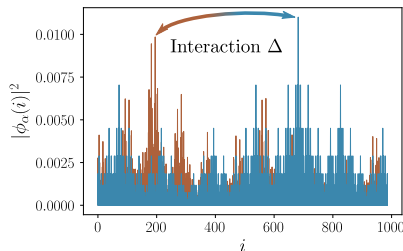


# Interacting electrons on a Fibonacci chain at high temperature

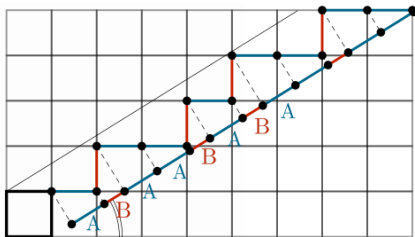
Nicolas Macé, Nicolas Laflorencie, Fabien Alet

Laboratoire de Physique Théorique, Université Paul Sabatier, Toulouse

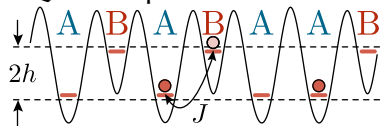


# QUANTUM PARTICLES IN A QUASIPERIODIC (QP) ENVIRONMENT

QP environment: Fibonacci chain

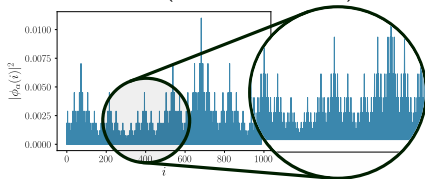


Quantum particles: fermions



$$H = \sum_{i=1}^L \left[ J(c_i^\dagger c_{i+1} + \text{h.c.}) - h_i n_i \right]$$

**Multifractal (scale invariant) states**



# MOTIVATION: QUASIPERIODICITY + INTERACTING ELECTRONS

## No interactions:

Periodic

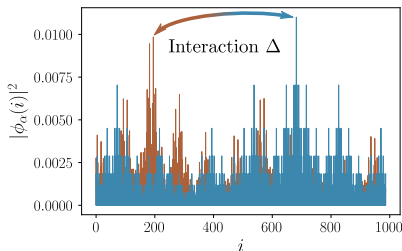
Quasiperiodic

Random

Complexity

- Extended states
- Fast transport
- Multifractal states
- Anomalous transport
- Localized states
- No transport

Quasiperiodicity (QP) + **interactions** between particles?



Naively: delocalisation, fast transport

## Results:

- weak QP: delocalisation, fast transport
- strong QP: **many-body localisation**, no transport

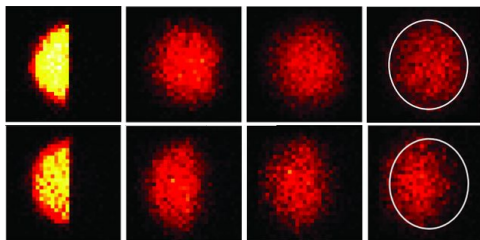
# OUTLINE

- 1 Many-body localisation
- 2 Free Fibonacci chain at high energy
- 3 Interacting Fibonacci chain

# MANY-BODY LOCALISATION

**Isolated** quantum system, **strong interactions**, disorder or quasiperiodicity

- 1 Usual: ergodic dynamics, transport, **eigenstate thermalisation hypothesis (ETH)**,
- 2 Unusual: non-ergodicity, no transport, **many-body localisation (MBL)**.



[Choi *et al* 16]

**Experiments:** cold ions/atoms [Schreiber *et al* 15; Smith *et al* 15; Bordia *et al* 17].

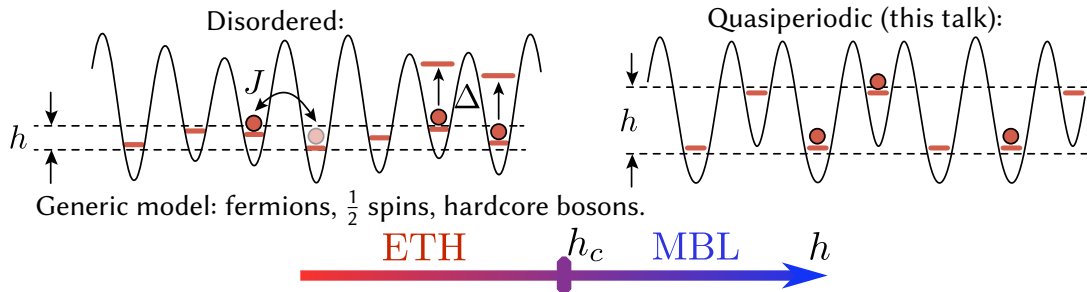
**Motivations:**

- ETH/MBL phase transition,
- MBL in more than 1D,
- Ingredients for MBL (this talk).

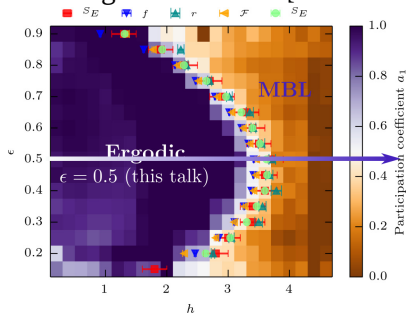
# A MODEL FOR MBL

Chain of interacting spinless fermions (nb: no phonons):

$$H = \sum_{i=1}^L \left[ J(c_i^\dagger c_{i+1} + \text{h.c.}) + \Delta n_i n_{i+1} - h_i n_i \right]$$

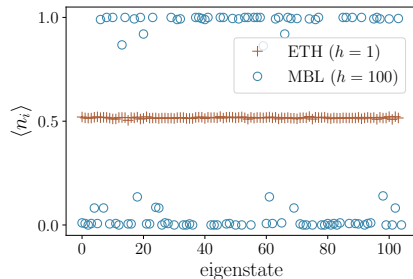


## MBL PHENOMENOLOGY

Phase diagram at  $\Delta = 1$  [Luitz *et al* 15]

ETH:

- Transport, thermal observables
- High entanglement
- Non-integrability

Fermion density at  $\epsilon = 0.5$ 

MBL:

- No transport, non-thermal observables
- Low entanglement
- Emergent integrability

# INGREDIENTS FOR MBL

Usually:

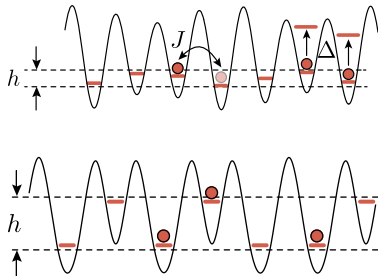
- 1  $\Delta = 0$ : **localized** (random, Aubry-André potential)
- 2  $\Delta \neq 0$ : localization persists

This talk

- 1  $\Delta = 0$ : **multifractal** (quasiperiodic potential)
- 2  $\Delta \neq 0$ : localization appears

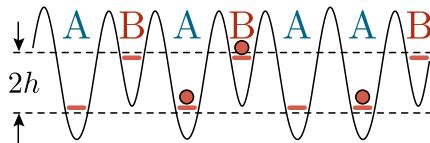
## Interest

- MBL is **generic**
- Interplay between quasiperiodicity and MBL





# INTERACTING FERMIONS ON THE FIBONACCI CHAIN



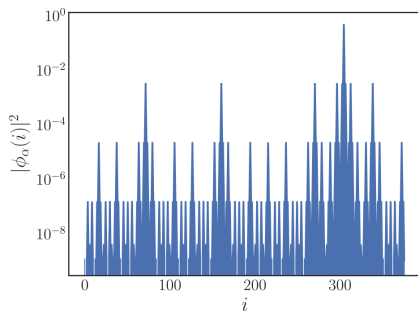
$$H = \sum_{i=1}^L \left[ J(c_i^\dagger c_{i+1} + \text{H.C.}) + \Delta n_i n_{i+1} - h_i n_i \right]$$

Method: numerical **exact diagonalization**

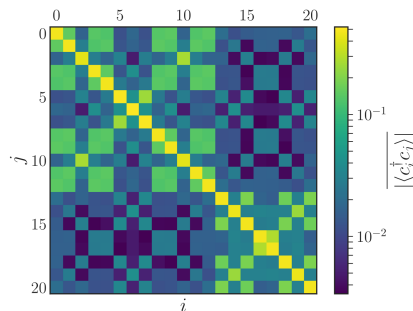
- High energy + non-integrable: **no analytical methods**
- $L/2$  fermions on  $L$  sites:  $\# \text{states} \sim 2^L / \sqrt{L} \rightarrow$  **memory is limiting**  
State-of-the art:  $L = 24$  [Pietracaprina *et al* 18]
- Fibonacci: **few samples**:  $L/2$  non-equivalent systems of size  $L$ .

# FREE FERMIONS PROPERTIES

- Multifractal single particle wavefunctions [Ostlund; Kalugin; Kohmoto; ...]
- Anomalous transport [Mayou; Schreiber; Varma & Žnidarič; ...]



Single particle wavefunction at the Fermi level

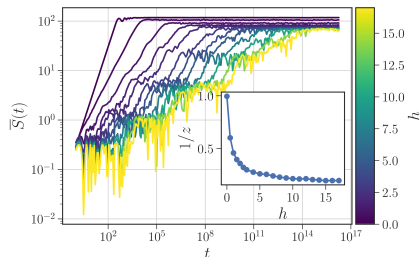


Correlations of highly excited states [Macé *et al* 19]

# FREE FERMIONS ENTANGLEMENT

Entanglement entropy  $S(\psi)$ : a many-body **locality** probe

- $S(\psi) = \#\{\text{bits of information recoverable by local measurements}\}$
- $S(\psi)$  large: extended (entangled) state,  $S(\psi)$  small: localized state.



Entanglement growth starting from localized fermions [Macé *et al* 19]

Fibonacci fermions: **anomalous** growth

$$S(t) \sim t^{\frac{1}{z}}, \quad z > 1$$

Compare with:

- Periodic system:  $z = 1$  (ballistic growth),
- Disordered system:  $z \rightarrow \infty$  (no growth).

## Conclusion

Anomalous, intermediate prop. even at high energy.

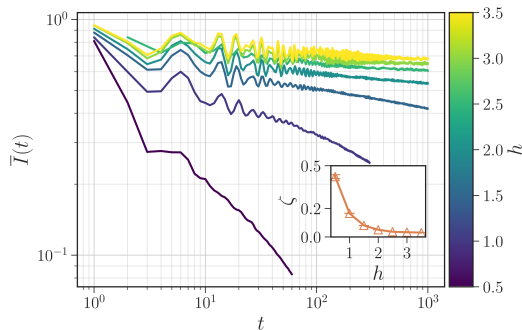
# THE ETH/MBL TRANSITION: 1) DYNAMICS

Imbalance experiment:

- Initially: fermions on even sites  
 $|\psi(t=0)\rangle = |0101\dots\rangle$
- Imbalance: **distance** from initial state  
 $I(t) = \frac{2}{L}(N_e(t) - N_o(t))$

Properties:

- **ETH: power-law decay**  $I(t) \sim t^{-\zeta}$
- **MBL: saturation**  $I(t \rightarrow \infty) = \text{cst} > 0$   
→ **memory** of the initial state  
[Luitz *et al* 15]



$\zeta(h \geq 2.5) = 0 \rightarrow$  **MBL phase transition.**

## THE ETH/MBL TRANSITION: 2) ENTANGLEMENT

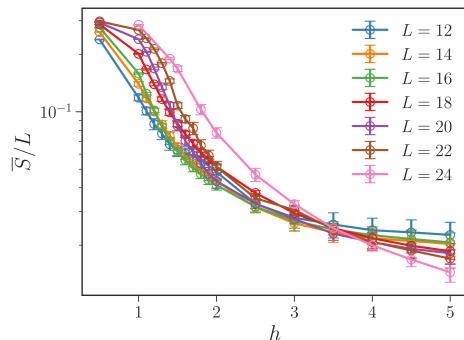
Entanglement entropy:

- **ETH:** coincides with thermodynamic entropy: **extensive**

$$\bar{S}_{\text{ETH}} \simeq L$$

- **MBL:** **sub-extensive**

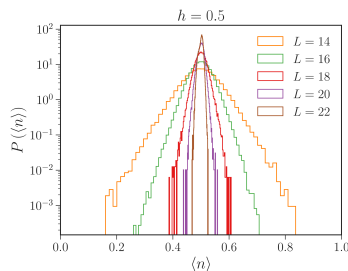
$$\bar{S}_{\text{MBL}}/L \rightarrow 0$$



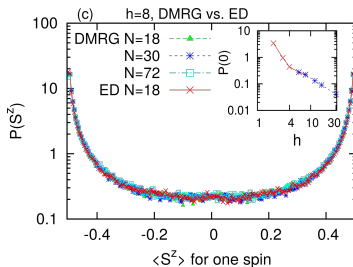
Compatible with **ETH/MBL transition**,  
 $h^* \simeq 3.5$ .

# THE ETH/MBL TRANSITION: 3) LOCAL OBSERVABLES

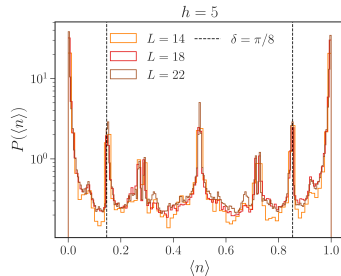
Expect:  $\langle n_i \rangle_{\text{ETH}} = \frac{1}{2}$ ,  $\langle n_i \rangle_{\text{MBL}} \simeq 0 \text{ or } 1$ .



ETH Fibonacci



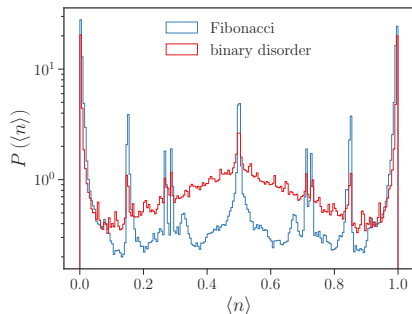
MBL random [Lim, Sheng 15]



MBL Fibonacci

Fibonacci MBL: **extra structure**  $\rightarrow$  link with QP geometry?

# FIBONACCI MBL: LOCAL ENTANGLEMENT



Scrambled potentials  $\rightarrow$  peak suppression

QP induces **local “cat states”**

$$|\psi\rangle \propto \alpha |01\rangle + \beta |10\rangle$$

Peak ingredients:

- **Binary** modulation: A/B letters
- **Correlated** modulation (quasiperiodic)

# CONCLUSION

## Non-interacting Fibonacci fermions

- Geometry: intermediate complexity
- **Multifractality** even at high energy
- Intermediate, **anomalous** transport, even at the many-body level

## Interacting Fibonacci fermions

- Weak quasiperiodicity: **thermal phase (ETH)**
- Strong quasiperiodicity: **localized phase (MBL)**
- MBL bears sign of **quasiperiodicity** (locally entangled states)

→ anomalous transport in the ETH phase: why?

→ vicinity of the non-interacting point?

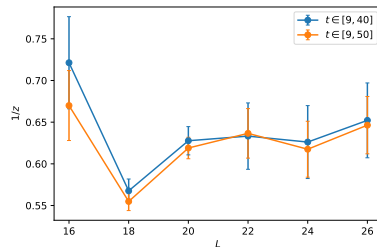
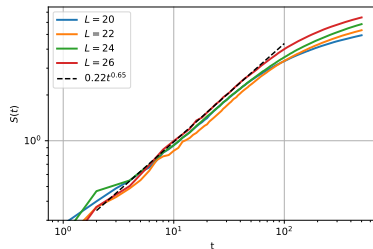
[*Macé, Laflorencie, Alet, SciPost Phys. 6, 050 (2019)*]



## DYNAMICS: ENTANGLEMENT

**ETH phase:** expect  $S(t) \propto t$

Fibonacci: **anomalous**  $S(t) \propto t^{1/z}$ ,  $z > 1$ .



Usual explanation: **rare regions** [Vosk; Potter 15]

Fibonacci: no rare regions ...

Finite size effect? [Setiawan *et al* 17], initial state fluctuations? [Lüschen *et al* 17]

# SPECTRAL PROPERTIES

Gap ratios [Oganesian, Huse]

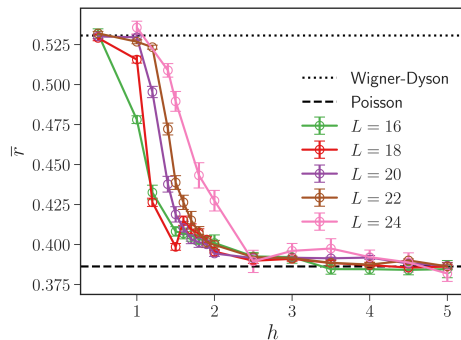
$$r_n = \min \left( \frac{g_{n+1}}{g_n}, \frac{g_n}{g_{n+1}} \right)$$

- **ETH: random matrix-like spectrum**

$$\bar{r}_{\text{ETH}} \simeq 0.53$$

- **MBL: independant levels**

$$\bar{r}_{\text{MBL}} \simeq 0.39$$



Compatible with **ETH/MBL transition**,  
 $h^* \simeq 2.5$ .