left(\Re(\tilde{g}^{\uparrow\downarrow}) \mathbf{e}_M \times \frac{d\mathbf{e}_M}{dt} - \Im(\tilde{g}^{\uparrow\downarrow}) \frac{d\mathbf{e}_M}{dt} \right)$$

Electrical spin injection theory

- Spin polarized current is injected from FM to semiconductor
- Current spin polarizability $P_I = \frac{I_{\uparrow} - I_{\downarrow}}{I_{\uparrow} + I_{\downarrow}}$
- Main problem :conductance mismatch
- Spin injection efficiency $P_j = \frac{R_F P_{\sigma F} + R_c P_{\Sigma}}{R_F + R_c + R_N}$



Different contact

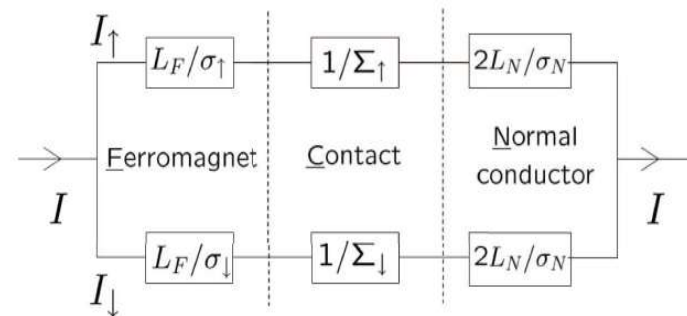
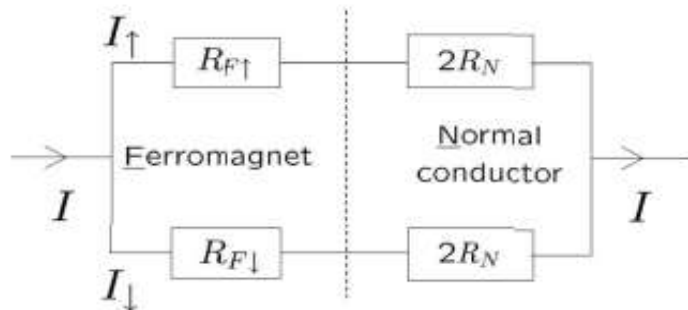
Spin injection efficiency: $P_j = \frac{R_F P_{\sigma F} + R_c P_{\Sigma}}{R_F + R_c + R_N}$

- Transparent contact: $R_c \ll R_N, R_F$

$$P_j = \frac{R_F}{R_N + R_F} P_{\sigma F}$$

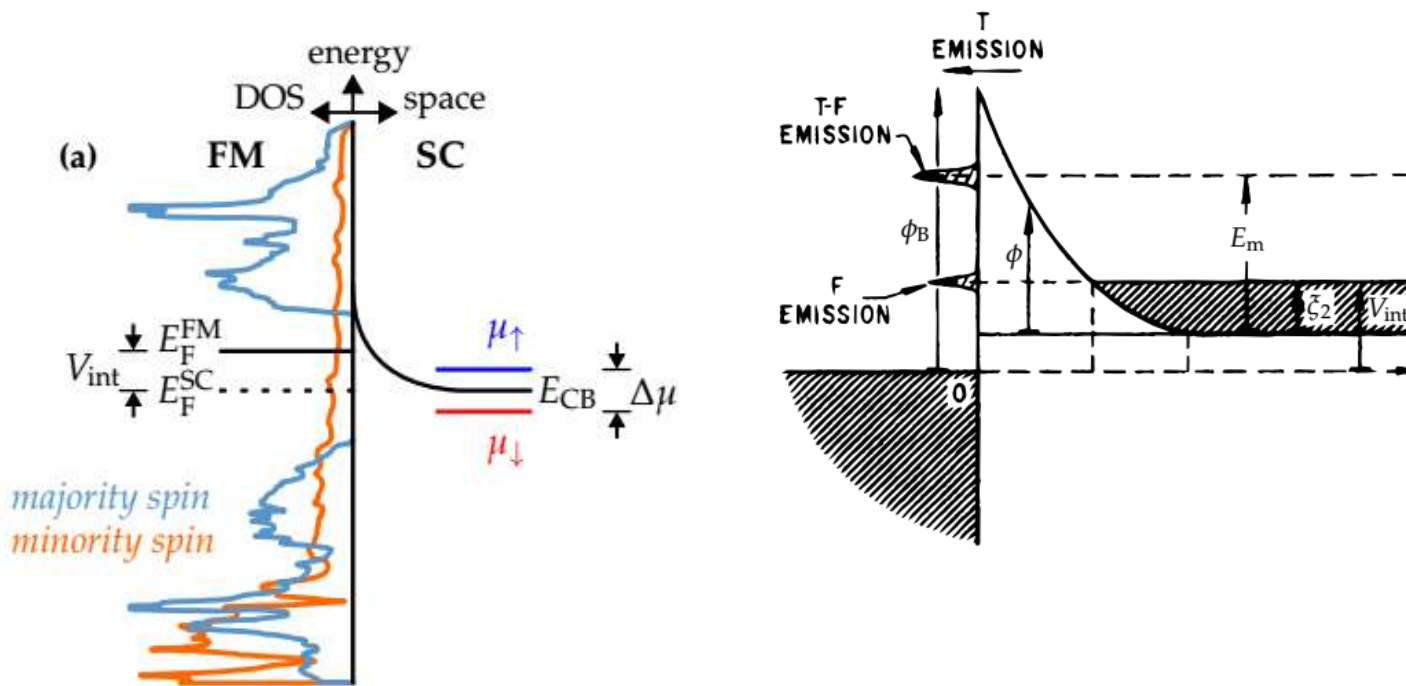
- Tunneling contact: $R_c \gg R_F, R_N$

$$P_j = P_{\Sigma}$$



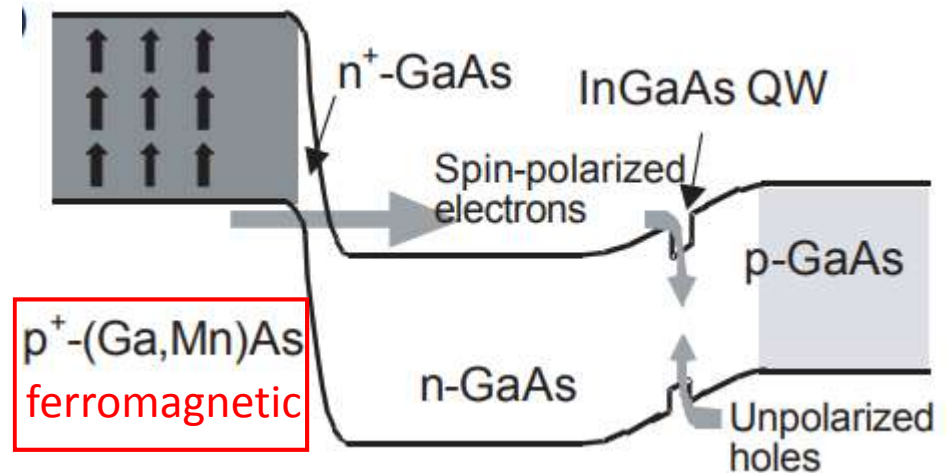
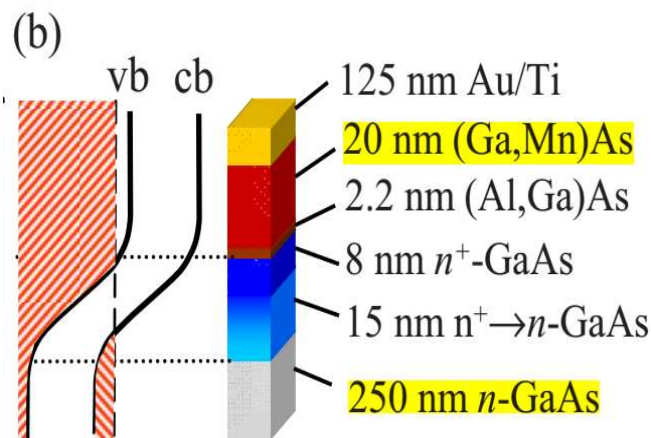
Schottky barrier

- ferromagnetic Heusler alloy Co_2FeSi /n-type GaAs



Spin Esaki diode

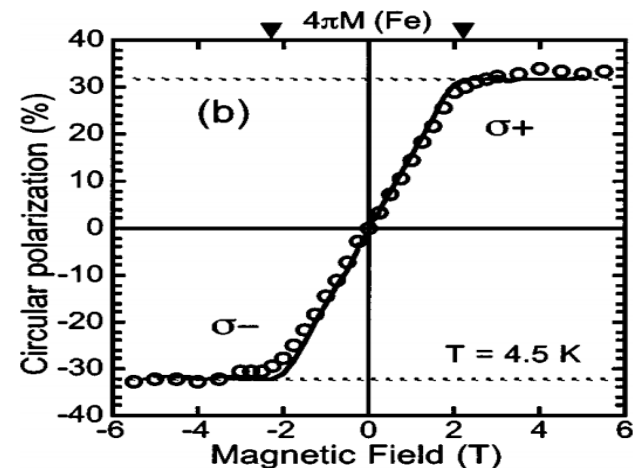
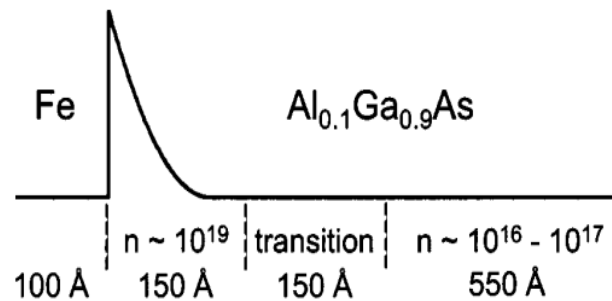
- Under a small reverse bias electrons from VB of (Ga, Mn)As tunnel to CB of GaAs.
- The conversion of spin-polarized electrons via Esaki tunneling leaves its mark in a bias dependence of the spin-injection efficiency, which at maximum reaches the value of 50%.



Improvement for depletion region

- The use of a thin, heavily doped surface region reduces the depletion width as well as the effective barrier height, significantly enhancing the probability for tunneling.

$$P = \exp\left(-\frac{4\pi\sqrt{2m\phi}}{h}d\right)$$

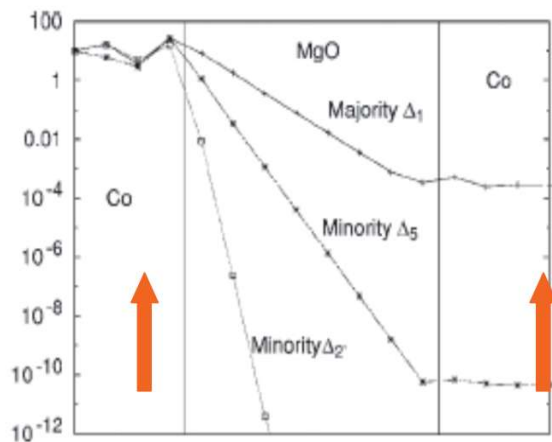


Oxide barrier

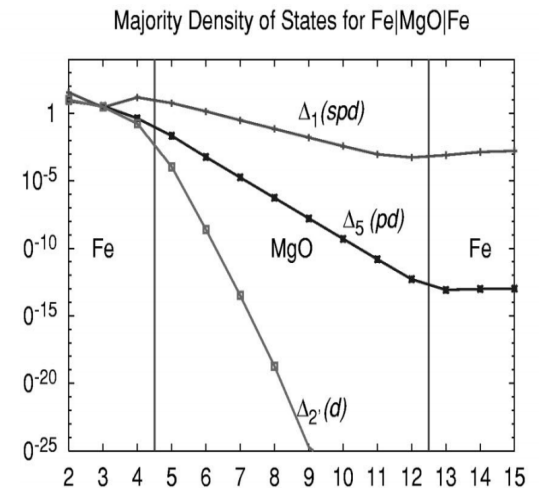
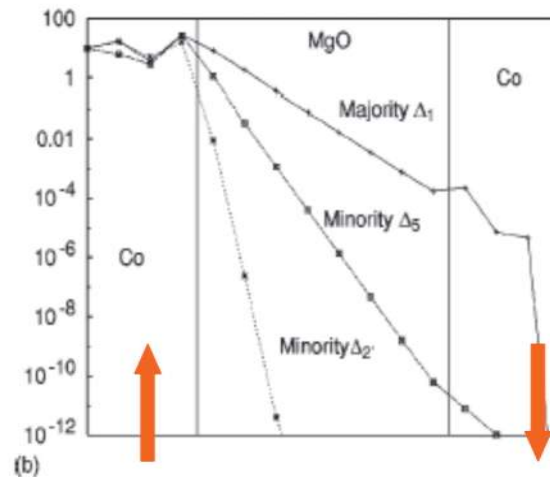
- Spin dependent conductance:
- Spin efficiency:

$$G = e^2 \rho(E_F) \tilde{D}$$

$$P_j = P_\Sigma$$

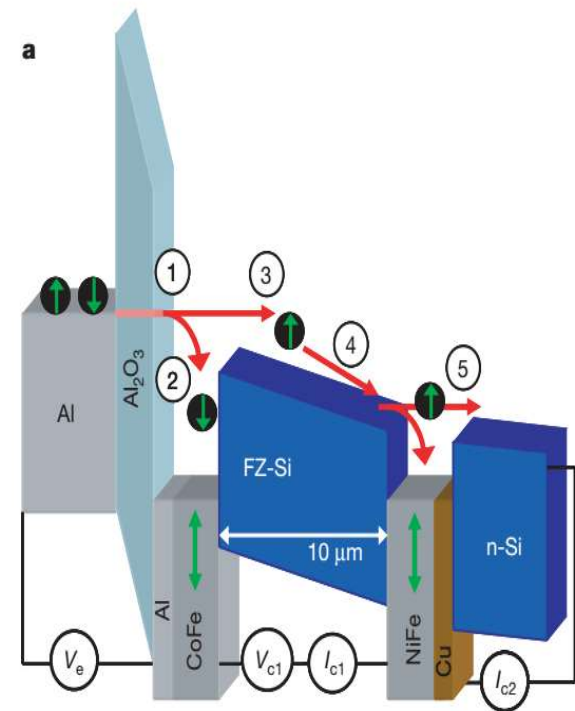


AP



Hot electron spin injection

- This method to achieve spin injection is by spin-dependent ballistic hot-electron filtering through ferromagnetic thin films.
- The exponential spin selective mean free path dependence in the ferromagnetic films create very large spin polarizations. In principle, this can approach 100%, allowing effective injection and detection at cryogenic and room temperature.



Spin detection

- ISHE
- Silsbee-Johnson spin-charge coupling

AHE, SHE and ISHE

- Charge current: $\mathbf{j}_c = \sigma \mathbf{E} + D \nabla n$
- Spin current: $j_{s,ij} = -\frac{\hbar}{2} \mu E_i s_j - \frac{\hbar}{2} D \frac{\partial s_j}{\partial x_i}$
- Including spin-orbit coupling and anomalous current density

$$j_{c,i} = \sigma E_i + eD \frac{\partial n}{\partial x_i} - \frac{2e}{\hbar} \zeta \epsilon_{ijk} \left(\frac{\hbar}{2} \mu E_j s_k + \frac{\hbar}{2} D \frac{\partial s_k}{\partial x_j} \right)$$

$$j_{s,ij} = \frac{\hbar}{2} \mu E_i s_j + \frac{\hbar}{2} D \frac{\partial s_j}{\partial x_i} - \frac{\hbar}{2e} \zeta \epsilon_{ijk} \left(\sigma E_k + eD \frac{\partial n}{\partial x_k} \right)$$

$$\mathbf{j}_c = \frac{2e}{\hbar} \alpha_{\text{SH}} \mathbf{j}_s \times \mathbf{e}_s$$

AHE and ISHE

$$\mathbf{j}_s = \frac{\hbar}{2e} \alpha_{\text{SH}} \mathbf{e}_s \times \mathbf{j}_c$$

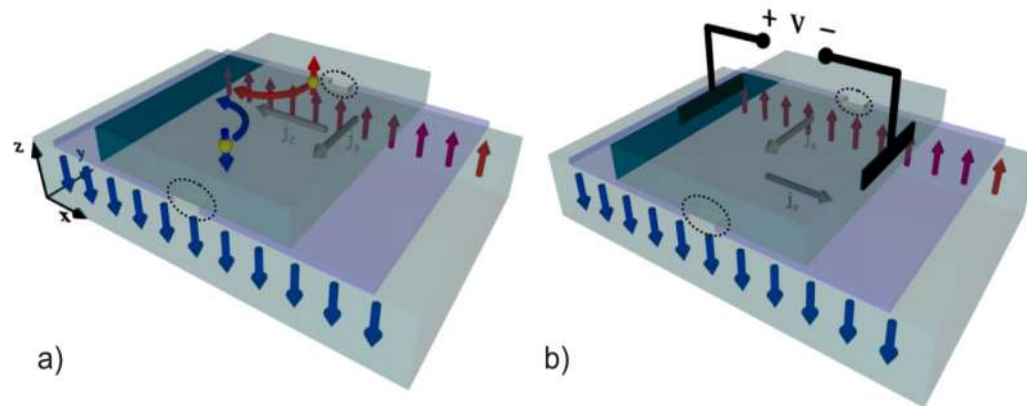
SHE

ISHE

- Spin current convert to charge current:

$$\mathbf{j}_c = \frac{2e}{\hbar} \alpha_{\text{SH}} \mathbf{j}_s \times \mathbf{e}_s$$

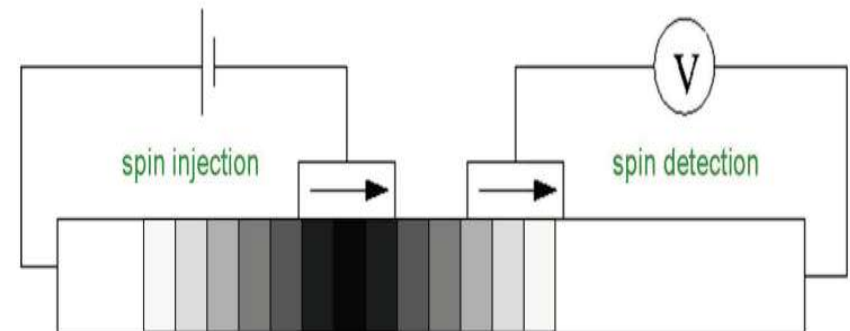
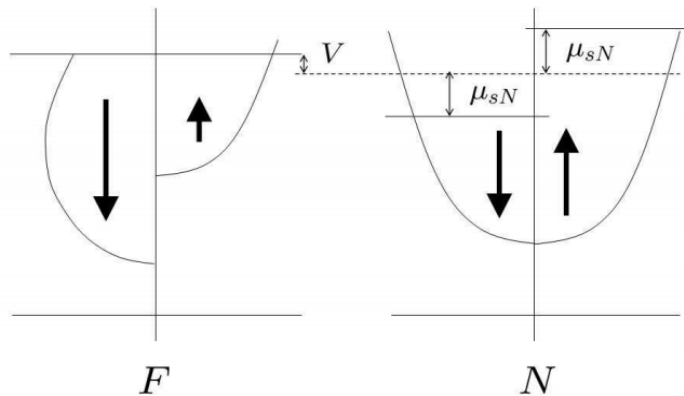
- Some challenges:
 - Voltage detectable is small (is proportional to the resistivity and α_{SH})
 - The device's dimensions are smaller than λ_{sd}



Silsbee-Johnson spin-charge coupling

- If a spin accumulation is generated in a nonmagnetic conductor that is in a proximity of a ferromagnet, a current flows in a closed circuit, or an electromotive force appears in an open circuit.

$$\text{emf} = -\frac{R_F P_{\sigma F} + R_c P_{\Sigma}}{R_F + R_c + R_N} \mu_{sN}(\infty) = -P_j \mu_{sN}(\infty)$$



Spin modulation

- Hanle spin precession frequency
- Spin diffusion length
- Conductivity
- Spin dependent barrier
- Magnetoelectric effect

Spin procession frequency

- Spin orbit coupling in 2DEG

$$H_R = \alpha_R (\boldsymbol{\sigma} \times \mathbf{k}) \cdot \hat{\mathbf{n}}$$

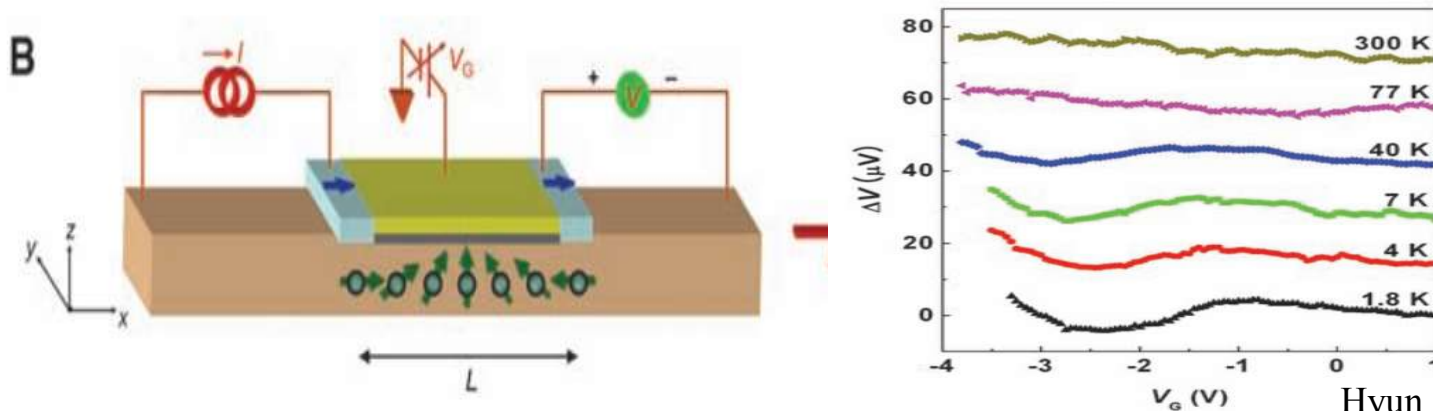
$$B_R(\vec{v}) = \frac{2}{\hbar} \alpha_R (\vec{k} \times \vec{n})$$

- Gate modulation of Rashba coefficient

$$\alpha_R = \alpha_0 \langle \mathcal{E}_v(\mathbf{r}) \rangle$$

- At low temperature, transport of electrons in 2DEG is ballistic(coherent).

$$V = A \cos(2m^* \alpha L / \hbar^2 + \varphi)$$



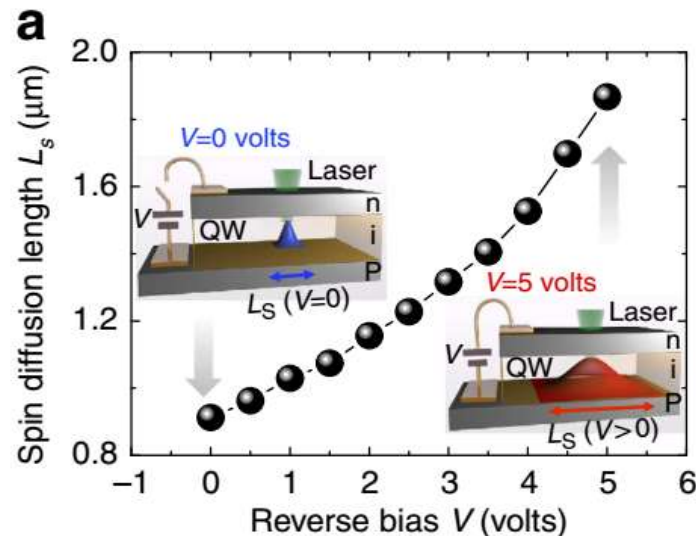
Spin diffusion length

- At RT, in semiconductor(GaAs,GaN),electron spin relaxation dominates by DP mechanism

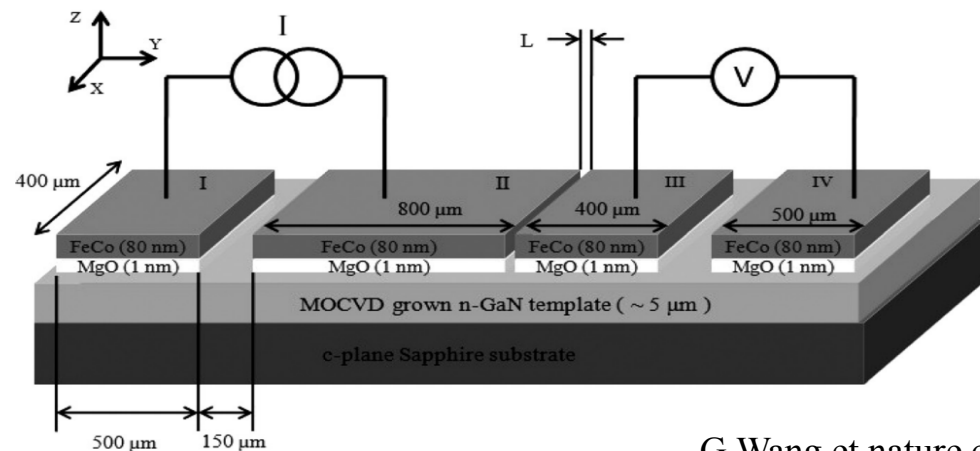
$$\begin{aligned}\Omega_{\text{Tot}}(\mathbf{k}_{//}) &= \Omega_{\text{BIA}}(\mathbf{k}_{//}) + \Omega_{\text{SIA}}(\mathbf{k}_{//}) \\ &= \left(\frac{\beta}{\hbar} + \frac{2r_{41}E}{\hbar} \right) (k_y, -k_x, 0)\end{aligned}$$

$$\frac{1}{\tau_s} = \langle \Omega_{\text{Tot}}^2 \rangle \tau_p^*$$

$$\bar{L}_s = \sqrt{D_s T_s}$$

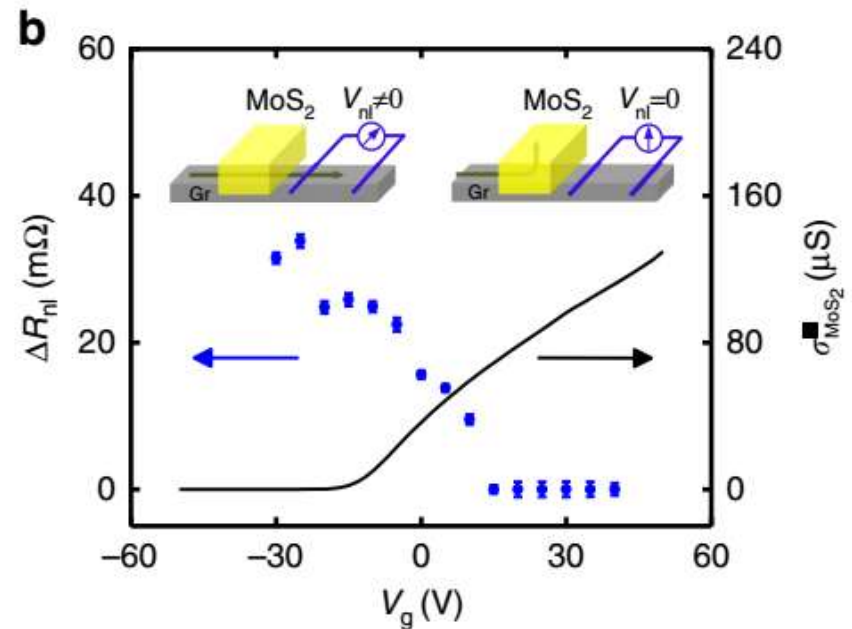
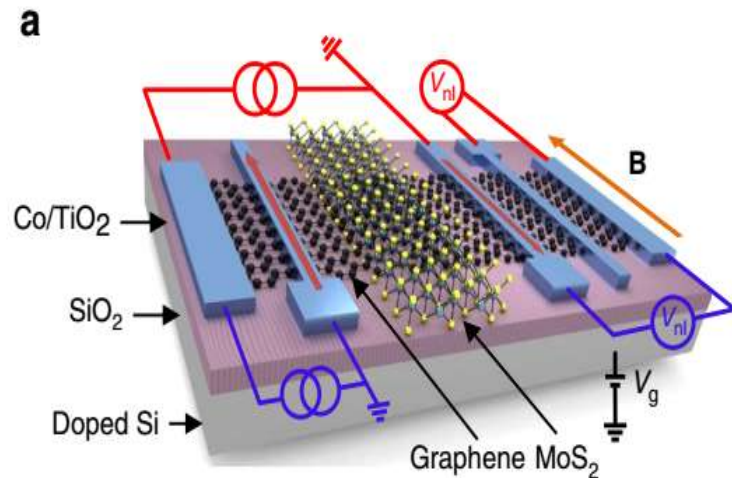


$$\mu_{sN} = \mu_{sN}(0) e^{-x/L_{sN}}$$



A spin field-effect switch

- Gate voltage modulate the conductivity of MoS₂



Spin-dependent barrier

- Using magnetic insulator rather than normal insulator as dielectric layer in FET could induce a split according to spin caused by ferromagnetic proximity.

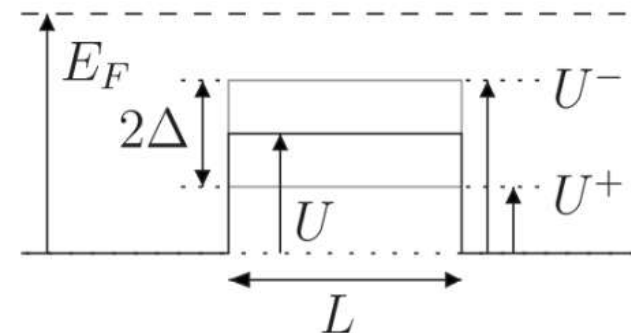
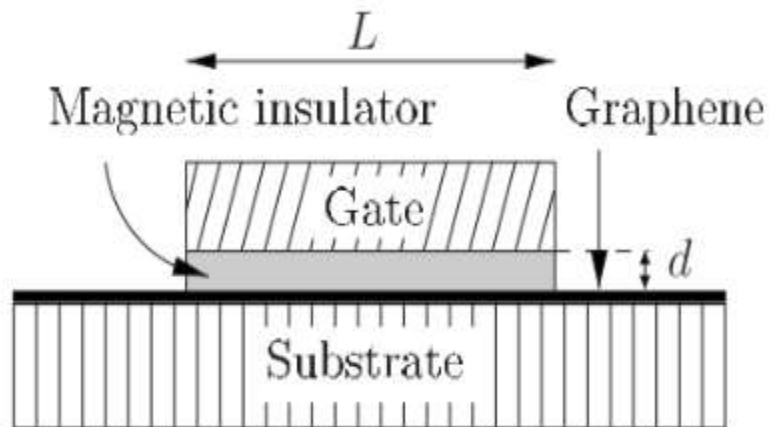
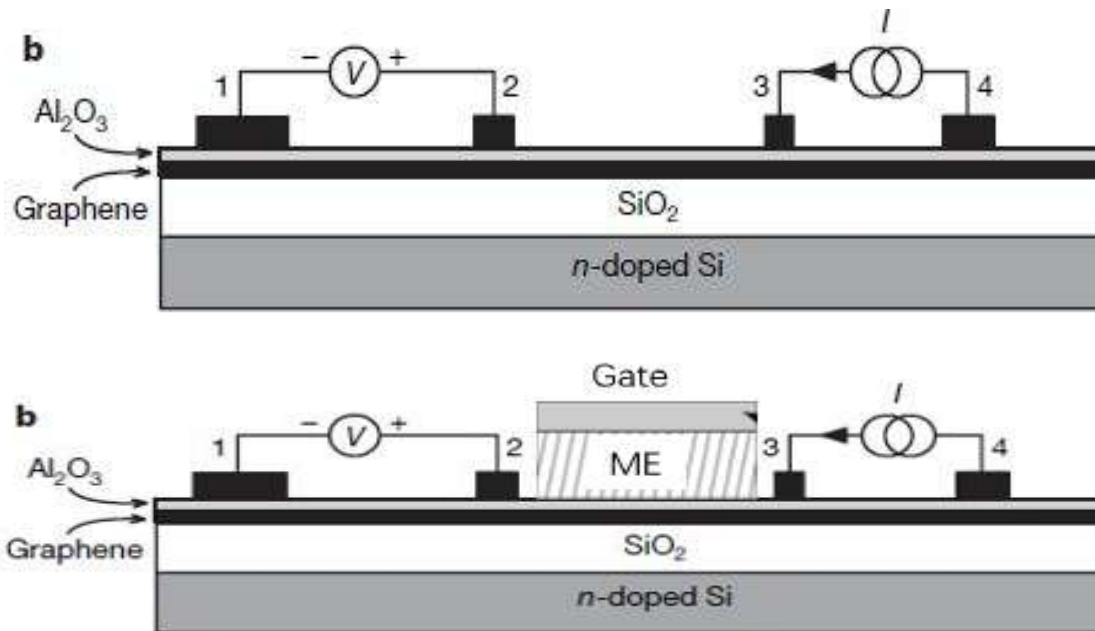


FIG. 4. Ferromagnetic proximity effect splits the barrier according to spin such that $U^\pm = U \mp \Delta$.

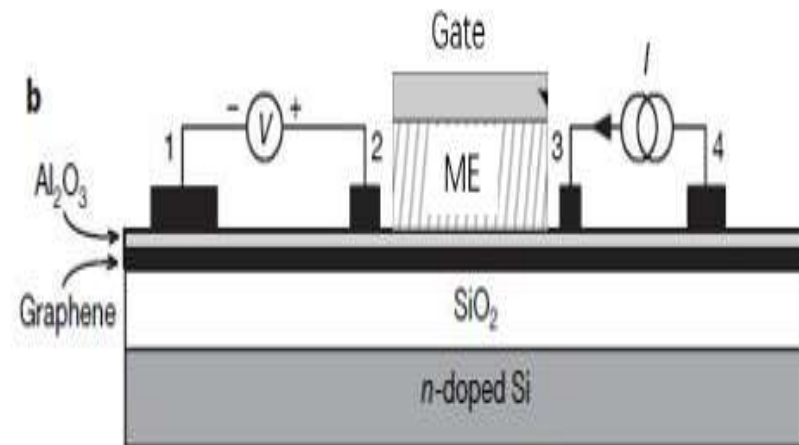
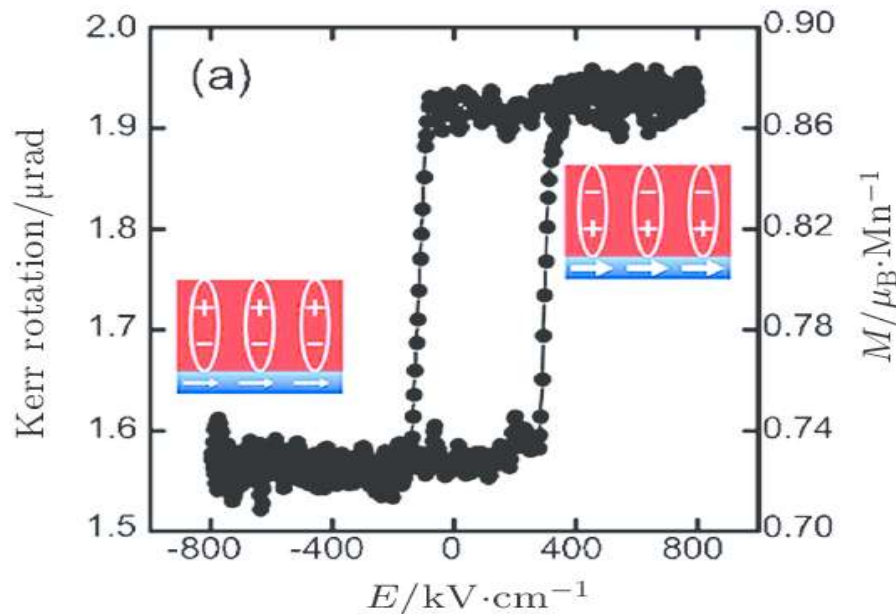
A proposal for spin FET

If such structure is deposited on the middle of a non-local spin-valve, the signal would be affected by the gate voltage.



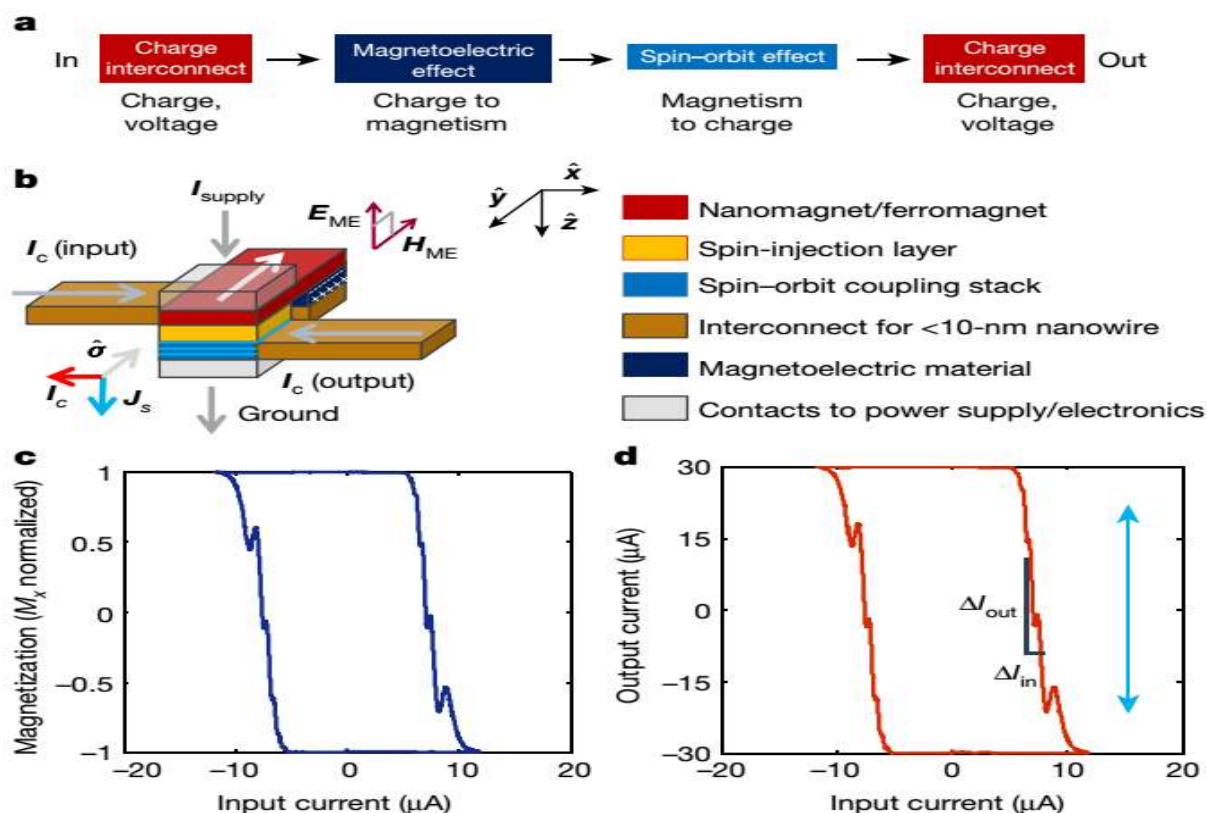
Magnetoelectric effect

- A dielectric material moving through an electric field would become magnetized. A material where such a coupling is intrinsically present is called a magnetoelectric.



Scalable energy-efficient magnetoelectric spin-orbit logic

Sasikanth Manipatruni^{1*}, Dmitri E. Nikonov¹, Chia-Ching Lin¹, Tanay A. Gosavi¹, Huichu Liu², Bhagwati Prasad³, Yen-Lin Huang^{3,4}, Everton Bonturim³, Ramamoorthy Ramesh^{3,4,5} & Ian A. Young¹



Thank you