Appendix A: Proof of the Inequality and the Expression for $B^2_{\rm COST}$ in Thm 1

This appendix proves the following inequality from the proof sketch of Thm 1 and derives the corresponding expression for B_{cost}^2 :

$$\mathbb{E}_{\theta_T' \sim \mathcal{D}_{\theta_T}} \left[e^{\lambda (C(K) - \hat{C}(K))} \right] \le e^{\frac{\lambda^2 B_{\text{cost}}^2}{8m}}, \tag{1}$$

where

$$B_{\text{cost}}^{2} := \sup_{K \in \mathcal{K}} \sigma_{w}^{4} d_{x} \cdot \rho_{S}(K)^{2} \cdot \left(128 \sum_{i=0}^{T} \rho_{M}(K)^{4(T-i)} + 64 \sum_{0 \le i < j \le T} \rho_{M}(K)^{4T-2(i+j)} \right), \tag{2}$$

for $|\lambda| \leq \inf_{K \in \mathcal{K}} \frac{1}{4\sigma_w^2 \cdot \rho_S(K) \cdot \rho_M(K)^{2T}}$.

Proof: For notational simplicity, we write \mathbb{E} in place of $\mathbb{E}_{\theta_T' \sim \mathcal{D}_{\theta_T}}$ throughout the proof. The closed-loop system

$$x_{t+1} = Ax_t + Bu_t + w_t,$$

$$u_t = -Kx_t,$$

can be written in the form

$$x_{t+1} = Mx_t + w_t,$$

where M := A - BK. We assume the system starts from the initial state $x_0 = 0$. The control performance is measured by the quadratic cost:

$$C_q(K, \theta_T) = \sum_{t=0}^T x_t^{\top} S x_t,$$

where $S := Q + K^{\top}RK$. We first expand x_t in terms of the noise variables:

$$x_t = \sum_{k=0}^{t-1} M^{t-1-k} w_k.$$

Substituting this into the cost expression, we obtain:

$$C_q(K, \theta_T) = \sum_{t=0}^{T} \sum_{i,j=0}^{t-1} w_i^{\top} M^{t-1-i\top} S M^{t-1-j} w_j.$$

Define $D_{t,ij} := M^{t-1-i\top} SM^{t-1-j}$. Then we have

$$C_q(K, \theta_T) = \sum_{t=0}^{T} \sum_{i=0}^{t-1} w_i^{\top} D_{t,ij} w_j.$$

By grouping terms over time, we define:

$$H_i = \sum_{t=i+1}^{T} D_{t,ii}, \quad H_{ij} = \sum_{t=\max(i,j)+1}^{T} D_{t,ij}, \quad i \neq j,$$

and we rewrite the cost as:

$$C_{q}(K, \theta_{T}) = \sum_{i=0}^{T-1} w_{i}^{\top} H_{i} w_{i} + \sum_{0 \le i < j \le T-1} \left(w_{i}^{\top} H_{ij} w_{j} + w_{j}^{\top} H_{ij}^{\top} w_{i} \right).$$

Using the spectral decomposition $H_i = U\Lambda U^{\top}$ and defining $\tilde{w}_i := U^{\top} w_i$, we have

$$w_i^{\top} H_i w_i = \tilde{w}_i^{\top} \Lambda \tilde{w}_i,$$

where $\Lambda = \text{diag}(\mu_1, \dots, \mu_{d_x})$ is a diagonal matrix of eigenvalues. Since Λ is diagonal and positive semidefinite, we can upper bound the quadratic form as

$$\tilde{w}_i^\top \Lambda \tilde{w}_i = \sum_{j=1}^{d_x} \mu_j \tilde{w}_{i,j}^2 \leq \|\Lambda\| \cdot \sum_{j=1}^{d_x} \tilde{w}_{i,j}^2 = \|\Lambda\| \cdot \|\tilde{w}_i\|_2^2,$$

due to $\|\Lambda\| = \max_j \mu_j$. The inequality follows from the fact that each $\mu_i \leq \|\Lambda\|$.

Since the noise vector w_i has independent sub-Gaussian entries with parameter σ_w^2 , and $U \in \mathbb{R}^{d_x \times d_x}$ is an orthogonal matrix, then for any $v \in \mathbb{R}^{d_x}$ with $||v||_2 = 1$, by Def. 1 and Def. 2 in the main body of the paper, which define sub-Gaussian random variables and sub-Gaussian random vectors, we have

$$\mathbb{E}\left[e^{\lambda v^{\top} \tilde{w}_i}\right] = \mathbb{E}\left[e^{\lambda (Uv)^{\top} w_i}\right] \leq \exp\left(\frac{\lambda^2 \sigma_w^2}{2}\right),$$

where we used that Uv is a unit vector and that linear combinations of independent sub-Gaussian variables remain sub-Gaussian. Hence, \tilde{w}_i is a sub-Gaussian vector with parameter σ_w^2 .

We note that the definition of a sub-Gaussian vector guarantees sub-Gaussian concentration along any unit direction $v \in \mathbb{R}^{d_x}$. In particular, taking $v = e_j$, the j-th standard basis vector, selects the j-th coordinate of \tilde{w}_i , i.e., $\tilde{w}_{i,j} = e_j^\top \tilde{w}_i$. Since $||e_j||_2 = 1$, by the Def. 1, applying the inequality above with $v = e_j$ gives:

$$\mathbb{E}\left[e^{\lambda \tilde{w}_{i,j}}\right] = \mathbb{E}\left[e^{\lambda e_j^\top \tilde{w}_i}\right] \leq \exp\left(\frac{\lambda^2 \sigma_w^2}{2}\right),$$

which shows that each entry $\tilde{w}_{i,j}$ is a sub-Gaussian random variable with parameter σ_w^2 .

Since each $\tilde{w}_{i,j}$ is a sub-Gaussian variable with parameter σ_w^2 , we apply the result from [1, Appendix B], which gives:

$$\mathbb{E}\left[e^{\lambda(\tilde{w}_{i,j}^2 - \mathbb{E}[\tilde{w}_{i,j}^2])}\right] \leq \exp(16\lambda^2 \sigma_w^4).$$

Let $\Lambda = \text{diag}(\mu_1, \dots, \mu_{d_x})$, where μ_i is the i^{th} eigenvalue of H_i . Then we have

$$\tilde{w}_i^{\top} \Lambda \tilde{w}_i = \sum_{j=1}^{d_x} \mu_j \tilde{w}_{i,j}^2,$$

and we can upper bound the centered sum as

$$\sum_{i=1}^{d_x} \mu_j(\tilde{w}_{i,j}^2 - \mathbb{E}[\tilde{w}_{i,j}^2]) \leq \|\Lambda\| \cdot \sum_{i=1}^{d_x} (\tilde{w}_{i,j}^2 - \mathbb{E}[\tilde{w}_{i,j}^2]),$$

due to $\|\Lambda\| = \max_j |\mu_j|$. Applying the bound for each $\tilde{w}_{i,j}$ and combining, we obtain:

$$\mathbb{E}\left[e^{\lambda(\tilde{w}_{i}^{\top}\Lambda\tilde{w}_{i}-\mathbb{E}[\tilde{w}_{i}^{\top}\Lambda\tilde{w}_{i}])}\right] = \mathbb{E}\left[e^{\lambda(w_{i}^{\top}H_{i}w_{i}-\mathbb{E}[w_{i}^{\top}H_{i}w_{i}])}\right]$$

$$\leq \exp\left(16\lambda^{2}\|\Lambda\|^{2}\sigma_{w}^{4}d_{x}\right),$$

for all λ such that $|\lambda| \le \frac{1}{4\sigma_w^2 \|\Lambda\|}$, by the result from [1, Appendix B]. That is,

$$\mathbb{E}\left[e^{\lambda(w_i^\top H_i w_i - \mathbb{E}[w_i^\top H_i w_i])}\right] \le \exp\left(16\lambda^2 ||H_i||^2 \sigma_w^4 d_x\right),\,$$

for all λ such that $|\lambda| \leq \frac{1}{4\sigma_w^2 ||H_i||}$, by the result from [1, Appendix B].

We now provide a detailed derivation for bounding the moment generating function (MGF) of the cross term $w_i^{\top} H_{ij} w_j$, where w_i and w_j are independent random vectors with sub-Gaussian entries and $i \neq j$. Our goal is to upper bound the centered MGF:

$$\mathbb{E}\left[e^{\lambda(w_i^\top H_{ij}w_j - \mathbb{E}[w_i^\top H_{ij}w_j])}\right].$$

We show how to express the bilinear form $w_i^{\top} H_{ij} w_j$ as a quadratic form, which allows us to apply the result from [1, Appendix B].

Let $\zeta_{ij} \in \mathbb{R}^{2d_x}$ be the concatenated vector

$$\zeta_{ij} := \begin{bmatrix} w_i \\ w_j \end{bmatrix},$$

and define the symmetric block matrix $\Psi_{ij} \in \mathbb{R}^{2d_x \times 2d_x}$ as

$$\Psi_{ij} := egin{bmatrix} 0 & H_{ij} \ H_{ij}^{ op} & 0 \end{bmatrix}.$$

We compute the quadratic form $\zeta_{ij}^{\top} \Psi_{ij} \zeta_{ij}$:

$$\zeta_{ij}^{\top} \Psi_{ij} \zeta_{ij} = \begin{bmatrix} w_i^{\top} & w_j^{\top} \end{bmatrix} \begin{bmatrix} 0 & H_{ij} \\ H_{ij}^{\top} & 0 \end{bmatrix} \begin{bmatrix} w_i \\ w_j \end{bmatrix}$$
$$= w_i^{\top} H_{ij} w_j + w_j^{\top} H_{ij}^{\top} w_i$$
$$= 2w_i^{\top} H_{ij} w_i,$$

where the last step uses that both terms are scalars and equal. Thus, we have:

$$w_i^{\top} H_{ij} w_j = \frac{1}{2} \zeta_{ij}^{\top} \Psi_{ij} \zeta_{ij}.$$

This shows that the centered bilinear form can be written as a centered quadratic form over the augmented vector ζ_{ij} .

$$w_i^{\top} H_{ij} w_j - \mathbb{E}[w_i^{\top} H_{ij} w_j] = \frac{1}{2} \left(\zeta_{ij}^{\top} \Psi_{ij} \zeta_{ij} - \mathbb{E}[\zeta_{ij}^{\top} \Psi_{ij} \zeta_{ij}] \right).$$

Then we have

$$\mathbb{E}\left[e^{\lambda(w_i^\top H_{ij}w_j - \mathbb{E}[w_i^\top H_{ij}w_j])}\right] = \mathbb{E}\left[e^{\frac{\lambda}{2}(\zeta_{ij}^\top \Psi_{ij}\zeta_{ij} - \mathbb{E}[\zeta_{ij}^\top \Psi_{ij}\zeta_{ij}])}\right].$$

We now perform a spectral decomposition of the symmetric matrix Ψ_{ij} : let

$$\Psi_{ij} = \mathcal{U}_{ij} \Sigma_{ij} \mathcal{U}_{ij}^{\top},$$

where $\Sigma_{ij} = \operatorname{diag}(v_1, \dots, v_{2d_x})$ is a diagonal matrix of eigenvalues and $\mathcal{U}_{ij} \in \mathbb{R}^{2d_x \times 2d_x}$ is orthogonal. Define the rotated noise vector

$$\tilde{\zeta}_{ij} := \mathcal{U}_{ij}^{\top} \zeta_{ij} \in \mathbb{R}^{2d_x}.$$

Then we can rewrite the quadratic form as

$$\zeta_{ij}^{\top} \Psi_{ij} \zeta_{ij} = \tilde{\zeta}_{ij}^{\top} \Sigma_{ij} \tilde{\zeta}_{ij} = \sum_{k=1}^{2d_x} v_k \tilde{\zeta}_{ij,k}^2.$$

Since w_i and w_j are independent and each has i.i.d. sub-Gaussian entries with parameter σ_w^2 , the concatenated vector ζ_{ij} is also sub-Gaussian with parameter σ_w^2 . Moreover,

similar to the analysis of \tilde{w}_i , since \mathcal{U}_{ij} is orthogonal, $\tilde{\zeta}_{ij} = \mathcal{U}_{ij}^{\top} \zeta_{ij}$ is sub-Gaussian vector with parameter σ_w^2 .

By the definition of a sub-Gaussian vector Def. 2, for any $v \in \mathbb{R}^{2d_x}$ with $||v||_2 = 1$, we have

$$\mathbb{E}\left[e^{\lambda v^{\top} \tilde{\zeta}_{ij}}\right] = \mathbb{E}\left[e^{\lambda (\mathcal{U}_{ij}v)^{\top} \zeta_{ij}}\right] \leq \exp\left(\frac{\lambda^2 \sigma_w^2}{2}\right).$$

Hence, $\tilde{\zeta}_{ij}$ is a sub-Gaussian vector with parameter σ_w^2 .

In particular, for each k, taking $v = e_k$ (the k-th standard basis vector in \mathbb{R}^{2d}), we obtain

$$\mathbb{E}\left[e^{\lambda \tilde{\zeta}_{ij,k}}\right] = \mathbb{E}\left[e^{\lambda e_k^\top \tilde{\zeta}_{ij}}\right] \leq \exp\left(\frac{\lambda^2 \sigma_w^2}{2}\right),$$

which shows that each coordinate $\tilde{\zeta}_{ij,k}$ is a sub-Gaussian random variable with parameter σ_w^2 .

Since each $\tilde{\zeta}_{ij,k}$ is sub-Gaussian with parameter σ_w^2 , we apply the scalar MGF bound from [1, Appendix B], which gives:

$$\mathbb{E}\left[e^{\lambda(\tilde{\zeta}_{ij,k}^2 - \mathbb{E}[\tilde{\zeta}_{ij,k}^2])}\right] \leq \exp(16\lambda^2 \sigma_w^4).$$

Then, the centered quadratic form becomes:

$$\tilde{\zeta}_{ij}^{\top} \Sigma_{ij} \tilde{\zeta}_{ij} - \mathbb{E}[\tilde{\zeta}_{ij}^{\top} \Sigma_{ij} \tilde{\zeta}_{ij}] = \sum_{k=1}^{2d_x} \nu_k (\tilde{\zeta}_{ij,k}^2 - \mathbb{E}[\tilde{\zeta}_{ij,k}^2]),$$

which can be bounded as

$$\sum_{k=1}^{2d_x} v_k(\tilde{\zeta}_{ij,k}^2 - \mathbb{E}[\tilde{\zeta}_{ij,k}^2]) \leq \|\Sigma_{ij}\| \cdot \sum_{k=1}^{2d_x} (\tilde{\zeta}_{ij,k}^2 - \mathbb{E}[\tilde{\zeta}_{ij,k}^2]),$$

where $\|\Sigma_{ij}\| := \max_k |v_k|$.

To bound the moment generating function of the centered bilinear form, we recall that

$$\mathbb{E}\left[e^{\lambda(w_i^\top H_{ij}w_j - \mathbb{E}[w_i^\top H_{ij}w_j])}\right] = \mathbb{E}\left[e^{\frac{\lambda}{2}(\zeta_{ij}^\top \Psi_{ij}\zeta_{ij} - \mathbb{E}[\zeta_{ij}^\top \Psi_{ij}\zeta_{ij}])}\right].$$

By the result from [1, Appendix B], applying the composition of independent sub-Gaussian quadratic bounds and combining the $2d_x$ terms, we obtain:

$$\mathbb{E}\left[\exp\left(\lambda\sum_{k=1}^{2d_x}(\tilde{\zeta}_{ij,k}^2-\mathbb{E}[\tilde{\zeta}_{ij,k}^2])\right)\right] \leq \exp\left(16\cdot 2d_x\lambda^2\sigma_w^4\right).$$

By replacing λ as $\frac{\lambda \|\Sigma_{ij}\|}{2}$ to account for the factor in front of the quadratic form, we conclude:

$$\mathbb{E}\left[e^{\frac{\lambda}{2}(\zeta_{ij}^{\top}\Psi_{ij}\zeta_{ij}-\mathbb{E}[\zeta_{ij}^{\top}\Psi_{ij}\zeta_{ij}])}\right] \leq \exp\left(8\lambda^{2}\|\Sigma_{ij}\|^{2}\sigma_{w}^{4}d_{x}\right),$$

which implies

$$\mathbb{E}\left[e^{\lambda(w_i^\top H_{ij}w_j - \mathbb{E}[w_i^\top H_{ij}w_j])}\right] \leq \exp\left(8\lambda^2 \|\Sigma_{ij}\|^2 \sigma_w^4 d_x\right),\,$$

for all λ such that $|\lambda| \leq \frac{1}{4\sigma_w^2 \|\Sigma_{ii}\|}$. That is,

$$\mathbb{E}\left[e^{\lambda\left(w_i^\top H_{ij}w_j - \mathbb{E}[w_i^\top H_{ij}w_j]\right)}\right] \leq \exp\left(8\lambda^2 \|H_{ij}\|^2 \sigma_w^4 d_x\right),\,$$

for all λ such that $|\lambda| \leq \frac{1}{4\sigma_w^2 \|H_{ij}\|}$.

We first consider the deviation of the quadratic cost for a single trajectory $\theta_T' \sim \mathcal{D}_{\theta_T}$. From the previous derivation

using the structure of the cost in terms of $\{H_i, H_{ij}\}$ and the sub-Gaussianity of w_t , we obtain:

$$\begin{split} & \mathbb{E}\left[e^{\lambda\left(\mathbb{E}\left[C_{q}(K, \theta_{T}^{\prime})\right] - C_{q}(K, \theta_{T}^{\prime})\right)}\right] \\ & = \mathbb{E}\left[\exp\left(\lambda\sum_{i=0}^{T-1}\left(\mathbb{E}\left[w_{i}^{\top}H_{i}w_{i}\right] - w_{i}^{\top}H_{i}w_{i}\right) \right. \\ & + \sum_{0 \leq i < j \leq T-1}\left(\mathbb{E}\left[w_{i}^{\top}H_{ij}w_{j}\right] - w_{i}^{\top}H_{ij}w_{j}\right)\right)\right], \\ & \leq \mathbb{E}\left[\exp\left(\lambda\sum_{i=0}^{T}\left(\mathbb{E}\left[w_{i}^{\top}H_{i}w_{i}\right] - w_{i}^{\top}H_{i}w_{i}\right) \right. \\ & + \sum_{0 \leq i < j \leq T}\left(\mathbb{E}\left[w_{i}^{\top}H_{ij}w_{j}\right] - w_{i}^{\top}H_{ij}w_{j}\right)\right)\right], \end{split}$$

for all $|\lambda| \leq \frac{1}{4\sigma_w^2 \max\left\{\max_i \|H_i\|, \max_{i < j} \|H_{ij}\|\right\}}$. We apply the previous bounds for the diagonal and off-diagonal terms. By Jensen's inequality and the fact that the moment generating function of a sum is upper bounded by the product of the individual MGFs (for sub-exponential variables), we obtain:

$$\mathbb{E}\left[e^{\lambda\left(\mathbb{E}\left[C_{q}(K,\theta_{T}')\right]-C_{q}(K,\theta_{T}')\right)}\right]$$

$$\leq \exp\left(\lambda^{2}\sigma_{w}^{4}d_{x}\left(\sum_{i=0}^{T}16\|H_{i}\|^{2}+\sum_{0\leq i< j\leq T}8\|H_{ij}\|^{2}\right)\right),$$

Let $\hat{C}(K) := \frac{1}{m} \sum_{i=1}^{m} C_q(K, \theta_{i,T})$ denote the empirical cost computed from m i.i.d. trajectories $\{\theta_{i,T}\}_{i=1}^{m}$.

Then the deviation between the expected and empirical cost can be written as

$$\mathbb{E}\left[e^{\lambda(C(K)-\hat{C}(K))}\right] = \mathbb{E}\left[e^{\lambda\left(\mathbb{E}\left[C_q(K,\theta_T)\right]-\frac{1}{m}\sum_{i=1}^{m}C_q(K,\theta_{i,T})\right)}\right].$$

We now use linearity of expectation to rewrite:

$$\mathbb{E}[C_q(K, \theta_T)] - \frac{1}{m} \sum_{i=1}^m C_q(K, \theta_{i,T})$$

$$= \frac{1}{m} \sum_{i=1}^m \left(\mathbb{E}[C_q(K, \theta_T)] - C_q(K, \theta_{i,T}) \right).$$

Therefore,

$$\begin{split} & \mathbb{E}\left[e^{\lambda(C(K)-\hat{C}(K))}\right] \\ & = \mathbb{E}\left[e^{\lambda\cdot\frac{1}{m}\sum_{i=1}^{m}\left(\mathbb{E}\left[C_{q}(K,\theta_{T})\right]-C_{q}(K,\theta_{i,T})\right)\right]} \\ & = \mathbb{E}\left[\prod_{i=1}^{m}e^{\frac{\lambda}{m}\left(\mathbb{E}\left[C_{q}(K,\theta_{T})\right]-C_{q}(K,\theta_{i,T})\right)}\right]. \end{split}$$

Since the $C_q(K,\theta_{i,T})$ are independent (due to the independent system parameters and noise sequence), the random variable $\mathbb{E}[C_q(K,\theta_T)] - C_q(K,\theta_{i,T})$ is also independent. Hence, the product of exponentials factorizes:

$$\begin{split} & \mathbb{E}\left[\prod_{i=1}^{m} e^{\frac{\lambda}{m}\left(\mathbb{E}\left[C_{q}(K,\theta_{T})\right] - C_{q}(K,\theta_{i,T})\right)}\right] \\ & = \prod_{i=1}^{m} \mathbb{E}\left[e^{\frac{\lambda}{m}\left(\mathbb{E}\left[C_{q}(K,\theta_{T})\right] - C_{q}(K,\theta_{i,T})\right)}\right]. \end{split}$$

Thus, we conclude:

$$\mathbb{E}\left[e^{\lambda(C(K)-\hat{C}(K))}\right] = \prod_{i=1}^{m} \mathbb{E}\left[e^{\frac{\lambda}{m}\left(\mathbb{E}\left[C_{q}(K,\theta_{T})\right]-C_{q}(K,\theta_{i,T})\right)}\right].$$

From the previous derivation for a single trajectory, we know that for all $\lambda' \in \left[0, \frac{1}{4\sigma_w^2 \max\left\{\max_i \|H_i\|, \max_{i < j} \|H_{ij}\|\right\}}\right]$, we have:

$$\mathbb{E}\left[e^{\lambda'\left(\mathbb{E}\left[C_q(K,\theta_T)\right]-C_q(K,\theta_{i,T})\right)}\right]$$

$$\leq \exp\left(\lambda'^2\sigma_w^4d_x\cdot\left(\sum_{i=0}^T 16\|H_i\|^2 + \sum_{0\leq i\leq j\leq T} 8\|H_{ij}\|^2\right)\right).$$

We apply this bound with $\lambda' = \lambda/m$, and use that all m terms are identical. This gives:

$$\mathbb{E}\left[e^{\lambda(C(K)-\hat{C}(K))}\right] = \prod_{i=1}^{m} \mathbb{E}\left[e^{\frac{\lambda}{m}\left(\mathbb{E}\left[C_{q}(K,\theta_{T})\right]-C_{q}(K,\theta_{i,T})\right)}\right]$$

$$\leq \left(\exp\left(\frac{\lambda^{2}}{m^{2}}\cdot\sigma_{w}^{4}d_{x}\cdot\left(\sum_{i=0}^{T}16\|H_{i}\|^{2}+\sum_{0\leq i< j\leq T}8\|H_{ij}\|^{2}\right)\right)\right)^{m}$$

$$= \exp\left(\frac{\lambda^{2}}{m}\cdot\sigma_{w}^{4}d_{x}\cdot\left(\sum_{i=0}^{T}16\|H_{i}\|^{2}+\sum_{0\leq i< j\leq T}8\|H_{ij}\|^{2}\right)\right).$$

This implies the quadratic cost is sub-exponential with variance proxy:

$$B^2 := \sigma_w^4 d_x \left(128 \sum_{i=0}^T \|H_i\|^2 + 64 \sum_{0 \le i < j \le T} \|H_{ij}\|^2 \right).$$

We now derive an explicit upper bound on B^2 in terms of system-dependent quantities $\rho_S(K)$ and $\rho_M(K)$. Recall the definitions:

$$\rho_S(K) := \|Q\|_F + \|K\|^2 \|R\|_F, \quad \rho_M(K) := \rho_A + \rho_B \|K\|,$$

where $\rho_A := \max\{|a_1|, |a_2|\}$ and $\rho_B := \max\{|b_1|, |b_2|\}$. These provide bounds on ||S|| and ||M||:

$$||S|| \le \rho_S(K), \quad ||M|| \le \rho_M(K).$$

Each H_i corresponds to the contribution of w_i to the quadratic cost:

$$H_i = M^{T-i\top} \cdots M^{\top} SM \cdots M^{T-i}$$

Hence, we can bound:

$$||H_i|| \le ||S|| \cdot ||M||^{2(T-i)} \le \rho_S(K) \cdot \rho_M(K)^{2(T-i)}.$$

The cross-term matrix H_{ij} arises from the interaction between w_i and w_j :

$$H_{ij} = M^{T-i\top} \cdots M^{\top} SM \cdots M^{T-j},$$

so that

$$||H_{ij}|| \le ||S|| \cdot ||M||^{T-i+T-j} = ||S|| \cdot ||M||^{2T-i-j}$$

 $\le \rho_S(K) \cdot \rho_M(K)^{2T-i-j}$.

Substituting into the expression:

$$B^{2} \leq \sigma_{w}^{4} d_{x} \cdot \rho_{S}(K)^{2} \cdot \left(128 \sum_{i=0}^{T} \rho_{M}(K)^{4(T-i)} + 64 \sum_{0 \leq i < j \leq T} \rho_{M}(K)^{4T-2(i+j)}\right).$$

Then we define

$$B_{\text{cost}}^{2} := \sup_{K \in \mathcal{K}} \sigma_{w}^{4} d_{x} \cdot \rho_{S}(K)^{2} \cdot \left(128 \sum_{i=0}^{T} \rho_{M}(K)^{4(T-i)} + 64 \sum_{0 \le i < j \le T} \rho_{M}(K)^{4T-2(i+j)} \right).$$

This gives a compact upper bound on B_{cost}^2 in terms of $\rho_S(K)$, $\rho_M(K)$, σ_w^2 , d_x , and T.

For the admissible range of λ , we need the maximum spectral norm across all H_i and H_{ij} . Since

$$\max_{i} ||H_i|| = ||H_0|| \le \rho_S(K) \cdot \rho_M(K)^{2T},$$

and

$$\max_{i < j} ||H_{ij}|| = ||H_{01}|| \le \rho_S(K) \cdot \rho_M(K)^{2T-1},$$

we get

$$|\lambda| \leq \inf_{K \in \mathcal{K}} \frac{1}{4\sigma_w^2 \cdot \rho_S(K) \cdot \rho_M(K)^{2T}}.$$

Hence, we have obtained Eq. (1) and the expression for B_{cost} Eq. (2) used in Thm 1, which completes the proof.

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